

# Hyperon-Nuclei/Nucleon Scattering at BESIII

Han Miao (妙晗) and Jianyu Zhang (张剑宇)

Institute of High Energy Physics  
University of Chinese Academy of Sciences

2024. 1. 16

arxiv: 2209.12601 [hep-ex], Phys.Rev.Lett. 130 (2023) 25, 251902, arXiv: 2310.00720 [nucl-ex]

# CONTENT

1

**Why to measure**

2

**How to measure**

3

**What have been measured**

4

**What will be studied**





# Why to measure



# From Hyperon to Neutron Stars



## Hyperons in neutron stars

Phys.Lett.B 747 (2015) 43-47

Tetsuya Katayama\*, Koichi Saito<sup>1</sup>

Department of Physics, Faculty of Science and Technology, Tokyo University of Science, Noda 278-8510, Japan

## The Hyperon Puzzle in Neutron Stars

Ignazio BOMBACI<sup>1,2,3</sup>

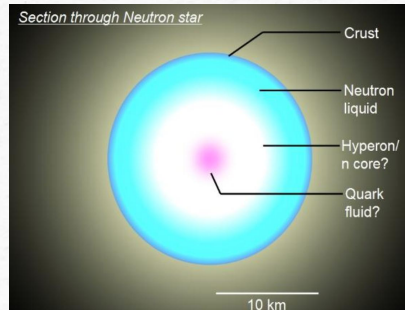
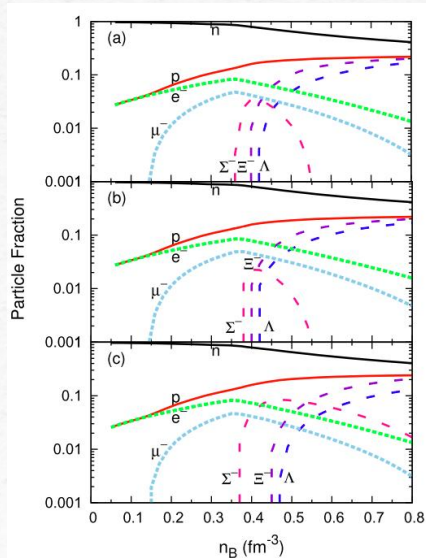
Nucl.Phys.News 31 (2021) 3, 17-21

Neutron stars need hyperon!  
(Pauli exclusion principle)

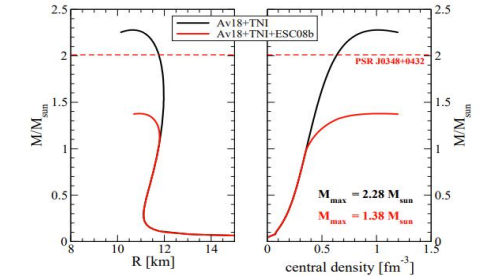
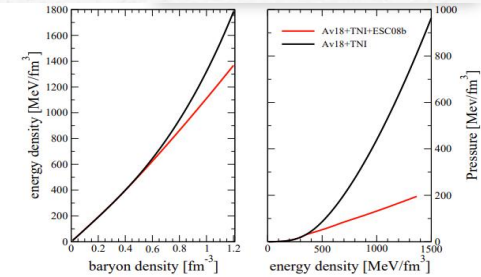
contradiction



The existence of hyperons will reduce the neutron star density.  
Inconsistent with observation!



Hyperon Puzzle





# Hyperon-nucleon Interaction

## The Hyperon Puzzle in Neutron Stars

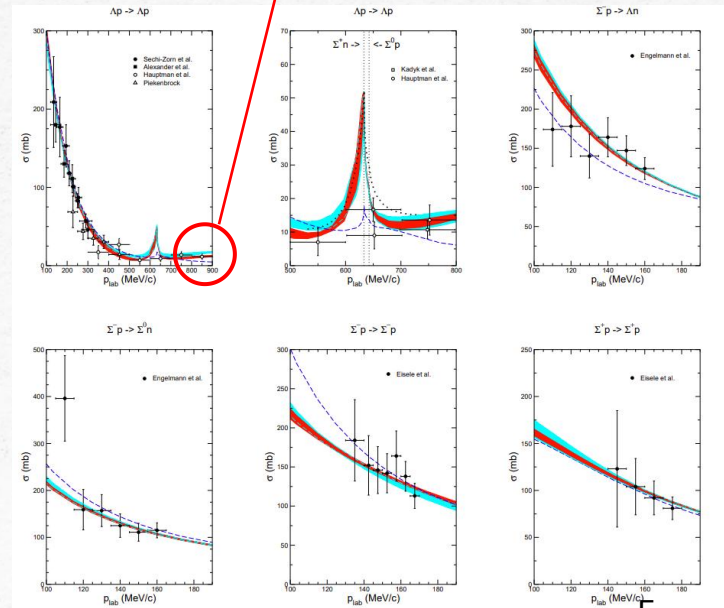
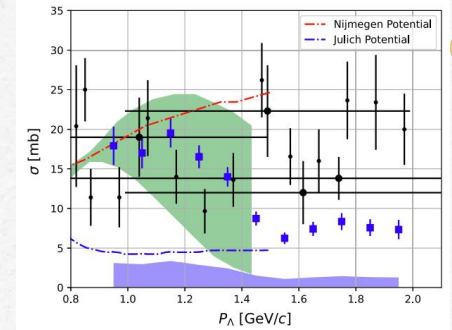
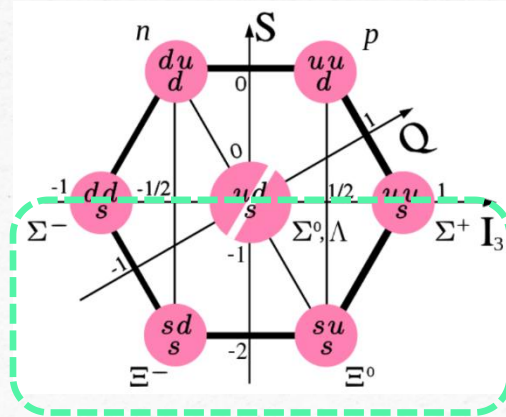
Ignazio BOMBACI<sup>1,2,3</sup>

Nucl.Phys.News 31 (2021) 3, 17-21

Clearly, one should try to trace back the origin of this problem to the underlying YN and YY two-body interactions or to conceivable hyperonic three-body interactions (YTBI) of the type YNN, YYN and YYY. Unfortunately, these two- and three-body strangeness  $S \neq 0$  baryonic interactions are rather uncertain and poorly known. Basically this is due to the scarce amount of experimental data and to the considerable difficulties in their theoretical analysis. This situation is in sharp contrast to the case of the NN interaction, which is satisfactorily well known mostly due to the large number of NN scattering data and to the huge amount of measured properties of stable and unstable nuclei. The study

The vast majority of experiments came from fixed-target experiments before the 1980s.

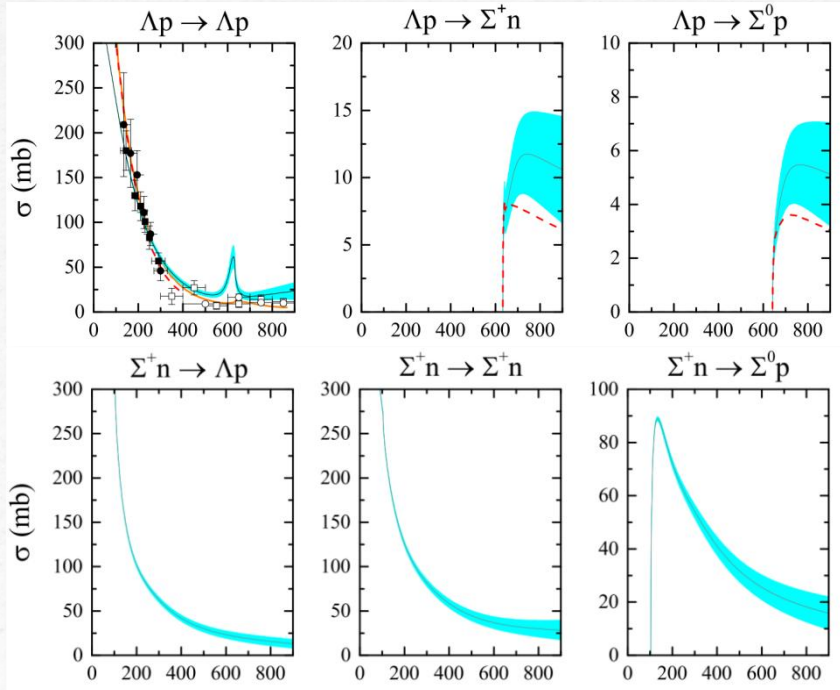
More studies of hyperon-hyperon (Y-Y) and hyperon-nucleon (Y-N) interactions are needed.



# Theoretical Work



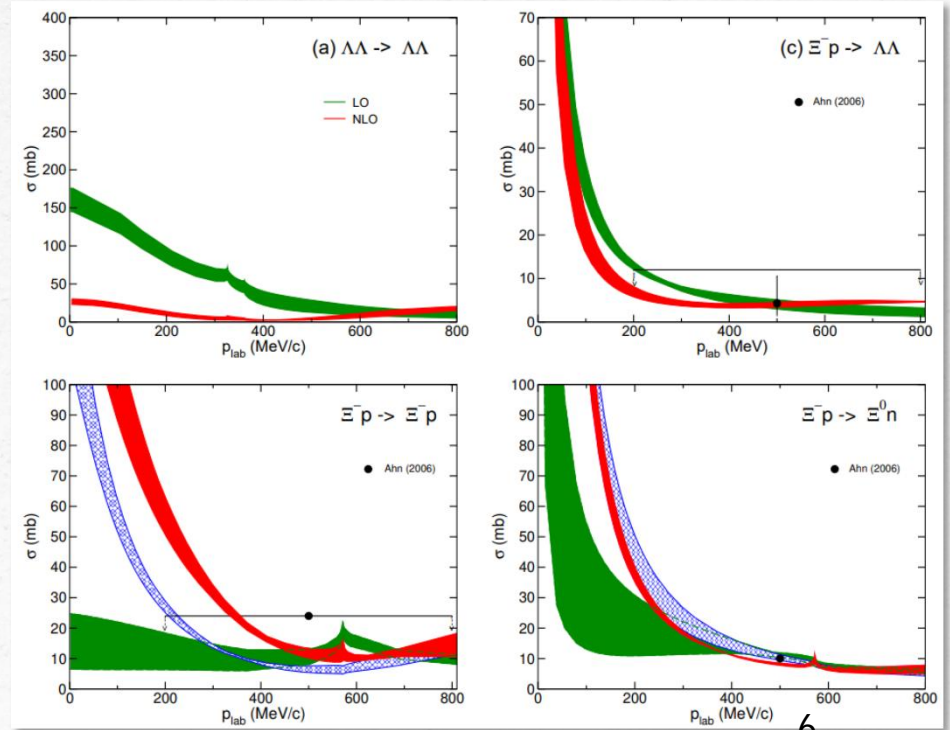
Phys. Rev. C 105, 035203 (2022)



LO : H. Polinder, J.H., U.-G. Meißner, PLB 653 (2007) 29

NLO16: J.H., U.-G. Meißner, S. Petschauer, NPA 954 (2016) 273

NLO19: J.H., U.-G. Meißner, EPJA 55 (2019) 23

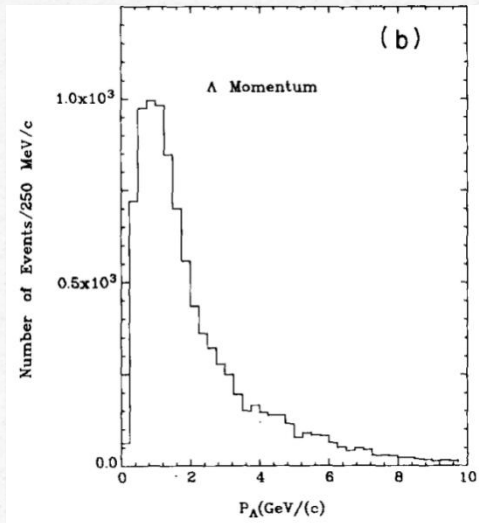




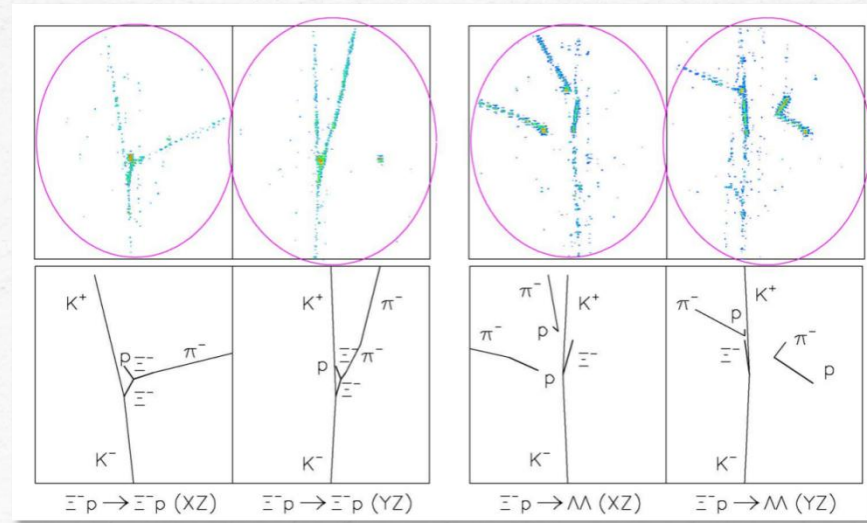
# Early Measurement



- Hyperons are obtained by bombarding hydrogen bubble chamber or scintillating fiber target with  $K^-$
- **Low statistics and high background**



Nucl. Phys. B 125, 29 (1977)



Phys. Lett. B 633, 214 (2006)



# Early Measurements

Phys. Lett. B 38, 123 (1972)

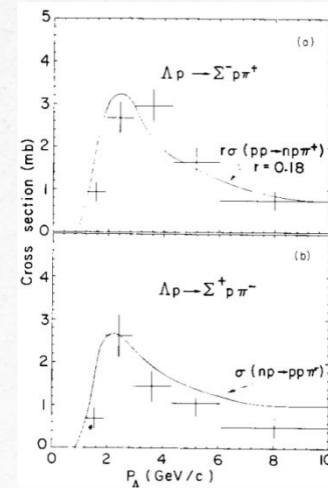
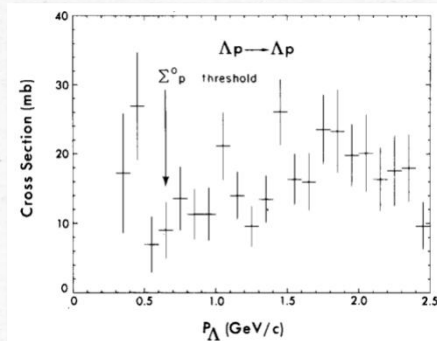
Reaction	Number of events	
$\Lambda p \rightarrow \Lambda p$ (elastic)	584	(1)
$\Lambda p \rightarrow \Sigma^- p \pi^+$	132	(2)
$\Lambda p \rightarrow \Sigma^+ p \pi^-$	60	(3)
$\Lambda p \rightarrow \Lambda p \pi^+ \pi^-$	181	(4)
$\Lambda p \rightarrow \Sigma^0 p$	35	(5)
various $\Xi^0 p$ interactions	25	

reaction	events *	signature	cross-section events **	cross-section (mb)
$\Xi^0 + p \rightarrow \Xi^0 + p$	2	K, $\Lambda$	1	8
$\Xi^0 + p \rightarrow \Lambda + \Sigma^+$	6	$\Lambda$	4	24
$\Xi^0 + p \rightarrow \Sigma^0 + \Sigma^+$	1	$\Lambda$	1	6
$\Xi^0 + p \rightarrow \pi^+ + \Lambda + \Lambda$	1	K, $\Lambda$	1	6
$\Xi^0 + p \rightarrow \pi^0 + \Lambda + \Sigma^+$	1	$\Lambda$	1	6
$\Xi^0 + p \rightarrow \pi^+ + \Xi^- + p$	1	K or $\Lambda$	1	5
$\Xi^0 + p \rightarrow \pi^+ + \pi^+ + \Xi^- + n$	1	K, $\Lambda$	1	6
$\Xi^0 + p \rightarrow p \Xi^- + p$	2	$\Lambda$	2	8
$\Xi^0 + p \rightarrow \Sigma^- + \Sigma^+$	1	K	1	4
$\Xi^0 + p \rightarrow \Sigma^- + K^0 + p$	1	K	1	4

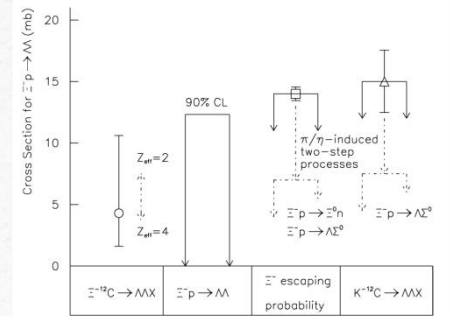
Phys. Lett. B 32, 720 (1970)

Reaction	Momentum interval (GeV/c)	Number of events	$\sigma$ (mb)
$\Lambda p \rightarrow$ all	0.5 - 1.0	25.8 ± 6.2	25.8 ± 6.2
	1.0 - 1.5	31.3 ± 6.5	31.3 ± 6.5
	1.5 - 2.0	42.8 ± 7.1	42.8 ± 7.1
	2.0 - 2.5	37.5 ± 7.2	37.5 ± 7.2
	2.5 - 3.0	34.1 ± 8.3	34.1 ± 8.3
	3.0 - 4.0	41.8 ± 10.0	41.8 ± 10.0
$\Lambda p \rightarrow \Lambda p$	0.5 - 1.0	20	22.2 ± 5.0
	1.0 - 1.5	21	12.9 ± 2.8
	1.5 - 2.0	37	22.0 ± 3.6
	2.0 - 2.5	28	16.1 ± 3.1
	2.5 - 3.0	12	11.0 ± 3.2
	3.0 - 4.0	13	12.5 ± 3.4
$\Lambda p \rightarrow \Sigma^0$	0.66 - 4.0	11	1.5 ± 0.5
	$\Lambda p \rightarrow \Lambda p \pi^0$	29	4.1 ± 0.8
	$\Lambda p \rightarrow \Lambda p \pi^+ \pi^-$	12	1.9 ± 0.6
$\Sigma^+ p \rightarrow \Sigma^+ p$	0.5 - 1.5	10	31.2 ± 10.1
	1.5 - 2.5	8	18.7 ± 6.6
	2.5 - 4.0	4	15.3 ± 7.8
$\Sigma^- p \rightarrow \Sigma^- p$	0.5 - 1.5	6	13.2 ± 4.7
	1.5 - 2.5	11	13.9 ± 4.1
	2.5 - 4.0	4	7.5 ± 3.8
$\Xi^- p \rightarrow \Xi^- p$	1.0 - 4.0	6	13 ± 6
	$\Xi^0 p \rightarrow \Xi^0 p$	4	19 ± 10

Nucl. Phys. B 125, 29 (1977)



Phys. Lett. B 633, 214 (2006)





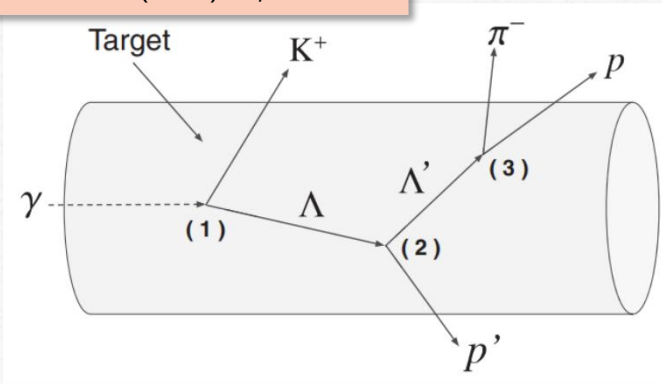


PHYSICAL REVIEW LETTERS **127**, 272303 (2021)

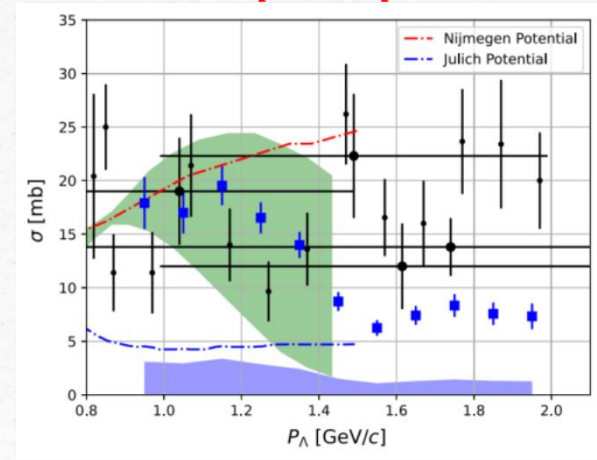
(CLAS Collaboration)

**Improved  $\Lambda p$  Elastic Scattering Cross Sections between 0.9 and 2.0 GeV/c  
as a Main Ingredient of the Neutron Star Equation of State**

Phys.Rev.Lett. 127 (2021) 27, 272303

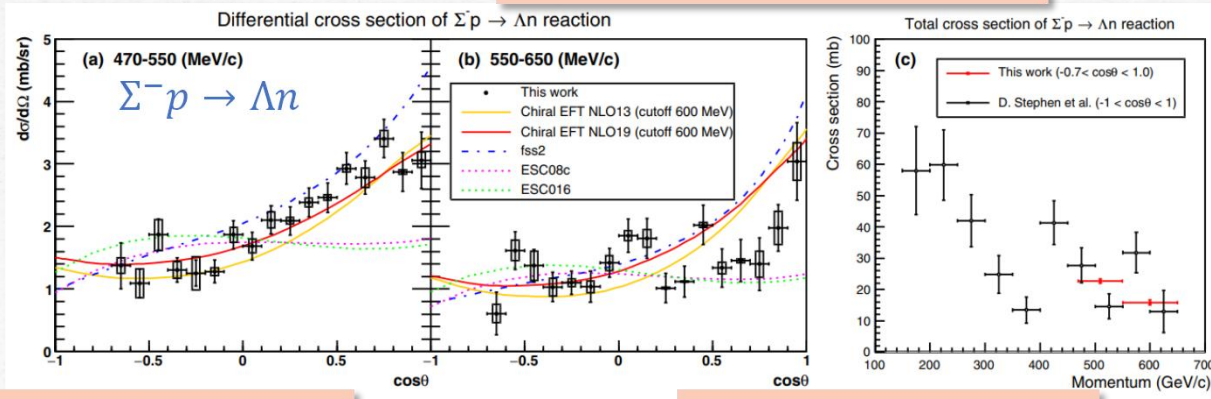
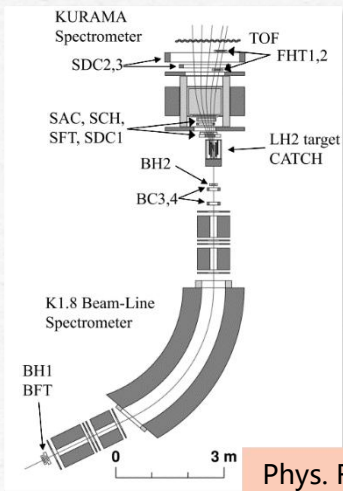


$\Lambda p \rightarrow \Lambda p$



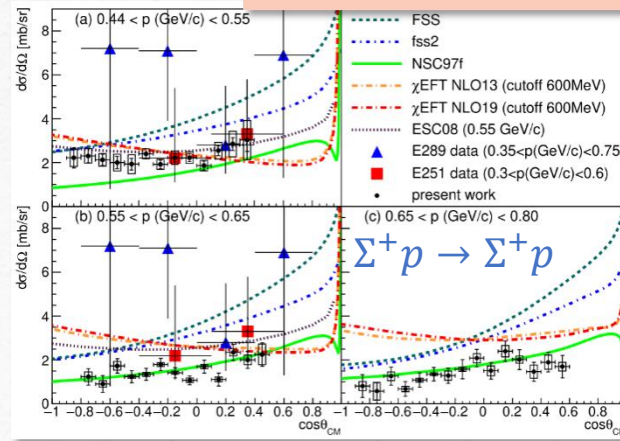
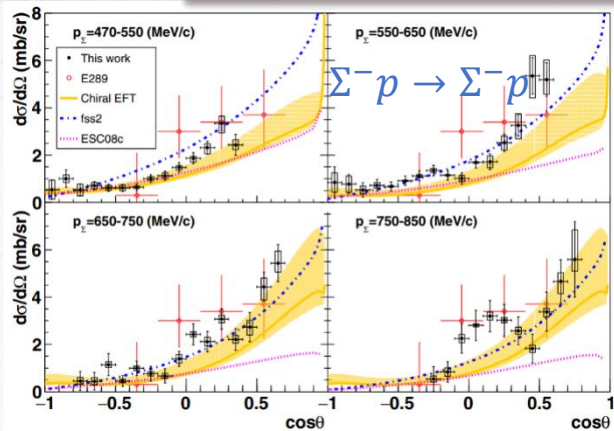
This is the first data on this reaction since the 1970s.

Phys. Rev. Lett. 128, 072501 (2022)



Phys. Rev. C 104, 045204 (2021)

PTEP 2022, 093D01 (2022)





# How to measure

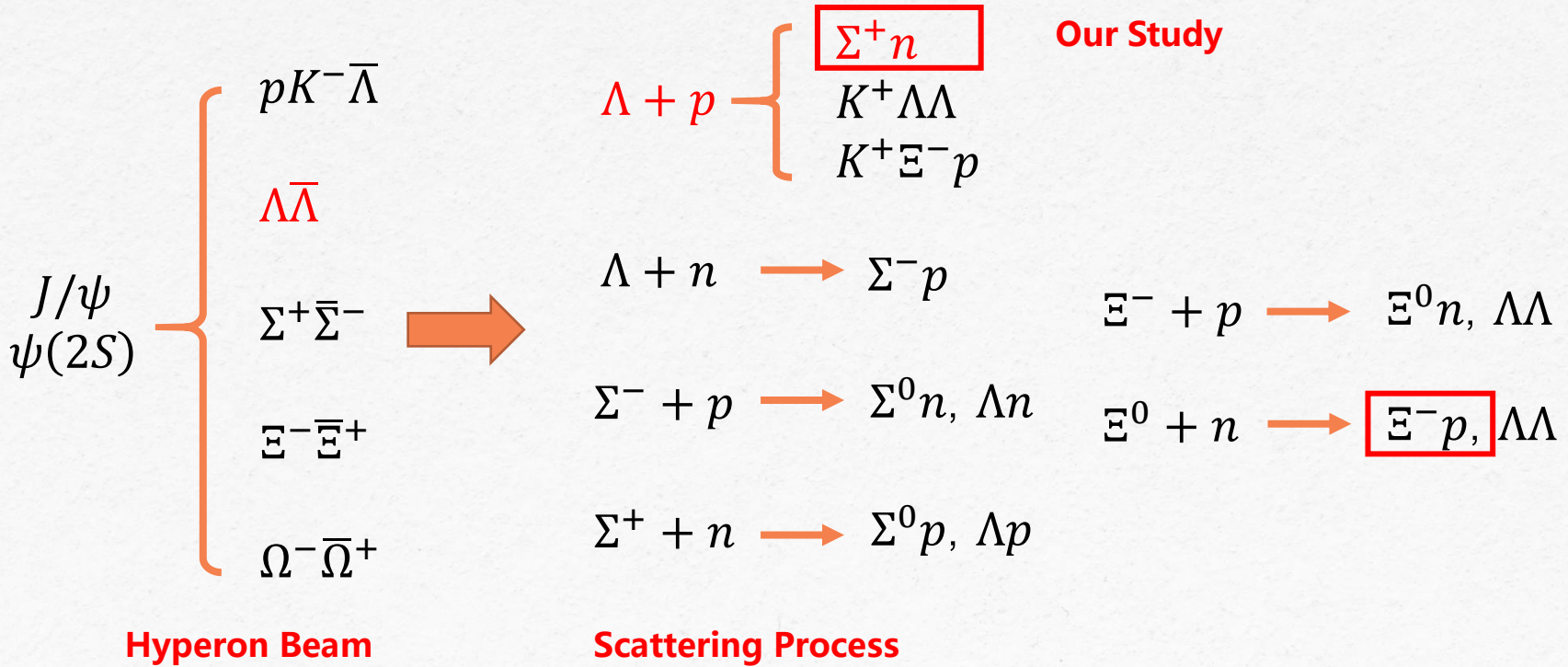








# Potential Channels





# Double-Tag Method



Double Tag Events:

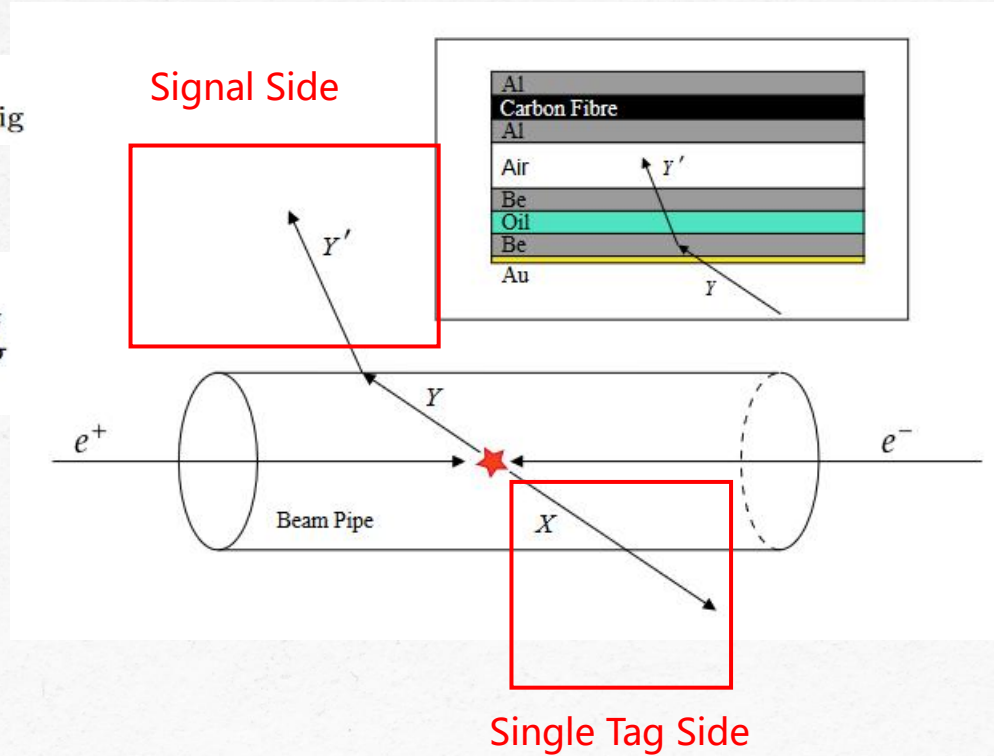
$$N_{\text{DT}} = \mathcal{L}_Y \cdot \sigma(YA \rightarrow Y'A') \cdot \mathcal{B}(Y') \cdot \epsilon_{\text{sig}}$$

Effective Luminosity of  $\Lambda$  beam:

$$\mathcal{L}_Y = N_{\text{ST}} \cdot \frac{N_A}{N_{\text{ST}}^{\text{MC}}} \cdot \sum_j^7 \sum_i^{N_{\text{ST}}^{\text{MC}}} \frac{\rho_T^j \cdot l^{ij}}{M^j} \cdot \mathcal{R}_\sigma^j$$

Cross section:

$$\sigma(YA \rightarrow Y'A') = \frac{N_{\text{DT}}}{\epsilon_{\text{sig}} \cdot \mathcal{L}_Y} \cdot \frac{1}{\mathcal{B}(Y')}$$





# Effective Luminosity

$$\mathcal{L}_Y = N_{ST} \cdot \frac{N_A}{N_{ST}^{MC}} \cdot \sum_j^7 \sum_i^{N_{ST}^{MC}} \frac{\rho_T^j \cdot l^{ij}}{M^j} \cdot \mathcal{R}_\sigma^j$$

$\rho_T^j$  - density of the  $j_{th}$  layer

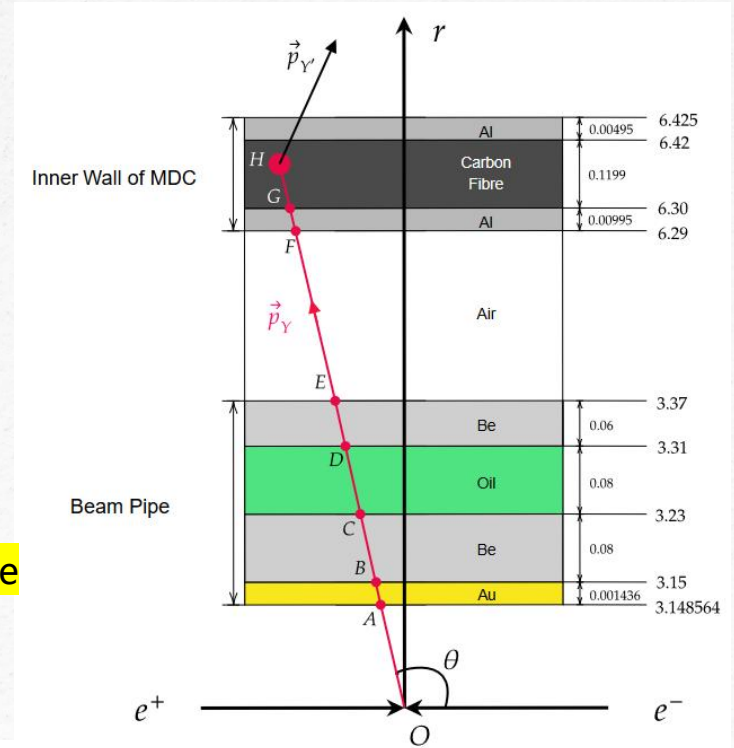
$M^j$  - molar mass of the  $j_{th}$  layer

$l^{ij}$  - path length of the  $i_{th}$  event in the  $j_{th}$  layer (will be 0 if the incident hyperon does not reach the  $j_{th}$  layer)

$\mathcal{R}_\sigma^j$  - the ratio of the cross sections between layers

- proportional to the number of nucleons in the nuclei surface
- proportional to the number of nucleons in the nuclei
- Eikonal Approximation

$$N_{\text{eff}}(Z_{\text{eff}}) = \frac{N(Z)}{A} \int \rho(\mathbf{r}) \exp \left\{ -\bar{\sigma}_i \int_{-\infty}^z \rho(x, y, z') dz' - \bar{\sigma}_f \int_z^{\infty} \rho(x, y, z') dz' \right\} d^3\mathbf{r}$$





What have been measured



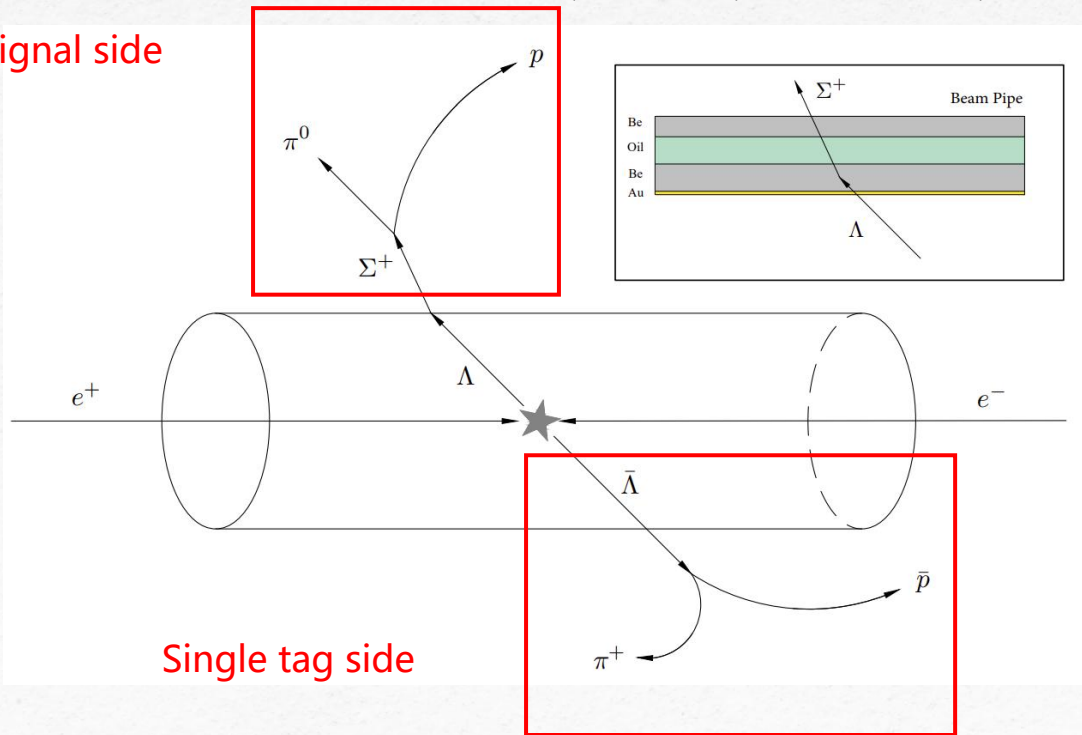
$$\Lambda N \rightarrow \Sigma^+ X$$



## Reaction chain :

$$J/\psi \rightarrow \Lambda \bar{\Lambda}, \bar{\Lambda} \rightarrow \bar{p} \pi^+, \Lambda + N(\text{nucleus}) \rightarrow \Sigma^+ + X(\text{anything}), \Sigma^+ \rightarrow p \pi^0, \pi^0 \rightarrow \gamma \gamma.$$

Signal side



Single tag side

Two-body decay,  $P_\Lambda \approx 1.074 \text{ GeV}/c$ ,

Very small horizontal crossing angle  
of 11 mrad for  $e^+$  and  $e^-$  beams

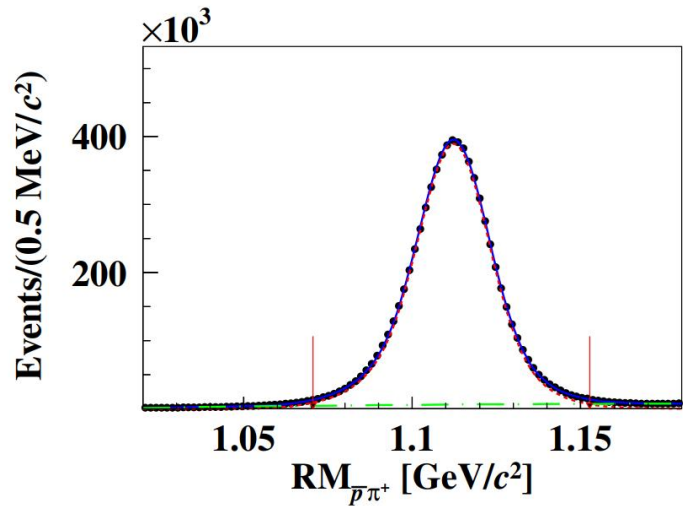


Very small range of  $0.017 \text{ GeV}/c$   
above and below  $1.074 \text{ GeV}/c$  for  $P_\Lambda$ .

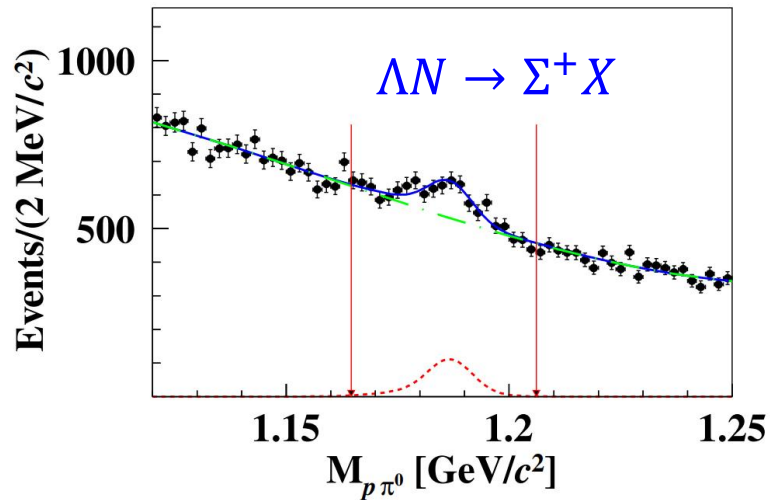




# $\Lambda N \rightarrow \Sigma^+ X$



$$N_{ST} = 7207565 \pm 3741$$



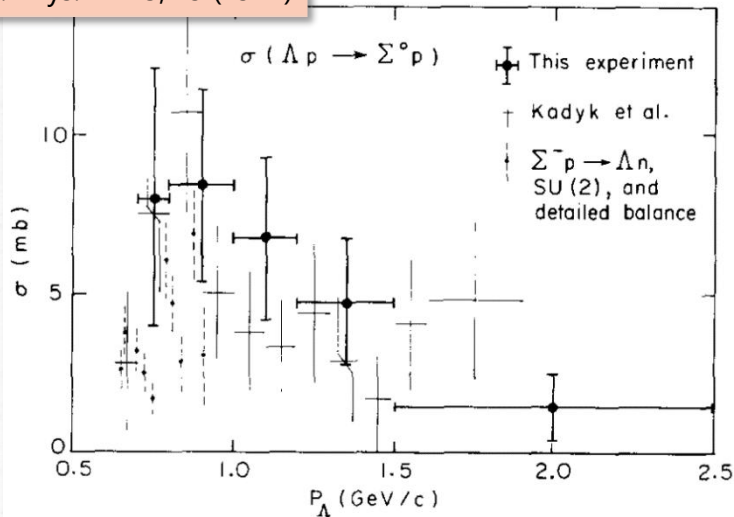
$$N_{DT} = 795 \pm 101$$

Parameter	Value
$N_{DT}$	$795 \pm 101$
$\epsilon_{sig}$	24.32%
$\mathcal{L}_\Lambda$	$(17.00 \pm 0.01) \times 10^{28} \text{ cm}^{-2}$
$\mathcal{B}(\Sigma^+ \rightarrow p\pi^0)$	$(51.57 \pm 0.30)\%$

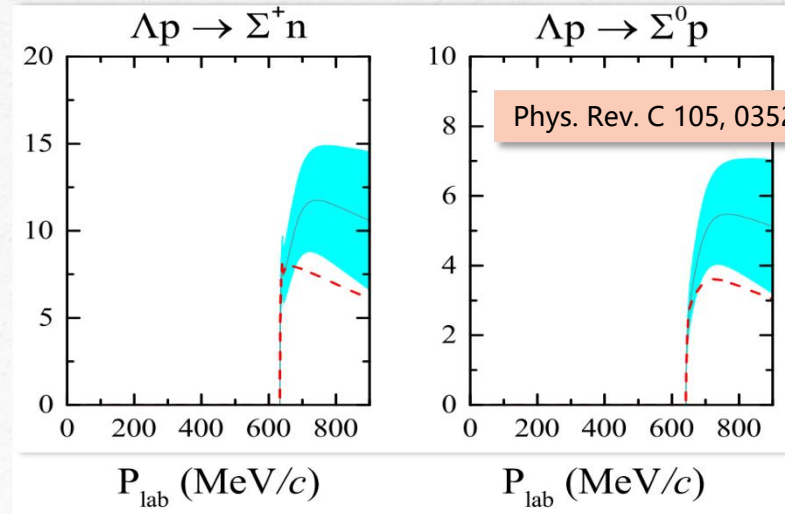
# $\Lambda N \rightarrow \Sigma^+ X$


- $\sigma(\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X) = (37.3 \pm 4.7_{\text{stat}} \pm 3.5_{\text{sys}})$  mb at  $P_\Lambda \approx 1.074$  GeV/c. The first attempt to investigate  $\Lambda$ -nucleus interaction at an  $e^+e^-$  collider.
- Taking the effective number of reaction protons in  ${}^9\text{Be}$  nucleus as 1.93, the cross section of  $\Lambda p \rightarrow \Sigma^+ X$  for single proton is  $\sigma(\Lambda p \rightarrow \Sigma^+ X) = (19.3 \pm 2.4_{\text{stat}} \pm 1.8_{\text{sys}})$  mb.

Nucl. Phys. B 125, 29 (1977)



$\sigma(\Lambda p \rightarrow \Sigma^+ n)$  is twice of  $\sigma(\Lambda p \rightarrow \Sigma^0 p)$



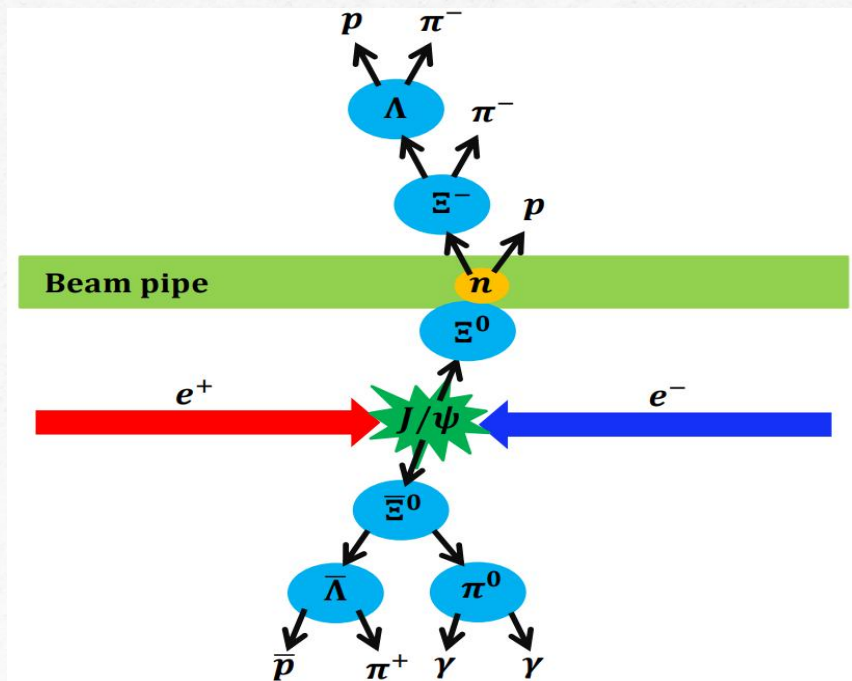


$$\Xi^0 n \rightarrow \Xi^- p$$



Reaction chain :

$$J/\psi \rightarrow \Xi^0 \bar{\Xi}^0, \bar{\Xi}^0 \rightarrow \bar{\Lambda} \pi^0, \bar{\Lambda} \rightarrow \bar{p} \pi^+, \pi^0 \rightarrow \gamma \gamma, \Xi^0 n \rightarrow \Xi^- p, \Xi^- \rightarrow \Lambda \pi^-, \Lambda \rightarrow p \pi^-.$$

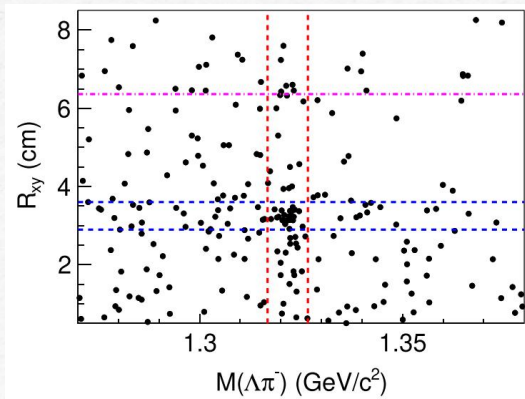


10 billion  $J/\psi$  data

Two-body decay,  $P_{\Xi^0} \approx 0.818 \text{ GeV}/c$



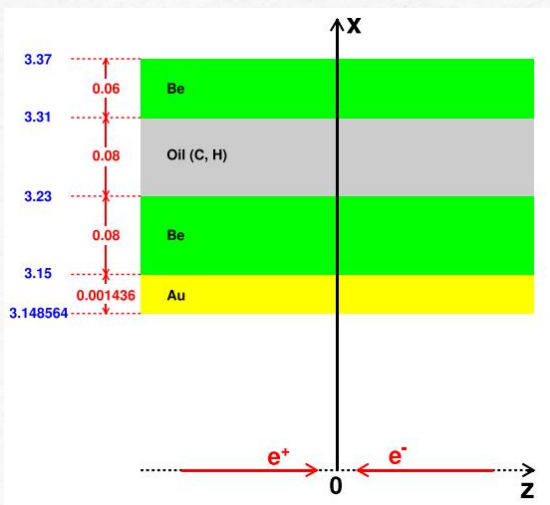
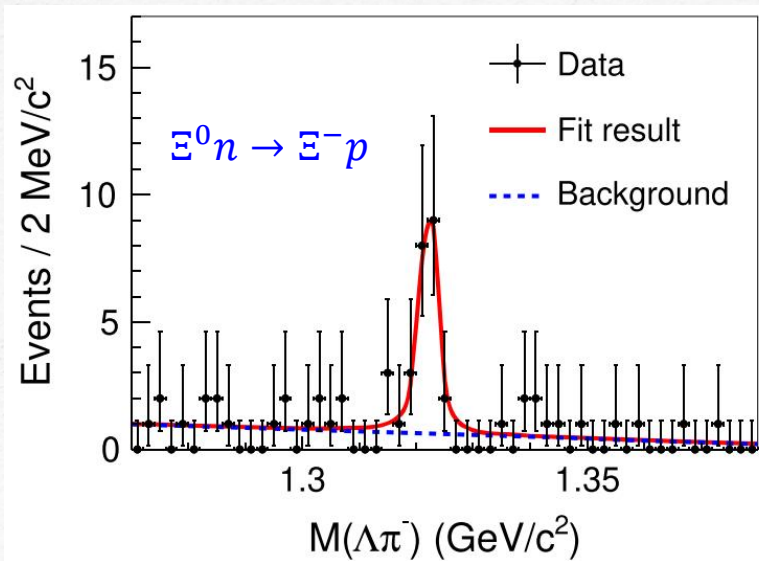
$\Xi^0 n \rightarrow \Xi^- p$



Inner wall of MDC

Beam pipe

$R_{xy}$  is distance from reconstructed  $\Xi^- p$  vertex to  $z$  axis



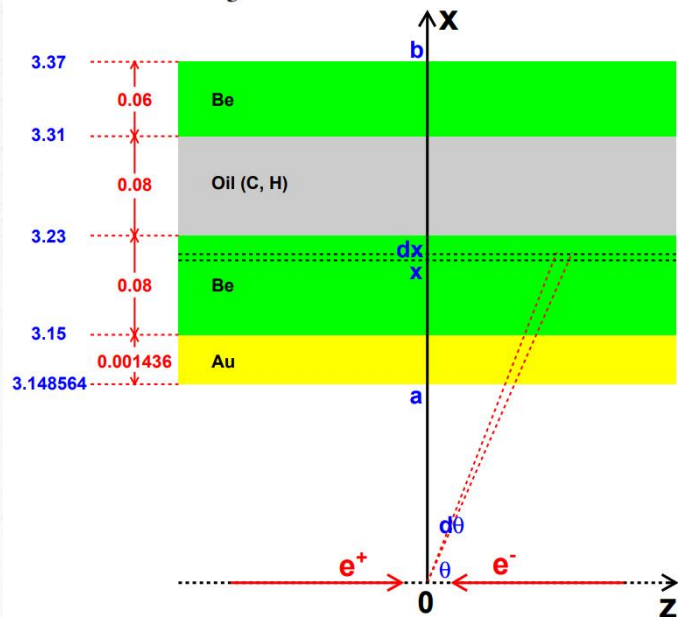
$$N = 22.9 \pm 5.5$$

$$S = 7.1\sigma$$

$$\Xi^0 n \rightarrow \Xi^- p$$

$$\sigma(\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}) = \frac{N^{\text{sig}}}{\epsilon \mathcal{B} \mathcal{L}_{\text{eff}}}$$

$$\mathcal{L}_{\text{eff}} = \frac{N_{J/\psi} \mathcal{B}_{J/\psi}}{2 + \frac{2}{3}\alpha} \int_a^b \int_0^\pi (1 + \alpha \cos^2 \theta) e^{-\frac{x}{\sin \theta \beta \gamma L}} N(x) C(x) d\theta dx$$



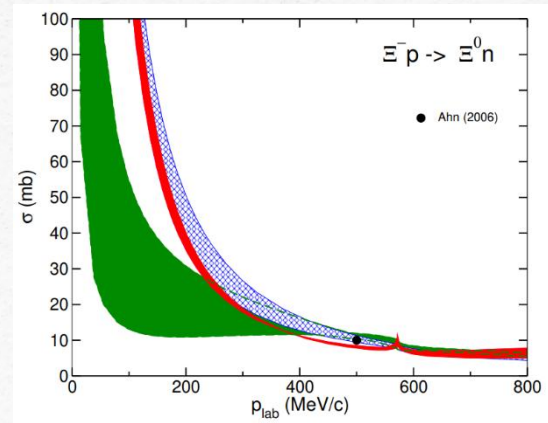
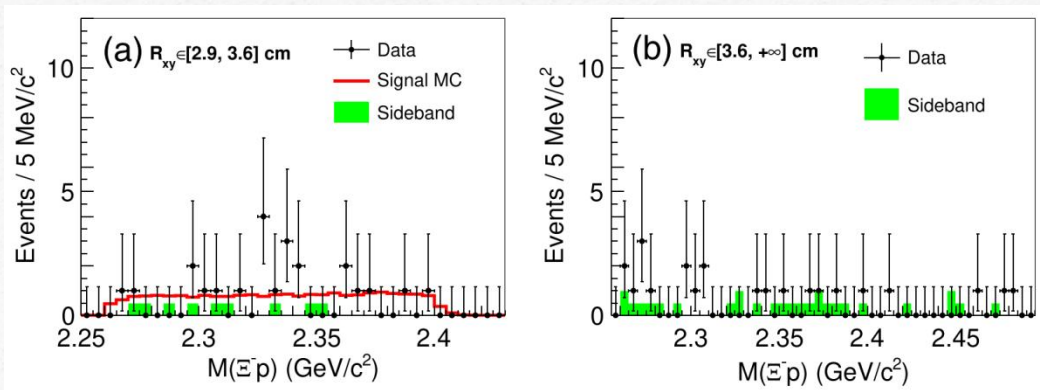
pure surface  
process  
assumption  
(proportional to  
number of  
neutrons)

Parameter	Result
$N^{\text{sig}}$	$22.9 \pm 5.5$
$\epsilon$	1.873%
$\mathcal{B}$	$(40.114 \pm 0.444)\%$ [53]
$N_{J/\psi}$	$(1.0087 \pm 0.0044) \times 10^{10}$ [46]
$\mathcal{B}_{J/\psi}$	$(0.117 \pm 0.004)\%$ [53]
$\alpha$	$0.514 \pm 0.016$ [56]
$L$	$(8.69 \pm 0.27)$ cm [53]
$E_{\text{beam}}$	1.5485 GeV
$m_{\Xi^0}/c^2$	$(1.31486 \pm 0.00020)$ GeV/ $c^2$ [53]
$a$	3.148564 cm [45]
$b$	3.37 cm [45]
$N(x)$	$\begin{cases} 5.91 \times 10^{22} \text{ cm}^{-3}, & 3.148564 \leq x \leq 3.15 \text{ cm} \\ 1.24 \times 10^{23} \text{ cm}^{-3}, & 3.15 < x \leq 3.23 \text{ cm} \\ 3.45 \times 10^{22} \text{ cm}^{-3}, & 3.23 < x \leq 3.31 \text{ cm} \\ 1.24 \times 10^{23} \text{ cm}^{-3}, & 3.31 < x \leq 3.37 \text{ cm} \end{cases}$
$C(x)$	$\begin{cases} 8.437(23.6), & 3.148564 \leq x \leq 3.15 \text{ cm} \\ 1.000(1.00), & 3.15 < x \leq 3.23 \text{ cm} \\ 1.090(1.20), & 3.23 < x \leq 3.31 \text{ cm} \\ 1.000(1.00), & 3.31 < x \leq 3.37 \text{ cm} \end{cases}$

# $\Xi^0 n \rightarrow \Xi^- p$



- $\sigma(\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}) = (22.1 \pm 5.3_{\text{stat}} \pm 4.5_{\text{sys}})$  mb at  $P_{\Xi^0} \approx 0.818$  GeV/c.
- Taking the effective number of reaction neutrons in  ${}^9\text{Be}$  nucleus as 3,  $\sigma(\Xi^0 n \rightarrow \Xi^- p) = (7.4 \pm 1.8_{\text{stat}} \pm 1.5_{\text{sys}})$  mb, consistent with theoretical predictions.



This work is the first study of hyperon-nucleon interaction in electron-positron collisions, and opens up a new direction for such research.

LO : H. Polinder, J.H., U.-G. Meißner, PLB 653 (2007) 29  
 NLO16: J.H., U.-G. Meißner, S. Petschauer, NPA 954 (2016) 273  
 NLO19: J.H., U.-G. Meißner, EPJA 55 (2019) 23

No significant H-dibaryon signals are seen

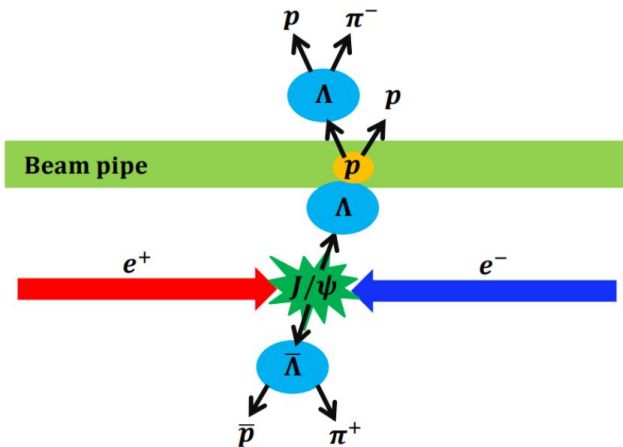




$$\Lambda(\bar{\Lambda})p \rightarrow \Lambda(\bar{\Lambda})p$$

Reaction chain :

$$J/\psi \rightarrow \Lambda\bar{\Lambda}, \Lambda p \rightarrow \Lambda p, \Lambda \rightarrow p\pi^-, \bar{\Lambda} \rightarrow \bar{p}\pi^+.$$



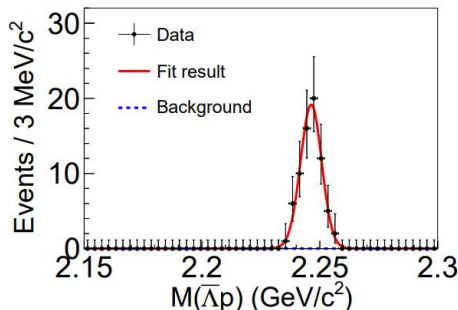
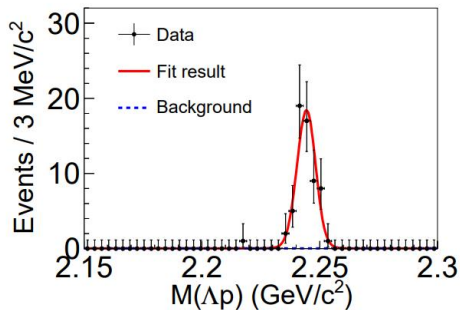
Two-body decay,  $P_{\Lambda} \approx 1.074 \text{ GeV}/c$ ,

Very small horizontal crossing angle  
of 11 mrad for  $e^+$  and  $e^-$  beams

Very small range of  $0.017 \text{ GeV}/c$   
above and below  $1.074 \text{ GeV}/c$  for  $P_{\Lambda}$ .

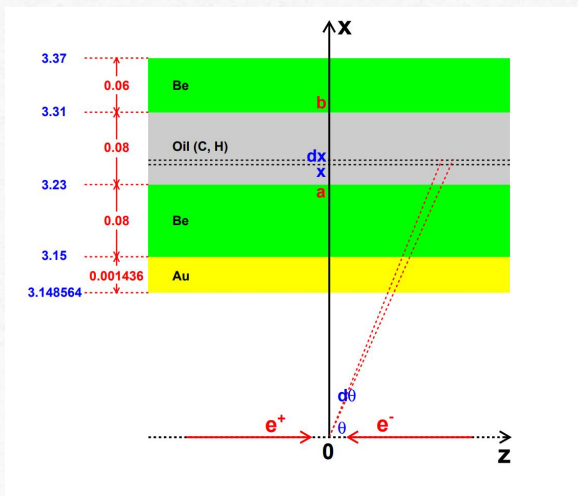


# $\Lambda(\bar{\Lambda})p \rightarrow \Lambda(\bar{\Lambda})p$



$$N(\Lambda p \rightarrow \Lambda p) = 60.9 \pm 7.8$$

$$N(\bar{\Lambda} p \rightarrow \bar{\Lambda} p) = 72.0 \pm 8.5$$



The center-of-mass energies for the incident  $\Lambda/\bar{\Lambda}$  and a static  $p$  are all  $2.243 \text{ GeV}/c^2$  within a range of  $\pm 0.005 \text{ GeV}/c^2$

Clear enhancements are seen around  $2.243 \text{ GeV}/c^2$ , corresponding to the reactions  $\Lambda(\bar{\Lambda})p \rightarrow \Lambda(\bar{\Lambda})p$ , respectively

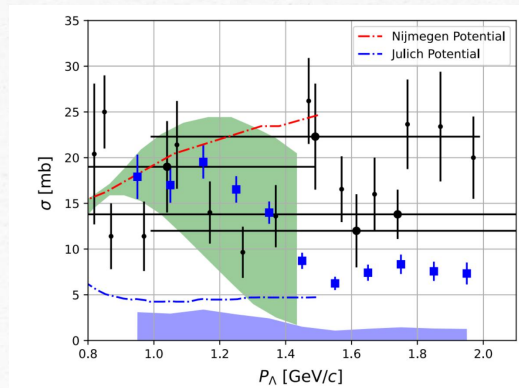
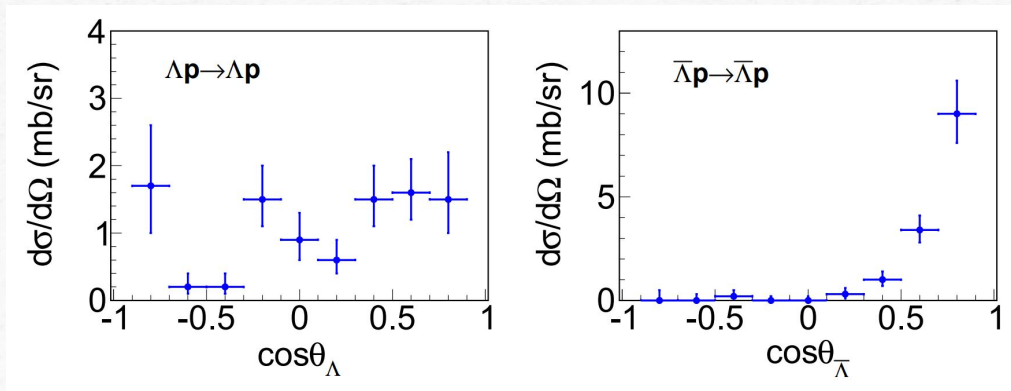


# $\Lambda(\bar{\Lambda})p \rightarrow \Lambda(\bar{\Lambda})p$

## First measurement of antihyperon-nucleon scattering!

- $\sigma(\Lambda + p \rightarrow \Lambda + p) = (12.2 \pm 1.6_{\text{stat}} \pm 1.1_{\text{sys}})$  mb and  $\sigma(\bar{\Lambda} + p \rightarrow \bar{\Lambda} + p) = (17.5 \pm 2.1_{\text{stat}} \pm 1.6_{\text{sys}})$  mb at  $P_{\Lambda} \approx 1.074$  GeV/c within  $-0.9 < \cos\theta_{\Lambda} < 0.9$ .
- The differential cross sections of the two reactions are measured within  $-0.9 < \cos\theta_{\Lambda} < 0.9$ , while there is a slight tendency of forward scattering for  $\Lambda p \rightarrow \Lambda p$ , and a strong forward peak for  $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$

## Consistent results from CLAS experiment







What will be studied





➤  $\Sigma^+ n \rightarrow \Lambda p, \Sigma^+ n \rightarrow \Sigma^0 p$

➤  $\Xi^0 n \rightarrow \Lambda \Lambda, \Xi^- p \rightarrow \Lambda \Lambda$

.....

**More results will come out soon !!!**



# Prospects at STCF



- Far more precise measurements thanks to the gaint statistics
- First measurement of the interaction between  $\Omega^-$  and nuclei/nucleon (three strange quarks, spin-3/2)
- Measurements of the differential cross sections and momentum-dependent cross section
- Search for potential hypernucleus

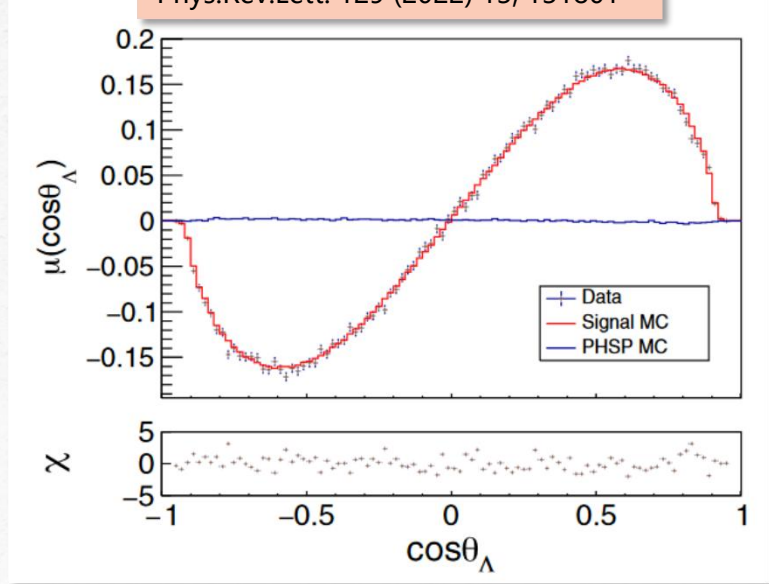
Hyperon	$c\tau$ (cm)	decay mode	$\mathcal{B}_{decay}$ [65] ( $\times 10^{-3}$ )	$p_{max}$ (MeV/c)	$n_{BP}^Y$ ( $\times 10^5$ for BESIII or $\times 10^8$ for STCF)	$\mathcal{B}_{tag}$ (%)	$\mathcal{L}_Y/N_{ST}$ ( $10^{21}$ . $cm^{-2}$ )	Estimated signal yield ( $\times 10^3$ for STCF)
$\Lambda$	7.89	$J/\psi \rightarrow \Lambda \bar{\Lambda}$	$1.89 \pm 0.09$	1074	26	64	23.59	5290
$\Sigma^+$	2.40	$J/\psi \rightarrow \Sigma^+ \bar{\Sigma}^-$	$1.07 \pm 0.04$	992	4	52	4.83	537
$\Xi^0$	8.71	$J/\psi \rightarrow \Xi^0 \bar{\Xi}^0$	$1.17 \pm 0.04$	818	7	64	15.81	2368
$\Xi^-$	4.91	$J/\psi \rightarrow \Xi^- \bar{\Xi}^+$	$0.97 \pm 0.08$	807	3	64	7.44	924
$\Omega^-$	2.46	$\psi(3686) \rightarrow \Omega^- \bar{\Omega}^+$	$0.056 \pm 0.003$	774	0.05	43	2.61	29



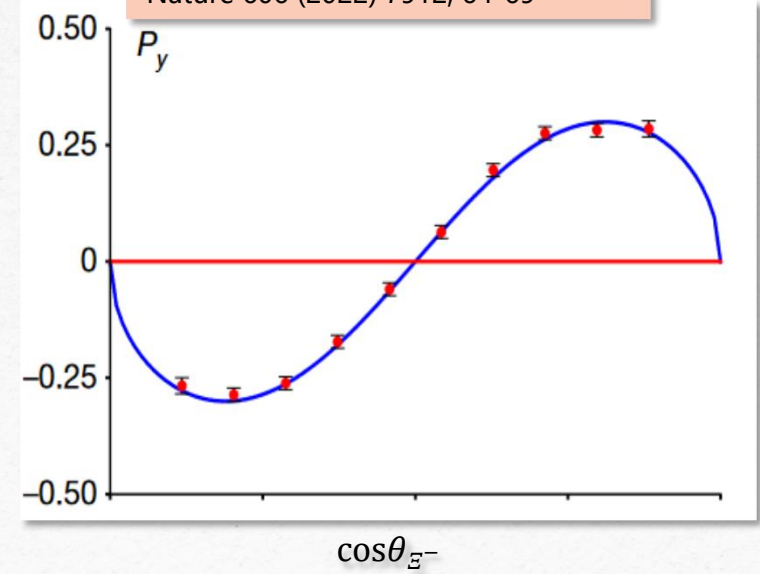


- Potential measurement of the differential cross sections with respect to the polarization of the incident hyperons

Phys.Rev.Lett. 129 (2022) 13, 131801



Nature 606 (2022) 7912, 64-69





# Summary





## Summary



- The cross section of  $\Xi^0 n \rightarrow \Xi^- p$  is firstly measured with  $\Xi^0$  beam from the decay  $J/\psi \rightarrow \Xi^0 \bar{\Xi}^0$  based on 10 billion  $J/\psi$  data at BESIII to be  $\sigma(\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}) = (22.1 \pm 5.3_{\text{stat}} \pm 4.5_{\text{sys}})$  mb. The first study of hyperon-nucleon interaction in electron-positron collisions, opening a new direction for such research.
- The cross section of  $\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X$  is studied with  $\Lambda$  from  $J/\psi \rightarrow \Lambda \bar{\Lambda}$  to be  $\sigma(\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X) = (37.3 \pm 4.7_{\text{stat}} \pm 3.5_{\text{sys}})$  mb. The first attempt to investigate  $\Lambda$ -nucleus interaction at an  $e^+e^-$  collider.
- With more statistics in future STCF, the momentum-dependent cross section and differential cross sections can also be studied.



Thank you !

**Han Miao (妙晗) and Jianyu Zhang (张剑宇)**

Institute of High Energy Physics  
University of Chinese Academy of Sciences

2024. 1. 16