

η/η' meson physics at STCF

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η/η' mesons

- The η and η' are **special**:
 - The η is a pNGB, with $m_\eta = 547.862$ MeV and $\Gamma_\eta = 1.31$ keV
 - The η' : not a pNGB due to $U(1)_A$ anomaly, $m_{\eta'} = 957.78(6)$ MeV, $\Gamma_{\eta'} = 196$ keV
 - Eigenstates of the C, P, CP and G operators: $I^G J^{PC} = 0^+ 0^- +$
 - Flavor **conserving** decays \Rightarrow laboratory for symmetry tests
 - All their EM and strong decays are **suppressed** at LO $\sim \mathcal{O}(\alpha_{\text{em}}^2)$ or $\mathcal{O}((m_u - m_d)^2)$
 - Window to **BSM** physics \Rightarrow Dark sector physics:
 - **Decays** to new light dark particles are 2-/or 3-body decays that mimic 3-, 4-, or 5-body final states (often very rare)
- Perfect **laboratory** to stress-test the SM in search for BSM physics

η/η' events: BESIII vs STCF

- $e^+e^- \rightarrow J/\psi$
 $\sqrt{s} \sim 2\text{--}4.95 \text{ GeV (BESIII), } 2\text{--}7 \text{ GeV (STCF)}$
- $10^{10} J/\psi$ (BESIII)
 $3.4 \times 10^{12} J/\psi/\text{year}$ (STCF)
- High production of η/η' in J/ψ decays

Decay mode	η/η' events (BESIII)	η/η' events (STCF)
$J/\psi \rightarrow \gamma\eta'$	5.2×10^7	1.8×10^{10}
$J/\psi \rightarrow \gamma\eta$	1.1×10^7	3.7×10^9
$J/\psi \rightarrow \phi\eta'$	2.5×10^6	2.5×10^9
$J/\psi \rightarrow \phi\eta$	4×10^6	1.6×10^9

STCF: M. Achasov, *et. al.*
Front. Phys. 19(1), 14701 (2024)

- KLOE: Total of $3 \times 10^8 \eta$ and $5 \times 10^5 \eta'$
- Future experiments:
 - JEF at JLab Hall D (approved): $6.5 \times 10^7 \eta$ and $4.9 \times 10^7 \eta'$ per 100 days
 - REDTOP (proposed, 2203.07651 [hep-exp]): $10^{13} \eta$ and $10^{11} \eta'$ per year

Decay Channel	Physics	Publication
$\eta' \rightarrow \pi^+\pi^-\pi^+\pi^- / \pi^+\pi^-\pi^0\pi^0$	First Observation - BR	PRL112, 251801 (2014)
$\eta' \rightarrow \gamma e^+e^-$	First Observation - BR - TFF	PRD92, 012001 (2015)
$\eta \rightarrow \pi^+\pi^-\pi^0 \& \eta/\eta' \rightarrow \pi^0\pi^0\pi^0$	Matrix Elements	PRD92, 012014 (2015)
$\eta' \rightarrow \omega e^+e^-$	First Observation - BR	PRD92, 051101 (2015)
$\eta' \rightarrow K\pi$	UL	PRD93, 072008 (2016)
$\eta' \rightarrow \rho\pi$	First Observation - BR	PRL118, 012001 (2017)
$\eta' \rightarrow \gamma\gamma\pi^0$	BR	PRD96, 012005 (2017)
$\eta' \rightarrow \gamma\pi^+\pi^-$	BR - Box Anomaly	PRL120, 242003 (2018)
$\eta' \rightarrow \pi^+\pi^-\eta/\pi^0\pi^0\eta$	Matrix Elements - Cusp Effect	PRD97, 012003 (2018)
$P \rightarrow \gamma\gamma$	BRs	PRD97, 072014 (2018)
$\omega \rightarrow \pi^+\pi^-\pi^0$	Dalitz Plot Analysis	PRD98, 112007 (2018)
Absolut BR of η' decays	BRs	PRL122, 142002 (2019)
$\eta' \rightarrow \gamma\gamma\eta$	UL	PRD100, 052015 (2019)
$\eta' \rightarrow \pi^0\pi^0\pi^0\pi^0$	UL	PRD101, 032001 (2020)
$\eta' \rightarrow \pi^+\pi^-\mu^+\mu^-$	First Observation - BR	PRD103, 072006 (2021)
Absolute BR of η decays	BRs	PRD104, 092004 (2021)

Rich physics program at η, η' factories

Standard Model highlights

- Theory input for light-by-light scattering for $(g-2)_\mu$
- Extraction of light quark masses
- QCD scalar dynamics

Fundamental symmetry tests

- P/CP violation
- C/CP violation

[Kobzarev & Okun (1964), Prentki & Veltman (1965), Lee (1965), Lee & Wolfenstein (1965), Bernstein et al (1965)]

Dark sectors (MeV—GeV)

- Vector bosons (dark photon, B boson, X boson)
- Scalars
- Pseudoscalars (ALPs)

(Plus other channels that have not been searched for to date)

Channel	Expt. branching ratio	Discussion
$\eta \rightarrow 2\gamma$	39.41(20)%	chiral anomaly, $\eta-\eta'$ mixing
$\eta \rightarrow 3\pi^0$	32.68(23)%	$m_u - m_d$
$\eta \rightarrow \pi^0\gamma\gamma$	$2.56(22) \times 10^{-4}$	χ PT at $O(p^6)$, leptophobic B boson, light Higgs scalars
$\eta \rightarrow \pi^0\pi^0\gamma\gamma$	$< 1.2 \times 10^{-3}$	χ PT, axion-like particles (ALPs)
$\eta \rightarrow 4\gamma$	$< 2.8 \times 10^{-4}$	$< 10^{-11}$ [52]
$\eta \rightarrow \pi^+\pi^-\pi^0$	22.92(28)%	$m_u - m_d$, C/CP violation, light Higgs scalars
$\eta \rightarrow \pi^+\pi^-\gamma$	4.22(8)%	chiral anomaly, theory input for singly-virtual TFF and $(g-2)_\mu$, P/CP violation
$\eta \rightarrow \pi^+\pi^-\gamma\gamma$	$< 2.1 \times 10^{-3}$	χ PT, ALPs
$\eta \rightarrow e^+e^-\gamma$	$6.9(4) \times 10^{-3}$	theory input for $(g-2)_\mu$, dark photon, protophobic X boson
$\eta \rightarrow \mu^+\mu^-\gamma$	$3.1(4) \times 10^{-4}$	theory input for $(g-2)_\mu$, dark photon
$\eta \rightarrow e^+e^-$	$< 7 \times 10^{-7}$	theory input for $(g-2)_\mu$, BSM weak decays
$\eta \rightarrow \mu^+\mu^-$	$5.8(8) \times 10^{-6}$	theory input for $(g-2)_\mu$, BSM weak decays, P/CP violation
$\eta \rightarrow \pi^0\pi^0\ell^+\ell^-$		C/CP violation, ALPs
$\eta \rightarrow \pi^+\pi^-e^+e^-$	$2.68(11) \times 10^{-4}$	theory input for doubly-virtual TFF and $(g-2)_\mu$, P/CP violation, ALPs
$\eta \rightarrow \pi^+\pi^-\mu^+\mu^-$	$< 3.6 \times 10^{-4}$	theory input for doubly-virtual TFF and $(g-2)_\mu$, P/CP violation, ALPs
$\eta \rightarrow e^+e^-e^+e^-$	$2.40(22) \times 10^{-5}$	theory input for $(g-2)_\mu$
$\eta \rightarrow e^+e^+\mu^+\mu^-$	$< 1.6 \times 10^{-4}$	theory input for $(g-2)_\mu$
$\eta \rightarrow \mu^+\mu^-\mu^+\mu^-$	$< 3.6 \times 10^{-4}$	theory input for $(g-2)_\mu$
$\eta \rightarrow \pi^+\pi^-\pi^0\gamma$	$< 5 \times 10^{-4}$	direct emission only
$\eta \rightarrow \pi^+\pi^-\nu_e$	$< 1.7 \times 10^{-4}$	second-class current
$\eta \rightarrow \pi^+\pi^-$	$< 4.4 \times 10^{-6}$ [53]	P/CP violation
$\eta \rightarrow 2\pi^0$	$< 3.5 \times 10^{-4}$	P/CP violation
$\eta \rightarrow 4\pi^0$	$< 6.9 \times 10^{-7}$	P/CP violation
		Gan, Kubis, Passemar, ST (2020)

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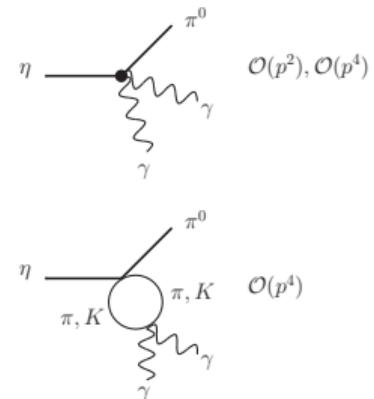
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$\eta \rightarrow \pi^0 \gamma\gamma$ decays: Theoretical motivation

- SM motivation:

Reference	$\Gamma(\eta \rightarrow \pi^0 \gamma\gamma)$ [eV]
$\mathcal{O}(p^2), \mathcal{O}(p^4)$ tree-level χ PT	0
$\pi + K$ loops at $\mathcal{O}(p^4)$	1.87×10^{-3}
Experimental value (pdg)	0.34(3)



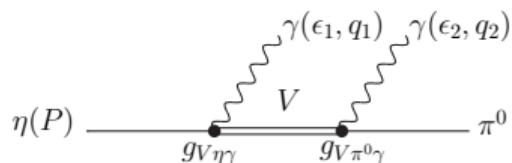
- 1st sizable contribution comes at $\mathcal{O}(p^6)$, but LEC's are not well known
- To test ChPT and a wide range of chiral models, *e. g.* VMD and L σ M



- BSM motivation: search for a B boson via $\eta \rightarrow B\gamma \rightarrow \pi^0 \gamma\gamma$

Vector meson exchange contributions

- Six diagrams corresponding to the exchange of $V = \rho^0, \omega, \phi$



$$\mathcal{A}_{\eta \rightarrow \pi^0 \gamma\gamma}^{\text{VMD}} = \sum_{V=\rho^0, \omega, \phi} g_{V\eta\gamma} g_{V\pi^0\gamma} \left[\frac{(P \cdot q_2 - m_\eta^2)\{a\} - \{b\}}{D_V(t)} + \left\{ \begin{array}{l} q_2 \leftrightarrow q_1 \\ t \leftrightarrow u \end{array} \right\} \right] ,$$

- Mandelstam variables and Lorentz structures given by:

$$t, u = (P - q_{2,1})^2 = m_\eta^2 - 2P \cdot q_{2,1} ,$$

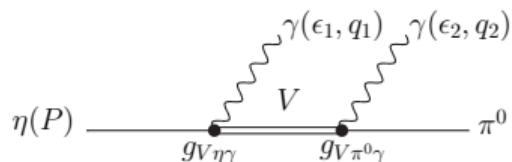
$$\{a\} = (\epsilon_1 \cdot \epsilon_2)(q_1 \cdot q_2) - (\epsilon_1 \cdot q_2)(\epsilon_2 \cdot q_1) ,$$

$$\begin{aligned} \{b\} = & (\epsilon_1 \cdot q_2)(\epsilon_2 \cdot P)(P \cdot q_1) + (\epsilon_2 \cdot q_1)(\epsilon_1 \cdot P)(P \cdot q_2) \\ & - (\epsilon_1 \cdot \epsilon_2)(P \cdot q_1)(P \cdot q_2) - (\epsilon_1 \cdot P)(\epsilon_2 \cdot P)(q_1 \cdot q_2) \end{aligned}$$

- The decays $\eta' \rightarrow \{\pi^0, \eta\}\gamma\gamma$ are formally identical: $g_{V\eta\gamma} g_{V\pi^0\gamma} \rightarrow g_{V\eta'\gamma} g_{V\{\pi^0, \eta\}\gamma}$

$\eta \rightarrow \pi^0 \gamma\gamma$ decays: VMD calculation

- Six diagrams corresponding to the exchange of $V = \rho^0, \omega, \phi$



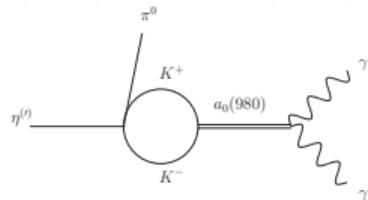
$$\mathcal{A}_{\eta \rightarrow \pi^0 \gamma\gamma}^{\text{VMD}} = \sum_{V=\rho^0, \omega, \phi} g_{V\eta\gamma} g_{V\pi^0\gamma} \left[\frac{(P \cdot q_2 - m_\eta^2)\{a\} - \{b\}}{D_V(t)} + \left\{ \begin{array}{l} q_2 \leftrightarrow q_1 \\ t \leftrightarrow u \end{array} \right\} \right],$$

- $g_{VP\gamma}$ couplings: $\Gamma_{V \rightarrow P\gamma}^{\text{exp}} = \frac{1}{3} \frac{g_{VP\gamma}^2}{32\pi} \left(\frac{m_V^2 - m_P^2}{m_V} \right)^3$, $\Gamma_{P \rightarrow V\gamma}^{\text{exp}} = \frac{g_{VP\gamma}^2}{32\pi} \left(\frac{m_P^2 - m_V^2}{m_P} \right)^3$,

Decay	Branching ratio (pdg)	$ g_{VP\gamma} \text{ GeV}^{-1}$
$\rho^0 \rightarrow \pi^0 \gamma$	$(4.7 \pm 0.6) \times 10^{-4}$	0.22(1)
$\rho^0 \rightarrow \eta\gamma$	$(3.00 \pm 0.21) \times 10^{-4}$	0.48(2)
$\eta' \rightarrow \rho^0 \gamma$	$(28.9 \pm 0.5)\%$	0.40(1)
$\omega \rightarrow \pi^0 \gamma$	$(8.40 \pm 0.22)\%$	0.70(1)
$\omega \rightarrow \eta\gamma$	$(4.5 \pm 0.4) \times 10^{-4}$	0.135(6)
$\eta' \rightarrow \omega\gamma$	$(2.62 \pm 0.13)\%$	0.127(4)
$\phi \rightarrow \pi^0 \gamma$	$(1.30 \pm 0.05) \times 10^{-3}$	0.041(1)
$\phi \rightarrow \eta\gamma$	$(1.303 \pm 0.025)\%$	0.2093(20)
$\phi \rightarrow \eta'\gamma$	$(6.22 \pm 0.21) \times 10^{-5}$	0.216(4)

L σ M for the scalar resonance contributions

- χ PT loops complemented by the exchange of scalar resonances, $a_0(980), \kappa, \sigma, f_0(980)$, e.g.:



$$\mathcal{A}_{\eta^{(\prime)} \rightarrow \pi^0 \gamma\gamma}^{\text{L}\sigma\text{M}} = \frac{2\alpha}{\pi} \frac{1}{m_{K^+}^2} L(s_K)\{a\} \times \mathcal{A}_{K^+ K^- \rightarrow \pi^0 \eta^{(\prime)}}^{\text{L}\sigma\text{M}},$$

- Scalar amplitudes:

$$\begin{aligned} \mathcal{A}_{K^+ K^- \rightarrow \pi^0 \eta^{(\prime)}}^{\text{L}\sigma\text{M}} = & \frac{1}{2f_\pi f_K} \left\{ (s - m_{\eta^{(\prime)}}^2) \frac{m_K^2 - m_{a_0}^2}{D_{a_0}(s)} \cos \varphi_P + \frac{1}{6} \left[(5m_{\eta^{(\prime)}}^2 + m_\pi^2 - 3s) \cos \varphi_P \right. \right. \\ & \left. \left. - \sqrt{2}(m_{\eta^{(\prime)}}^2 + 4m_K^2 + m_\pi^2 - 3s) \sin \varphi_P \right] \right\}, \end{aligned}$$

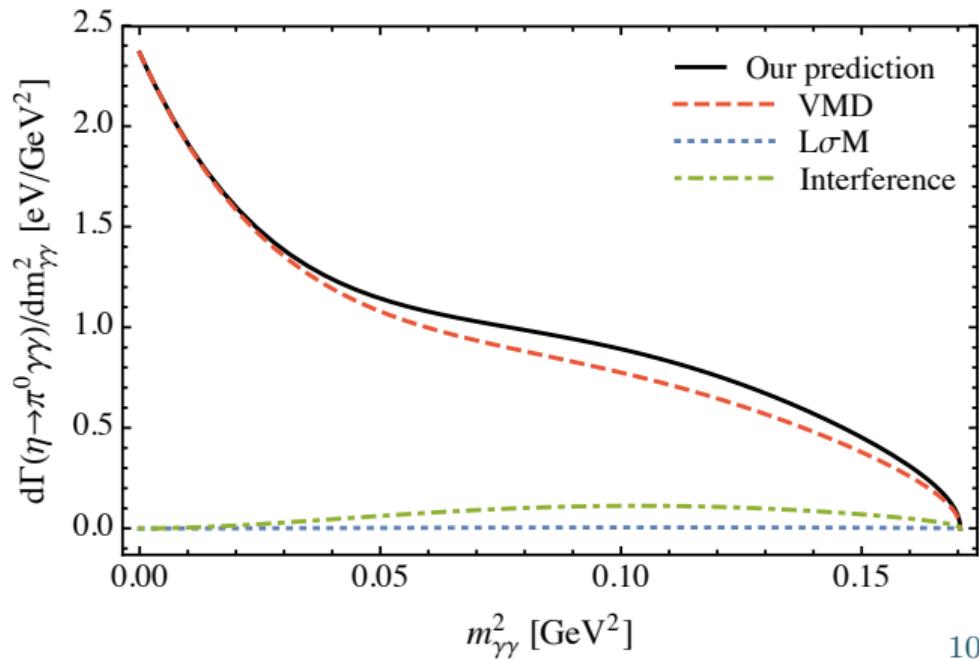
- Complete one-loop propagator for the scalar resonances:

$$D_R(s) = s - m_R^2 + \text{Re}\Pi(s) - \text{Re}\Pi(m_R^2) + i\text{Im}\Pi(s),$$

$\eta \rightarrow \pi^0 \gamma\gamma$ predictions

- Our theoretical prediction $BR = 1.35(8) \times 10^{-4}$
(Escribano, SGS, Jora, Royo, Phys.Rev.D 102, 034026 (2020))

- VMD dominates:
- ρ : 27% of the signal
- ω : 21% of the signal
- ϕ : 0% of the signal
- interference between ρ - ω - ϕ : 52%
- interference between scalar and vector mesons: 7%



$\eta \rightarrow \pi^0 \gamma\gamma$ predictions

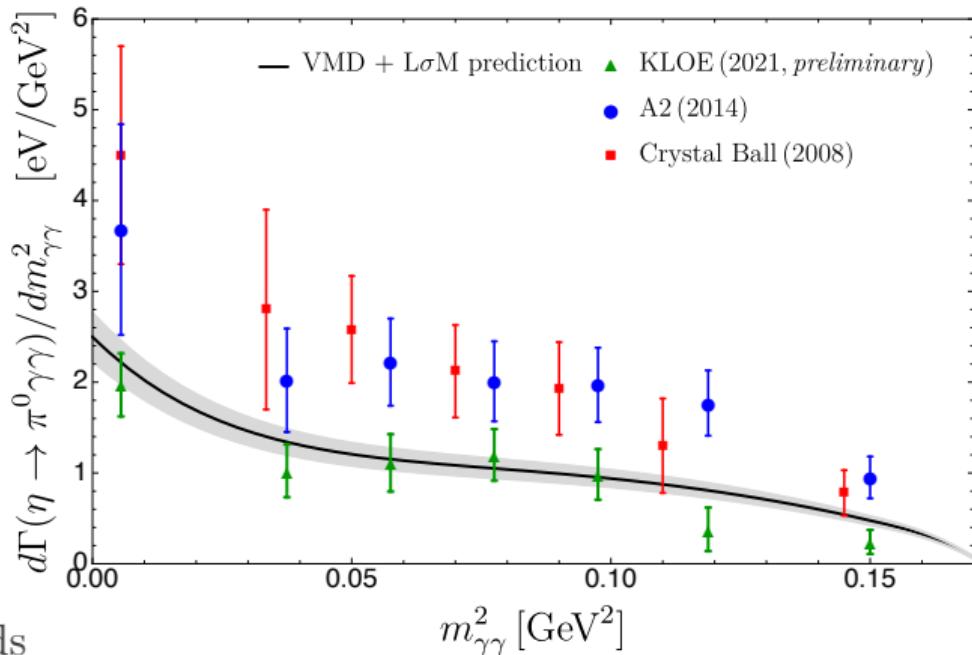
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- $BR_{\text{exp}} = 2.55(22) \times 10^{-4}$ (pdg)

- Comparison with data

- Shape of the A2 and Crystal Ball spectra is captured well (normalization offset)

- Good agreement with (preliminary) KLOE data
(B. Cao, PoS EPS-HEP2021 (2022) 409)

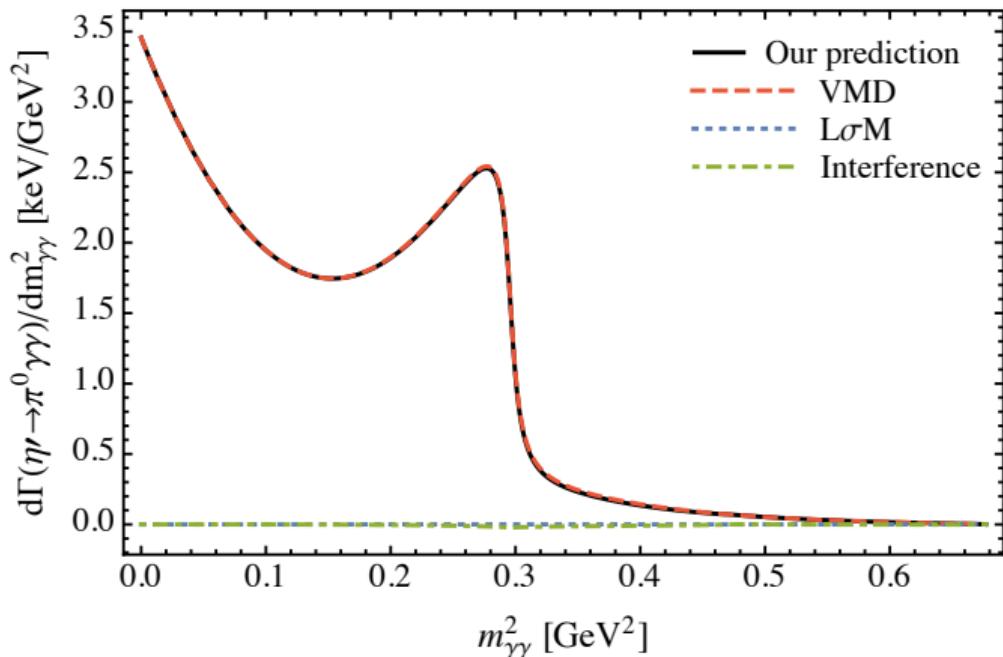


- The experimental situation needs to be **clarified** (A2, BESIII, JEF, REDTOP, SCTF)

$\eta' \rightarrow \pi^0 \gamma\gamma$ predictions

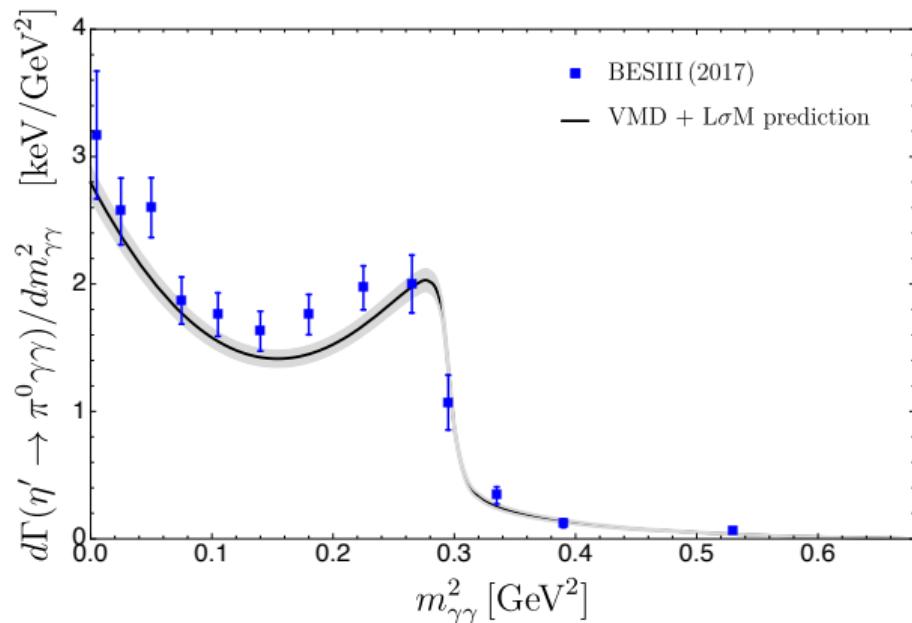
- Our theoretical prediction $BR = 2.91(21) \times 10^{-3}$
(Escribano, SGS, Jora, Royo, Phys.Rev.D 102, 034026 (2020))

- VMD completely dominates:
- ω : 78% of the signal
- ρ : 5% of the signal
- ϕ : 0% of the signal
- interference: 17%



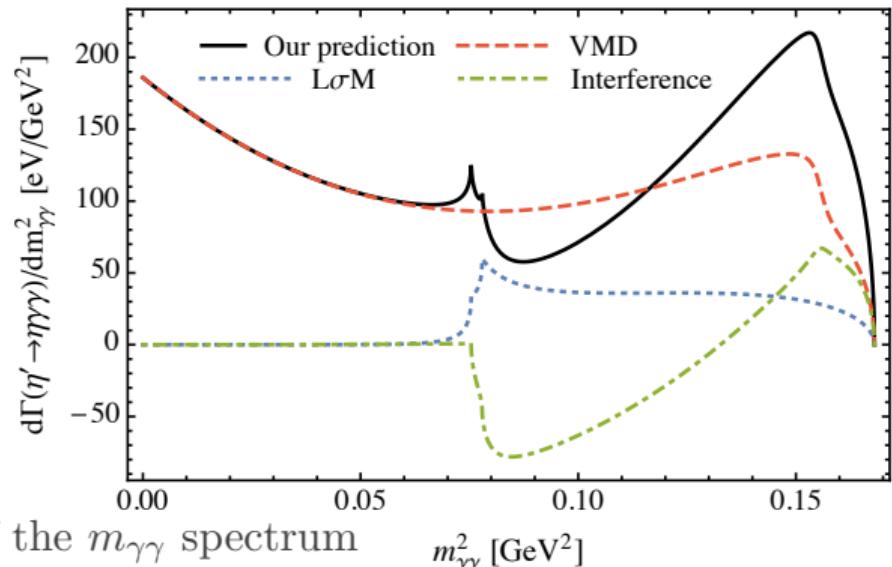
$\eta' \rightarrow \pi^0 \gamma\gamma$ predictions

- Our theoretical prediction $BR = 2.91(21) \times 10^{-3}$ ([Phys.Rev.D 102, 034026 \(2020\)](#))
- First time $m_{\gamma\gamma}$ invariant mass distribution by BESIII;
 $BR = 3.20(7)(23) \times 10^{-3}$ ([Ablikim *et. al.* Phys.Rev.D 96, 012005 \(2017\)](#))



$\eta' \rightarrow \eta\gamma\gamma$ predictions

- 1st BR measurement by BESIII, $BR = 8.25(3.41)(0.72) \times 10^{-5}$ or $BR < 1.33 \times 10^{-4}$ at 90% C.L. ([Ablikim *et. al.* Phys.Rev.D 100, 052015 \(2019\)](#))
- Our theoretical predictions $BR = 1.17(8) \times 10^{-4}$
([R. Escribano, S. G-S, R. Jora, E. Royo, Phys.Rev.D 102, 034026 \(2020\)](#))
 - VMD predominates (91% of the signal)
 - Substantial scalar meson effects (16%)
 - Interference between scalar and vector mesons (7%)



- We look forward to the release of the $m_{\gamma\gamma}$ spectrum

Leptophobic B boson model

- New boson arising from a new $U(1)_B$ gauge symmetry

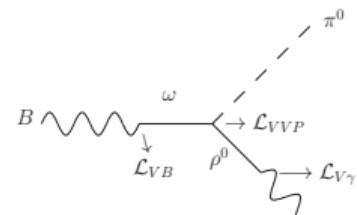
$$\mathcal{L}_{\text{int}} = \left(\frac{1}{3} \mathbf{g}_B + \varepsilon Q_q e \right) \bar{q} \gamma^\mu q B_\mu - \varepsilon e \bar{\ell} \gamma^\mu \ell B_\mu ,$$

- Couples (predominantly) to quarks
- \mathbf{g}_B new gauge (universal?) coupling, $\alpha_B = \mathbf{g}_B^2 / 4\pi$
- Preserves QCD symmetries (C, P, T)
- B is a singlet under isospin:
 - $I^G(J^{PC}) = 0^-(1^{--}) \Rightarrow B$ is ω **meson** like
 - $\varepsilon = eg_B/(4\pi)^2$: (subleading) γ -like coupling to fermions
- Searches depend on the mass m_B and decay channels
- Searches on meson decays are gaining attention
 - $\phi \rightarrow \eta B \rightarrow \eta \pi^0 \gamma$ (KLOE-II), $\eta \rightarrow \pi^0 \gamma \gamma$ (JEF), $\eta \rightarrow \pi^+ \pi^- \gamma$ (Belle-II)

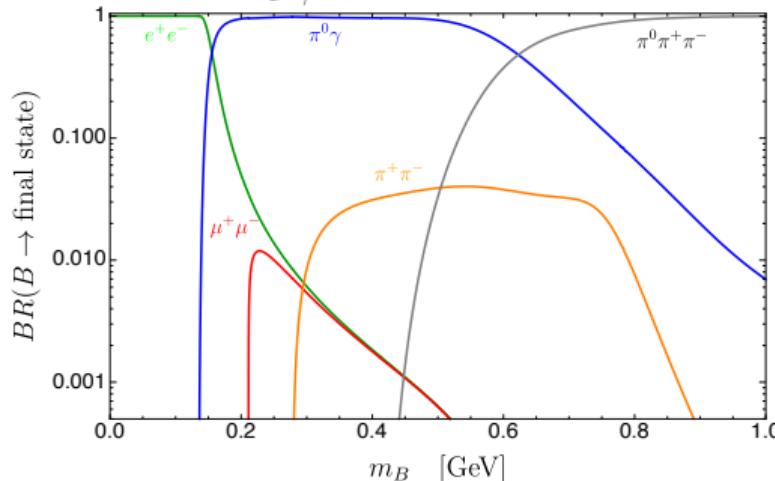
Calculation of hadronic processes and limits

- Following the conventional VMD picture, $\mathcal{L}_{V\gamma} \rightarrow \mathcal{L}_{VB}$

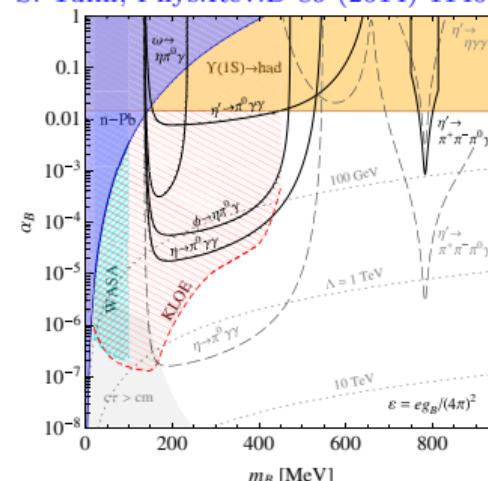
— $A^\mu \rightarrow B^\mu$, $e \rightarrow g_B$ and $Q = 1/3$, $\mathcal{L}_{VB} = -2\frac{1}{3}\textcolor{violet}{g}_B g f_\pi^2 B^\mu \text{tr}[V^\mu]$,



$$\Gamma_{B \rightarrow \pi^0 \gamma} = \frac{\alpha_B \alpha_{em} m_B^3}{96 \pi^3 f_\pi^2} \left(1 - \frac{m_\pi^2}{m_B^2}\right)^3 |F_\omega(m_B^2)|^2,$$



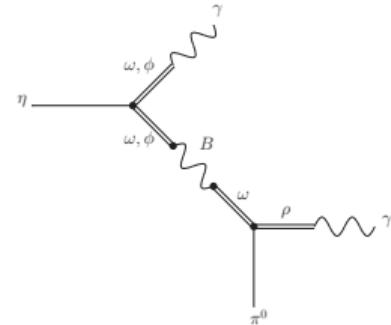
S. Tulin, Phys.Rev.D 89 (2014) 114008



$\eta \rightarrow \pi^0 \gamma \gamma$ decays: *B* boson calculation

- Two diagrams corresponding to the exchange of a ***B* boson**

$$\mathcal{A}_{\eta \rightarrow \pi^0 \gamma \gamma}^{B \text{ boson}} = g_{B\eta\gamma}(t) g_{B\pi^0\gamma}(t) \left[\frac{(P \cdot q_2 - m_\eta^2)\{a\} - \{b\}}{m_B^2 - t - i\sqrt{t}\Gamma_B(t)} + \left\{ \begin{array}{l} q_2 \leftrightarrow q_1 \\ t \leftrightarrow u \end{array} \right\} \right],$$

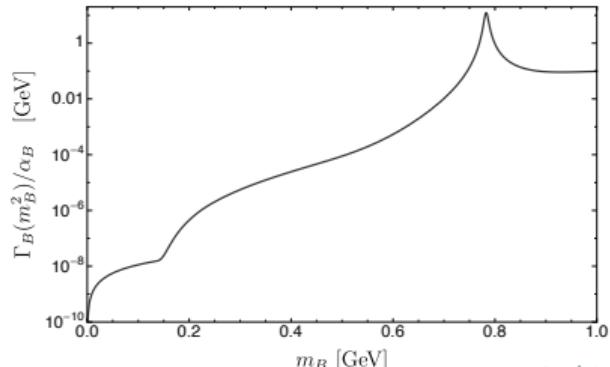


- $g_{BP\gamma}$ couplings:

$$g_{B\pi^0\gamma}(t) = \frac{\sqrt{2}eg_B}{4\pi^2 f_\pi} F_\omega(t), \quad g_{B\eta\gamma}(t) = \frac{eg_B}{12\pi^2 f_\pi} \frac{1}{\sqrt{3}} [(c_\theta - \sqrt{2}s_\theta)F_\omega(t) + (2c_\theta + \sqrt{2}s_\theta)F_\phi(t)],$$

- Energy-dependent width

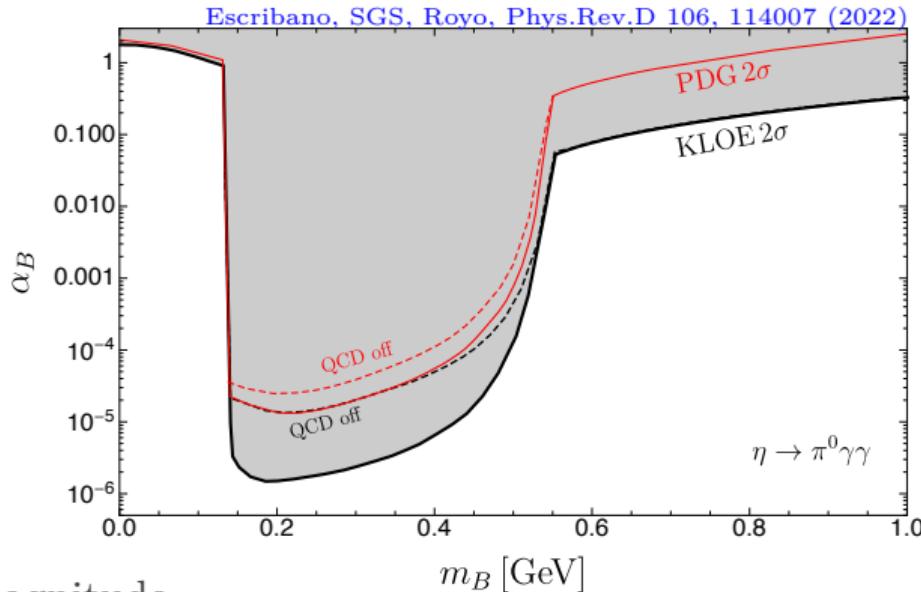
$$\begin{aligned} \Gamma_B(q^2) &= \frac{\gamma_{B \rightarrow \ell^+ \ell^-}(q^2)}{\gamma_{B \rightarrow \ell^+ \ell^-}(m_B^2)} \Gamma_{B \rightarrow \ell^+ \ell^-} \theta(q^2 - 4m_\ell^2) \\ &+ \frac{\gamma_{B \rightarrow \pi^0 \gamma}(q^2)}{\gamma_{B \rightarrow \pi^0 \gamma}(m_B^2)} \Gamma_{B \rightarrow \pi^0 \gamma} \theta(q^2 - m_{\pi^0}^2) \\ &+ \frac{\gamma_{B \rightarrow \pi \pi}(q^2)}{\gamma_{B \rightarrow \pi \pi}(m_B^2)} \Gamma_{B \rightarrow \pi \pi} \theta(q^2 - 4m_\pi^2) \\ &+ \frac{\gamma_{B \rightarrow 3\pi}(q^2)}{\gamma_{B \rightarrow 3\pi}(m_B^2)} \Gamma_{B \rightarrow 3\pi} \theta(q^2 - 9m_\pi^2) \end{aligned}$$



New limits on α_B and m_B

- Not assuming the NWA
- QCD contribution on
- $BR_{\text{VMD+Bboson}} < BR_{\text{exp}}$ at 2σ
 - $BR(\eta \rightarrow \pi^0 \gamma\gamma)^{\text{pdg}}_{\text{exp}} = 2.56(22) \times 10^{-4}$
 - $BR(\eta \rightarrow \pi^0 \gamma\gamma)^{\text{KLOE}}_{\text{exp}} = 1.23(14) \times 10^{-4}$

B. Cao [KLOE], PoS EPS-HEP2021 (2022) 409

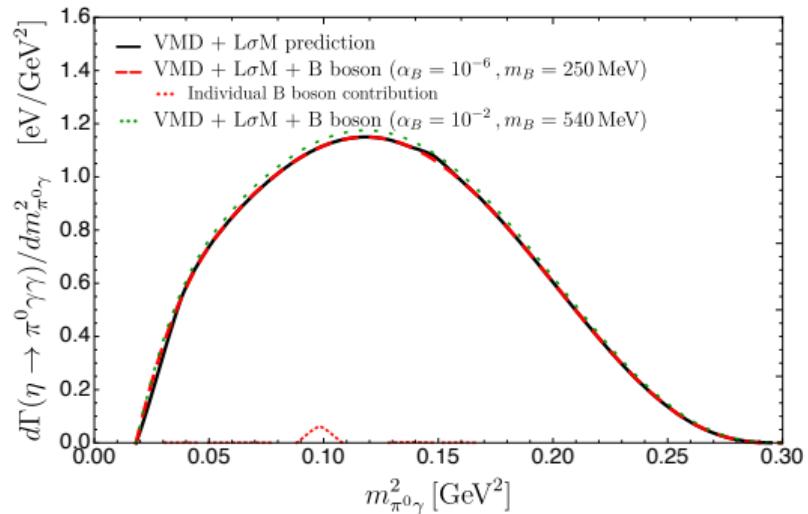
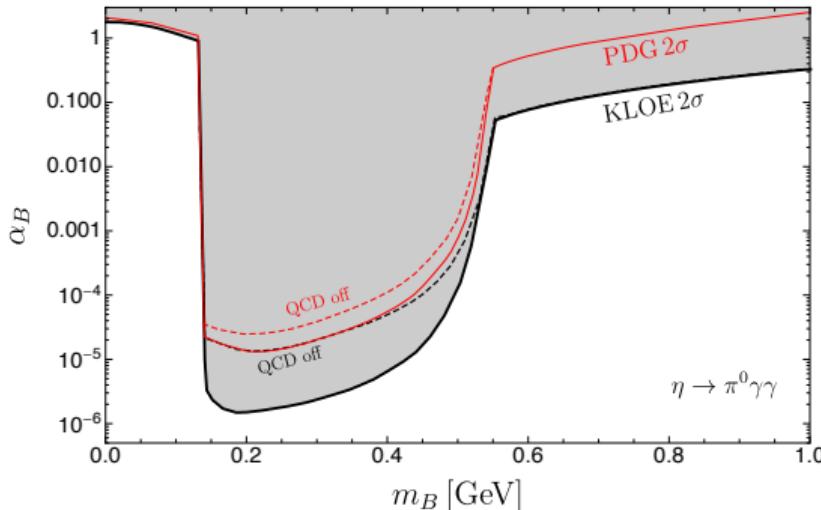


- Limits **strengthened** by one order of magnitude

$\pi^0\gamma$ mass distribution

- These constraints would make a B boson signature suppressed

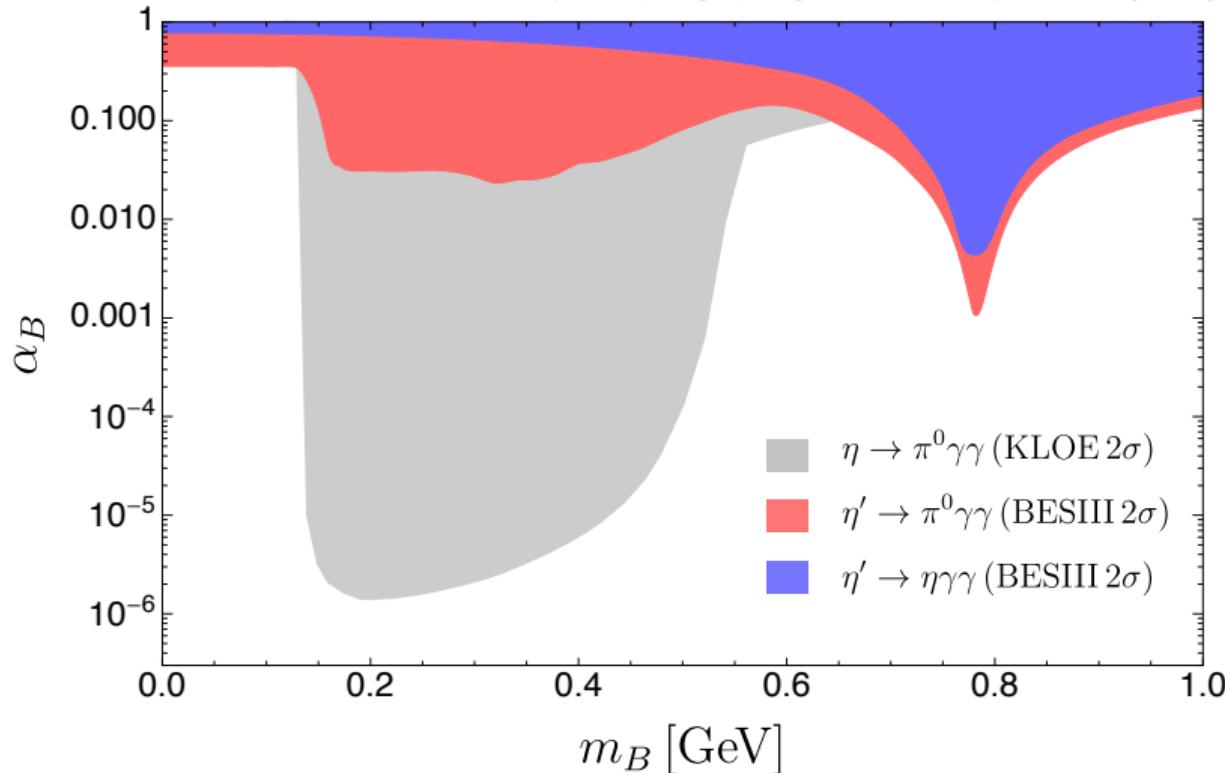
$$\Gamma(\eta \rightarrow \pi^0 \gamma \gamma) \propto \int \frac{\alpha_B^2 dt}{|\mathcal{D}_B(t)|^2} \rightarrow \frac{\alpha_B^2 \pi}{m_B \Gamma_B(m_B^2)} .$$



- Bump-hunt search: accumulate statistics in a certain energy range
- Experimental $\pi^0\gamma$ distribution will be very welcome (JEF?)

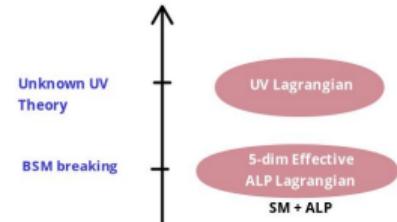
New limits on α_B and m_B

Escribano, SGS, Royo, Phys.Rev.D 106, 114007 (2022)



Axion-Like Particles (ALPs)

- Pseudo-Goldstone bosons from global symmetry breaking
- Example: QCD axion (Peccei-Quinn) \rightarrow solution of the CP problem, potential dark matter candidate. QCD axion mass: $m_a^2 \propto \frac{1}{f_a^2}$.
- “Yukawa basis” (at GeV scale): ALP with gluon and mass couplings

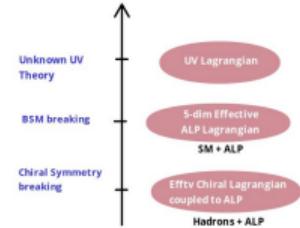


$$\mathcal{L}_{\text{ALP}} = \mathcal{L}_{\text{QCD}} + \frac{1}{2} (\partial_\mu a) (\partial^\mu a) - \frac{1}{2} \textcolor{red}{M}_a^2 a^2 - \textcolor{blue}{Q}_G \frac{\alpha_s}{8\pi} \frac{a}{\textcolor{red}{f}_a} G_{\mu\nu} \tilde{G}^{\mu\nu} + \sum_{q=u,d,s} m_q \bar{q} \left(e^{i \textcolor{blue}{Q}_q \frac{a}{\textcolor{red}{f}_a} \gamma_5} \right) q ,$$

$\textcolor{red}{M}_a^2$: PQ contribution to the mass, $\textcolor{red}{f}_a$: axion decay constant, $\textcolor{blue}{Q}_{q,G}$: PQ charges

- Equivalent to the “usual” derivative basis (related via chiral rotations of the quarks) if weak interactions are neglected
- The heavy-flavor c, b, t quarks contributions are absorbed in $\textcolor{blue}{Q}_G \rightarrow Q_G + Q_{t,b,c}$

Lagrangian for ALPs coupled to mesons



- Step 1: map \mathcal{L}_{ALP} into χPT at leading order

$$\mathcal{L}_{\text{ALP}}^{\chi\text{PT}@LO} = \frac{f_\pi^2}{4} \text{Tr} \left[\partial_\mu U^\dagger \partial^\mu U \right] + \frac{f_\pi^2}{4} \left[2B_0(M_q(a)U + M_q(a)^\dagger U^\dagger) \right] - \frac{1}{2} m_0^2 \left(\eta_0 - \frac{Q_G}{\sqrt{6}} \frac{f_\pi}{f_a} a \right)^2 + \frac{1}{2} \partial_\mu a \partial^\mu a - \frac{1}{2} \mathbf{M}_a^2 a^2$$

$$M_q(a) = \text{diag}(m_u e^{i\mathbf{Q}_u a/\mathbf{f}_a}, m_d e^{i\mathbf{Q}_d a/\mathbf{f}_a}, m_s e^{i\mathbf{Q}_s a/\mathbf{f}_a}),$$

$$U = \exp \left(\frac{i\sqrt{2}\Phi}{f} \right), \quad \Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi_3 + \frac{1}{\sqrt{6}}\eta_8 + \frac{1}{\sqrt{3}}\eta_0 & \pi^+ & K^+ \\ \pi^- & -\frac{1}{\sqrt{2}}\pi_3 + \frac{1}{\sqrt{6}}\eta_8 + \frac{1}{\sqrt{3}}\eta_0 & K^0 \\ K^- & K^0 & -\frac{2}{\sqrt{6}}\eta_8 + \frac{1}{\sqrt{3}}\eta_0 \end{pmatrix}.$$

- Step 2: **diagonalization** of the mass matrix \Rightarrow mixing angles $\theta_{\pi_3 a}$, $\theta_{\eta_8 a}$, $\theta_{\eta_0 a}$

$$\tilde{M}^2 = \begin{pmatrix} m_{\pi_3}^2 & m_{\pi_3 \eta_8}^2 & m_{\pi \eta_0}^2 & m_{\pi_3 a}^2 \\ & m_{\eta_8}^2 & m_{\eta_8 \eta_0}^2 & m_{\eta_8 a}^2 \\ & & m_{\eta_0}^2 & m_{\eta_0 a}^2 \\ & & & m_a^2 \end{pmatrix},$$

$\eta/\eta' \rightarrow \pi\pi a$ decay amplitudes at LO

- Step 3: re-express $\mathcal{L}_{\text{ALP}}^{\chi\text{PT@LO}}$ in terms of the **physical states**

$$\pi_3 \rightarrow \pi_3 + \theta_{\pi_3 a} a^{\text{phys}}, \quad \eta_8 \rightarrow \cos \theta \eta + \sin \theta \eta' + \theta_{\eta_8 a} a^{\text{phys}}, \quad \eta_0 \rightarrow -\sin \theta \eta + \cos \theta \eta' + \theta_{\eta_0 a} a^{\text{phys}},$$

- Physical ALP mass: $m_{a^{\text{phys}}}^2 = \frac{(Q_u + Q_d + Q_s + Q_G)^2 B_0 m_u m_d m_s}{\left(m_u m_d + m_u m_s + m_d m_s + \frac{6 B_0 m_u m_d m_s}{m_0^2} \right)} \frac{f_\pi^2}{f_a^2} + M_a^2,$

$$\mathcal{A}(\eta \rightarrow 2\pi^0 a)|_{\text{LO}} = 2! \frac{m_\pi^2}{f_\pi^2} (\cos \theta - \sqrt{2} \sin \theta) \left[\frac{f_\pi}{2\sqrt{3} f_a} \frac{Q_u m_u + Q_d m_d}{m_u + m_d} - \frac{1}{2\sqrt{3}} \frac{m_d - m_u}{m_u + m_d} \theta_{\pi_3 a} + \frac{1}{6} \theta_{\eta_8 a} + \frac{\sqrt{2}}{6} \theta_{\eta_0 a} \right],$$

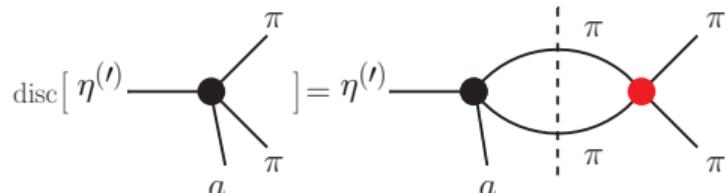
$$\mathcal{A}(\eta \rightarrow \pi^+ \pi^- a)|_{\text{LO}} = \frac{m_\pi^2}{f_\pi^2} (\cos \theta - \sqrt{2} \sin \theta) \left[\frac{f_\pi}{\sqrt{3} f_a} \frac{Q_u m_u + Q_d m_d}{m_u + m_d} - \frac{1}{3\sqrt{3}} \frac{m_d - m_u}{m_u + m_d} \theta_{\pi_3 a} + \frac{1}{3} \theta_{\eta_8 a} + \frac{\sqrt{2}}{3} \theta_{\eta_0 a} \right],$$

$$\mathcal{A}(\eta' \rightarrow 2\pi^0 a)|_{\text{LO}} = 2! \frac{m_\pi^2}{f_\pi^2} (\sqrt{2} \cos \theta + \sin \theta) \left[\frac{f_\pi}{2\sqrt{3} f_a} \frac{Q_u m_u + Q_d m_d}{m_u + m_d} - \frac{1}{2\sqrt{3}} \frac{m_d - m_u}{m_u + m_d} \theta_{\pi_3 a} + \frac{1}{6} \theta_{\eta_8 a} + \frac{\sqrt{2}}{6} \theta_{\eta_0 a} \right],$$

$$\mathcal{A}(\eta' \rightarrow \pi^+ \pi^- a)|_{\text{LO}} = \frac{m_\pi^2}{f_\pi^2} (\sqrt{2} \cos \theta + \sin \theta) \left[\frac{f_\pi}{\sqrt{3} f_a} \frac{Q_u m_u + Q_d m_d}{m_u + m_d} - \frac{1}{3\sqrt{3}} \frac{m_d - m_u}{m_u + m_d} \langle \pi^0 a \rangle + \frac{1}{3} \theta_{\eta_8 a} + \frac{\sqrt{2}}{3} \theta_{\eta_0 a} \right],$$

Effects of pion-pion final-state interactions (FSI)

- Unitarity:



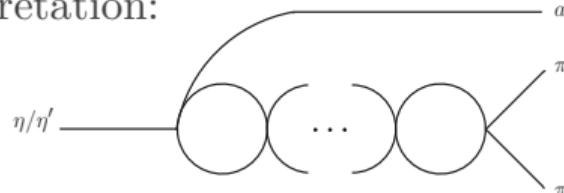
$$\text{disc}\mathcal{A}(s) = 2i\mathcal{A}(s)\sigma_\pi(s)T_0^{0*}(s) = 2i\mathcal{A}(s)\sin\delta_0^0(s)e^{-i\delta_0^0(s)},$$

$$\mathcal{A}(s) = \frac{1}{2i\pi} \int_{4M_\pi^2}^\infty ds' \frac{\text{disc}\mathcal{A}(s')}{s' - s - i\varepsilon},$$

- Analytic solution:

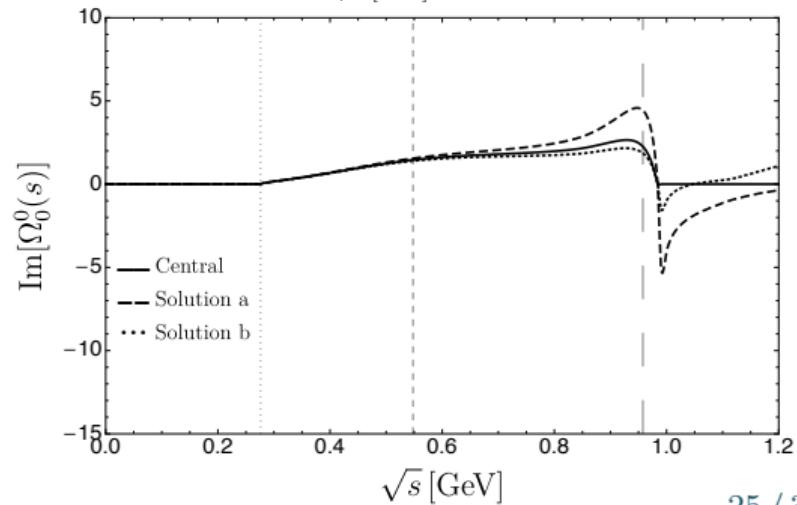
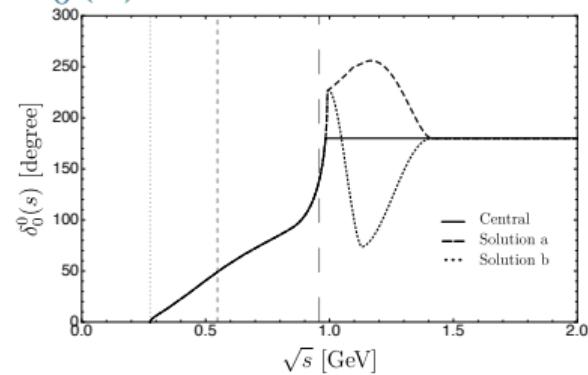
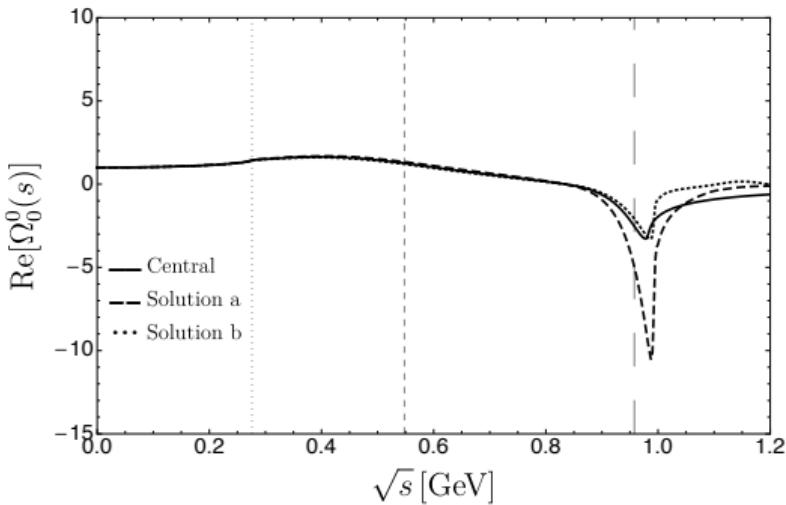
$$\mathcal{A}(s) = \mathcal{A}(\eta \rightarrow 2\pi a)|_{\text{LO}} \times \Omega_0^0(s), \quad \Omega_0^0(s) = \exp \left\{ \frac{s}{\pi} \int_{4M_\pi^2}^\infty ds' \frac{\delta_0^0(s')}{s'(s' - s - i\varepsilon)} \right\},$$

- Diagrammatic interpretation:



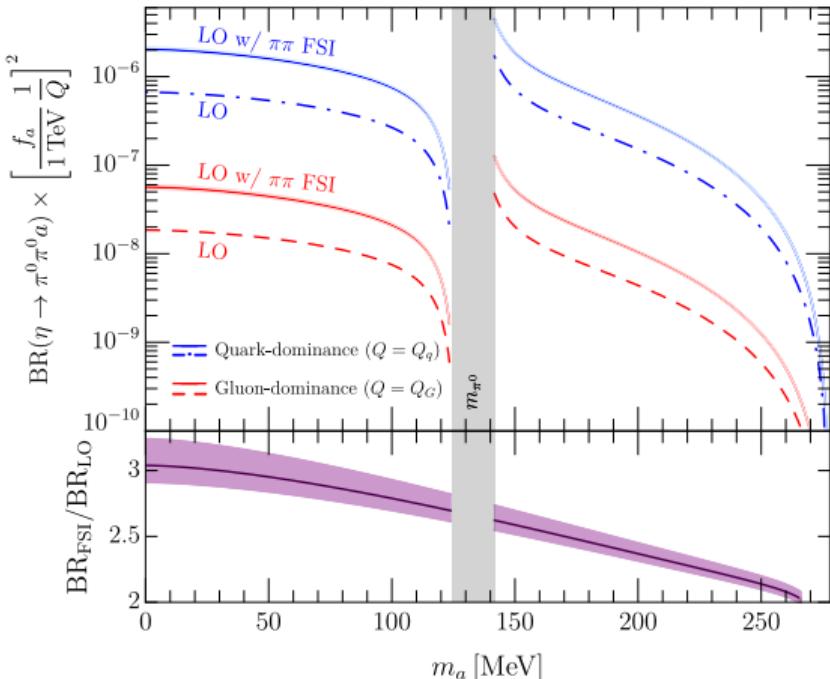
Solution of the Omnès function $\Omega_0^0(s)$

$$\Omega_0^0(s) = \exp \left\{ \frac{s}{\pi} \int_{4M_\pi^2}^\infty ds' \frac{\delta_0^0(s')}{s'(s' - s - i\varepsilon)} \right\},$$

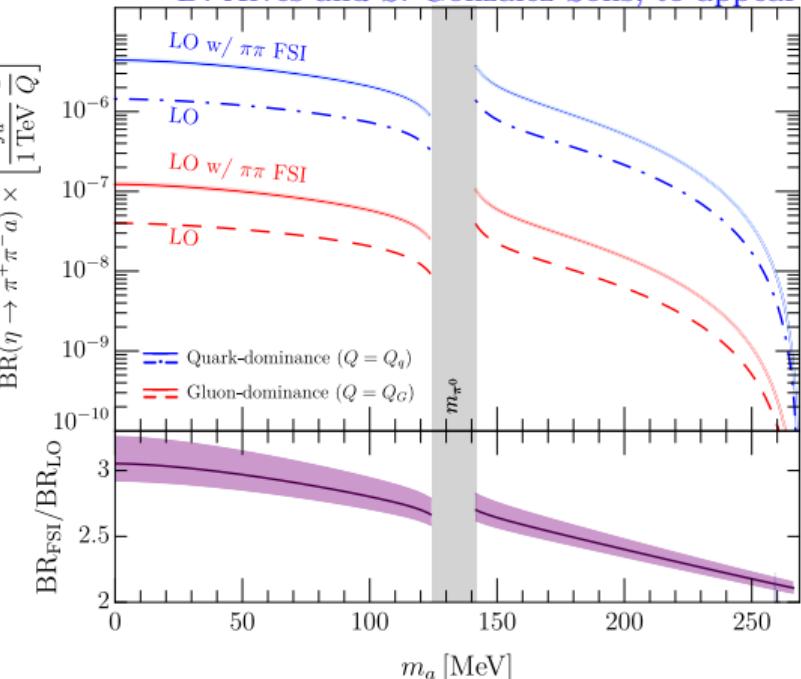


Branching ratio predictions for $\eta \rightarrow \pi\pi a$

- Two scenarios: **Quark-dominance** $Q_G = 0$ or **Gluon-dominance** $Q_q = 0$



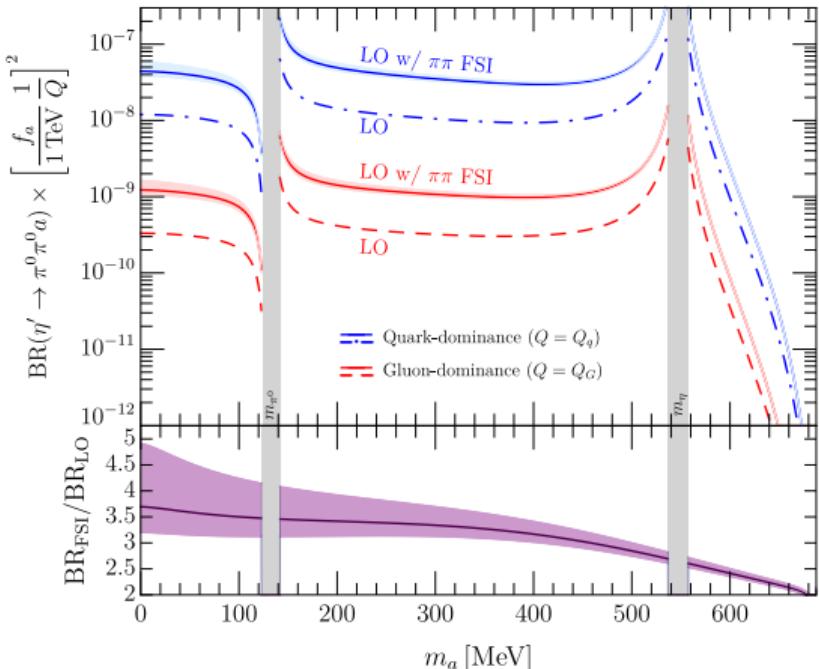
D. Alves and S. González-Solís, to appear



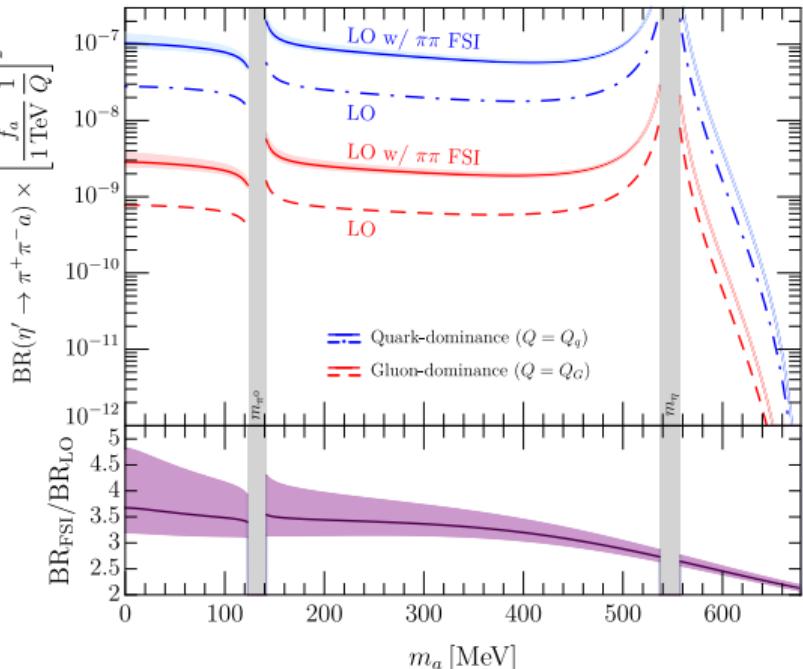
- We encourage searches in $\eta \rightarrow \pi\pi a \rightarrow \pi\pi\gamma\gamma$ and $\eta \rightarrow \pi\pi a \rightarrow \pi\pi\ell^+\ell^-$ (BESIII, KLOE, CMS, JEF, REDTOP)

Branching ratio predictions for $\eta' \rightarrow \pi\pi a$

- Two scenarios: **Quark-dominance** $Q_G = 0$ or **Gluon-dominance** $Q_q = 0$



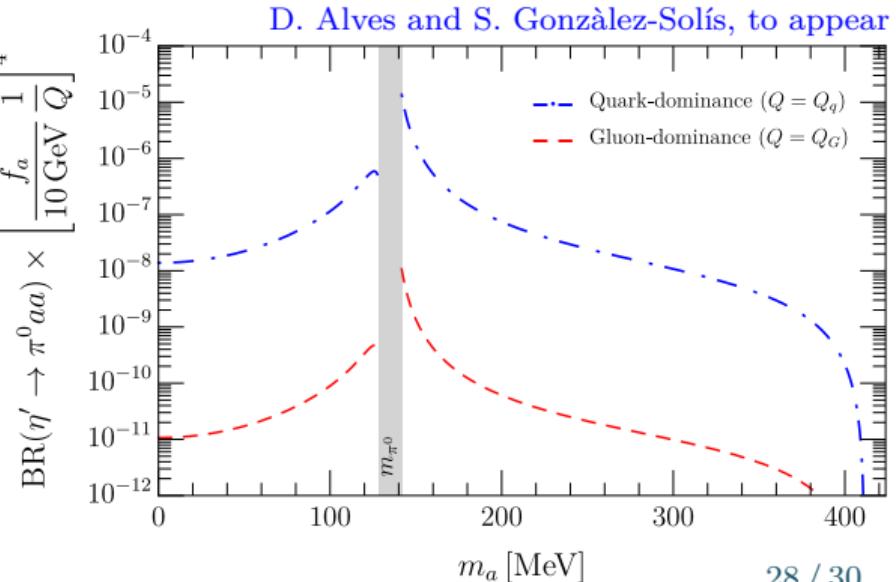
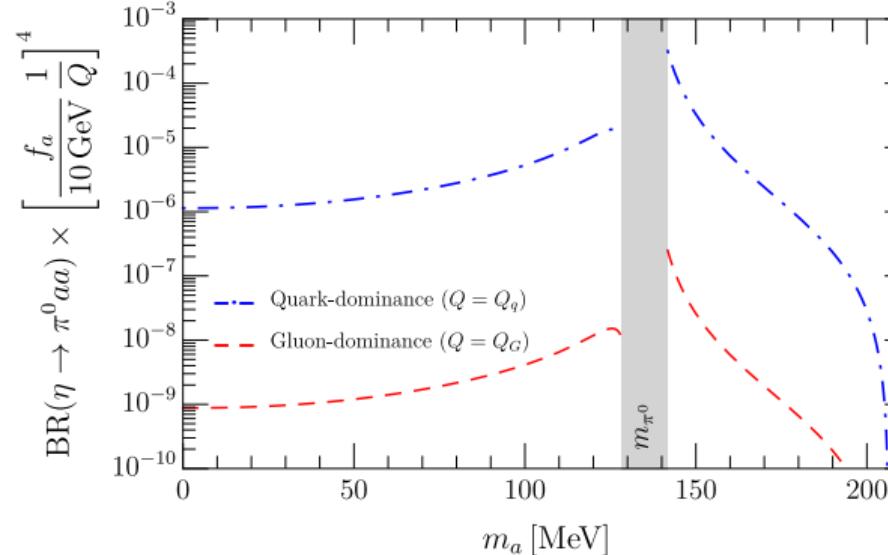
D. Alves and S. González-Solís, to appear



- We encourage searches in $\eta' \rightarrow \pi\pi a \rightarrow \pi\pi\gamma\gamma$ and $\eta' \rightarrow \pi\pi a \rightarrow \pi\pi\ell^+\ell^-$ (BESIII, KLOE, CMS, JEF, REDTOP)

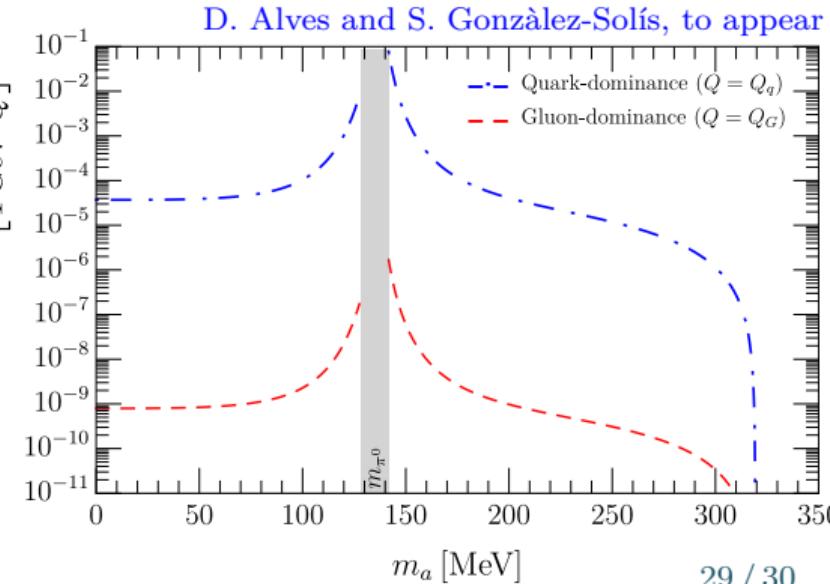
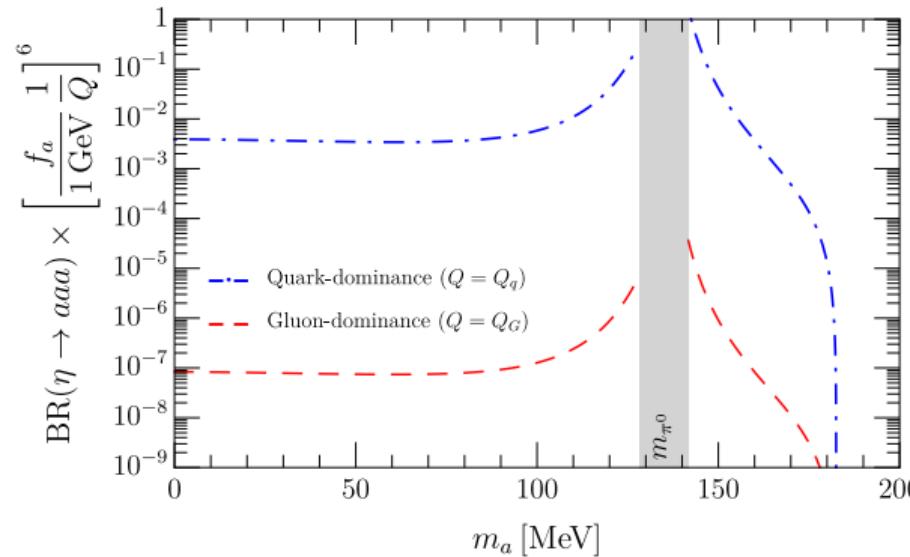
Double production of ALPs in η/η' decays

- $\eta/\eta' \rightarrow \pi^0 aa$ decays
- One extra power of $1/f_a$ suppression, $\text{BR} \sim \mathcal{O}(1/f_a^4)$
- $f_a \sim \mathcal{O}(1 - 10)$ GeV to be sensitive probes of ALPs



Triple production of ALPs in η/η' decays

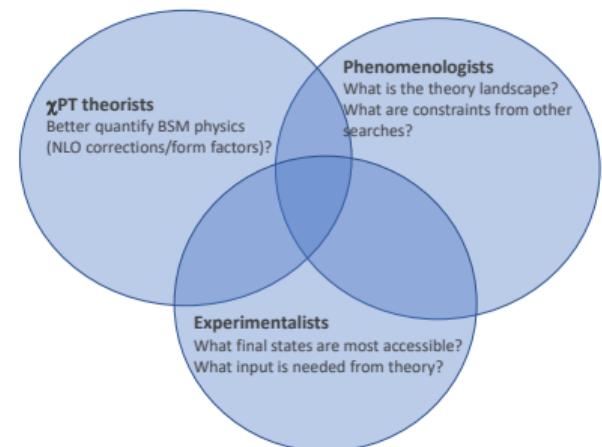
- $\eta/\eta' \rightarrow aaa$ decays
- $\text{BR} \sim \mathcal{O}(1/f_a^6)$
- $f_a \sim \mathcal{O}(1) \text{ GeV}$ to be sensitive probes of ALPs



D. Alves and S. González-Solís, to appear

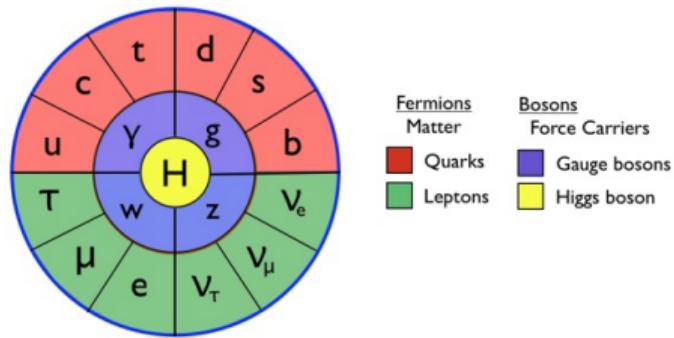
Conclusions

- Exploring **dark sectors** is an important and growing element of BSM physics
- A wealth of exciting ongoing/future **experiments** to search for dark sector particle signatures exist/planned
- η/η' mesons are an interesting place to look for dark particles because probe coupling to light quarks and gluons
- BSM searches in parallel with SM η/η' decay studies
- Progress on this front requires **collaboration!**



The Standard Model

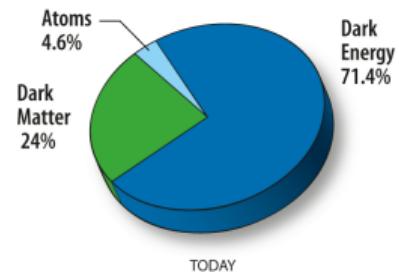
- Matter particles (quarks and leptons) in three families and mediator particles (bosons) of three interactions: electromagnetic, strong and weak
- Provides a consistent **description** of Nature's fundamental constituents and their interactions
- **Predictions** tested and confirmed by numerous experiments
- Experimental **completion** in 2012 (Higgs discovery)



Particles of the Standard Model

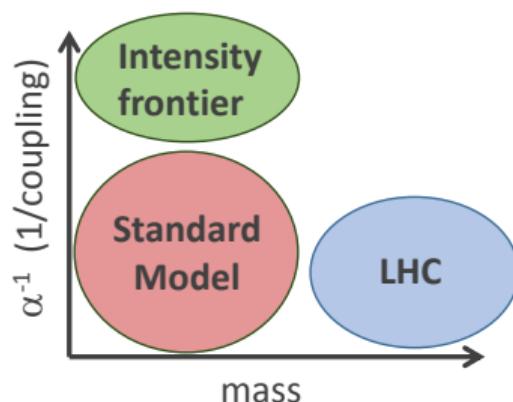
Beyond the Standard Model

- However, the SM **fails** to explain several observed phenomena in particle physics, astrophysics and cosmology:
 - **Dark matter:** what is the most prevalent kind of matter in our Universe?
 - **Dark Energy:** what drives the accelerated expansion of the Universe?
 - **Neutrino** masses and oscillations: why do neutrinos have mass? what makes neutrinos disappear and then re-appear in a different form?
 - **Baryon asymmetry** of the Universe: what mechanism created the tiny matter-antimatter imbalance in the early Universe?
 - Several **anomalies in data:** $(g - 2)_\mu$, B -physics anomalies, KOTO anomaly ($K_L \rightarrow \pi^0 \nu \bar{\nu}$), ${}^8\text{Be}$ excited decay, ...



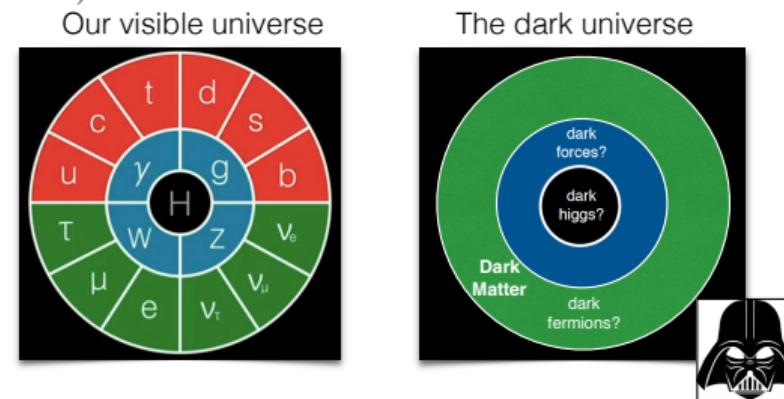
Energy and Intensity Frontier Research

- New Physics would be needed to explain observed phenomena
- Why have **not** new particles yet been observed?
 - Hypothetical new particles are **heavy** and require even higher collision energy to be observed \Rightarrow **Energy Frontier** research (LHC@CERN, Tevatron@FermiLab)
 - Another possibility is that our inability to observe new particles lies not in their heavy mass, but rather in their extremely **feeble interactions** \Rightarrow **Intensity Frontier** research
- We don't know in which **direction** BSM physics might be



Dark sector physics

- Why a dark sector?
 - Many open problem in particle physics, *e.g.* dark matter, neutrino mass generation or anomalies in data, let us think about dark particles
- What is a dark sector particle?
 - Any particle that does not interact through the SM forces (not charged under the SM symmetries)

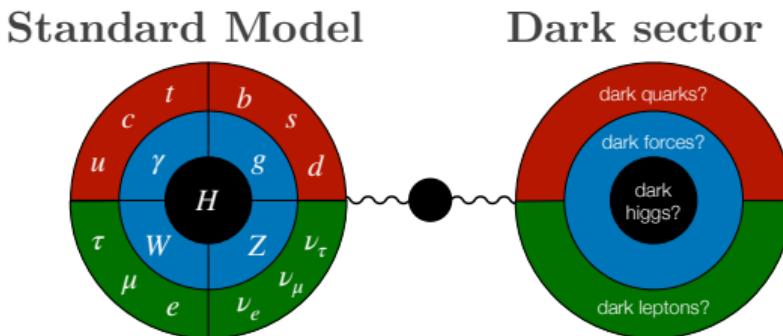


- We live in the SM world, how can we **access** (and test) the **dark sector**?

Dark sector portals to the Standard Model

We live in the Standard Model world, how can we access/test the **dark sector**?

⇒ **Portal** interactions with the SM, only a few are allowed by the SM symmetries



Portal

Vector

Scalar

Neutrino

Axion

Mediators

Dark photon

Dark scalar

Sterile Neutrino

Axion

Portal interactions

$$\epsilon B^{\mu\nu} A'_{\mu\nu}$$

$$\kappa |H|^2 |S|^2$$

$$y H L N$$

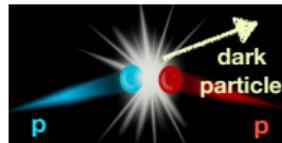
$$\frac{a}{f_a} \tilde{G}_{\mu\nu} G^{\mu\nu}$$

A broad program of searches of dark particles

- Vigorous effort of the community proposing new experiments & measurements

Energy frontier

LHC



Novel search strategies
are needed!

Flavor-factories

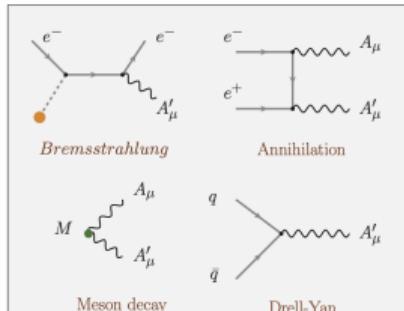
High-luminosity e^+e^- colliders



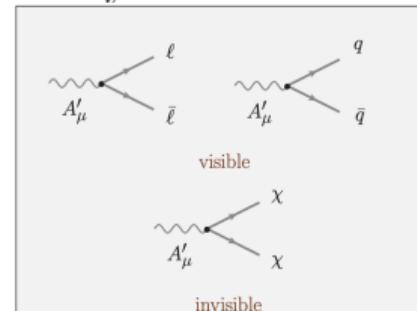
Unique access to dark
sectors!

- Plenty of dark particles can be produced from meson decays!!

Production modes



Decay modes



Large η/η' samples at future facilities

- Previous experiments:

Experiment	Total η	Total η'
CB at AGS	10^7	-
CB MAMI-B	2×10^7	-
CB MAMI-C	6×10^7	10^6
WASA-COSY	$\sim 3 \times 10^7$ (p+d), $\sim 5 \times 10^8$ (p+p)	-
KLOE-II	3×10^8	5×10^5
BESIII	$\sim 10^7$	$\sim 5 \times 10^7$

- Upcoming experiments:
 - Jefferson Lab Eta Factory Experiment (JEF) at JLab Hall D (approved)
 - 6.5×10^7 η and 4.9×10^7 η' per 100 days
 - Rare Eta Decays with a TPC for Optical Photons (REDTOP) possibly at Fermilab/GSI (proposed)
 - 10^{13} η and 10^{11} η' per year

Selected η/η' decays

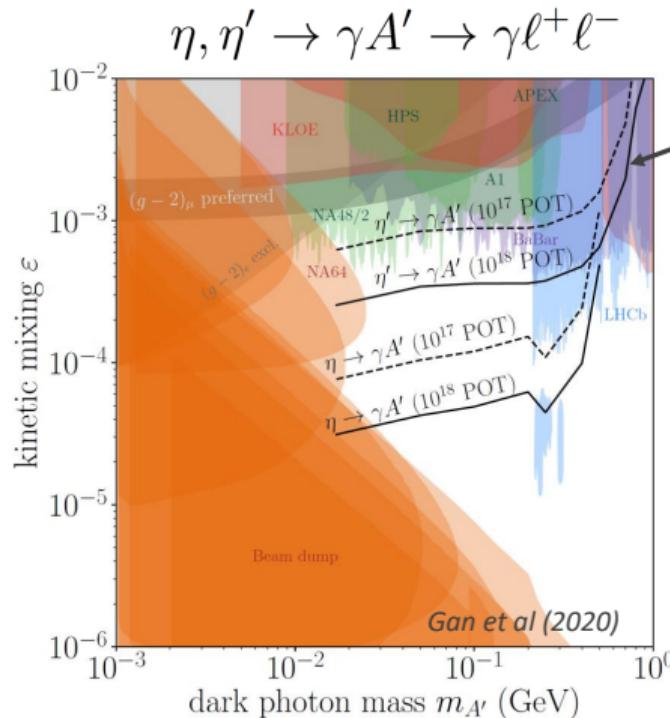
BSM particle	Decay mode	Signal channel	Search strategy
Dark photon (A')	$\eta/\eta' \rightarrow \gamma^{(*)} A'$	$A' \rightarrow \ell^+ \ell^-$ $A' \rightarrow \pi^+ \pi^-$	Bump-hunt in $d\Gamma/dm_{\ell\ell}$ Bump-hunt in $d\Gamma/dm_{\pi\pi}$
Leptophobic boson (B)	$\eta \rightarrow \gamma B$	$B \rightarrow \gamma \pi^0$	Enhancement in $m_{\pi^0 \gamma}$
		$B \rightarrow \pi^+ \pi^-$	Isospin suppressed
	$\eta' \rightarrow \gamma B$	$B \rightarrow \gamma \pi^0, \pi^+ \pi^-, \pi^+ \pi^- \pi^0, \gamma \eta$	Enhancement in $m_{\pi^0 \gamma}$
ALPs (a)	$\eta \rightarrow \pi \pi a$	$a \rightarrow \gamma \gamma, \ell^+ \ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma}$
	$\eta' \rightarrow \pi \pi a$	$a \rightarrow \gamma \gamma, \ell^+ \ell^-, \pi^+ \pi^- \gamma, 3\pi$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma}$
	$\eta^{(\prime)} \rightarrow \ell^+ \ell^-$		$\eta^{(\prime)}-a$ mixing
Scalar boson (S)	$\eta/\eta' \rightarrow \pi^0 S$	$S \rightarrow \gamma \gamma, \ell^+ \ell^-, \pi \pi$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma}$
	$\eta' \rightarrow \eta S$	$S \rightarrow \gamma \gamma, \ell^+ \ell^-, \pi \pi$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma}$

Other meson decays

BSM particle	Decay mode	Signal channel	Search strategy
ALPs (a)	$K^\pm \rightarrow \pi^\pm a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$K^\pm \rightarrow \pi^\pm \pi^0 a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$K_L \rightarrow \pi^0 a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$K_L \rightarrow \pi^0 \pi^0 a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$K_L \rightarrow \pi^+ \pi^- a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$B^\pm \rightarrow \pi^\pm a$	$a \rightarrow \ell^+\ell^-, 3\pi, \eta\pi\pi, KK\pi$	Higher ALP masses
	$B^\pm \rightarrow K^\pm a$	$a \rightarrow \ell^+\ell^-, 3\pi, \eta\pi\pi, KK\pi$	Higher ALP masses
	$B \rightarrow K^* a$	$a \rightarrow \ell^+\ell^-, 3\pi, \eta\pi\pi, KK\pi$	Higher ALP masses
	$\omega/\phi/J/\psi \rightarrow \pi^0 \pi^0 a$	$a \rightarrow \gamma\gamma, \ell^+\ell^- (\ell = e, \mu)$	Bump-hunt in $d\Gamma/dm_{\gamma\gamma, \ell\ell}$
	$\omega/\phi/J/\psi \rightarrow \pi^0 \pi^0 a$	$a \rightarrow \pi^+ \pi^- \gamma, 3\pi$	
Dark photon (A')	$\pi^0 \rightarrow \gamma A'$	$A' \rightarrow e^+e^-$	e^+e^- resonance
	$\pi^0 \rightarrow \gamma^* A'$	$\gamma^* \rightarrow e^+e^-, A' \rightarrow e^+e^-$	e^+e^- resonance
	$\omega/\phi/J/\psi \rightarrow \pi^0 A'$	$A' \rightarrow \ell^+\ell^- (\ell = e, \mu)$	$\ell^+\ell^-$ resonance
	$\omega/\phi/J/\psi \rightarrow \pi^0 A'$	$A' \rightarrow \pi^+ \pi^-$	$\pi^+ \pi^-$ resonance
Leptophobic boson (B)	$\omega/\phi \rightarrow \eta B$	$B \rightarrow \gamma\pi^0$	Enhancement in $m_{\pi^0\gamma}$

Dark photon searches

- Broad worldwide effort to search for dark photons (A')
- Most searches are for A' coupling to leptons, *i.e.* in $A' \rightarrow \ell^+ \ell^- (\ell = e, \mu)$



REDTOP sensitivities projected for
FNAL/BNL (10^{18}) or CERN (10^{17}) POT

Gatto (2019)

Many other experiments targeting
same dark photon parameter space

Worthwhile to also consider

$$\eta' \rightarrow \pi^+ \pi^- A' \rightarrow \pi^+ \pi^- \ell^+ \ell^-$$

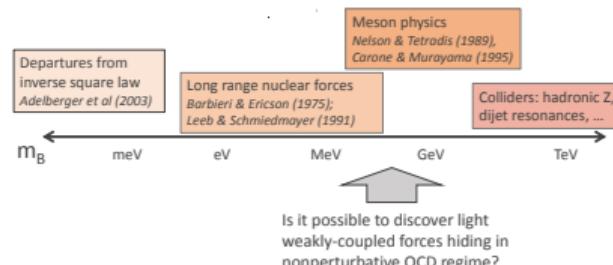
since $\mathcal{B}(\eta' \rightarrow \pi^+ \pi^- \gamma) \approx 10 \times \mathcal{B}(\eta' \rightarrow \gamma \gamma)$

Searches of a leptophobic dark photon in rare $\eta^{(\prime)}$ decays

- What if a **new force** couples mainly to quarks?
- **Simplest model:** gauge boson (B) coupled to baryon number

$$\mathcal{L}_{\text{int}} = \frac{1}{3} \mathbf{g}_B \bar{q} \gamma^\mu q B_\mu ,$$

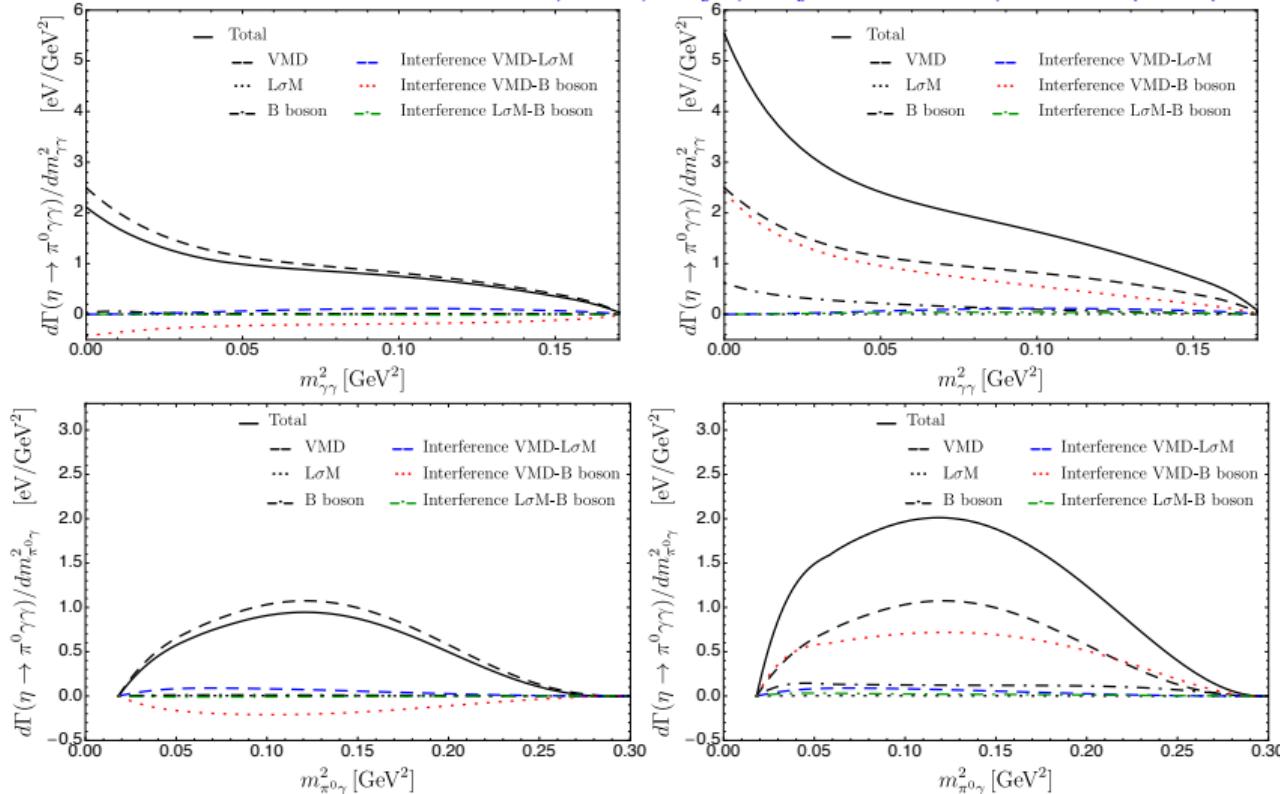
- \mathbf{g}_B : flavor-universal coupling to all quarks, preserves QCD symmetries (C, P, T)
- Also known as: B boson, leptophobic Z' or baryonic photon γ_B
- Discovery signals depend on the mass m_B and decay channels



- Searches in meson factories are gaining attention
 - $\eta \rightarrow \gamma B \rightarrow \gamma\gamma\pi^0$ (JEF), $\phi \rightarrow \eta B \rightarrow \eta\pi^0\gamma$ (KLOE-II), $\eta \rightarrow B\gamma \rightarrow \pi^+\pi^-\gamma$ (Belle-II)

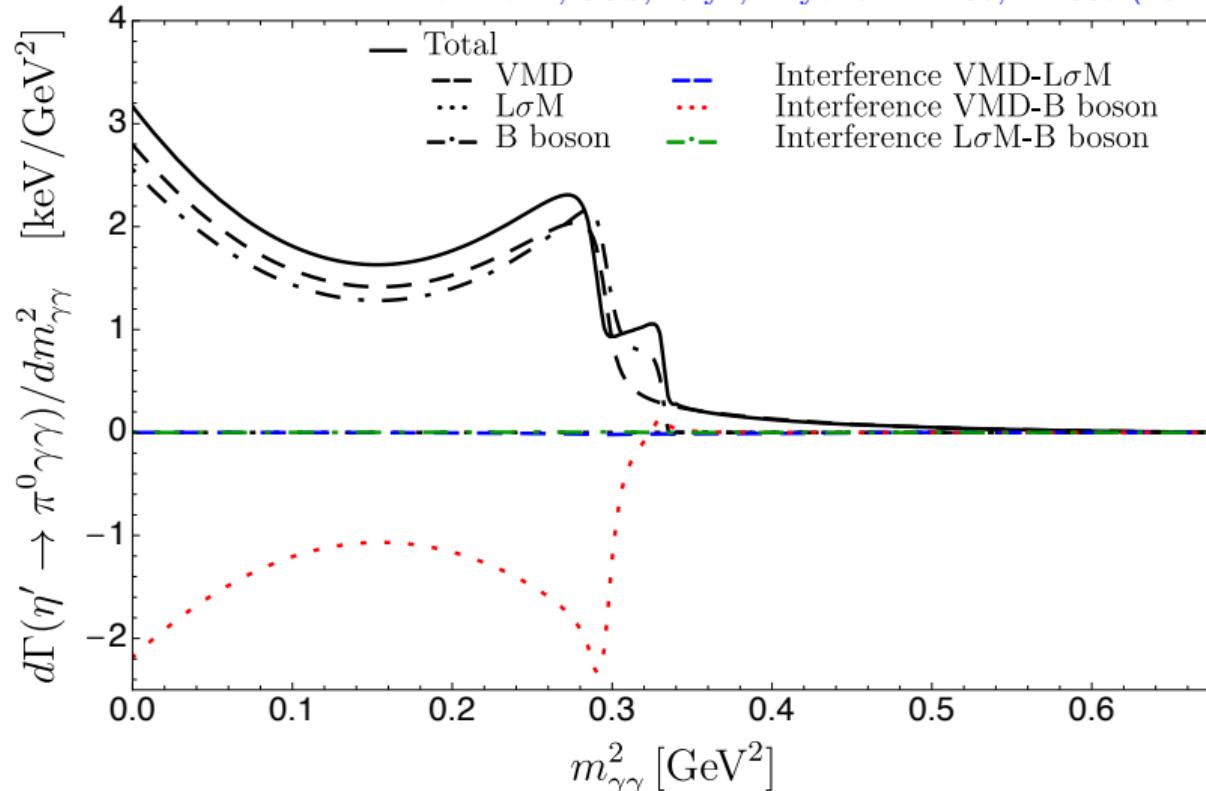
Individual contributions

Escribano, SGS, Royo, Phys.Rev.D 106, 114007 (2022)

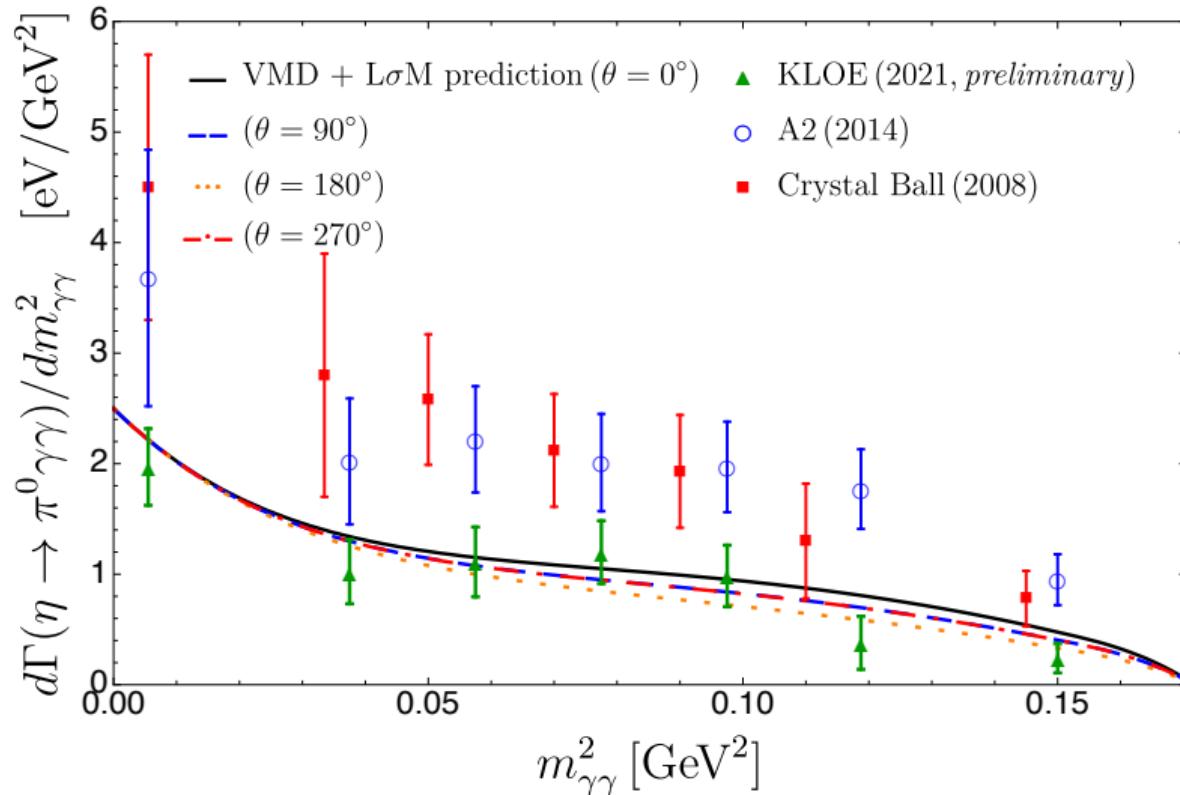


Individual contributions

Escribano, SGS, Royo, Phys.Rev.D 106, 114007 (2022)



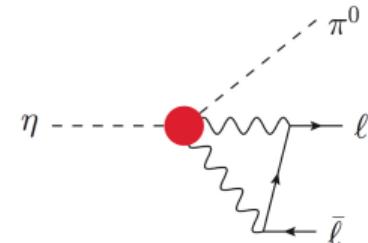
Interference phase between VMD and L σ M



$\eta^{(\prime)} \rightarrow \{\pi^0, \eta\} \ell^+ \ell^-$ decays ($\ell = e, \mu$)

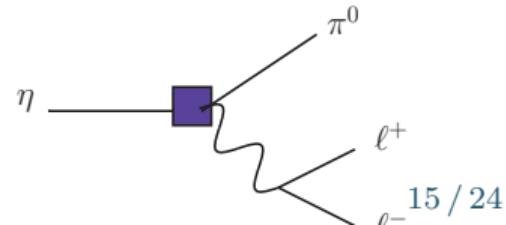
- In the SM:

- $\eta \rightarrow \pi^0 \gamma^* \rightarrow \pi^0 \ell^+ \ell^-$ forbidden by C and CP
- $\eta \rightarrow \pi^0 \ell^+ \ell^-$ proceed via C -conserving two-photon intermediate state



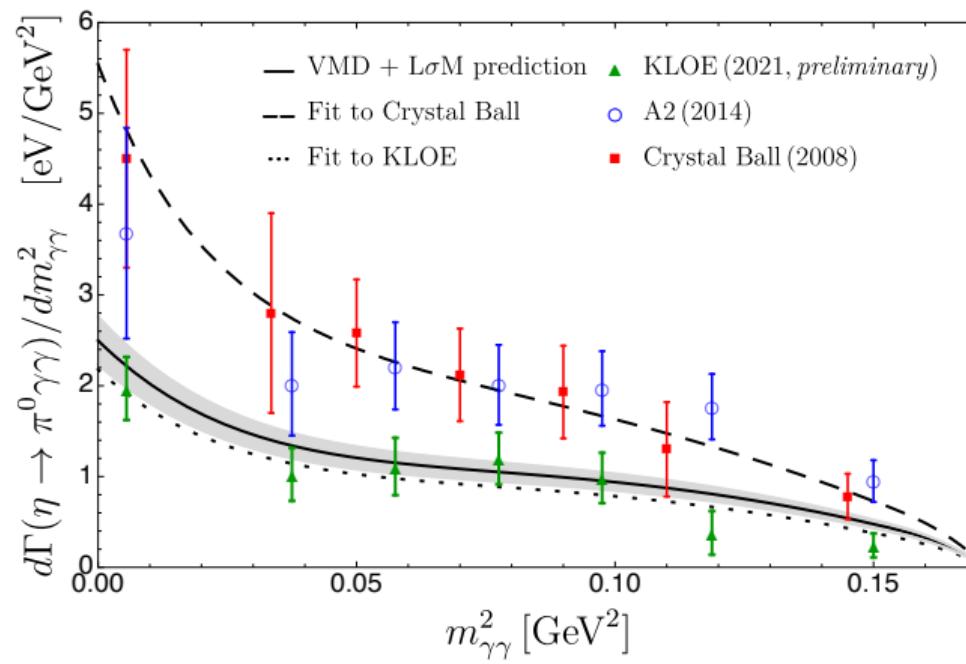
Decay channel	BR_{th} (Escribano&Royo 2007.12467)	BR_{exp} (pdg)
$\eta \rightarrow \pi^0 e^+ e^-$	$2.1(1)(2) \times 10^{-9}$	$< 7.5 \times 10^{-6}$ (CL=90%)
$\eta \rightarrow \pi^0 \mu^+ \mu^-$	$1.2(1)(1) \times 10^{-9}$	$< 5 \times 10^{-6}$ (CL=90%)
$\eta' \rightarrow \pi^0 e^+ e^-$	$4.6(3)(7) \times 10^{-9}$	$< 1.4 \times 10^{-3}$ (CL=90%)
$\eta' \rightarrow \pi^0 \mu^+ \mu^-$	$1.8(1)(2) \times 10^{-9}$	$< 6.0 \times 10^{-5}$ (CL=90%)
$\eta' \rightarrow \eta e^+ e^-$	$3.9(3)(4) \times 10^{-10}$	$< 2.4 \times 10^{-3}$ (CL=90%)
$\eta' \rightarrow \eta \mu^+ \mu^-$	$1.6(1)(2) \times 10^{-10}$	$< 1.5 \times 10^{-5}$ (CL=90%)

- Background for BSM searches, *e.g.* C -violating virtual photon exchange or new scalar mediators
- REDTOP can improve the experimental state



Fits to the $\eta \rightarrow \pi^0 \gamma\gamma$ decays

- Crystal Ball: $\alpha_B = 0.40^{+0.07}_{-0.08}$, $m_B = 583^{+32}_{-20}$ MeV, $\chi^2_{\text{dof}} = 0.4/5 = 0.1$
- KLOE: $\alpha_B = 0.049^{+40}_{-27}$, $m_B = 135^{+1}_{-135}$ MeV, $\chi^2_{\text{dof}} = 4.5/5 = 0.9$
- signatures outside $m_{\pi^0} \lesssim m_B \lesssim m_\eta$ may be visible

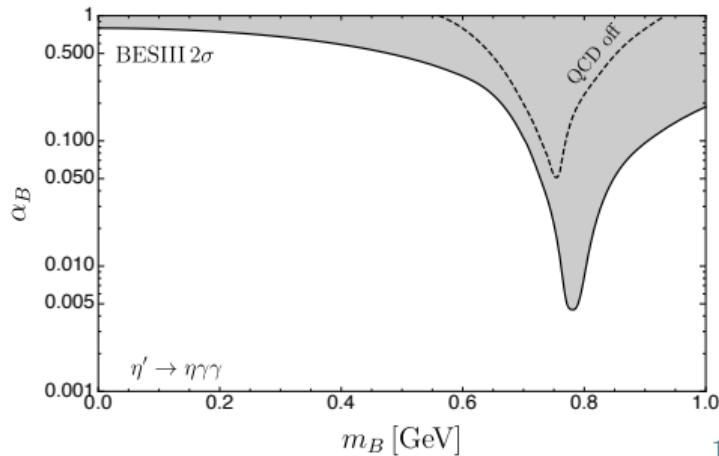
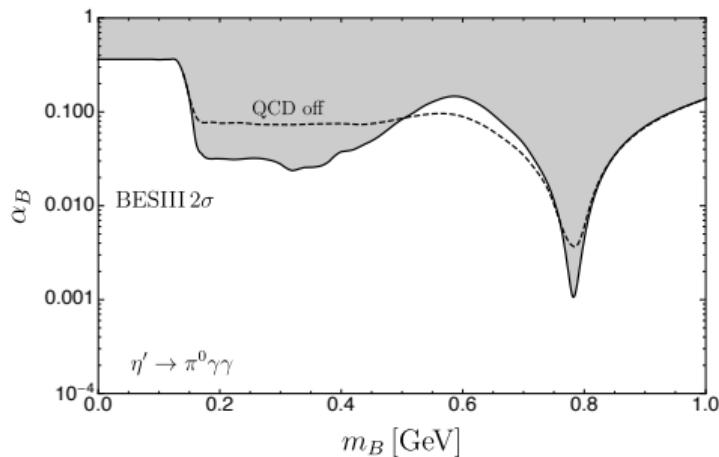


Conclusions

- Exploring **dark sectors** is an important and growing element of BSM physics
- A wealth of exciting ongoing **experiments** exist
- **Meson decays** offer a unique opportunity to look for New Physics
- Within the VMD and $L\sigma M$ frameworks **we have described**
 - $\eta \rightarrow \pi^0 \gamma\gamma$: the situation is **not conclusive**
$$BR = 1.35(8) \times 10^{-4} \left\{ \begin{array}{ll} \sim 1/2 \text{ of } BR = 2.54(27) \times 10^{-4} & (\text{A2, 2014}) \\ \sim 1.6\sigma \text{ from } BR = 2.21(24)(47) \times 10^{-4} & (\text{CB, 2008}) \\ \text{agrees with } BR = 1.23(14) \times 10^{-4} & (\text{KLOE prel., 2022}) \end{array} \right.$$
 - $\eta' \rightarrow \pi^0(\eta)\gamma\gamma$: **in line** with BESIII data
- **Constraints** on α_B, m_B have been strengthened by one order of magnitude from $\eta \rightarrow \pi^0 \gamma\gamma$
- We have tested ALPs with $\eta/\eta' \rightarrow \pi\pi a$ decays
 - We encourage searches in $\eta/\eta' \rightarrow \pi\pi a \rightarrow \pi\pi\gamma\gamma, \pi\pi\ell^+\ell^-$ (BESIII, KLOE,CMS, REDTOP)

New limits on α_B and m_B

- Not assuming the **NWA**
- QCD contribution **on**
- $BR < BR_{\text{exp}}$ at 2σ
- $BR(\eta' \rightarrow \pi^0 \gamma \gamma)_{\text{exp}} = 3.20(7)(23) \times 10^{-3}$
M. Ablikim *et.al* [BESIII], Phys.Rev. D 96 (2017) 012005
- $BR(\eta' \rightarrow \eta \gamma \gamma)_{\text{exp}} = 8.25(3.41)(72) \times 10^{-5}$
M. Ablikim *et.al* [BESIII], Phys.Rev. D 100 (2019) 052015
- Sharp dip when $m_B \sim m_\omega$
- Bounds 4 orders of magnitude weaker than $\eta \rightarrow \pi^0 \gamma \gamma$



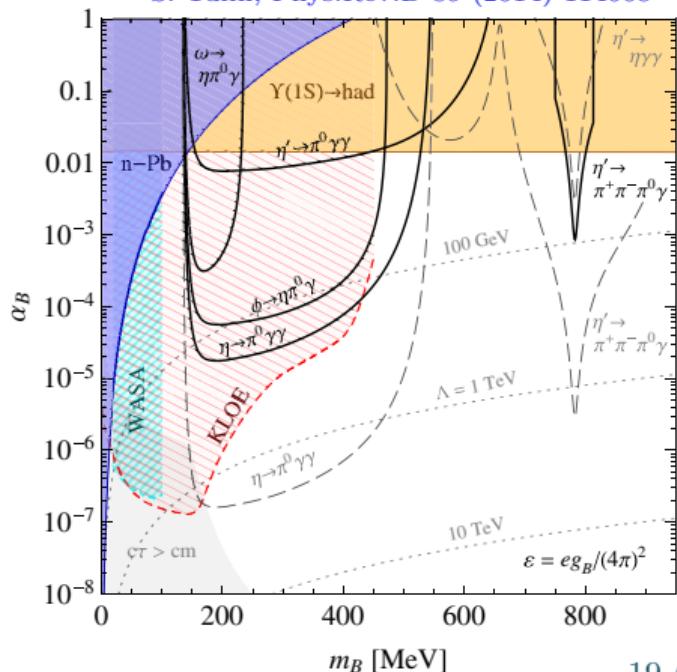
Previous limits on α_B and m_B

- Assuming the **Narrow-Width Approximation (NWA)**

$$BR(\eta \rightarrow \pi^0 \gamma \gamma) = BR(\eta \rightarrow B\gamma) \times BR(B \rightarrow \pi^0 \gamma),$$

S. Tulin, Phys.Rev.D 89 (2014) 114008

- QCD contribution off
- $BR(\eta \rightarrow \pi^0 \gamma \gamma) < BR_{\text{exp}}$ at 2σ
 - $BR(\eta \rightarrow \pi^0 \gamma \gamma)_{\text{exp}} = 2.21(53) \times 10^{-4}$
 - $BR(\eta' \rightarrow \pi^0 \gamma \gamma)_{\text{exp}} < 8 \times 10^{-4}$ (90% C.L.)
 - $BR(\eta' \rightarrow \eta \gamma \gamma)_{\text{exp}}$ no data



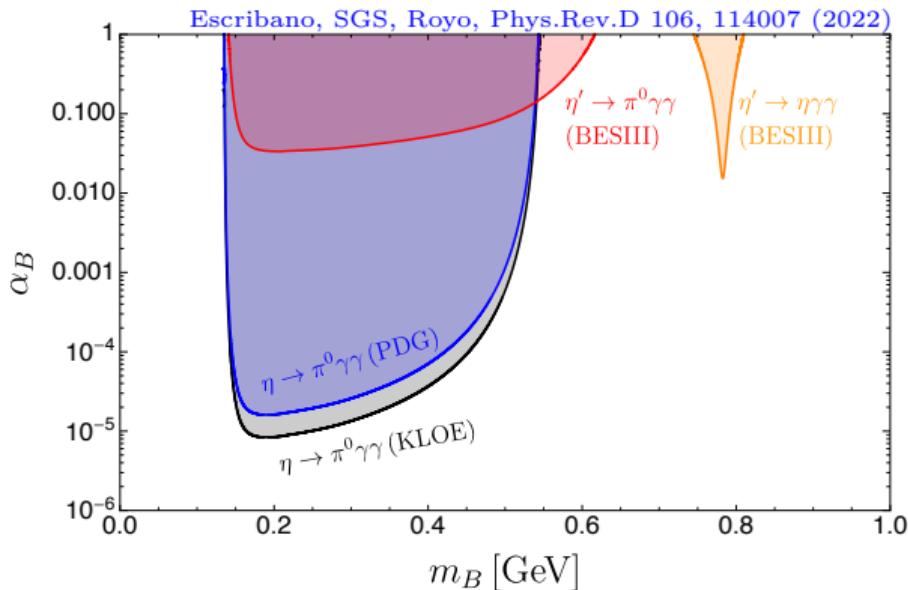
Present limits on α_B and m_B

- Assuming the **Narrow-Width** Approximation (NWA)

$$BR(\eta \rightarrow \pi^0 \gamma\gamma) = BR(\eta \rightarrow B\gamma) \times BR(B \rightarrow \pi^0 \gamma),$$

- QCD contribution off
- $BR(\eta \rightarrow \pi^0 \gamma\gamma) < BR_{\text{exp}}$ at 2σ

- $BR(\eta \rightarrow \pi^0 \gamma\gamma)_{\text{exp}}^{\text{pdg}} = 2.56(22) \times 10^{-4}$
- $BR(\eta \rightarrow \pi^0 \gamma\gamma)_{\text{exp}}^{\text{KLOE}} = 1.23(14) \times 10^{-4}$
B. Cao [KLOE], PoS EPS-HEP2021 (2022) 409
- $BR(\eta' \rightarrow \pi^0 \gamma\gamma)_{\text{exp}} = 3.20(7)(23) \times 10^{-3}$
M. Ablikim *et.al* [BESIII], Phys.Rev. D 96 (2017) 012005
- $BR(\eta' \rightarrow \eta \gamma\gamma)_{\text{exp}} = 8.25(3.41)(72) \times 10^{-5}$
M. Ablikim *et.al* [BESIII], Phys.Rev. D 100 (2019) 052015



Lagrangian for ALPs coupled to QCD

- “Derivative basis”: ALPs with gluon and derivative couplings

$$\begin{aligned}\mathcal{L}_{\text{ALP}} = & \mathcal{L}_{\text{QCD}} + \frac{1}{2} (\partial_\mu a) (\partial^\mu a) - \frac{1}{2} M_a^2 a^2 \\ & - \left(Q_G + \sum_{q=u,d,s} Q_q \right) \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{\partial_\mu a}{f_a} \sum_{q=u,d,s} \frac{Q_q}{2} \bar{q} \gamma^\mu \gamma^5 q ,\end{aligned}$$

M_a^2 : PQ contribution to the mass, f_a : axion decay constant, $Q_{q,G}$: PQ charges

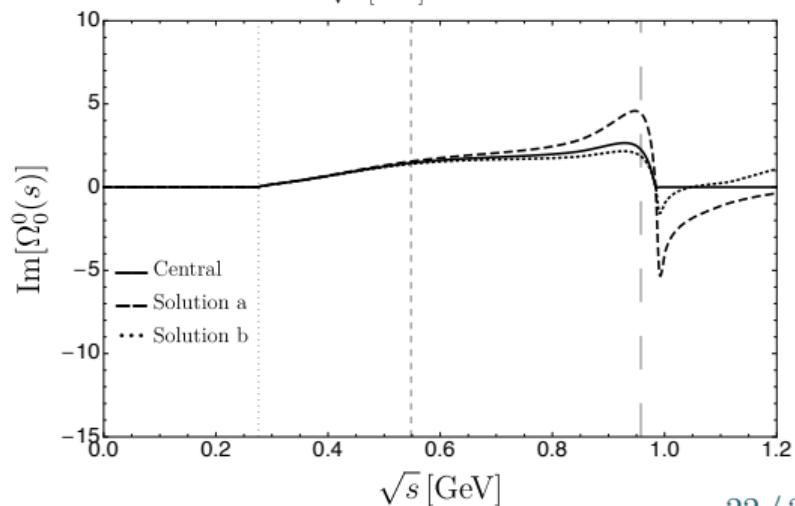
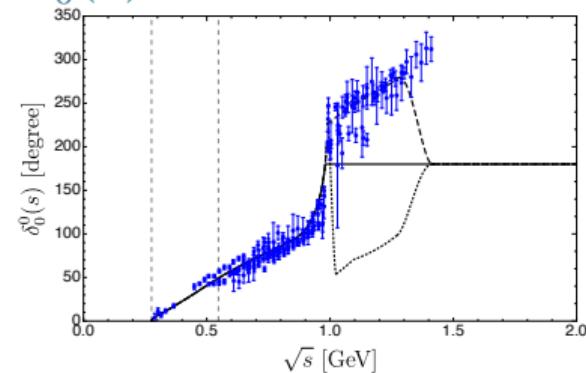
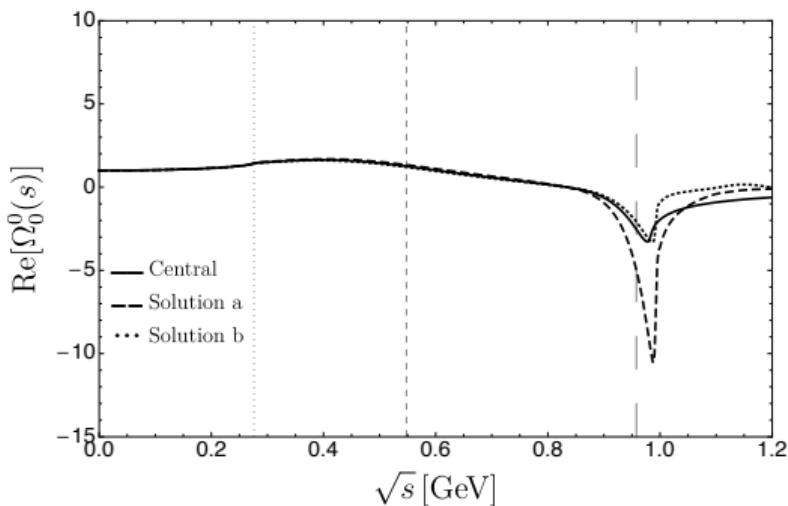
- “Yukawa basis” (this work, at GeV scale): ALP with gluon and mass couplings

$$\mathcal{L}_{\text{ALP}} = \mathcal{L}_{\text{QCD}} + \frac{1}{2} (\partial_\mu a) (\partial^\mu a) - \frac{1}{2} M_a^2 a^2 - Q_G \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu} \tilde{G}^{\mu\nu} + \sum_{q=u,d,s} m_q \bar{q} \left(e^{i Q_q \frac{a}{f_a} \gamma_5} \right) q ,$$

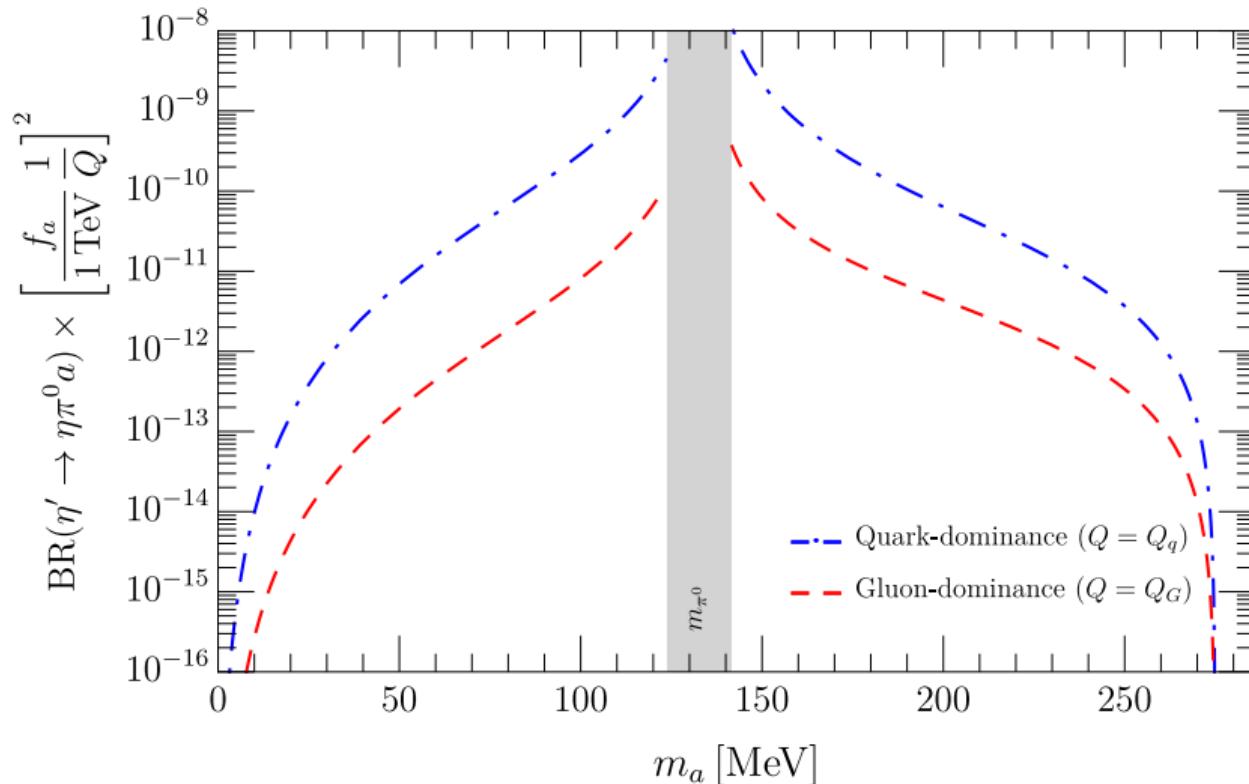
- Equivalent bases (related via chiral rotations of the quarks) if weak interactions are neglected
- The heavy-flavor c, b, t quarks contributions are absorbed in $Q_G \rightarrow Q_G + Q_{t,b,c}$

Solution of the Omnès function $\Omega_0^0(s)$

$$\Omega_0^0(s) = \exp \left\{ \frac{s}{\pi} \int_{4M_\pi^2}^\infty ds' \frac{\delta_0^0(s')}{s'(s' - s - i\varepsilon)} \right\},$$



$$\eta' \rightarrow \eta\pi^0 a$$



Multi production of ALPs

