

Λ 部分子光锥分布振幅的格点计算

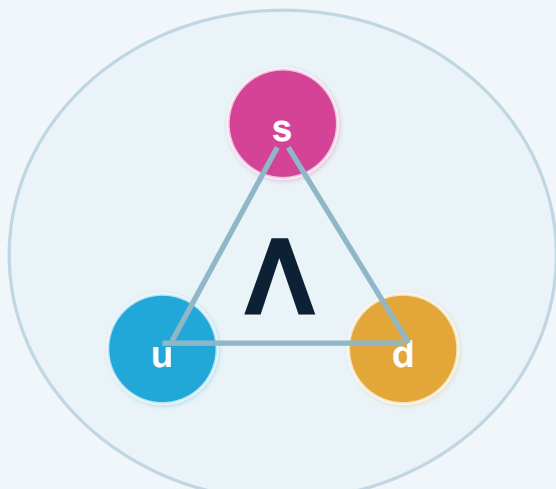
Complete leading-twist Λ -baryon LCDAs
from lattice QCD

超子-核子谱仪物理研讨会 南开大学 2026/08/16

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On behalf of the Lattice Parton Collaboration



为什么是 Λ ? 一个奇异重子实验室



uds

最轻奇异重子

isosinglet 轻夸克对 + strange valence quark

同一个强子同时连接
polarization · form factors · CP/EDM

$\Lambda \rightarrow p \pi^-$ 遍举弱衰变

$\Lambda \rightarrow p \pi^-$ 的角分布直接分析 Λ 极化, 并提供离散对称性检验。



奇异味结构标杆

最简单的 uds 系统之一, 可研究轻夸克对与 s 夸克之间的动量分配。



电磁结构

精确的 time-like form factors 正在提高对 Λ 内部结构的分辨能力。

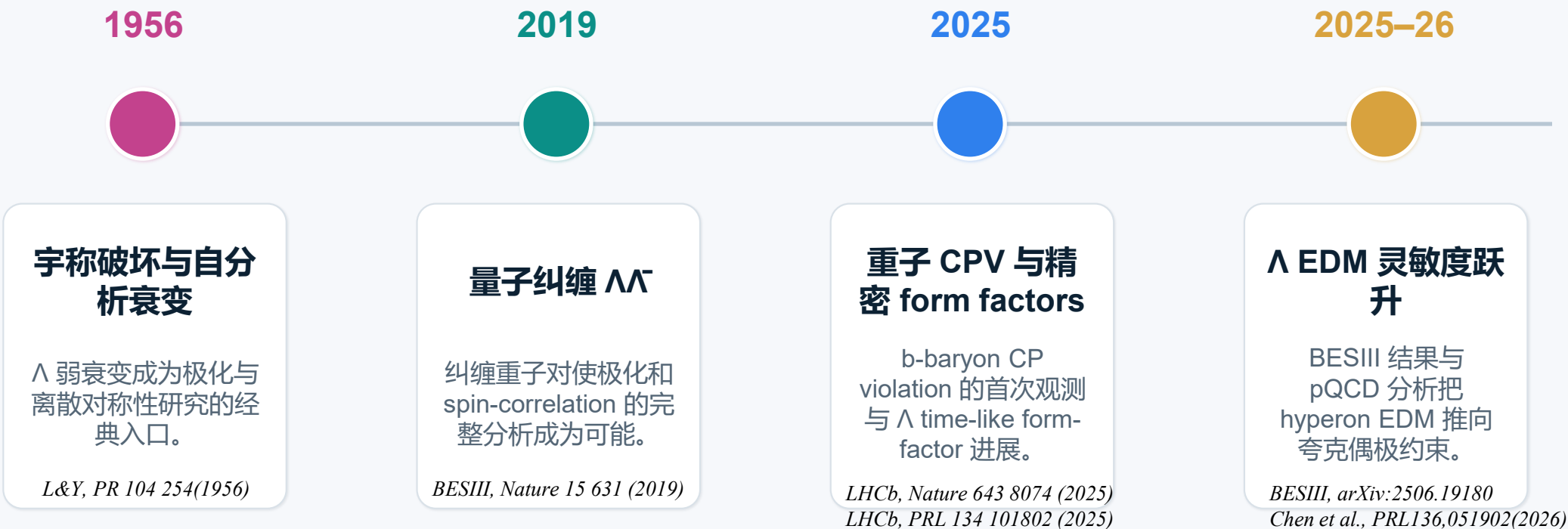


EDM / CP 探针

Λ EDM 对 CP-odd quark dipoles, 尤其 strange-quark CEDM, 提供互补灵敏度。

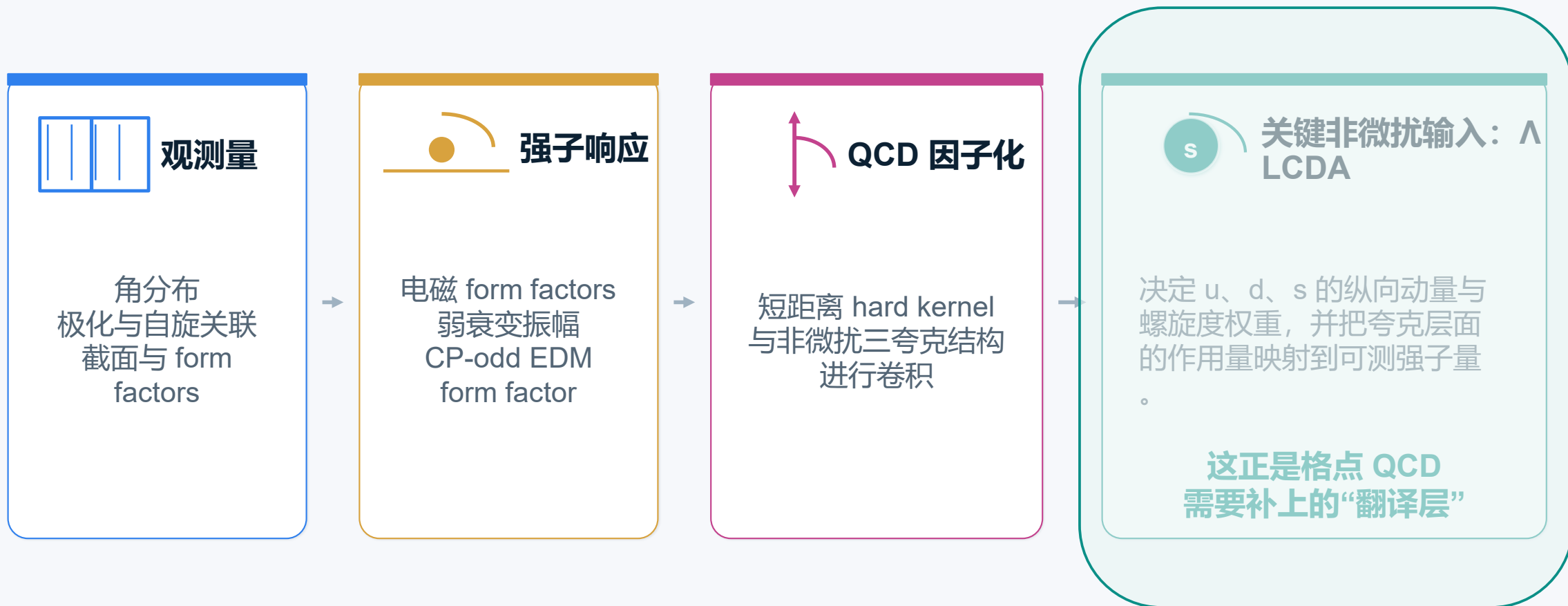
为什么是现在？超子精密物理进入新阶段

实验端的进展，正在把“缺少受控强子结构输入”变成主要瓶颈。



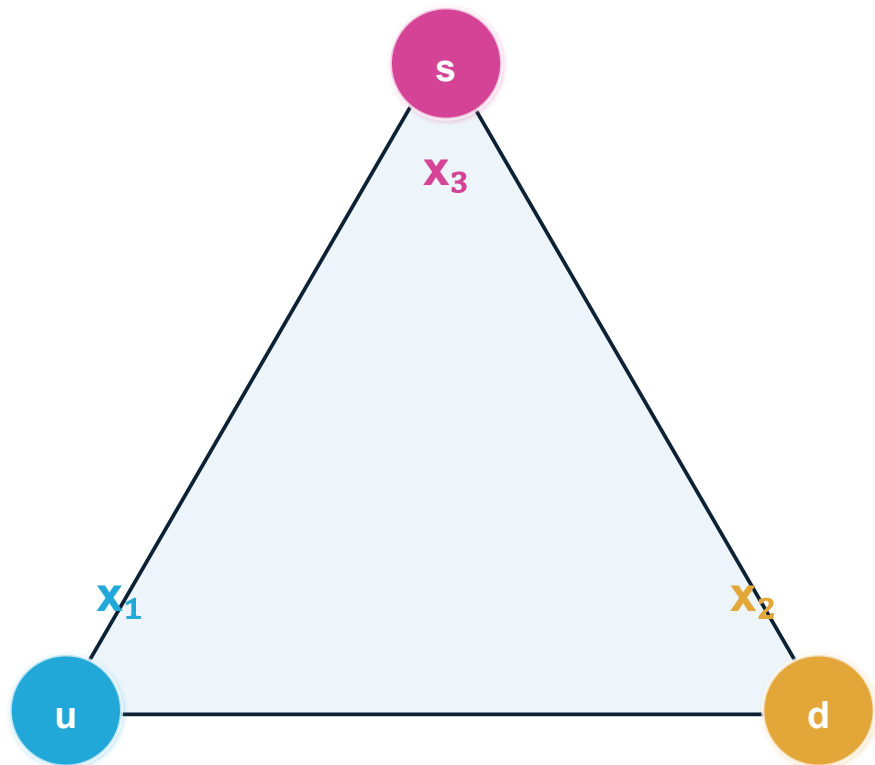
实验已经能更精细地问问题；理论需要提供完整、可卷积、带误差的 Λ 三夸克结构。

实验测什么？理论还缺什么？



LCDA 不是额外的模型参数；它是从夸克算符到谱仪观测量的非微扰传递函数。

LCDA: Λ 的三夸克纵向动量图像



$$x_1 + x_2 + x_3 = 1$$

物理支撑区域是一块二维三角形，而不是一条一维线

三种 leading-twist 结构

A

dominant ·
symmetric

V

antisymmetric

T

antisymmetric

V、A、T 同时编码动量分配、螺旋度与 flavor/permutation 结构。

● **完整二维函数** 分辨三夸克之间的相关性与内部峰形。

● **卷积输入**

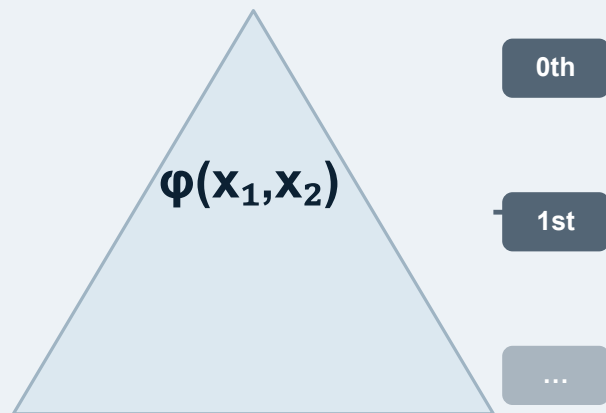
hard-exclusive observables 需要函数，而不只是几个数。

● **端点结构**

高阶 power corrections 和因子化可靠性与 $x_i \rightarrow 0$ 直接相关。

为什么必须直接求完整二维分布？

传统局域矩方法

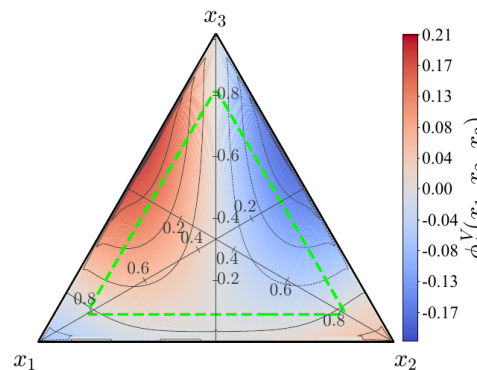


只得到少量积分约束
无法重建二维相关性与端点

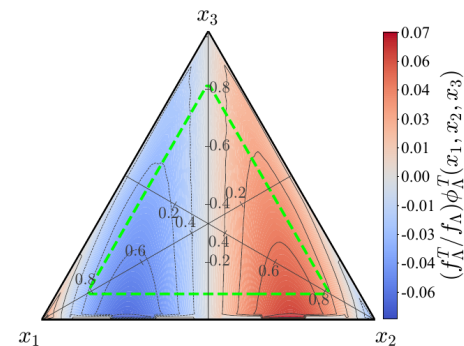
“几个数” ≠ “完整函数”

Λ 的特殊性：零矩为零，但函数并不为零

$$\int [dx] \phi_{\Lambda}^A = 1, \quad \int [dx] \phi_{\Lambda}^{V,T} = 0$$



V: sign-changing



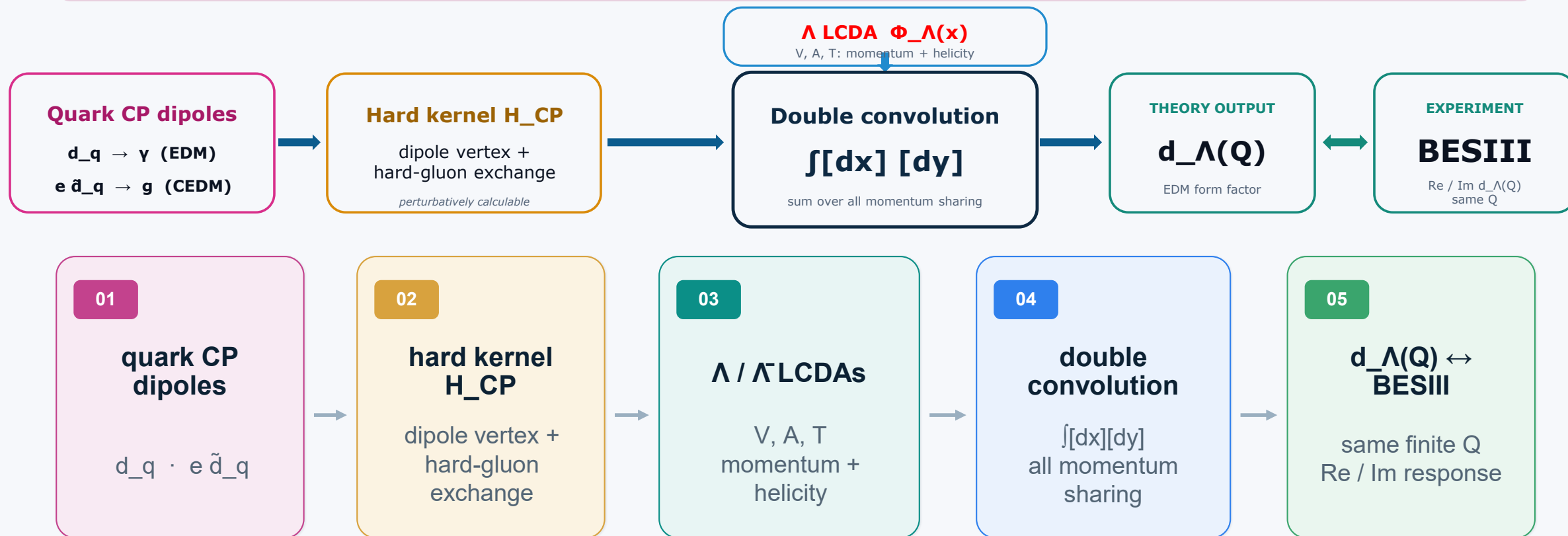
T: sign-changing

V、T 在 $x_1 = x_2$ 上为零，但在非局域二维平面上清晰非零；局域 decay constant 无法独立表征它们。

Λ LCDA 如何把夸克偶极“翻译”为 hyperon EDM?

$$d_{\Lambda}(Q) = \sum_q \int [dx][dy] \Phi_{\Lambda}(x) [d_q H_{\text{EDM}}^{(q)} + e \tilde{d}_q H_{\text{CEDM}}^{(q)}] \Phi_{\bar{\Lambda}}(y)$$

Chen et al., PRL 136,051902(2026)



LCDA 不是 CP 破坏的来源；它决定不同 flavor 的 QCD 传递系数和最终信号大小。

LCDA 形状会改变 flavor 权重与新物理约束

同一硬过程，不同 LCDA 输入

s-quark EDM coefficient

QCD SR	lattice moments
1.94×10^{-3}	$5.29 \times 10^{-4} \approx 3.7 \times$

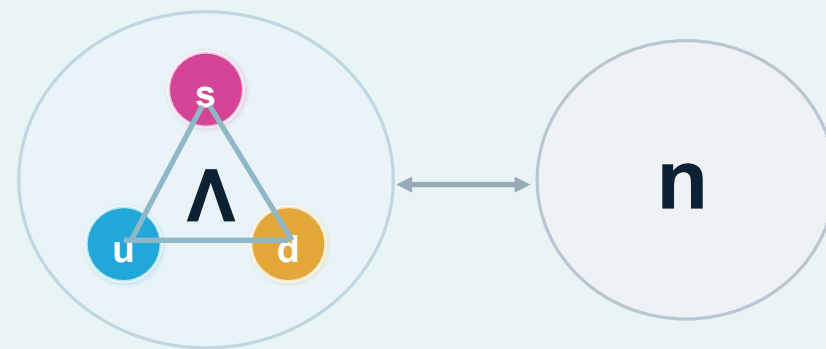
s-quark CEDM coefficient

QCD SR	lattice moments
1.78×10^{-4}	$6.21 \times 10^{-5} \approx 2.9 \times$

Chen et al., PRL 136,051902(2026)

LCDA 是当前数值计算中的主导强子不确定性最重要输入。

Λ EDM 与 neutron EDM 的互补性



s-CEDM sensitive

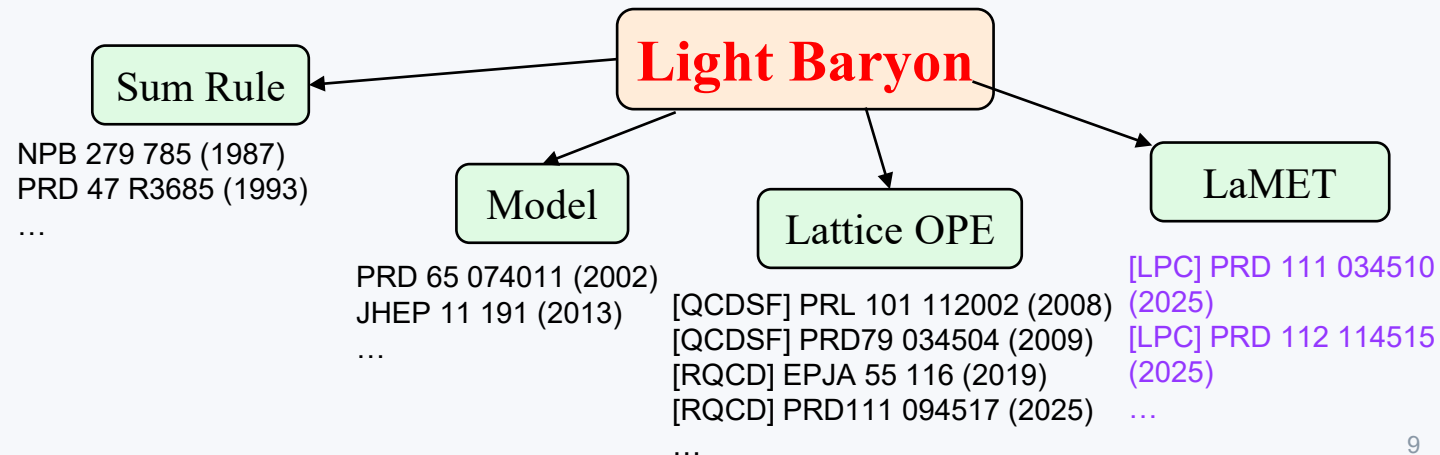
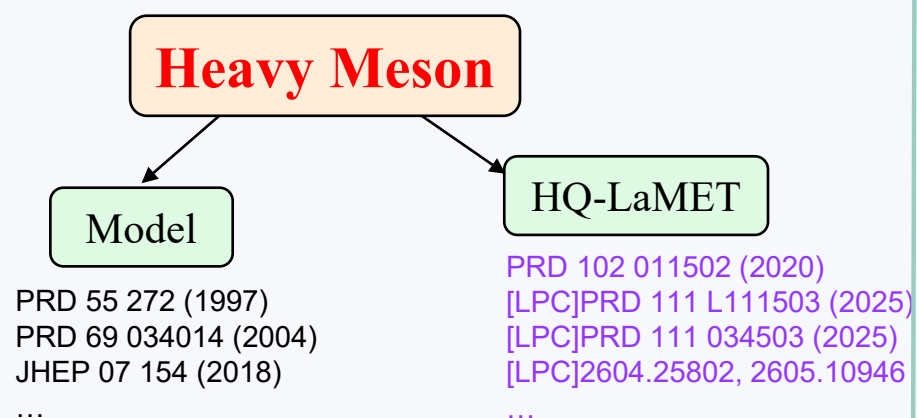
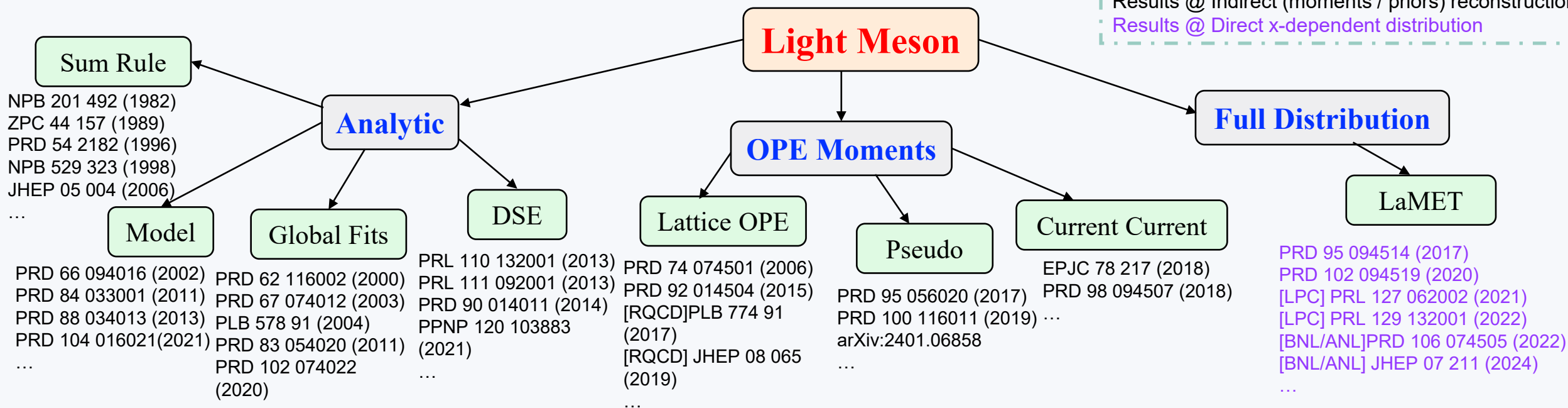
no $\tilde{d}_s$ term

example bound

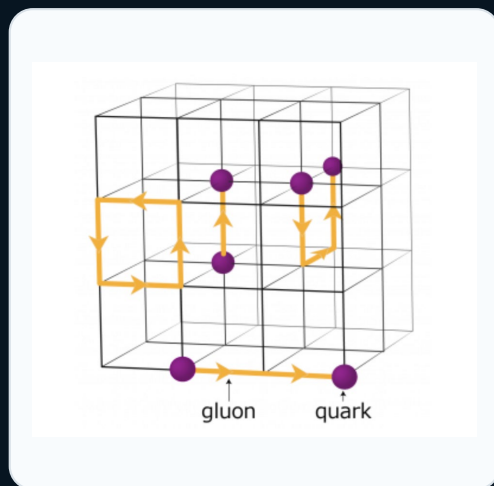
$$|\tilde{d}_s| \lesssim 1.4 \times 10^{-14} \text{ cm}$$



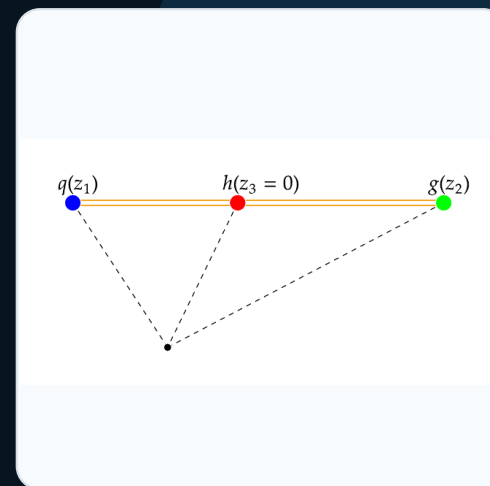
Results @ Indirect (moments / priors) reconstruction
Results @ Direct x-dependent distribution



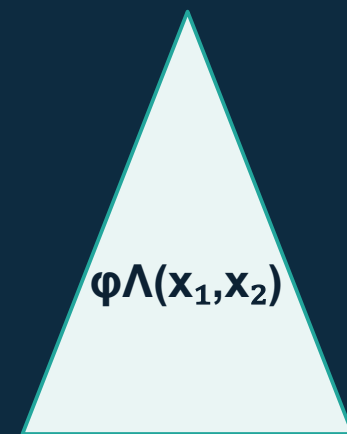
如何从 Euclidean lattice 算到 light-cone structure?



Euclidean grid

large P^z boost

equal-time quasi-DA



physical LCDA

$$\phi_\Lambda = C \otimes \tilde{\phi}_\Lambda + \mathcal{O}(\Lambda_{\text{QCD}}^2 / (x_i P^z)^2)$$

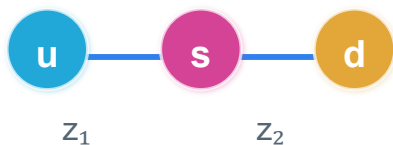
LaMET: 用大动量下的等时空间关联替代不可直接计算的光锥关联, 再通过微扰 matching 恢复物理 LCDA。

从 meson 到 baryon: 不是简单复制

Baryon LCDA 的难点来自“二维 + 三夸克 + 多结构”。

01

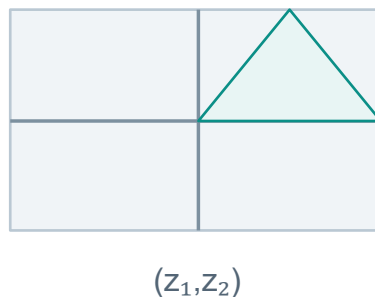
真正二维的非局域算符



两个独立空间分离 z_1 、 z_2 ；二维 Fourier transform；三个 valence quarks 的相关性不能降维。

02

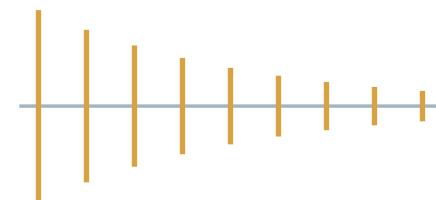
二维重整化与 matching



Wilson-line UV 发散依赖 $|z_1|$ 、 $|z_2|$ 、 $|z_1 - z_2|$ ；需要在整个坐标平面上一致处理。

03

信号与大距离完成

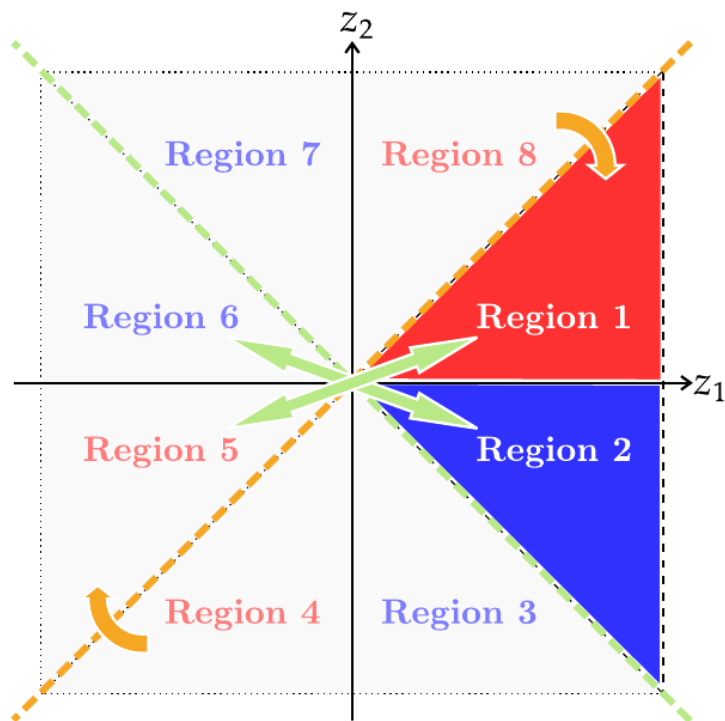


large $|z| \rightarrow$ noise / truncation

boosted baryon 信噪比快速恶化；有限坐标范围会在二维 Fourier transform 中放大系统误差。

这三类问题决定了：baryon LCDA 需要一套专门的完整分析框架。

Λ 对称性：把完整坐标平面压缩到两个独立区域



轻夸克交换： $z_1 \leftrightarrow z_2$

$$\tilde{\Phi}_{\Lambda}^A(z_1, z_2) = +\tilde{\Phi}_{\Lambda}^A(z_2, z_1)$$

$$\tilde{\Phi}_{\Lambda}^{V,T}(z_1, z_2) = -\tilde{\Phi}_{\Lambda}^{V,T}(z_2, z_1)$$

$$\tilde{\Phi}_{\Lambda}^X(-z_1, -z_2) = [\tilde{\Phi}_{\Lambda}^X(z_1, z_2)]^*$$

A

symmetric

V, T

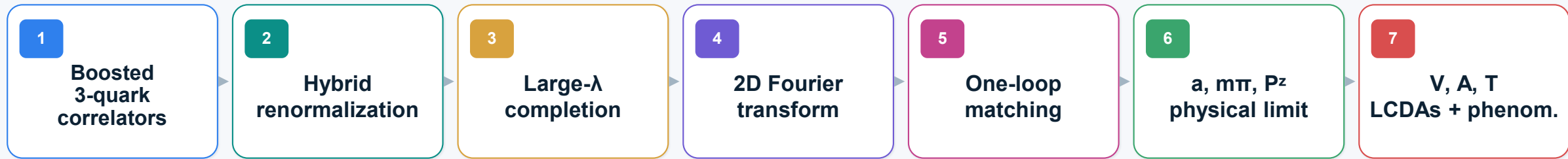
antisymmetric

Hermiticity

real momentum-space quasi-DAs

只需在 $z_1 > |z_2|$ 的独立四分之一区域分析，再用对称性重建整个平面。

两篇文章：一条从形式到物理结果的完整链条



arXiv:2606.30387v2 [hep-lat] 20 Jul 2026

Companion article

Formalism · Renormalization Extrapolation · Matching

(LPC) arXiv:2606.30387

建立 baryon-LaMET 计算分析流程, 并系统量化 end-to-end 误差来源。

Baryon Light-Cone Distribution Amplitudes from Lattice QCD: Formalism, Renormalization, Extrapolation, and Matching
Karin Bauer (Colorado UPRC)



Minlin Zhang¹, Baoping Bao^{2,3}, Shihong Chen^{4,5}, Jun Bao^{6,7}, Xinggang Li⁸, Xinggang Wang⁹, Jian Liang¹⁰, Yu-Liang Li¹¹, Jiahua Shao¹², Wei Wang^{13,14}, Yuhui Yang^{15,16,17,18,19,20,21,22,23,24,25,26,27,28,29,30,31,32,33,34,35,36,37,38,39,40,41,42,43,44,45,46,47,48,49,50,51,52,53,54,55,56,57,58,59,60,61,62,63,64,65,66,67,68,69,70,71,72,73,74,75,76,77,78,79,80,81,82,83,84,85,86,87,88,89,90,91,92,93,94,95,96,97,98,99,100,101,102,103,104,105,106,107,108,109,110,111,112,113,114,115,116,117,118,119,120,121,122,123,124,125,126,127,128,129,130,131,132,133,134,135,136,137,138,139,140,141,142,143,144,145,146,147,148,149,150,151,152,153,154,155,156,157,158,159,160,161,162,163,164,165,166,167,168,169,170,171,172,173,174,175,176,177,178,179,180,181,182,183,184,185,186,187,188,189,190,191,192,193,194,195,196,197,198,199,200,201,202,203,204,205,206,207,208,209,210,211,212,213,214,215,216,217,218,219,220,221,222,223,224,225,226,227,228,229,230,231,232,233,234,235,236,237,238,239,240,241,242,243,244,245,246,247,248,249,250,251,252,253,254,255,256,257,258,259,260,261,262,263,264,265,266,267,268,269,270,271,272,273,274,275,276,277,278,279,280,281,282,283,284,285,286,287,288,289,290,291,292,293,294,295,296,297,298,299,300,301,302,303,304,305,306,307,308,309,310,311,312,313,314,315,316,317,318,319,320,321,322,323,324,325,326,327,328,329,330,331,332,333,334,335,336,337,338,339,340,341,342,343,344,345,346,347,348,349,350,351,352,353,354,355,356,357,358,359,360,361,362,363,364,365,366,367,368,369,370,371,372,373,374,375,376,377,378,379,380,381,382,383,384,385,386,387,388,389,390,391,392,393,394,395,396,397,398,399,400,401,402,403,404,405,406,407,408,409,410,411,412,413,414,415,416,417,418,419,420,421,422,423,424,425,426,427,428,429,430,431,432,433,434,435,436,437,438,439,440,441,442,443,444,445,446,447,448,449,450,451,452,453,454,455,456,457,458,459,460,461,462,463,464,465,466,467,468,469,470,471,472,473,474,475,476,477,478,479,480,481,482,483,484,485,486,487,488,489,490,491,492,493,494,495,496,497,498,499,500,501,502,503,504,505,506,507,508,509,510,511,512,513,514,515,516,517,518,519,520,521,522,523,524,525,526,527,528,529,530,531,532,533,534,535,536,537,538,539,540,541,542,543,544,545,546,547,548,549,550,551,552,553,554,555,556,557,558,559,560,561,562,563,564,565,566,567,568,569,570,571,572,573,574,575,576,577,578,579,580,581,582,583,584,585,586,587,588,589,590,591,592,593,594,595,596,597,598,599,600,601,602,603,604,605,606,607,608,609,610,611,612,613,614,615,616,617,618,619,620,621,622,623,624,625,626,627,628,629,630,631,632,633,634,635,636,637,638,639,640,641,642,643,644,645,646,647,648,649,650,651,652,653,654,655,656,657,658,659,660,661,662,663,664,665,666,667,668,669,670,671,672,673,674,675,676,677,678,679,680,681,682,683,684,685,686,687,688,689,690,691,692,693,694,695,696,697,698,699,700,701,702,703,704,705,706,707,708,709,710,711,712,713,714,715,716,717,718,719,720,721,722,723,724,725,726,727,728,729,730,731,732,733,734,735,736,737,738,739,740,741,742,743,744,745,746,747,748,749,750,751,752,753,754,755,756,757,758,759,760,761,762,763,764,765,766,767,768,769,770,771,772,773,774,775,776,777,778,779,780,781,782,783,784,785,786,787,788,789,790,791,792,793,794,795,796,797,798,799,800,801,802,803,804,805,806,807,808,809,810,811,812,813,814,815,816,817,818,819,820,821,822,823,824,825,826,827,828,829,830,831,832,833,834,835,836,837,838,839,840,841,842,843,844,845,846,847,848,849,850,851,852,853,854,855,856,857,858,859,860,861,862,863,864,865,866,867,868,869,870,871,872,873,874,875,876,877,878,879,880,881,882,883,884,885,886,887,888,889,890,891,892,893,894,895,896,897,898,899,900,901,902,903,904,905,906,907,908,909,910,911,912,913,914,915,916,917,918,919,920,921,922,923,924,925,926,927,928,929,930,931,932,933,934,935,936,937,938,939,940,941,942,}

[illegible]

Lattice setup: 覆盖 a 、 m_π 与 P^z 三个方向

CLQCD · $N_f = 2+1$

7

lattice ensembles

0.052–0.105
fm

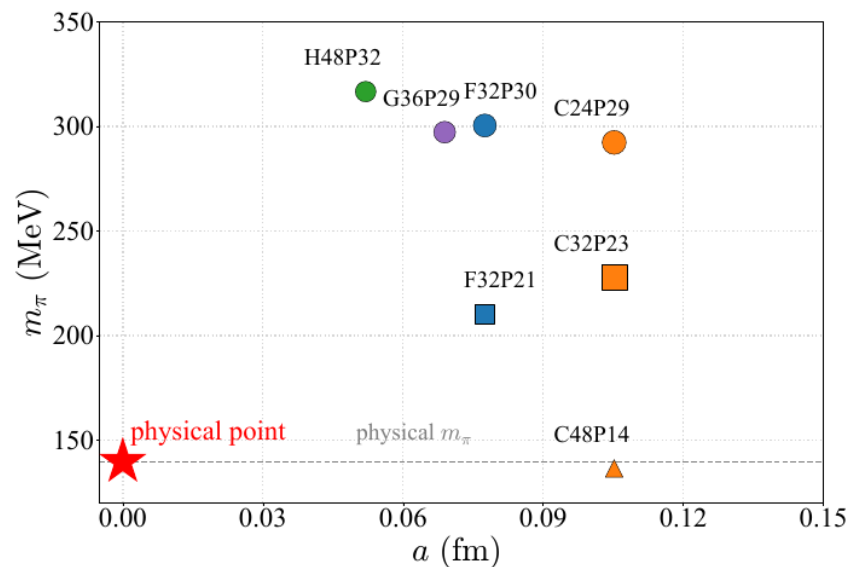
lattice spacing a

136–317 MeV

pion mass m_π

$P^z \leq 3 \text{ GeV}$

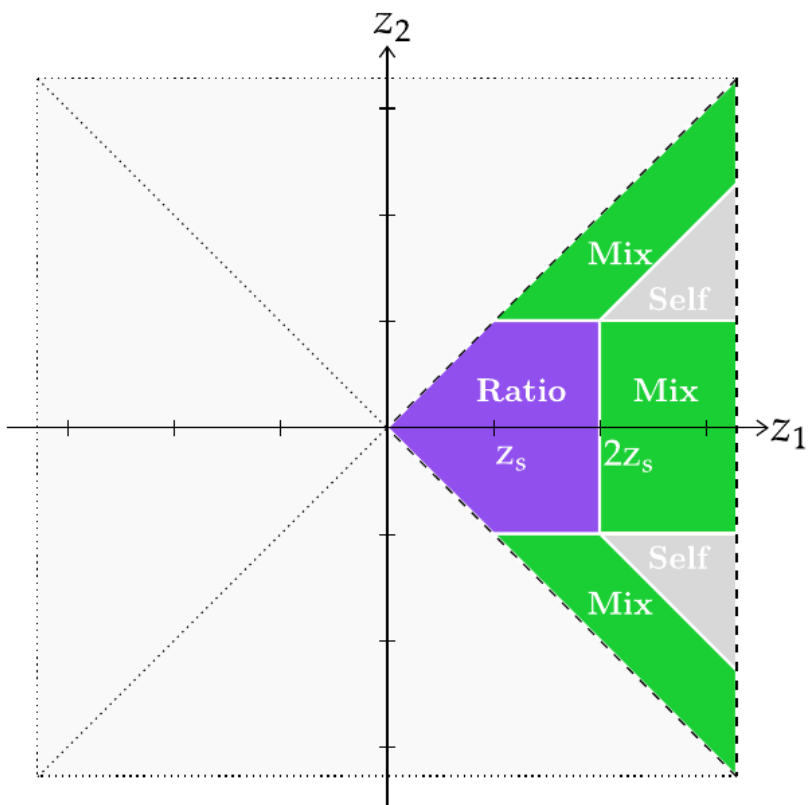
boosted Λ momenta



(CLQCD) PRD 109 054507 (2024)

多格距、多 pion mass、多 momentum 为 continuum、physical-mass 与 infinite-momentum 极限提供 lever arm。

Hybrid renormalization: 在二维坐标平面上分区处理



三种nonlocal距离划分

$$|z_1| \cdot |z_2| \cdot |z_1 - z_2|$$

SHORT

**Ratio
prescription**

用 zero-momentum matrix elements 消除短距离 UV dependence。

LONG

**Self-
renormalization**

长距离避免 ratio denominator 的红外污染。

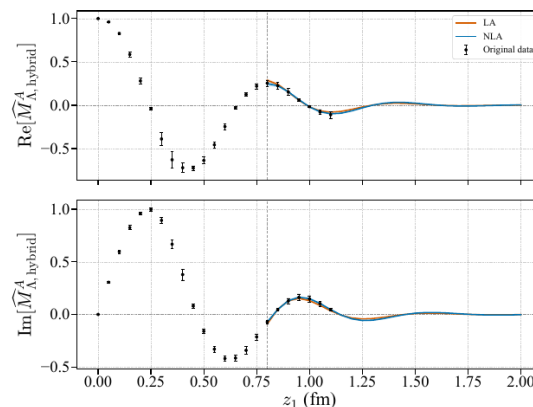
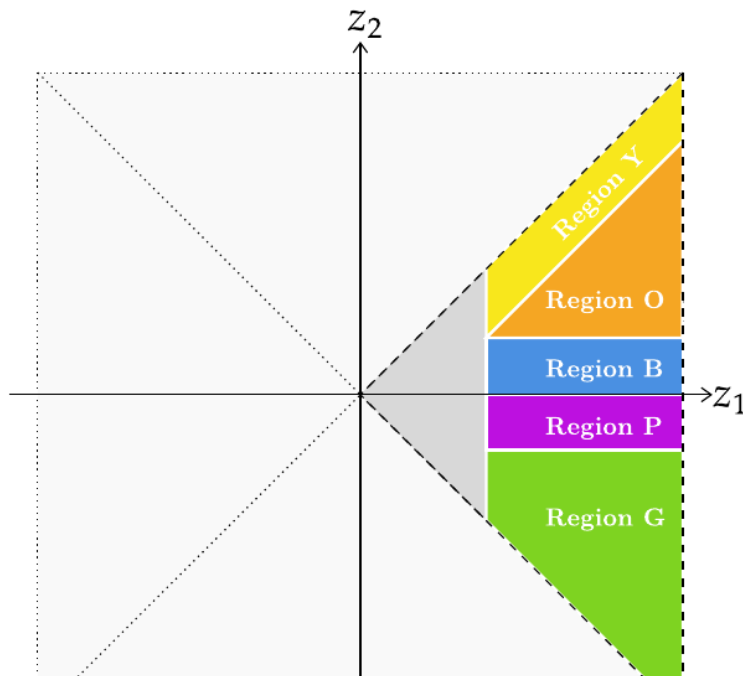
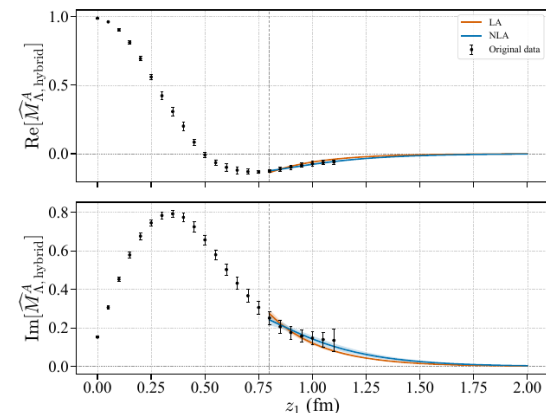
MIX

**Continuous
mixed regions**

对长、短距离成分分别处理，并在边界保持连续。

不是“一个 renormalization factor 乘遍全平面”，而是依据物理距离自治拼接。

Large- λ completion: 让二维 Fourier transform 可控

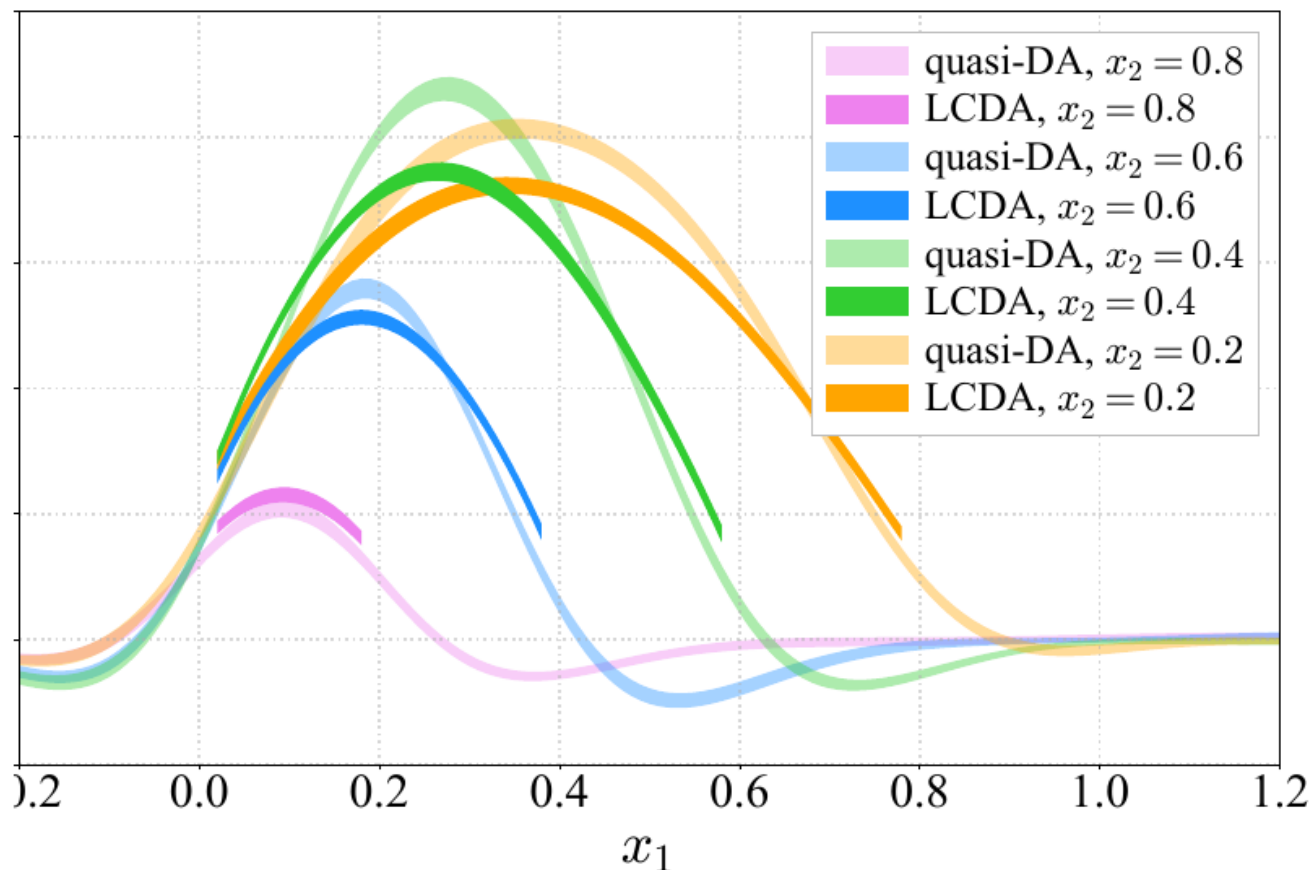
(a) $z_1 - z_2 = 0, z_1 + z_2 \geq 0$ (b) $z_2 = 1 \text{ grid unit}, z_1 \geq 0$

新的 Euclidean long-distance expansion *arXiv:2601.12189*

- 依据 $|z_1|$ 、 $|z_2|$ 、 $|z_1 - z_2|$ 哪个进入大距离，为不同区域构造 baryon-specific ansätze。
- LA / NLA 稳定性差异纳入系统误差。
- 结果平滑延拓到有限格点数据之外。

有限坐标数据 → 物理约束的平滑完成 → 稳定的二维 momentum-space quasi-DA。

LaMET matching: 改变的是形状，不是一个整体因子



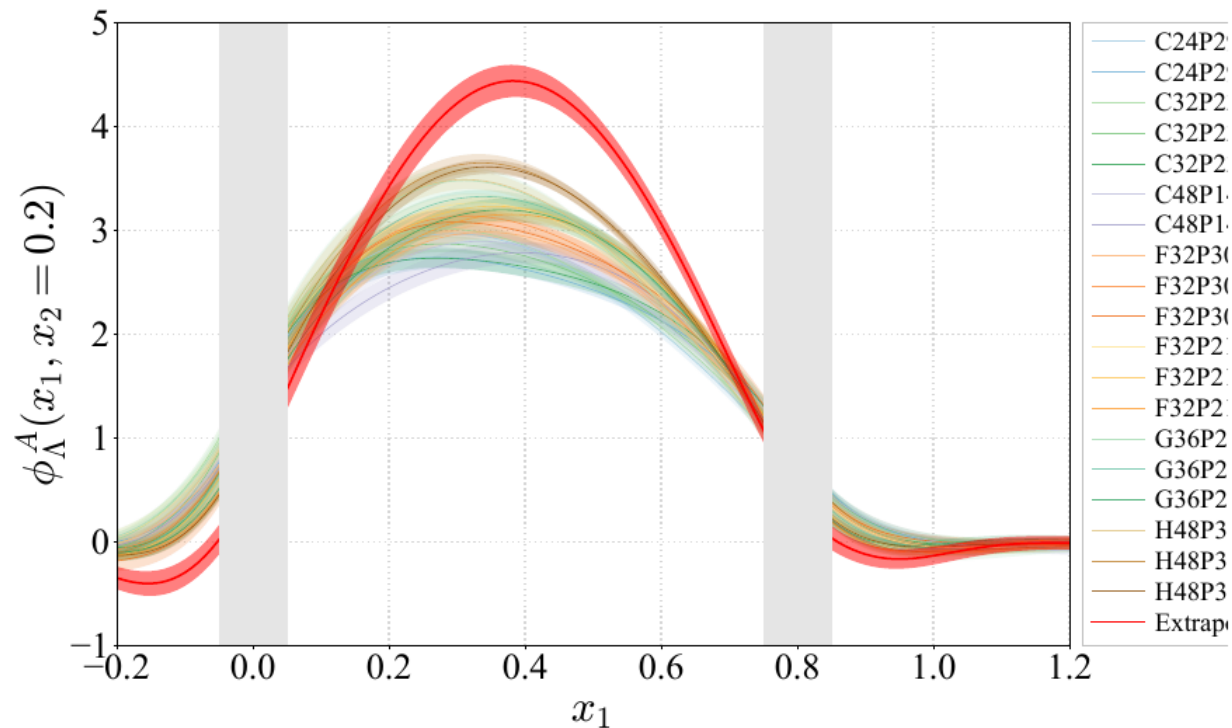
one-loop matching

$$\phi^X = C^X \otimes \tilde{\phi}^X$$

- 把 finite- P_z quasi-DA 转换为 light-cone LCDA
- x_1, x_2 依赖的 kernel 重新分配支撑与峰形
- hybrid counterterms 与 renormalization scheme 一致

visible, momentum-fraction-dependent modification

联合外推: continuum · physical pion mass · $P^z \rightarrow \infty$



point-by-point in (x_1, x_2)

$$\phi_{a,m_\pi,P^z} = \phi_{\text{phys}} + \frac{A}{(P^z)^2} + (m_\pi^2 - m_{\pi,\text{phys}}^2)B + a^2[D_1 + (P^z)^2 D_2]$$

$1/(P^z)^2$

leading LaMET power correction

$(m_\pi)^2$

physical-pion-mass dependence

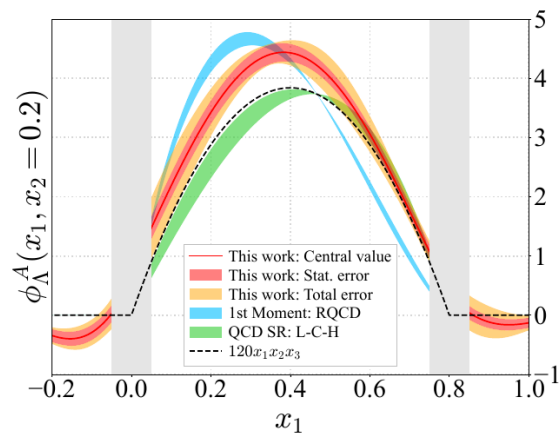
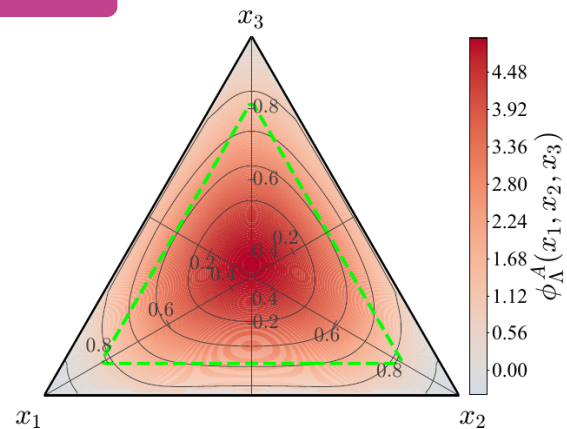
$a^2 + a^2 P^z^2$

static and boosted discretization

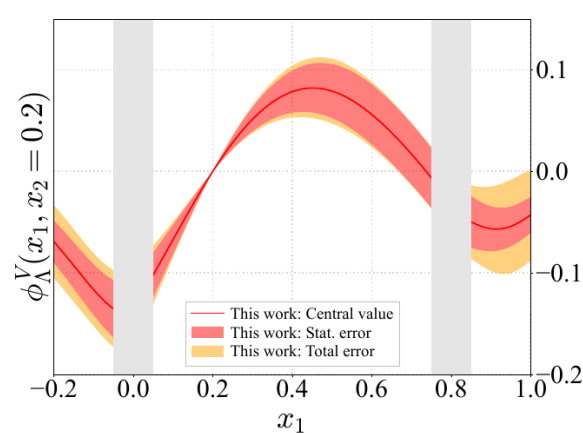
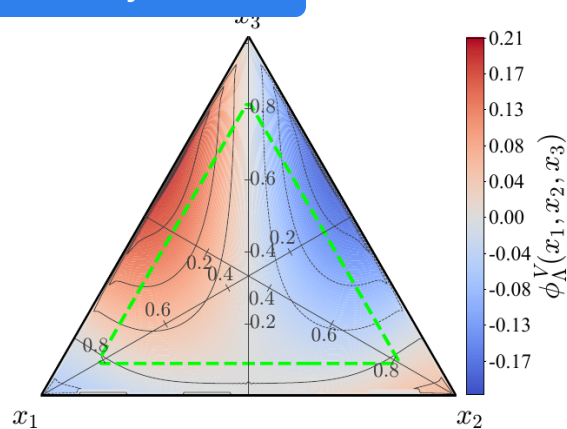
red band : physical-limit LCDA

第一性原理获得完整的 Λ leading-twist V、A、T

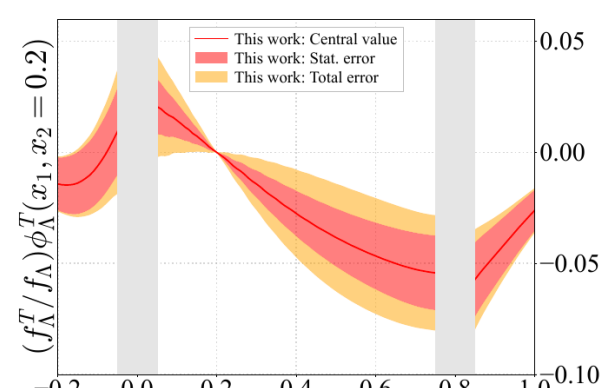
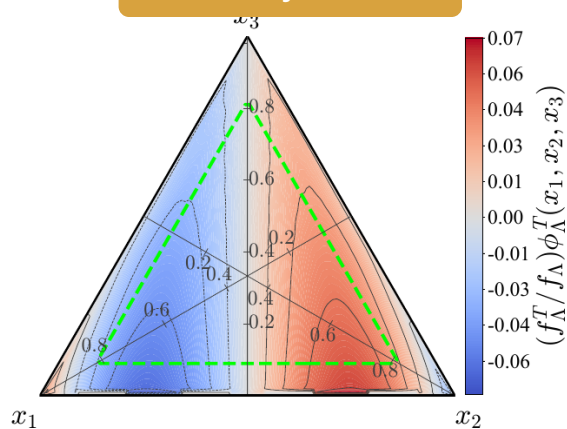
A · symmetric



V · antisymmetric



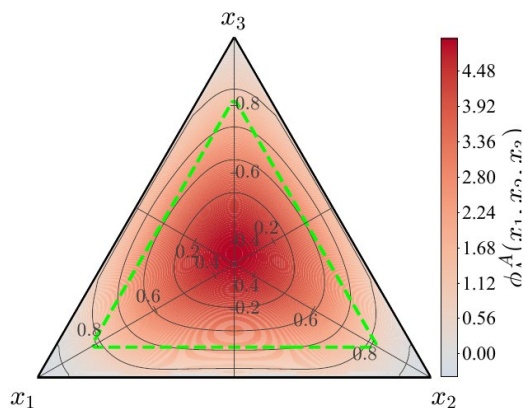
T · antisymmetric



first complete 2D lattice-QCD
determination

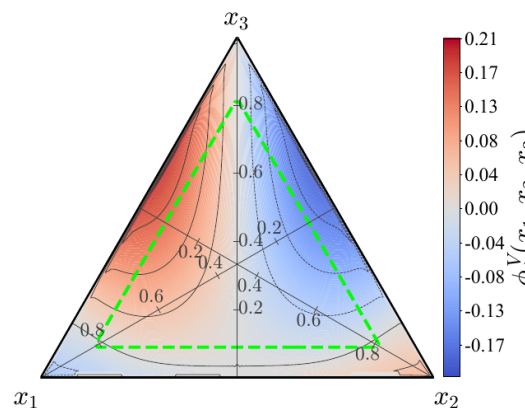
二维分布告诉我们什么？

A: 主导且对称



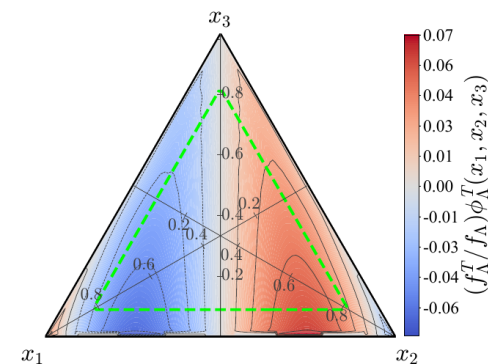
内部宽峰；峰形轻微向 x_3 方向偏移。
→ s 夸克携带的纵向动量份额略大。

V: 反对称



跨 $x_1 = x_2$ 变号，并在对角线上为零。
→ 非局域轻夸克交换结构被直接解析。

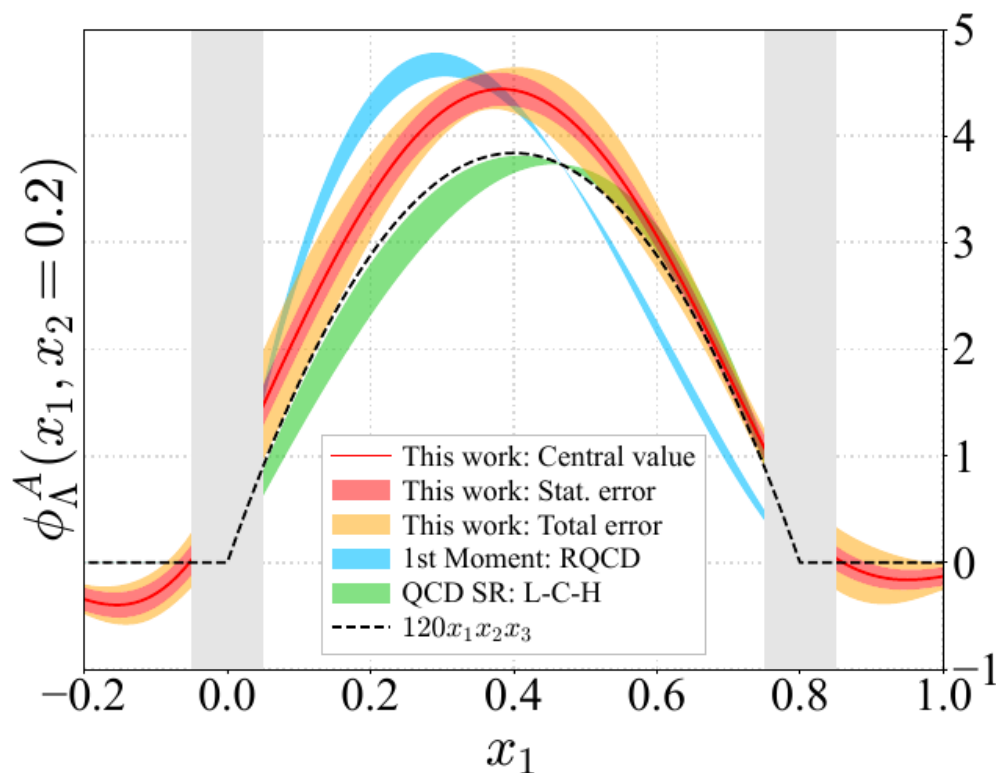
T: 反对称



同样具有 sign-changing pattern，
幅度较小但显著非零。
→ 超出 local-moment 能力范围。

结果不仅给出“有多大”，更首次在 momentum-fraction plane 中直接看见对称与反对称结构。

超越渐近形状与低阶矩：可观测量出现 $O(10\%)$ 改变



Λ Dirac form factor at $Q^2 = 20 \text{ GeV}^2$

LO	asymptotic	lattice LCDA	12.7%
	0.00425	→ 0.00371	
NLO	asymptotic	lattice LCDA	11.7%
	0.00913	→ 0.00806	

full 2D shape becomes a leading phenomenological input once the hard kernel is under control.

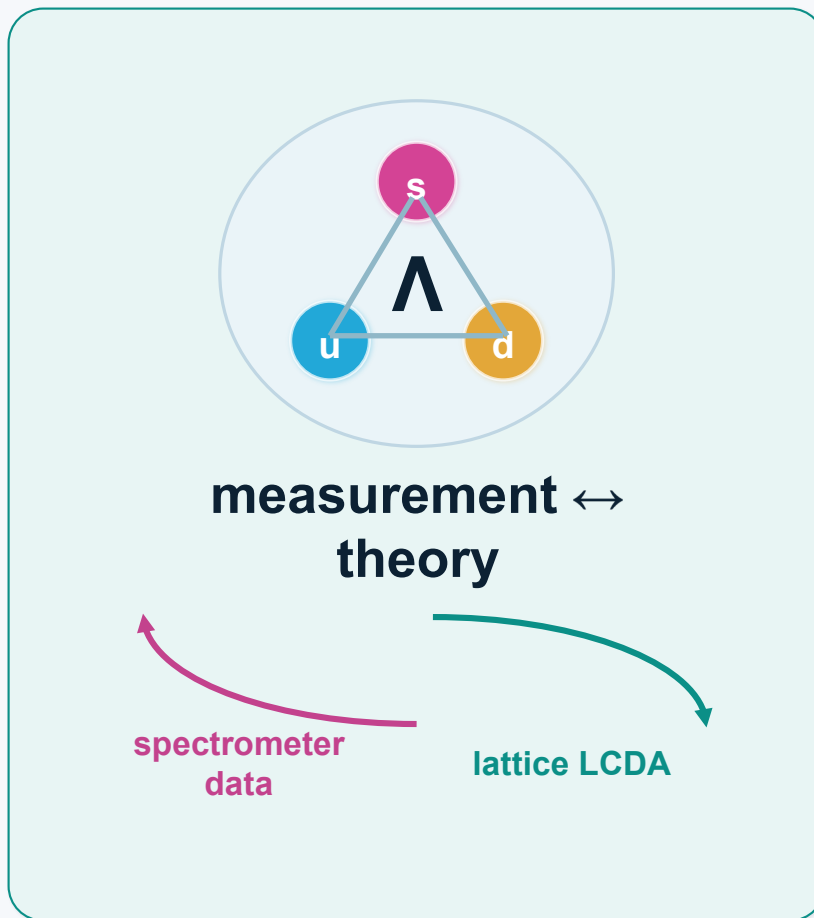
对超子-核子谱仪物理意味着什么？

time-like EM form factors

为高 Q^2 短距离 hard-scattering 部分提供第一性原理三夸克输入。

hyperon EDM / CP-odd FF

减少从 $d\Lambda(Q)$ 到 quark EDM/CEDM 的强子映射不确定性。



极化与弱衰变 CP

为 baryonic amplitudes 与 CP asymmetries 的因子化分析提供结构输入。

octet flavor systematics

同一框架可推广到 p 、 n 、 Σ 、 Ξ ，建立 strange/non-strange 对照。

最直接的近期用法：把文中给出的二维参数化接入 hyperon form-factor 与 EDM 分析。

三点总结

01

Λ 是连接奇异味结构、极化与 CP/EDM 的关键实验室

自分析弱衰变、精密 form factors 与 EDM 搜索使 Λ 成为谱仪物理中的高杠杆对象。

02

我们首次从格点 QCD 完整得到物理点 V、A、T 二维 LCDAs

完整流程包含二维 hybrid renormalization、large- λ completion、matching 与联合物理极限外推。

03

完整形状对可观测量和 BSM 解释是定量输入

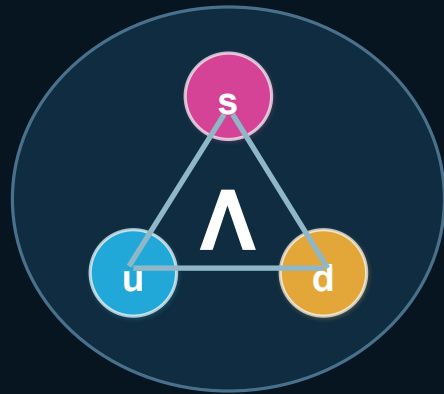
在 Λ form factor 中产生 O(10%) 效应，并显著影响 hyperon EDM 对 quark dipoles 的 flavor 权重。

谢谢 · Thank you

Discussion

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Lattice Parton Collaboration



complete 2D Λ light-cone
structure



arXiv:2606.29597

Complete Access to Leading- Twist Λ -Baryon LCDAs

physical V, A, T + phenomenology

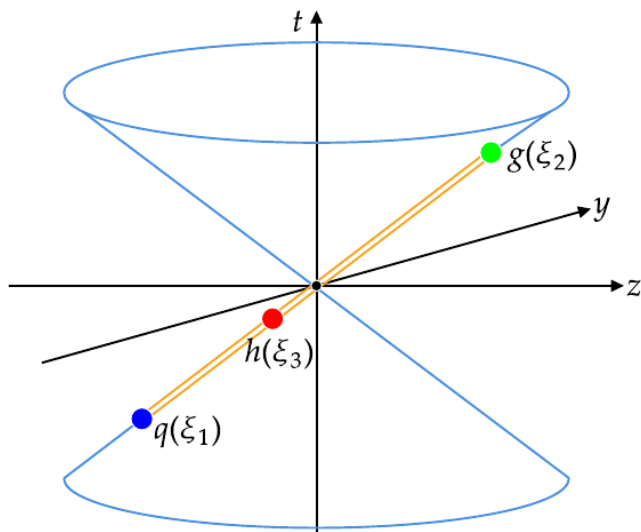


arXiv:2606.30387

Formalism · Renormalization · Extrapolation · Matching

complete baryon-LaMET infrastructure

定义、支撑区域与归一化



$$\phi_{\Lambda}^X(x_1, x_2; \mu) = (P^+)^2 \int \frac{d\xi_1}{2\pi} \frac{d\xi_2}{2\pi} e^{i(x_1\xi_1 + x_2\xi_2)P^+} \Phi_{\Lambda}^X(\xi_1 P^+, \xi_2 P^+; \mu)$$

$$0 \leq x_1, x_2, x_3 \leq 1, \quad x_1 + x_2 + x_3 = 1$$

$$\int [dx] \phi_{\Lambda}^A = 1, \quad \int [dx] \phi_{\Lambda}^{V,T} = 0$$

A has a nonzero local limit; V and T vanish locally and use the local A matrix element as the reference normalization.

$$X = V, A, T$$

推荐的二维参数化 (phenomenology input)

TABLE II. Recommended second-order parameterization of the lattice-determined Λ LCDAs. The coefficients $\phi_{\Lambda,nm}$ determine the A and V amplitudes, while $\tilde{\pi}_{\Lambda,nm}$ determine the T amplitude. To suppress the impact of endpoint-enhanced higher-power corrections, the fit is restricted to the region $x_{1,2,3} \geq 0.1$.

$\phi_{\Lambda,10}$	$\phi_{\Lambda,11}$	$\phi_{\Lambda,20}$	$\phi_{\Lambda,21}$	$\phi_{\Lambda,22}$	$\tilde{\pi}_{\Lambda,10}$	$\tilde{\pi}_{\Lambda,21}$
0.0343(12)	0.00543(40)	0.0012(14)	-0.0267(18)	-0.0087(14)	0.00750(10)	0.00043(14)

Recommended default: second-order conformal polynomial fit

fit region

$x_1, x_2, x_3 \geq 0.1$

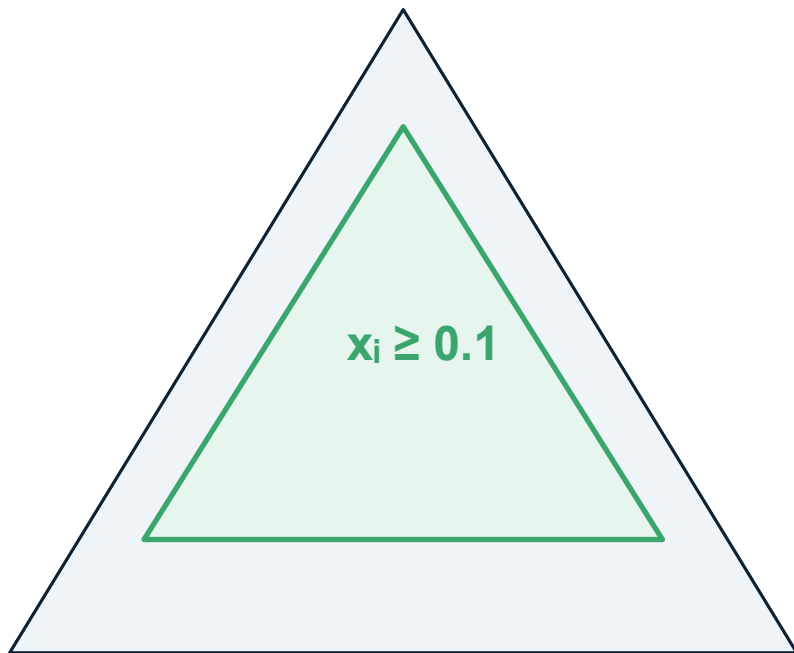
purpose

captures resolved 2D shape without weakly constrained higher components

normalization

absorb residual $\phi_{\Lambda,0}$ into $f_{\Lambda} / \phi_{\Lambda,0}$ for applications

系统误差与 LaMET reliable region



endpoint-enhanced
power corrections

Error budget

- renormalization / matching scale
- large- λ extrapolation
- a , $m\pi$, P^z physical-limit extrapolation
- statistics propagated through full pipeline

central values are quoted with statistical and total uncertainties point-by-point in (x_1, x_2)

EDM interpretation: 适用域与未计入修正

$$d_{\Lambda}(Q) \sim Q^{-4} \Phi_{\Lambda} \otimes H_{\text{CP}} \otimes \Phi_{\bar{\Lambda}}$$

KINEMATICS

BESIII probes $Q = M(J/\psi) \approx 3.1$ GeV; this is a finite- Q EDM form factor, not $d_{\Lambda}(0)$.

PERTURBATIVE

Current example uses leading-power / leading-order hard kernel; higher-order α_s is not included.

TWIST

Higher-twist contributions may be relevant at this moderate scale and require future study.

Use the lattice LCDA to reduce hadronic uncertainty — not to erase the remaining perturbative and power-correction systematics.