

# Validating Wilson Flow Renormalization on the Lattice: Quark Bilinear and Glue Operators

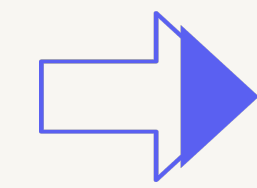
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April 22, 2026 第一届中国电子离子对撞机相关物理年会

# Outline



Wilson flow & SFTX

Gluon momentum fraction

Quark bilinear

# Wilson flow

$$B_\mu \Big|_{t_f=0} = A_\mu \quad \chi \Big|_{t_f=0} = \psi \quad \bar{\chi} \Big|_{t_f=0} = \bar{\psi}$$

$$\partial_{t_f} B_\mu = D_\nu G_{\nu\mu} \quad \partial_{t_f} \chi = \Delta \chi \quad \partial_{t_f} \bar{\chi} = \bar{\chi} \overleftarrow{\Delta}$$

$$G_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu + [B_\mu, B_\nu] \quad D_\mu = \partial_\mu + [B_\mu, \cdot]$$

$$\Delta = D_\mu D_\mu \quad D_\mu = \partial_\mu + B_\mu$$

After renormalization of the fundamental parameters of QCD, the flowed gluon field is finite and does not require renormalization

The flowed quark fields require additional multiplicative field renormalization

M. Lüscher, JHEP08 071 (2010)

M. Lüscher and P. Weisz, JHEP02 051 (2011)

M. Lüscher, JHEP04 123 (2013)

# Short flow time expansion (SFTX)

$$\tilde{O}_g(t_f) \equiv \lim_{a \rightarrow 0} [O_g^L(t_f, a)] = \zeta_O(\mu, t_f) O_g^R(\mu) + \mathcal{O}(t_f)$$

$$\tilde{O}_q(t_f) \equiv \lim_{a \rightarrow 0} [\dot{Z}_\chi(t_f, a) O_q^L(t_f, a)] = \zeta_O(\mu, t_f) O_q^R(\mu) + \mathcal{O}(t_f)$$

$$O^R(\mu) = \lim_{t_f \rightarrow 0} \frac{\lim_{a \rightarrow 0} [\dot{Z}_\chi(t_f, a) O^L(t_f, a)]}{\zeta_O(\mu, t_f)} = \lim_{t_f \rightarrow 0} \lim_{a \rightarrow 0} \left[ \frac{\dot{Z}_\chi(t_f, a) O^L(t_f, a)}{\zeta_O(\mu, t_f)} \right]$$

The 1st limit  $a \rightarrow 0$  should be taken with all  $t_f$  satisfying  $\sqrt{8t_f} \gg a$ . The 2nd limit  $t_f \rightarrow 0$  is usually done linearly.

The golden window:  $a \ll \sqrt{8t_f} \ll \Lambda_{\text{QCD}}^{-1}$ . Similarly, the golden window of RI/MOM is

$$\Lambda_{\text{QCD}} \ll p \ll \pi/a.$$

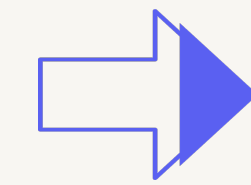
Gauge invariant, chiral symmetry, no mixing, smearing...

M. Lüscher and P. Weisz, JHEP02 051 (2011)

H. Suzuki, PTEP 083B03 (2013)

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# Formalism

$$T_{\mu\nu}^g = \frac{1}{g_0^2} \left[ O_1 - \frac{1}{4} O_2 \right]$$

$$O_{1,\mu\nu} = G_{\mu\rho}^a G_{\nu\rho}^a$$

$$O_{2,\mu\nu} = \delta_{\mu\nu} G_{\rho\sigma}^a G_{\rho\sigma}^a$$

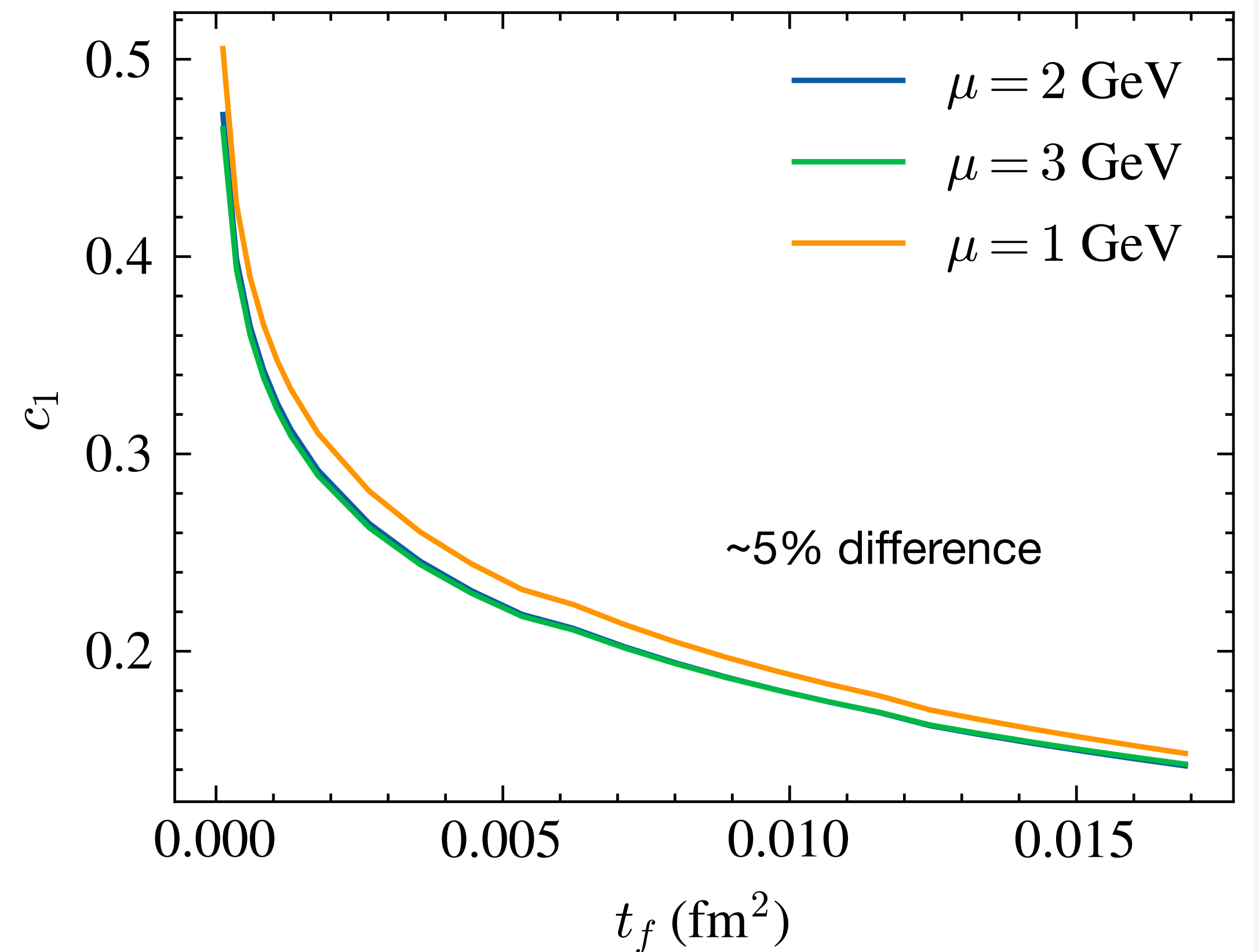
$$O_i^R(\mu) = \lim_{t_f \rightarrow 0} \lim_{a \rightarrow 0} M_{ij}(\mu, t_f) O_j^L(t_f),$$

$$O = O_{1,44} - \frac{1}{3} O_{1,ii}$$

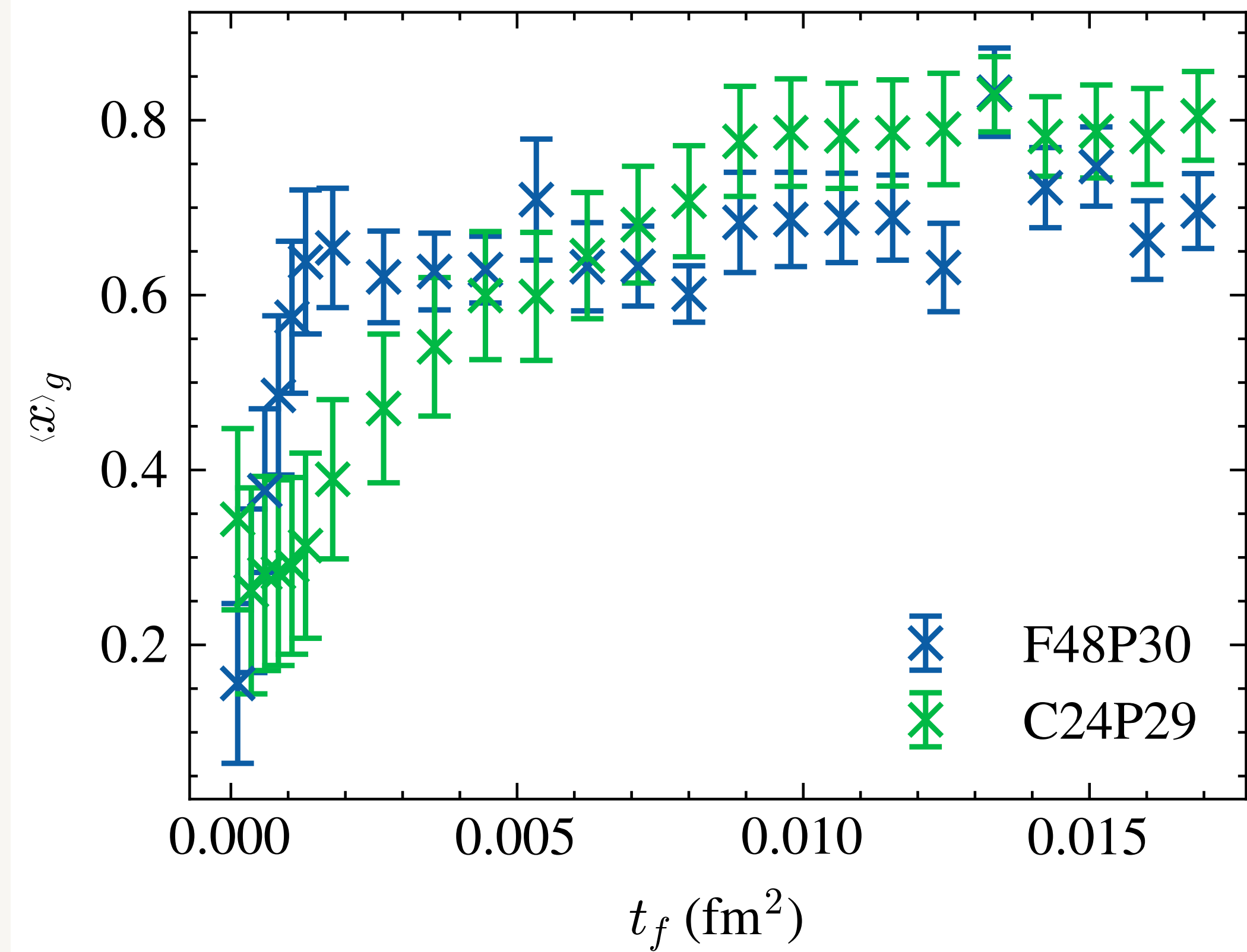
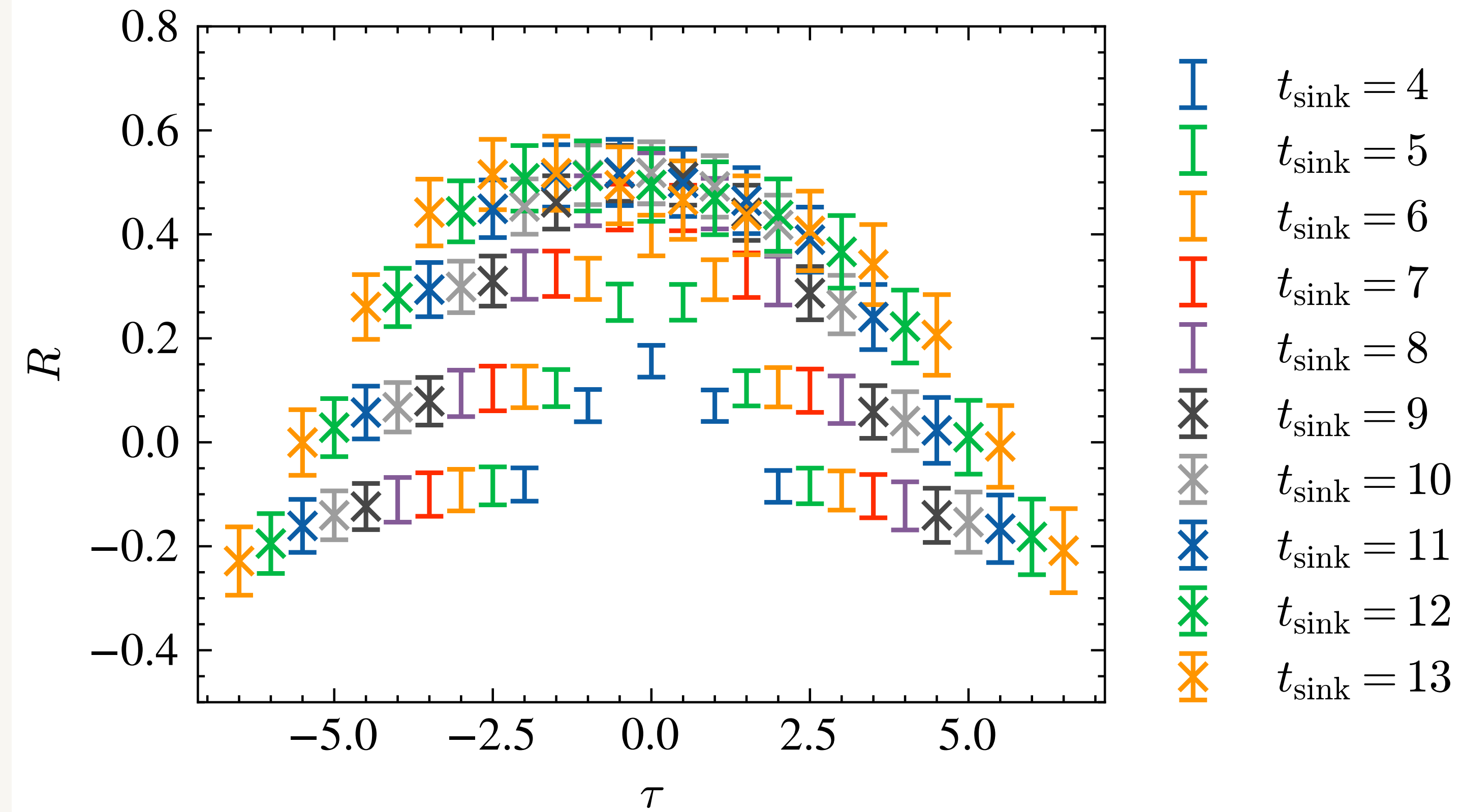
$$\langle N | O | N \rangle = 4E^2 \langle x \rangle_g$$

$$T_{\mu\nu}(x) = c_i(t) \tilde{\mathcal{O}}_{i,\mu\nu}(t, x)$$

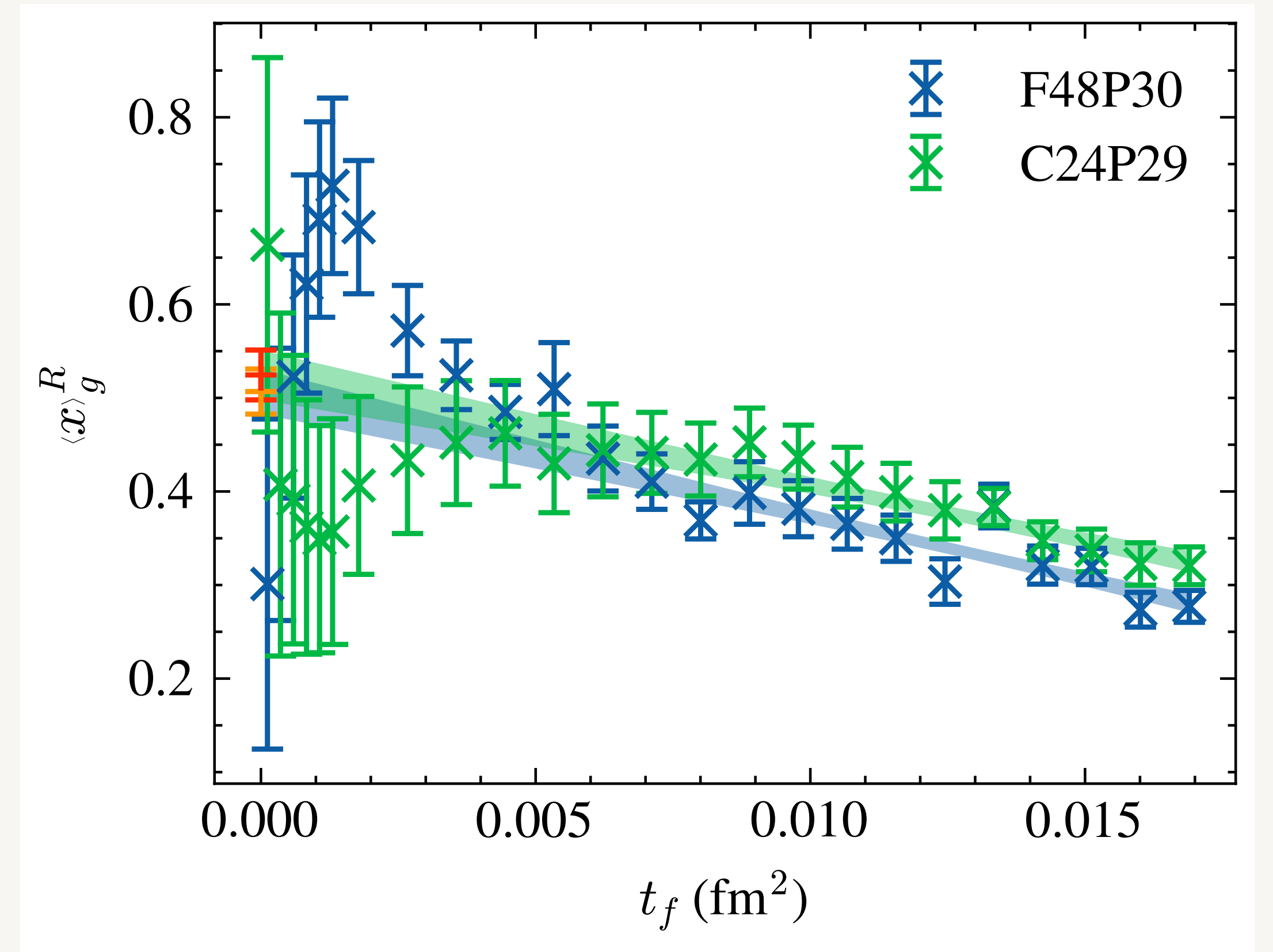
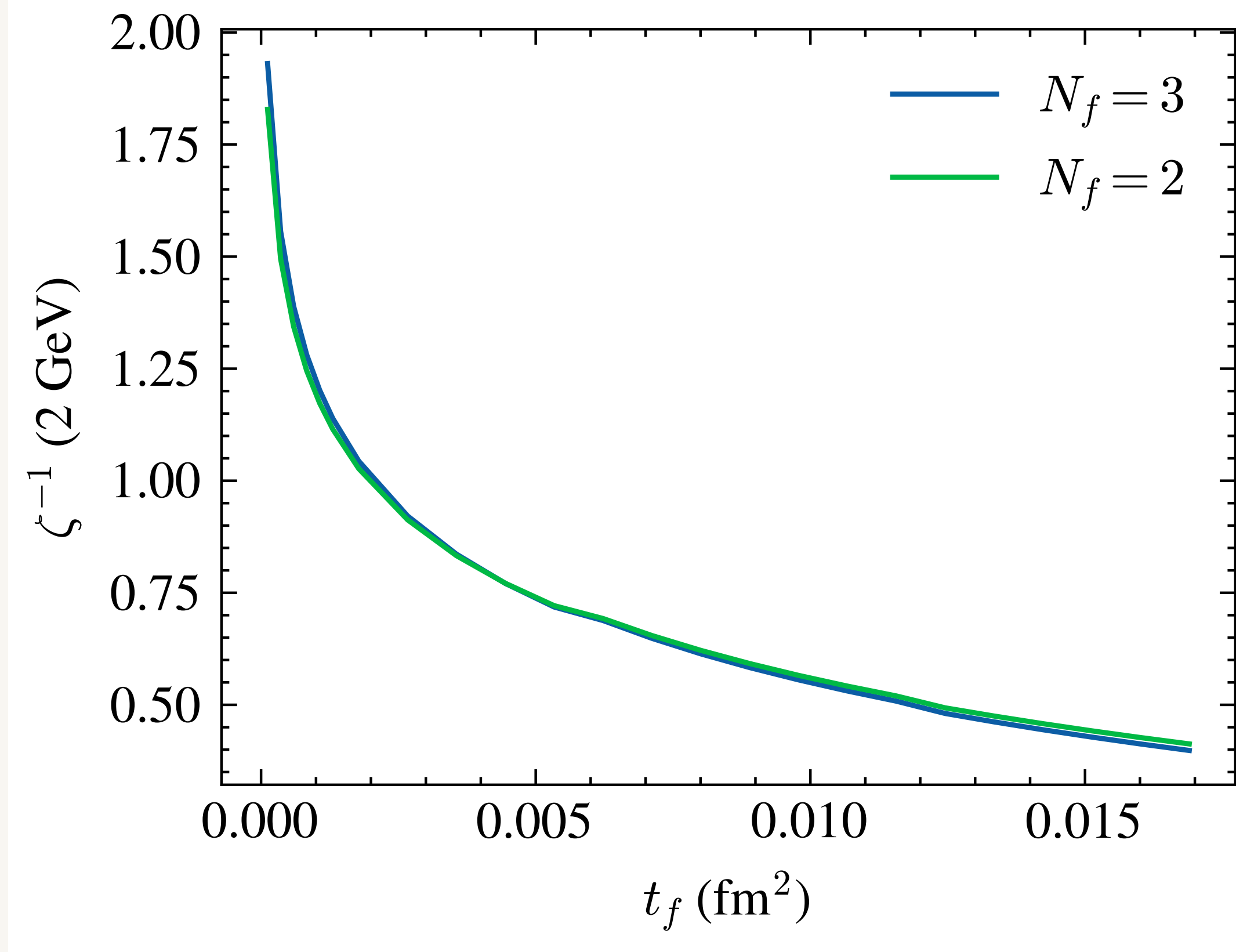
R. V. Harlander, Y. Kluth and F. Lange, EPJC78 944 (2018)



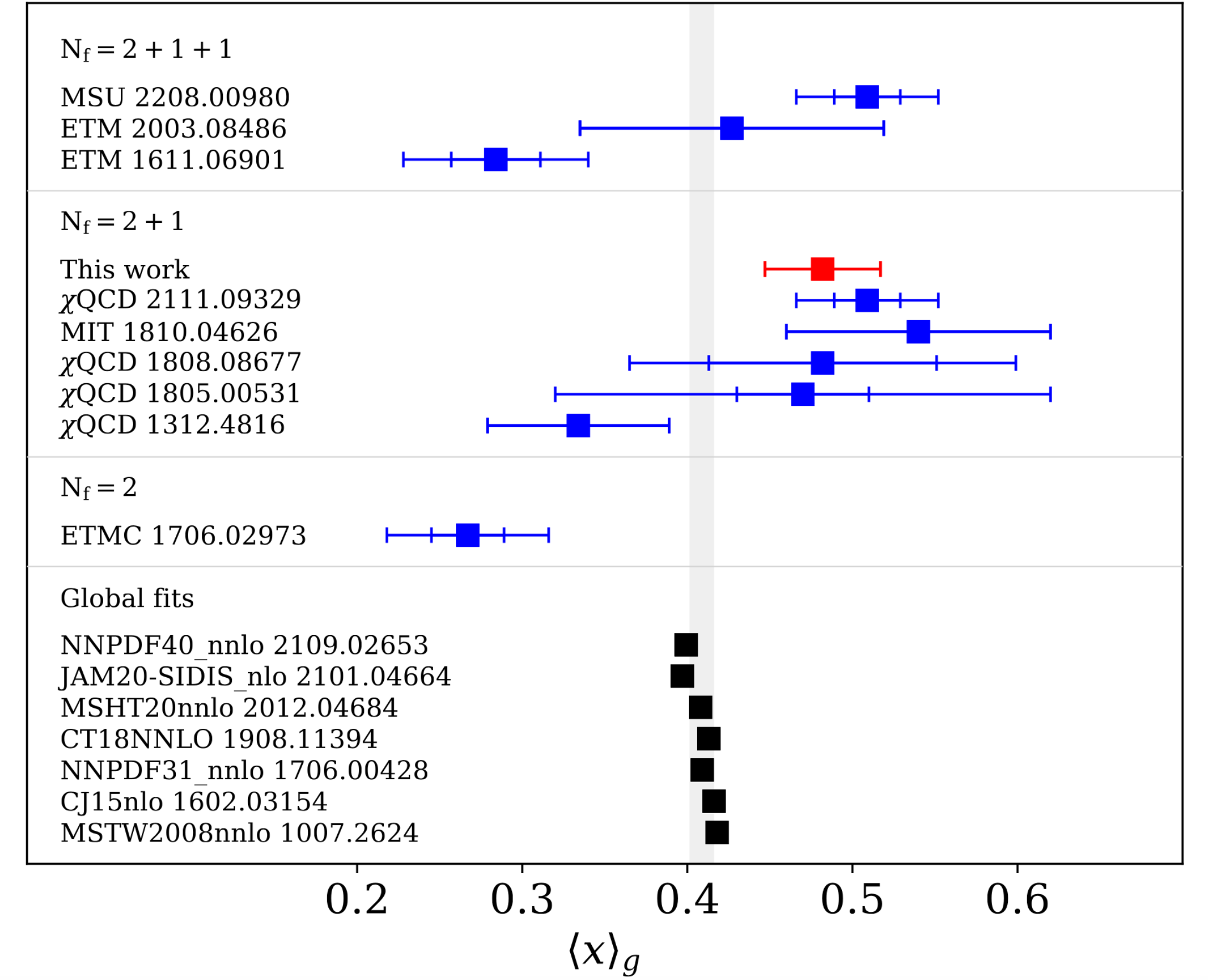
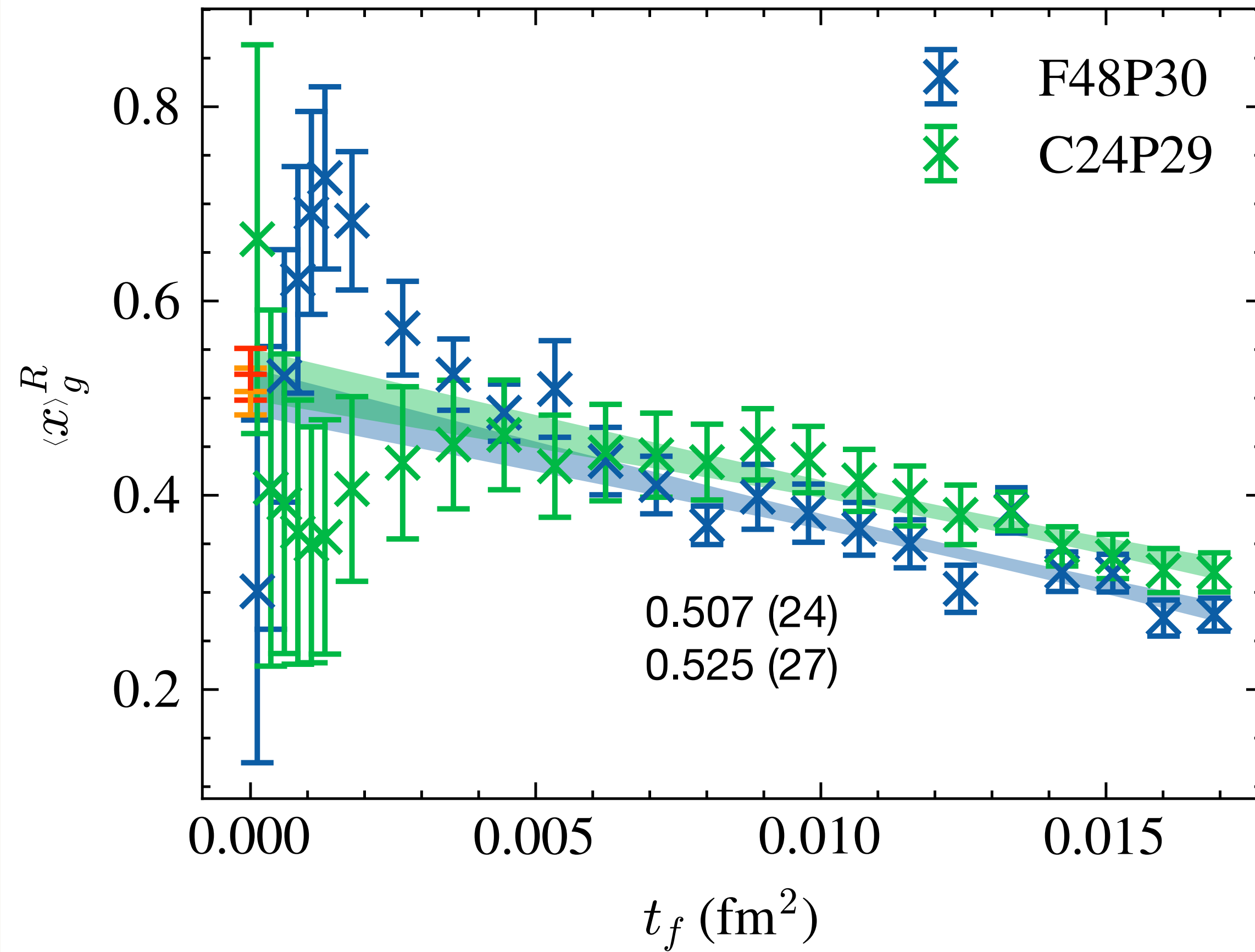
# Lattice matrix elements



# Preliminary results



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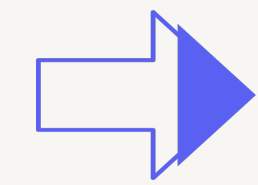


arXiv:2602.14260

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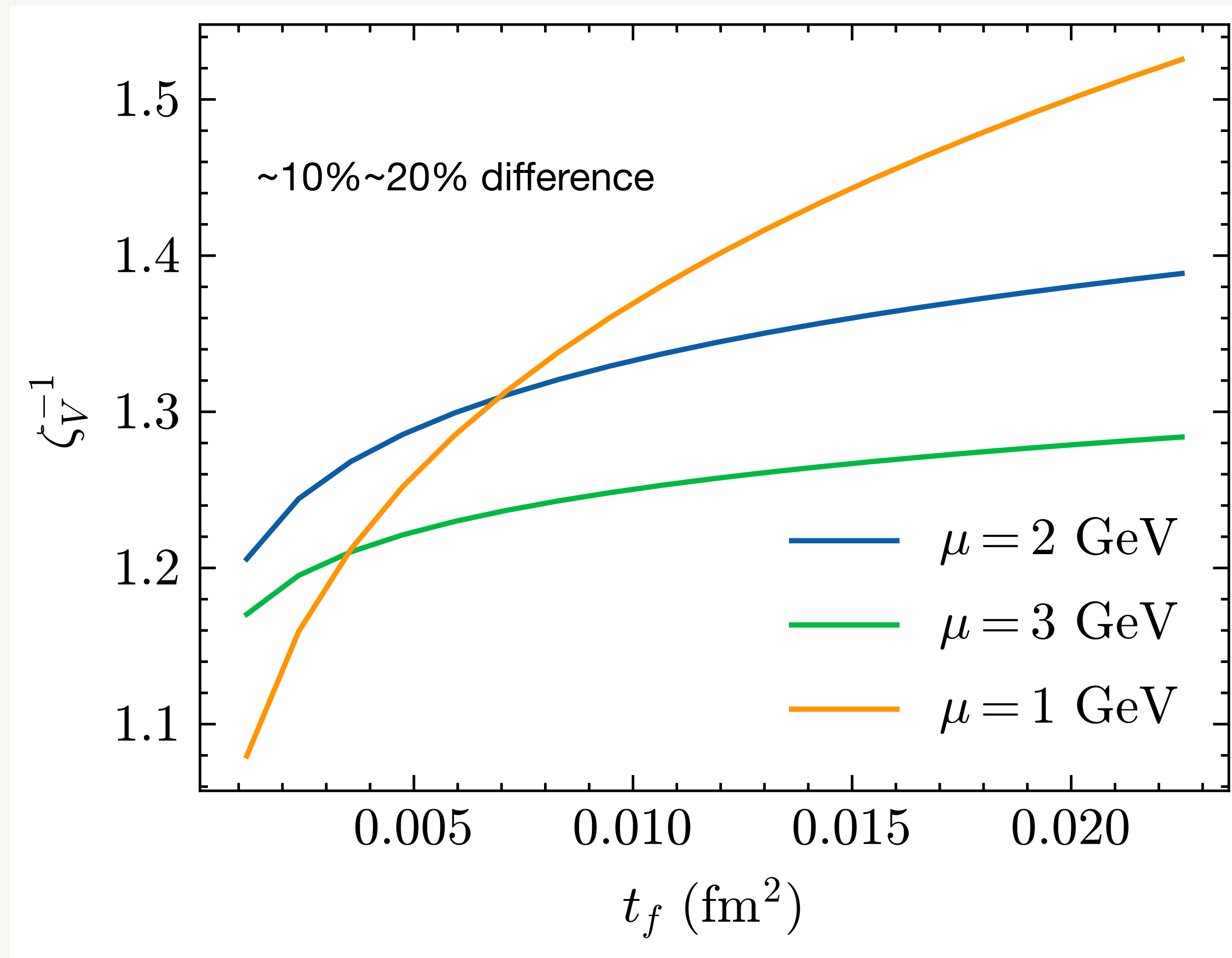


Quark bilinear

# Check on matching

$$\tilde{O}_q(t_f) \equiv \lim_{a \rightarrow 0} [\dot{Z}_\chi(t_f, a) O_q^L(t_f, a)] = \zeta_O(\mu, t_f) O_q^R(\mu) + \mathcal{O}(t_f)$$

J. Borgulat, R. V. Harlander, J. T. Kohnen and F. Lange, JHEP05 179 (2024)



# Renormalization of flowed quarks

$$\mathring{Z}_\chi(t_f, a) = \frac{-2N_c}{(4\pi t_f)^2 \left\langle \bar{\chi}(t_f, x) \overleftrightarrow{D} \chi(t_f, x) \right\rangle}$$

$$\mathring{Z}_\chi(t_f, a) \langle V_\mu^L(t_f, a) \rangle = \zeta_V(\mu, t_f) \langle V_\mu^R(\mu) \rangle + \mathcal{O}(t_f, a)$$

$$\frac{\mathring{Z}_\chi(t_f, a) \langle V_\mu^L(t_f, a) \rangle}{Z_V(\mu, a) \langle V_\mu^L(t_f = 0, a) \rangle} = \zeta_V(\mu, t_f) + \mathcal{O}(t_f, a)$$

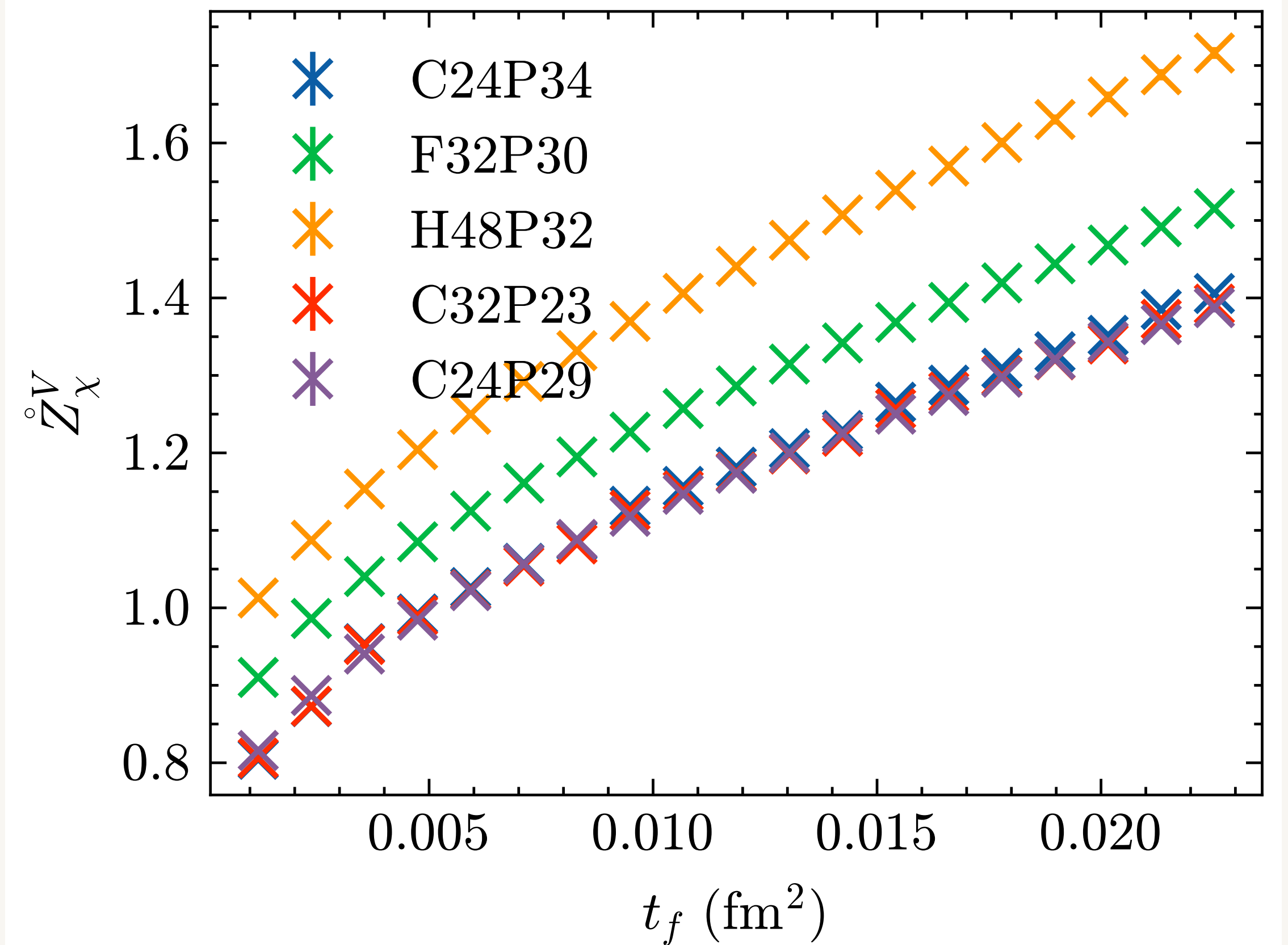
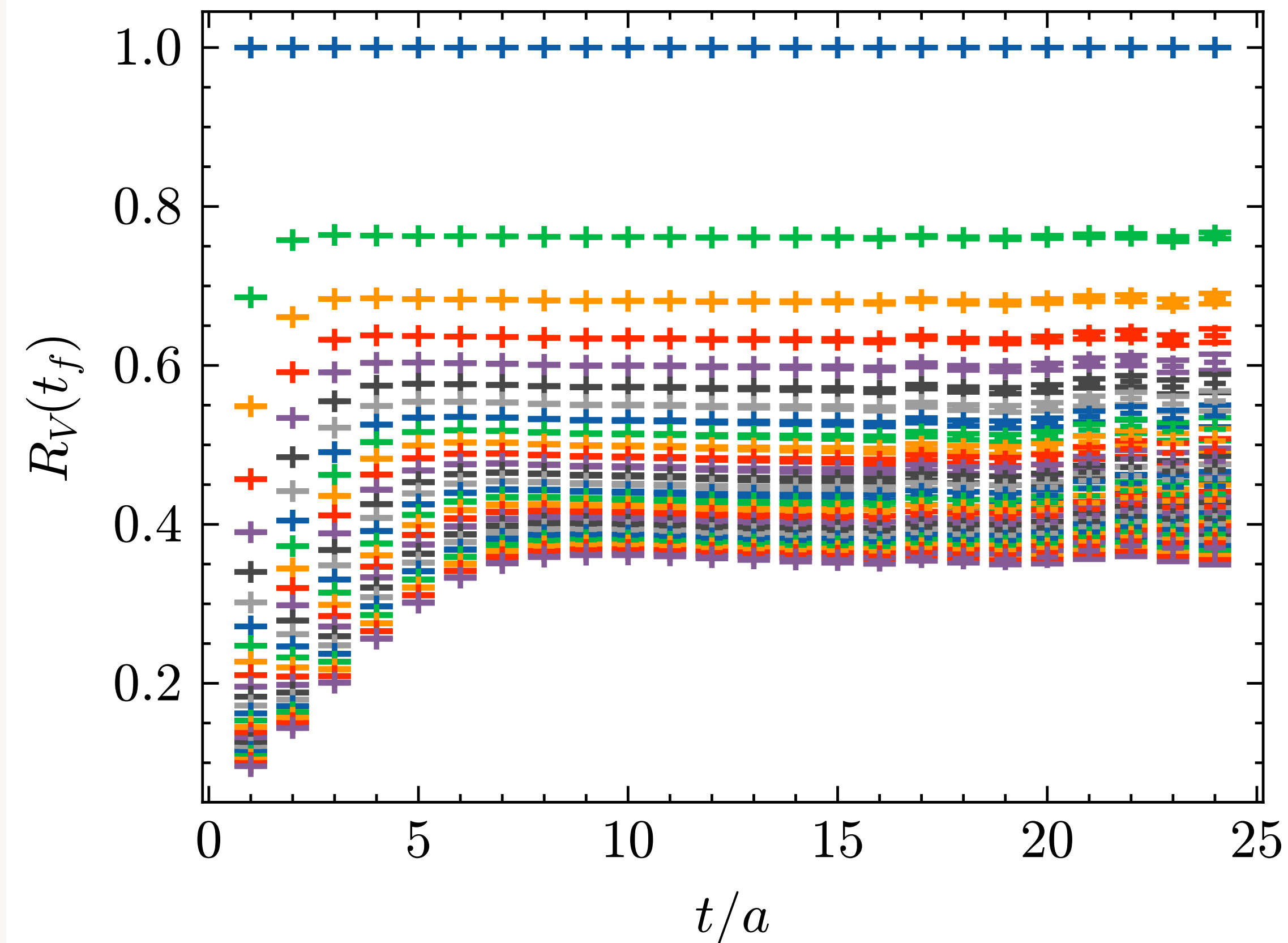
$$Z_V(\mu, a) \langle V_\mu^L(t_f = 0, a) \rangle = \langle V_\mu^R(\mu) \rangle$$

$$\mathring{Z}_\chi^V(t_f, a) \equiv \frac{\zeta_V(\mu, t_f) \langle V_\mu^L(t_f = 0, a) \rangle}{\langle V_\mu^L(t_f, a) \rangle} Z_V(\mu, a) \equiv \frac{\zeta_V(\mu, t_f)}{\langle R_{V,\mu}^L(t_f, a) \rangle} Z_V(\mu, a)$$

Similar idea of fixing  $Z_V$ : for the vector case,  $\mathring{Z}_\chi^V(t_f, a)$  absorbs all the discrete effects.

# Renormalization of flowed quarks

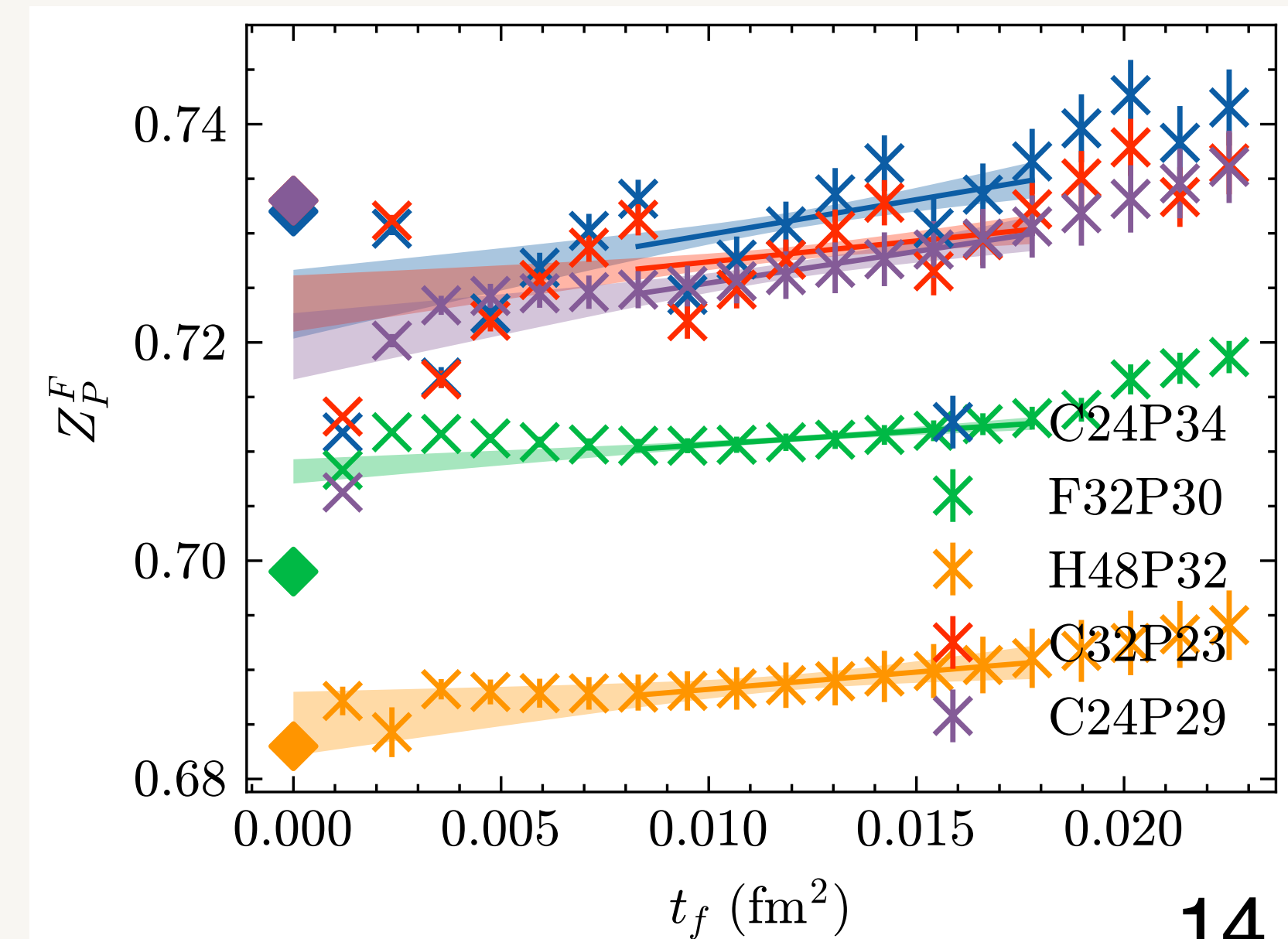
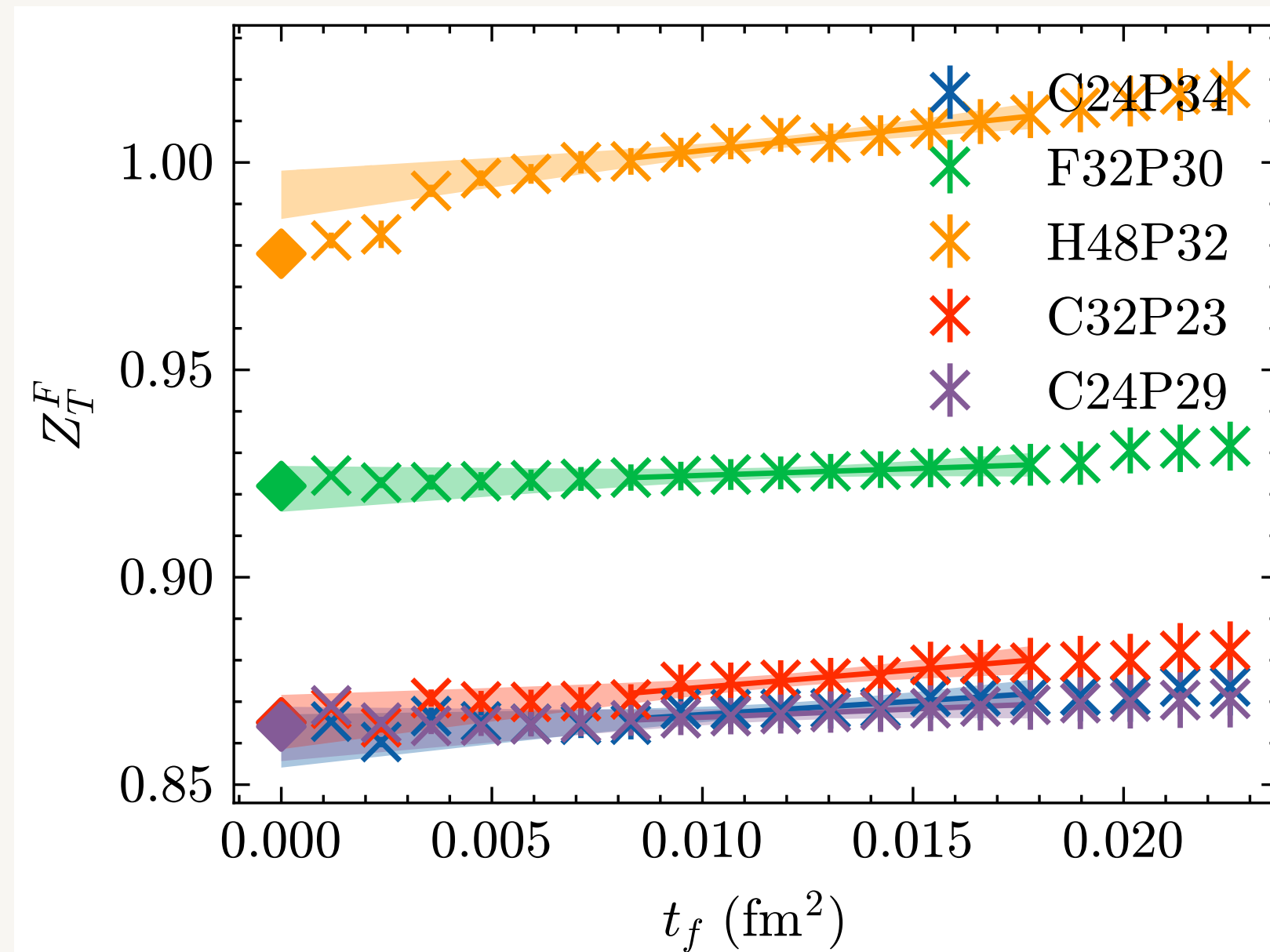
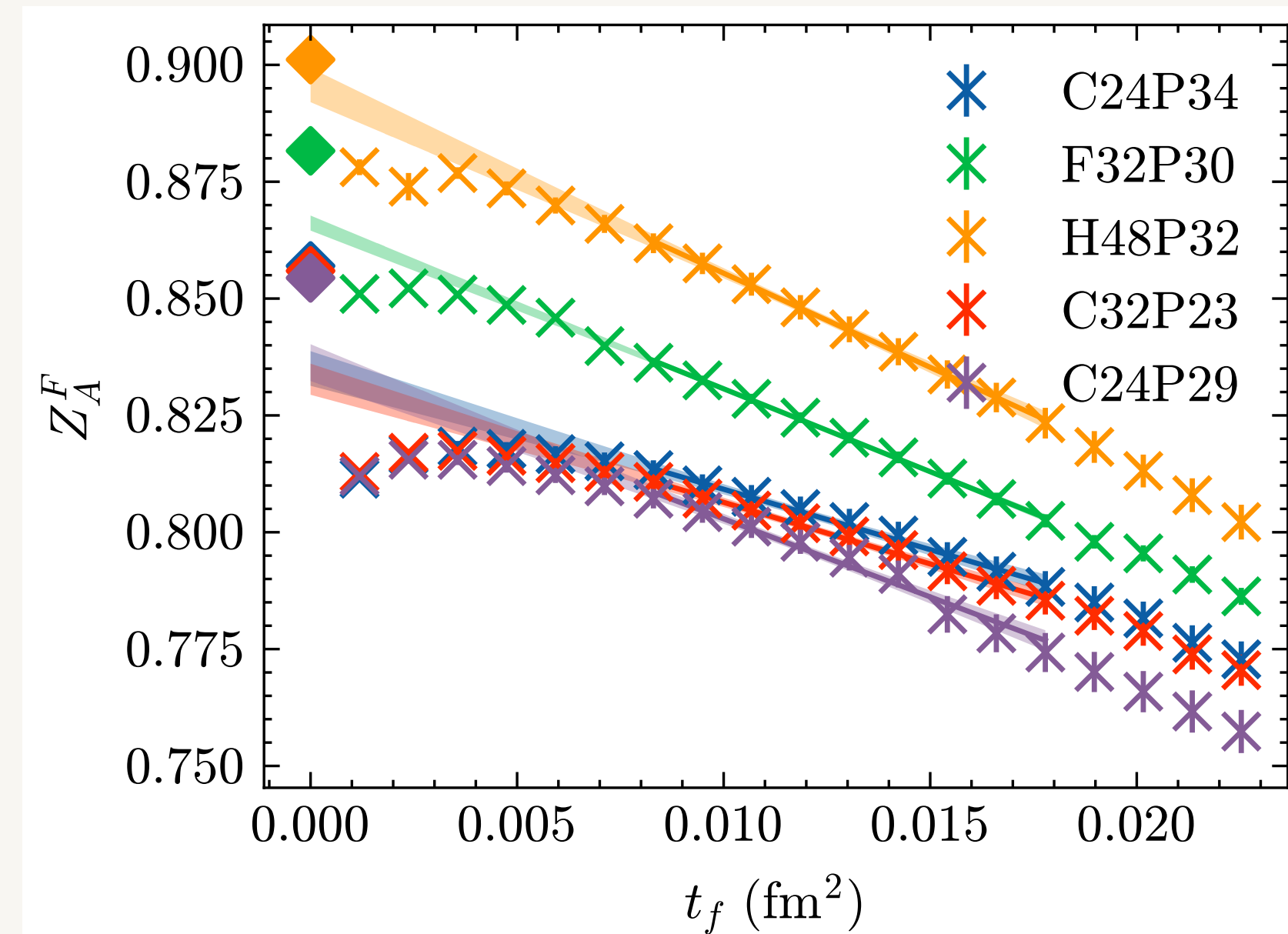
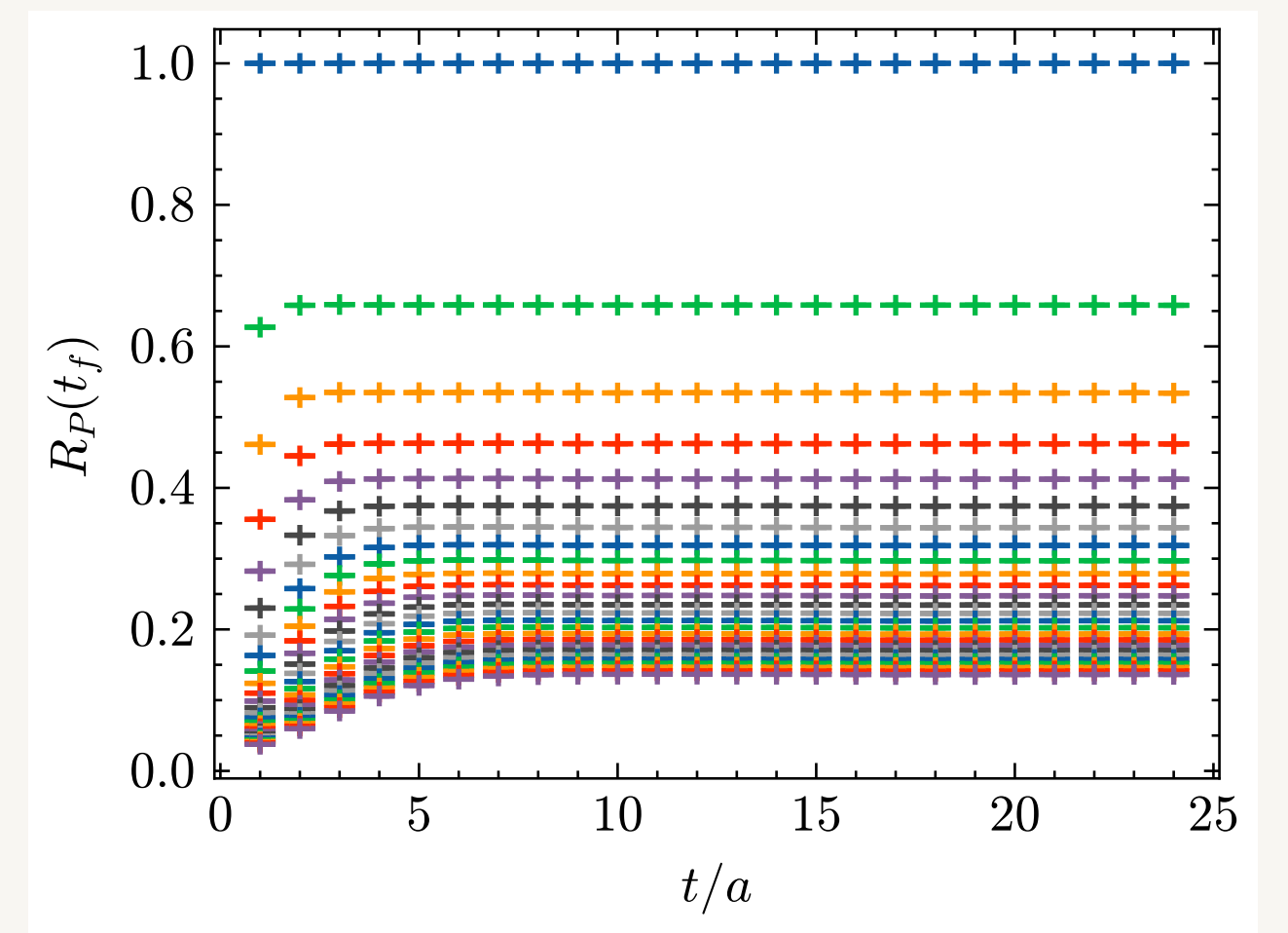
$$\mathring{Z}_\chi^V(t_f, a) \equiv \frac{\zeta_V(\mu, t_f) \langle V_\mu^L(t_f = 0, a) \rangle}{\langle V_\mu^L(t_f, a) \rangle} Z_V(\mu, a) \equiv \frac{\zeta_V(\mu, t_f)}{\langle R_{V,\mu}^L(t_f, a) \rangle} Z_V(\mu, a)$$



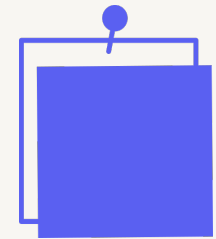
# Check against RI/MOM

$$Z_O^{\text{RI/MOM}}(a, \mu) \sim \frac{O^R(\mu)}{O^L(t_f = 0, a)}$$

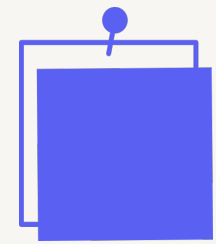
$$Z_O^{\text{F}}(a, \mu) = \frac{\lim_{t_f \rightarrow 0} \lim_{a \rightarrow 0} \left[ \frac{\dot{Z}_\chi^V(t_f, a) O^L(t_f, a)}{\zeta_O(\mu, t_f)} \right]}{O^L(t_f = 0, a)} \sim \lim_{t_f \rightarrow 0} \left[ \frac{\dot{Z}_\chi^V(t_f, a) R_O^L(t_f, a)}{\zeta_O(\mu, t_f)} \right] + \mathcal{O}(a)$$



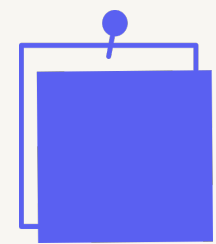
# Summary and Take-home



**Validation:** Rigorous numerical tests confirm the viability of Wilson flow for renormalizing quark bilinear and gluonic operators.



**Performance:** The scheme yields consistent results within current systematic uncertainties.



**Outlook:** Refinement of systematic controls and extensions to physical observables are currently in progress.

**Thank You!**