

Group Meeting Presentations

Measurements of inclusive D^0 -meson production in isobar collisions at $\sqrt{s_{NN}} = 200$ GeV

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Sep. 12th, 2025

Outline

- ⊙ Motivation and STAR Experiment
- ⊙ Energy loss
 - Measurements of D^0 p_T spectra and R_{AA} in isobar collisions at $\sqrt{s_{NN}} = 200$ at RHIC
- ⊙ Collectivity
 - Measurement of D^0 elliptic flow v_2 at mid-rapidity in isobar collisions
- ⊙ Summary

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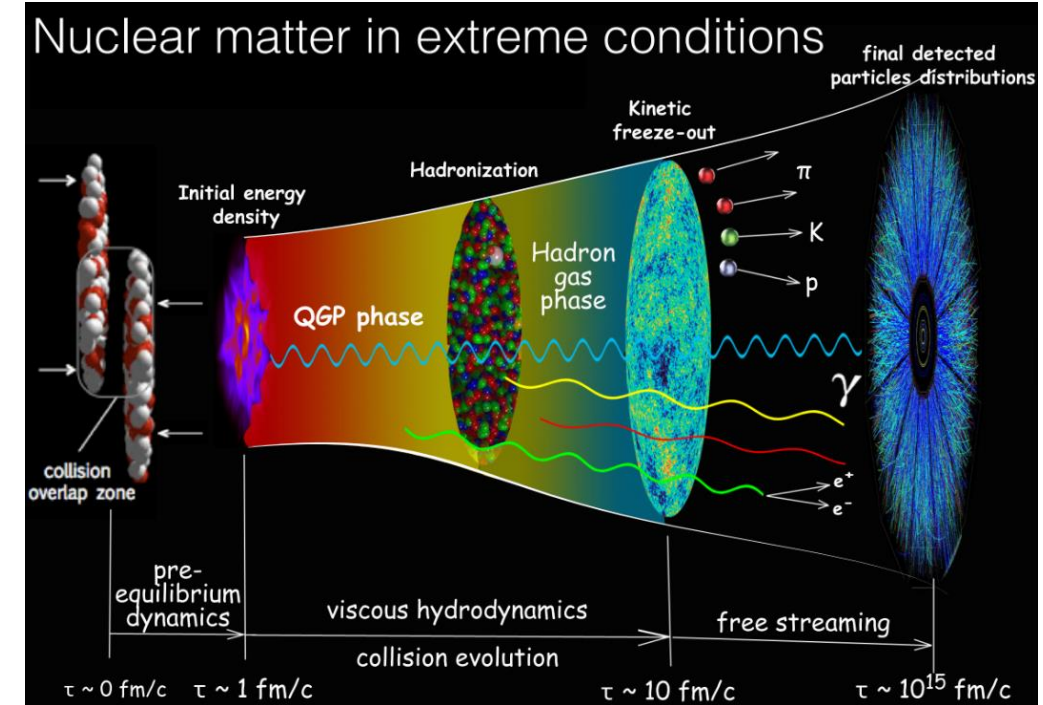
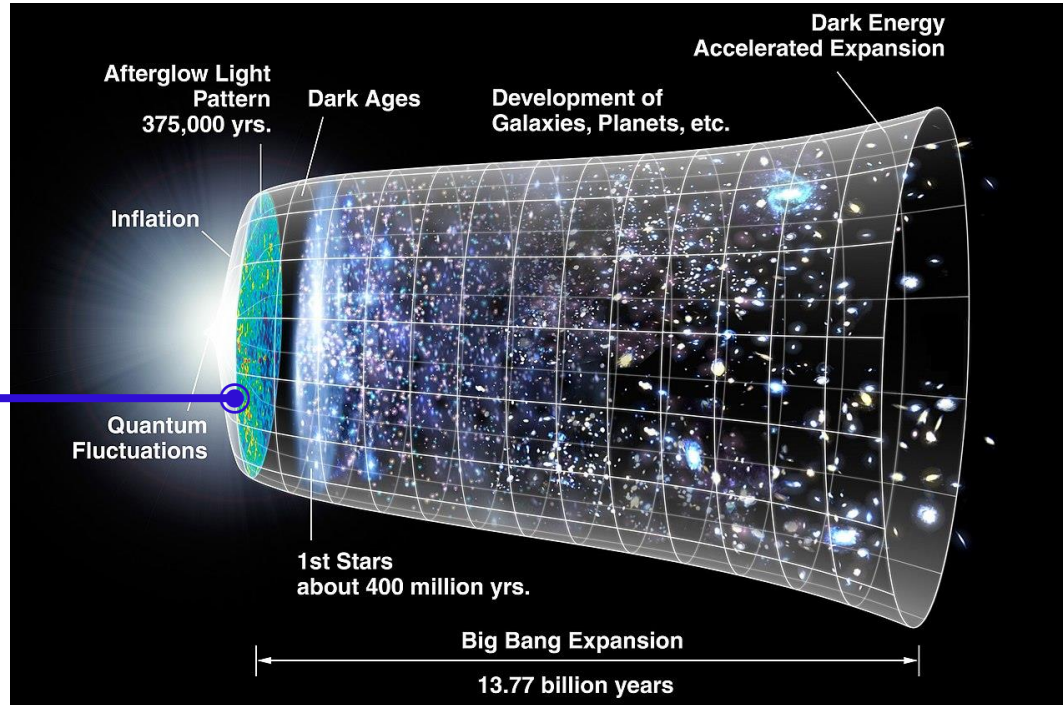
- Collectivity

 - Measurement of D^0 elliptic flow v_2 at mid-rapidity in isobar collisions

- Summary

Relativistic Heavy Ion Collision

Energies involved in the Quark epoch are **directly accessible** in RHIC (“Little” Big Bang)



● QGP is the hottest ($\sim 10^{12} \text{ K}$), smallest, and the most perfect fluid ($0.08 \sim 0.2 \hbar/k_B$) in nature!

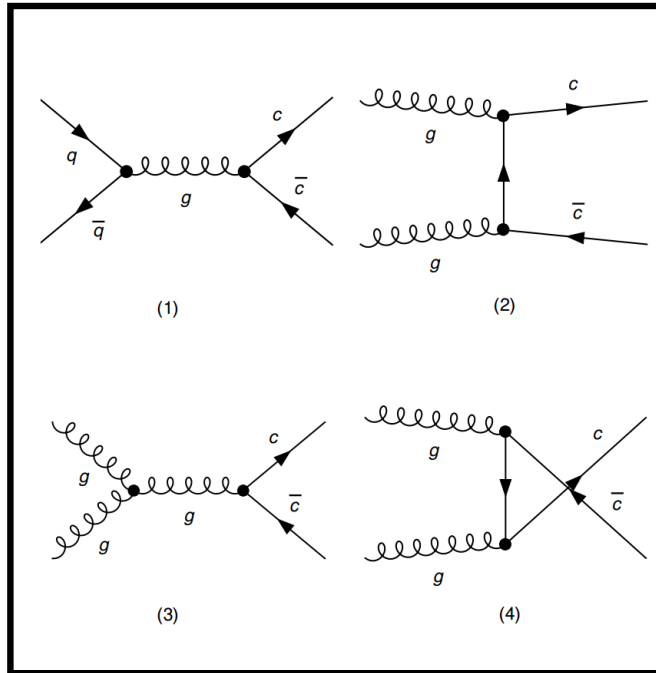
● Hadronization and chemical freeze-out stages

- Confined again & inelastic interactions & The relative yields of different hadron species are fixed.

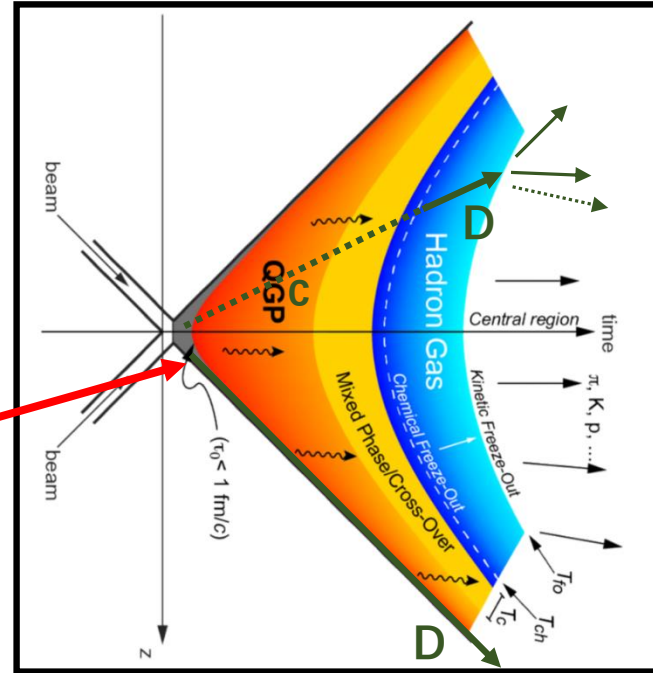
● Kinetic freeze-out stages

- elastic interactions & The momenta of particles get fixed.

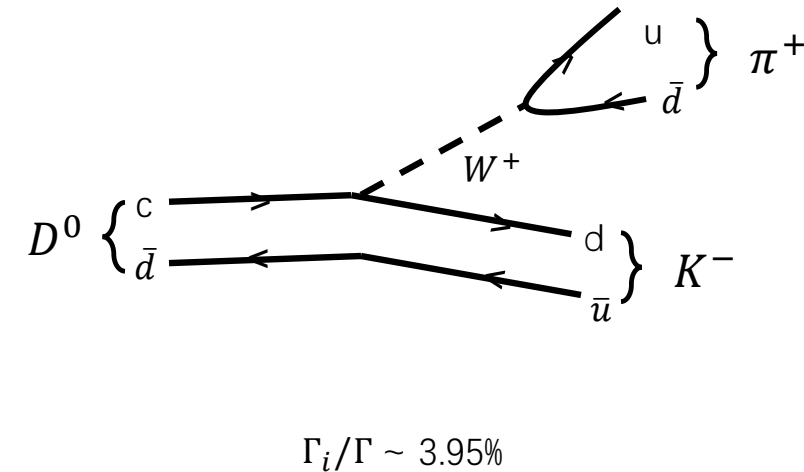
Open heavy-flavour production



$< 1 \text{ fm}/c$, initial partonic hard scattering



$m_c > \Lambda_{QCD}$, an ideal probe with a high probability of penetrating QGP

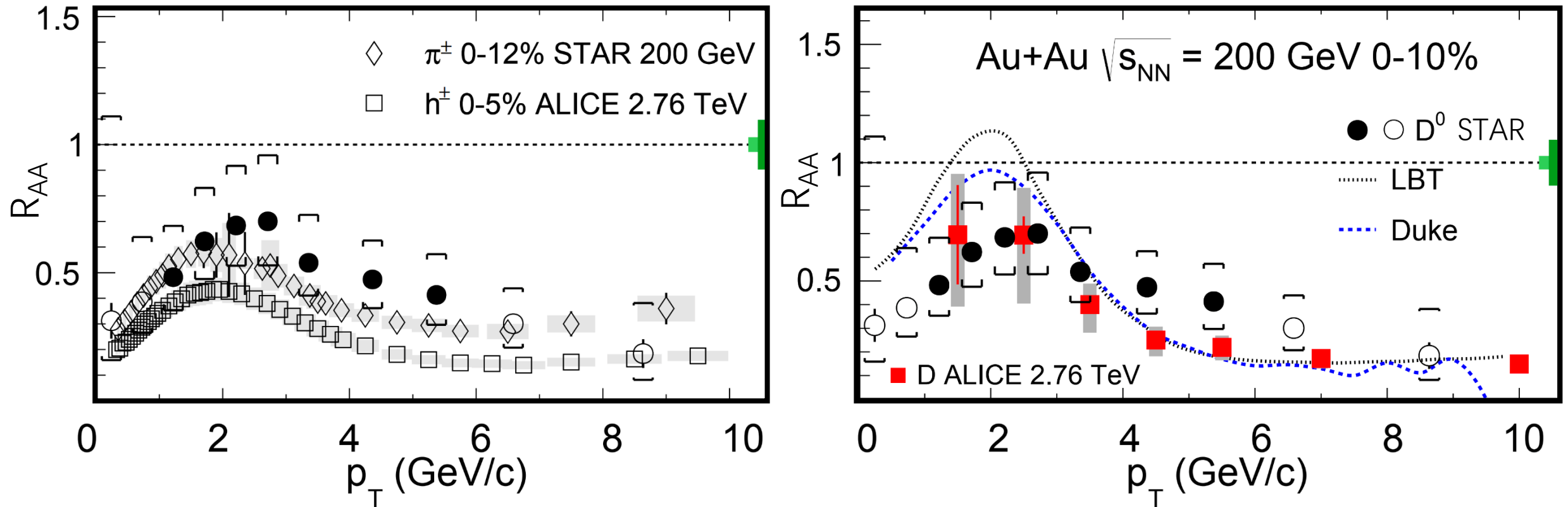


Open heavy-flavour production plays **a unique role** in heavy-ion physics:

- Production restricted to **early collision stages (mainly gluon fusion)** and **retain a “memory”** of their evolution through the QGP
- Consistent with the theoretical calculation (**pQCD in p+p collisions**) and partons have energy loss in the medium
 - Gluon radiation (dead cone effect; suppressed at $\theta < m_Q/E_Q$)
 - Collisional energy loss (**Brownian motion**)
- Heavy quarks **retain their flavour and mass identity**; can be “tagged” by the measurement of heavy-flavour hadrons

D^0 R_{AA} in Au+Au collisions

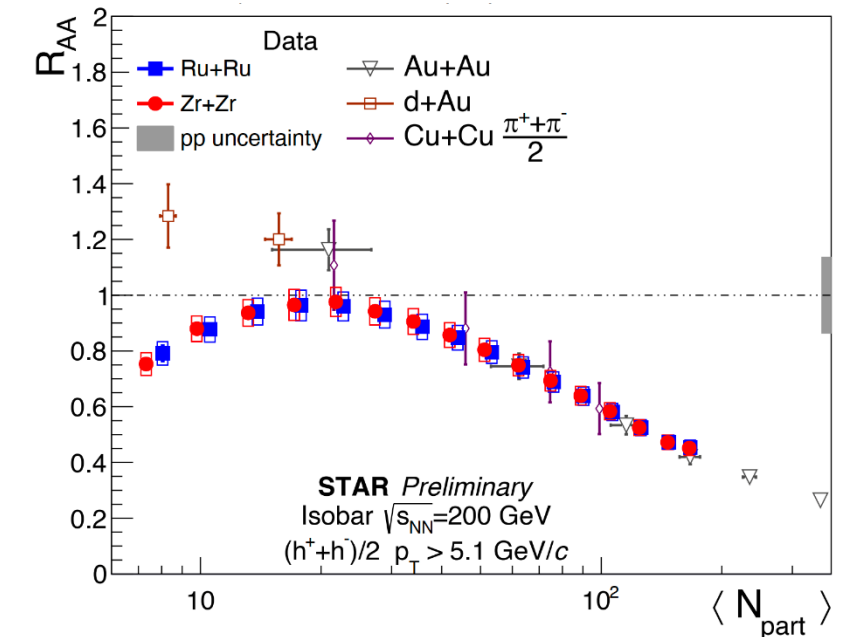
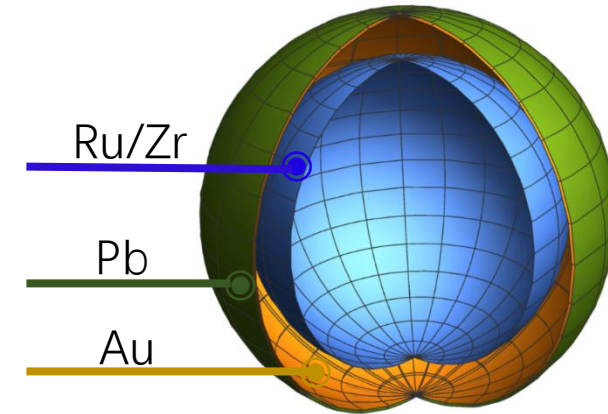
STAR, Phys. Rev. C 99, 034908



- Suppression level comparable to that of light hadrons in 0-10%
- Similar suppression in Au+Au and Pb+Pb
- What is the system size dependence of charm quark energy loss and hadronization?

How about system size dependence?

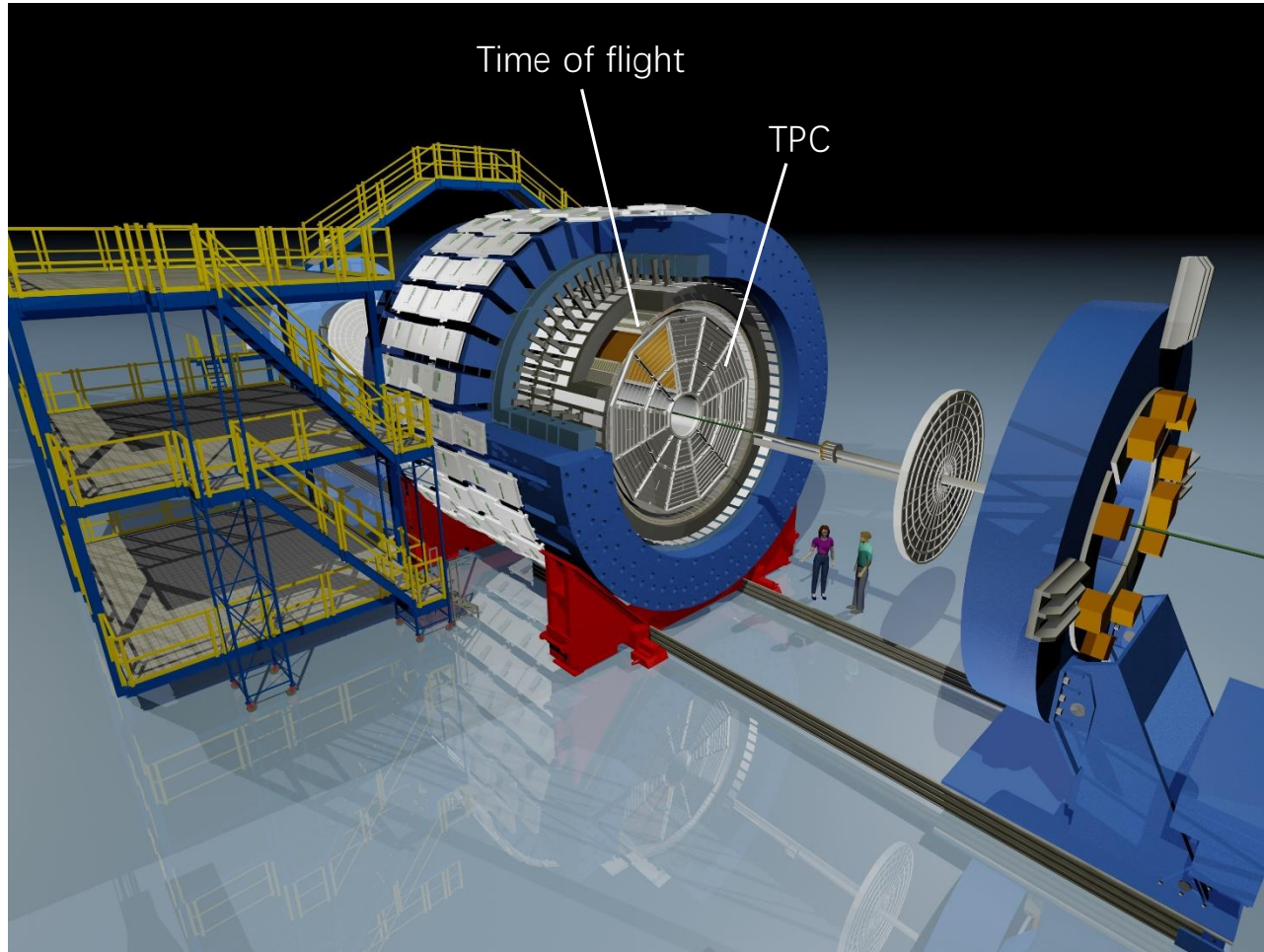
- STAR recorded large sample of **isobar collisions (Ru+Ru & Zr+Zr) at $\sqrt{s_{NN}} = 200$ GeV** in 2018
- Provide a great opportunity to study the system size dependence
- High- p_T charged hadrons: R_{AA} scales with $\langle N_{part} \rangle$ in different systems $\rightarrow$ driven by average energy density and volume
- How about D^0 ?



STAR Phys. Rev. Lett. 91, 172302 (2003) STAR Phys. Rev. C 81, 054907 (2010)

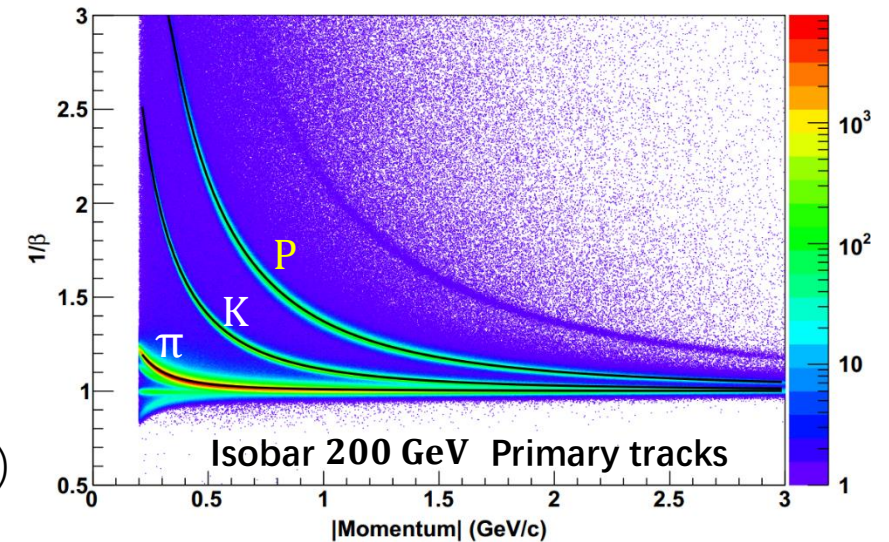
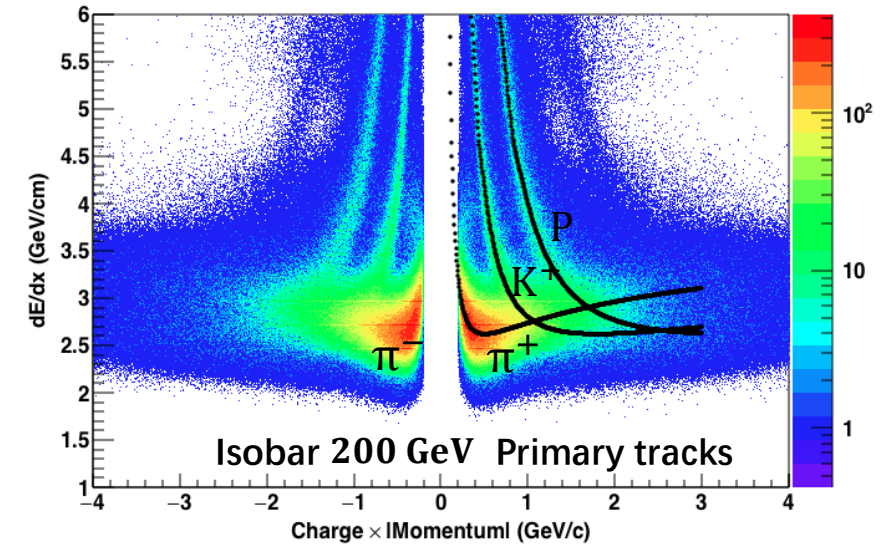
STAR Phys. Rev. Lett. 91, 072304 (2003)

STAR Detectors



- Time Projection Chamber (TPC)
~ 350 μm , vertex resolution with more than 1000 tracks

- Time Of Flight detector (TOF)
the precise $1/\beta$ extends PID reach



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Data set and Analysis cuts

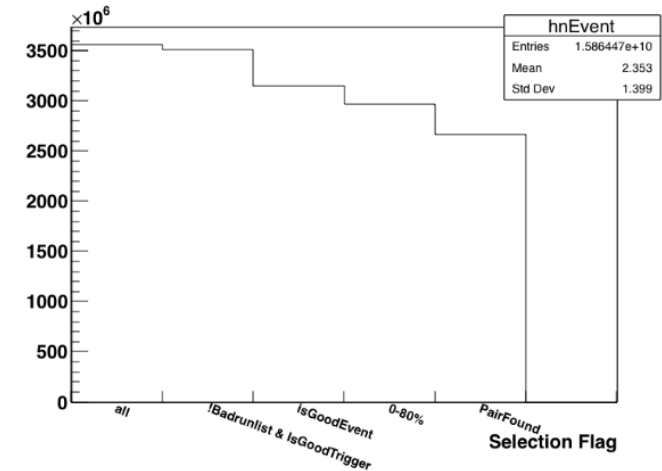
- **Dataset:** production_isobar_2018
- **Year:** 2018
- **Production tag:** P20ic
- **Triggers used:** 600001, 600011, 600021, 600031 (MB)
- **Embedding request id:** 20201503

hybrid PID functions

$$f_{\pi}^{max}(p) = \begin{cases} 5.1 - 2.25p, & p \leq 1.6 \\ 1.5, & p > 1.6 \end{cases}$$

$$f_K^{max}(p) = \begin{cases} 6.129 - 1.9316p, & p \leq 2.5 \\ 1.3, & p > 2.5 \end{cases}$$

$$f_K^{min}(p) = \begin{cases} -7.54 + 5.83p - 1.31p^2, & p \leq 1.7 \\ -1.4149, & p > 1.7 \end{cases}$$



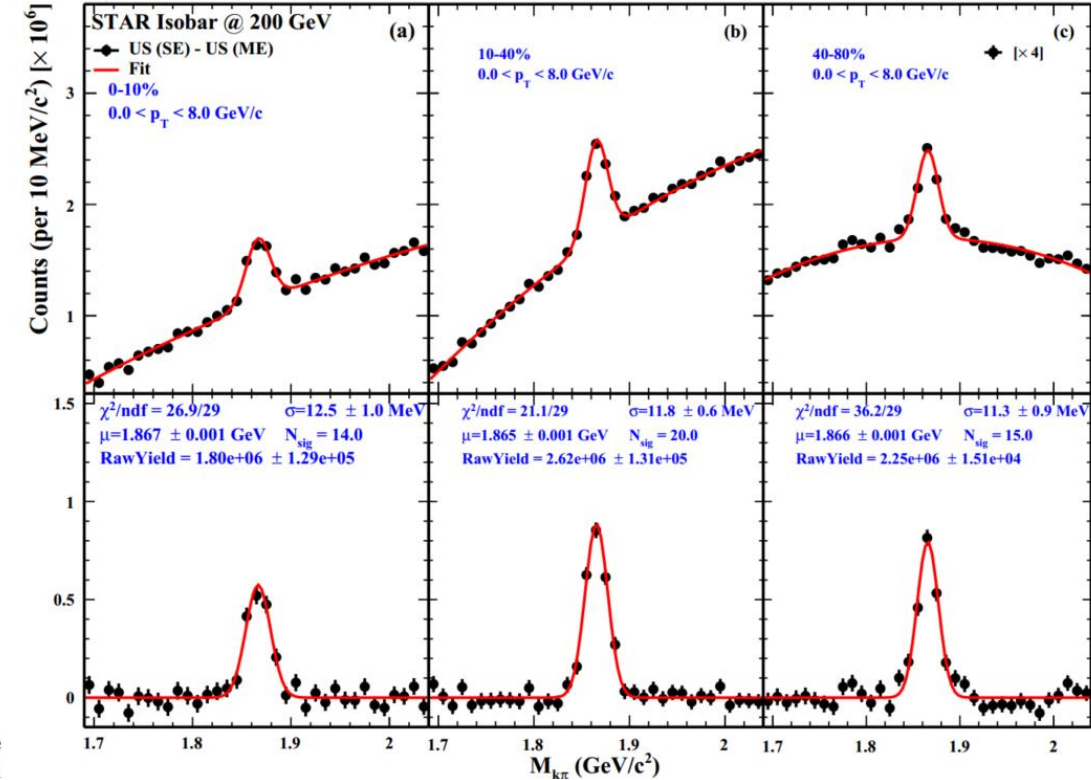
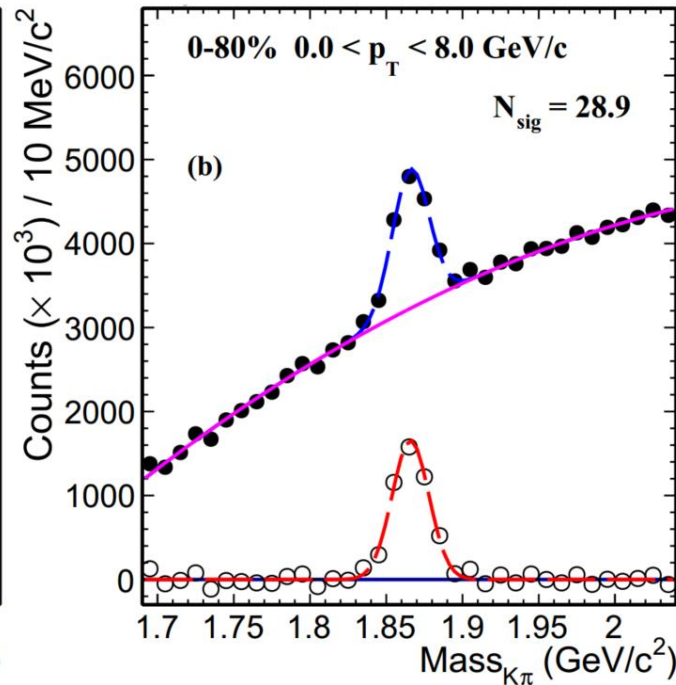
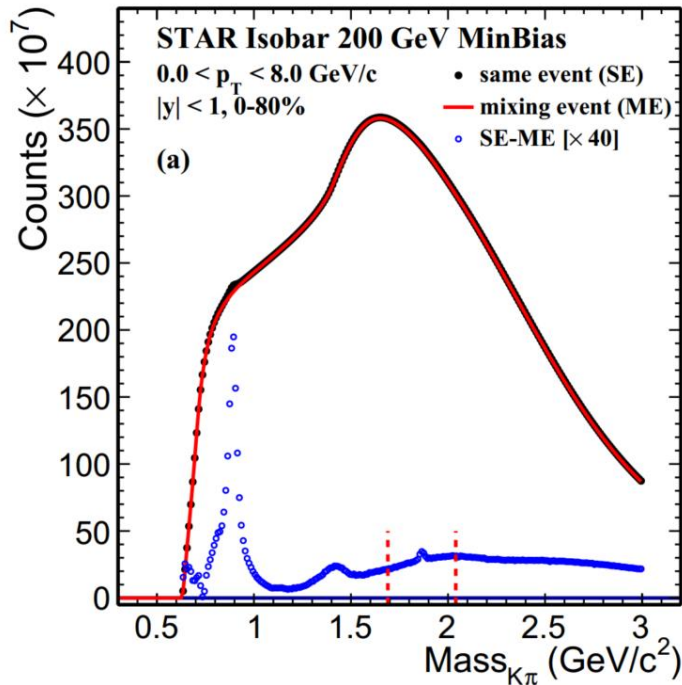
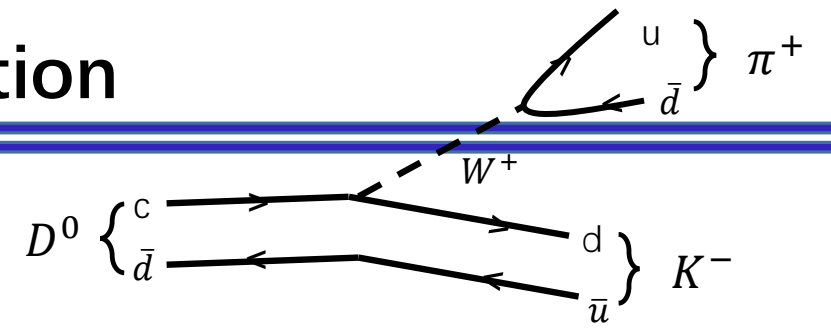
Event Level cuts	
Cuts	Value
$ V_z $	(-35,25) cm
V_r	< 2. cm
VzDiff	< 3. cm
!Badrun_list	
GoodTrigger	
!isBadRun()	
passnTofMatchRefmultCut()	
0-80%	
# Events	2.7B

Track quality cuts	
Cuts	Value
p_T (GeV/c)	> 0.6
TPC Nhits	≥ 20
nHits/HitsMax	[0.52, 1.2]
gDCA	≤ 2.0 cm
$ \eta $	< 1.
nHitsDedx	> 10.
charge	± 1
Primary tracks	

PID cuts	
Cuts	Value
btofYLocal (TOF matched)	[-1.8, 1.8]
$p < 1.6 \ \&\& \ \beta > 0$	hybrid PID functions
$p \geq 1.6 \ \&\& \ \beta > 0$ (Pion)	hybrid PID functions
$p \geq 1.6 \ \&\& \ !\text{TofisAvailble}$	$ \sigma_K < 2. \ \&\& \ \sigma_{\pi} < 2.$
$ y _{D^0}$	< 1.0

D^0 signal reconstruction

hadronic modes: $D^0 \rightarrow K^- + \pi^+$; $\bar{D}^0 \rightarrow K^+ + \pi^-$ ($\Gamma_i/\Gamma \sim 3.95\%$)



- The K, π invariant mass distribution with centrality 0-80% and p_T range 0-8 GeV/c at midrapidity.
- The mix-event method can well reproduce the combination background (solid red line).

- D^0 signal at 0-10% (a), 10-40% (b), 40-80% (c) centrality bins with transverse momentum range 0-8 GeV/c at midrapidity.

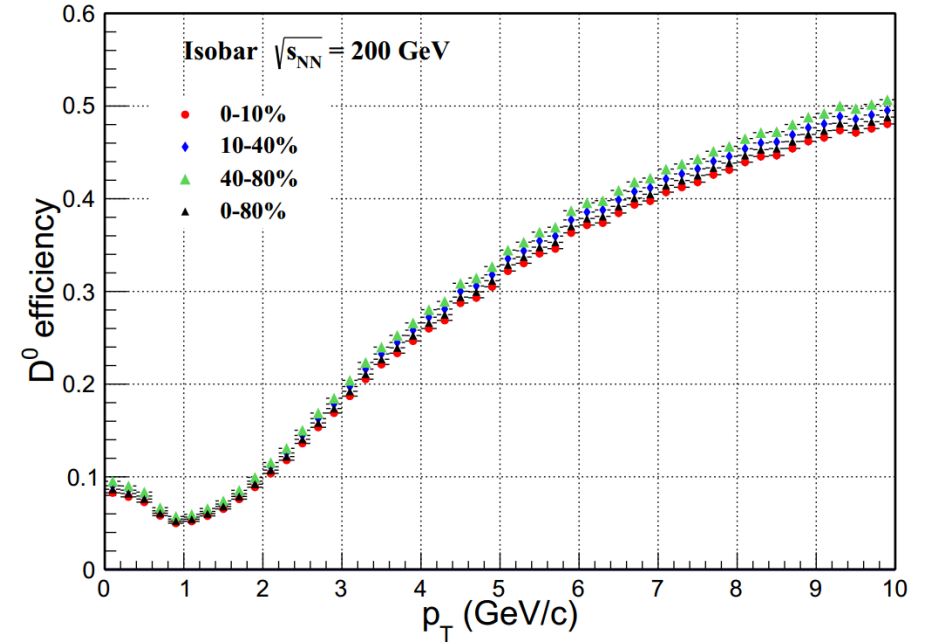
Efficiency correction procedures

$$\frac{d^2 N}{2\pi p_T dp_T dy} = \frac{\Delta N^{raw} / \epsilon_{D^0}^{reco} / 2}{2\pi p_T \Delta p_T \Delta y \times N_{events} \times B.R.}$$

$$\epsilon_{D^0}^{reco} = \epsilon_{Accept} \times \epsilon_{TPC} \times \epsilon_{PID}$$

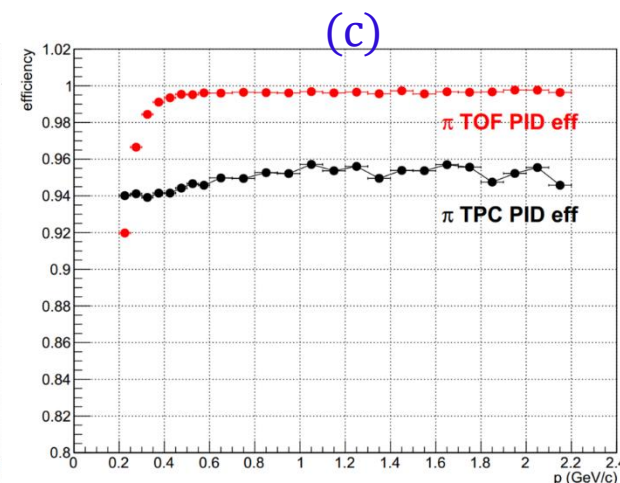
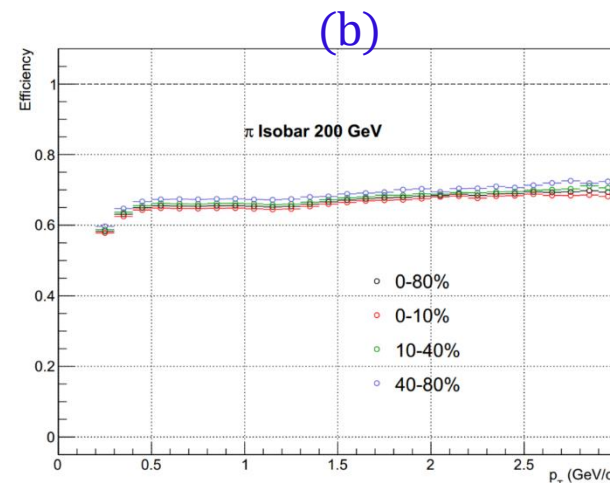
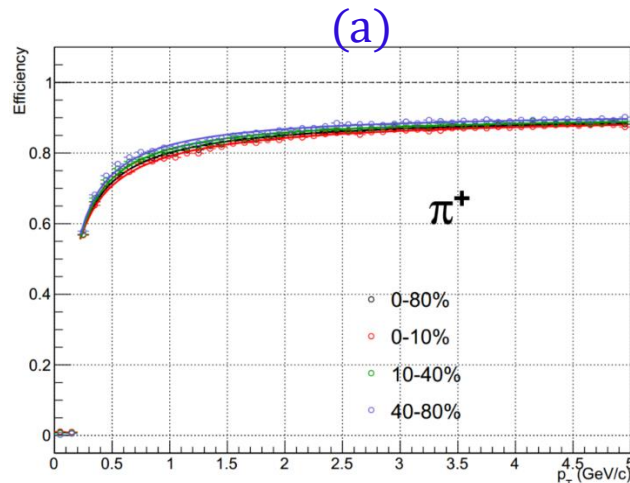
$$\epsilon_{PID} = \epsilon_{n\sigma_X} \cdot \epsilon_{TOF} \cdot \epsilon_{n\sigma_X}^{TOF} + \epsilon_{n\sigma_X} \cdot (1 - \epsilon_{TOF})$$

- ΔN^{raw} : the raw yield measured in the bin $\Delta p_T \Delta y$;
- $\epsilon_{Accept} \times \epsilon_{TPC}$: TPC acceptance and tracking efficiency (embedding);
- ϵ_{PID} : particle identification efficiency (data, K_S^0 , ϕ , Λ^0).



• D^0 efficiency as a function of D^0 transverse momentum at different centrality bins.

- (a) TPC tracking efficiency for Pion
- (b) TOF matching efficiency for Pion
- (c) Pion TPC PID efficiency (black circles) and TOF PID efficiency (red circles)



Systematic uncertainties

● Signal extraction

- The difference between the fitting and counting methods (1.69, 2.04) GeV/c^2 ;
- The order of polynomial function to depict the residual background (Pol3);
- Signal fit range (1.73, 2.00) GeV/c^2 ;
- p_T cut variation for daughter particles;
- Mix-event like-sign normalization factor

● TPC tracking

- DCA: 2cm (default);
- nHitsFit: 20 (default)

● PID cuts 3%

● B.R. 0.5%

● p + p inelastic scattering cross section 8%

		0–10%	10–40%	40–80%
spectra	Raw yield	3.9-16.5%	3.7-10.5%	4.0-8.5 %
	Track p_T variation	10.6%	9.3%	7.3%
	ϵ_{PID}	3%	3%	3%
	TRE	2-6%	2-6%	2-6%
	BR	0.5%	0.5%	0.5%
R_{AA}	$\langle N_{\text{bin}} \rangle$	1.6%	0.6%	0.4%
	pp base	20.6-71.8%	20.6-71.8%	20.6-71.8%
Rcp (/40–80%)		0–10%		10–40%
	Raw yield	13.6-20.7%		12.4-16.5 %
Integrated cross section		$p_T > 0$		$p_T > 4 \text{ GeV}/c$
	Total	12.7-15.8%		12.0-15.2 %

p_T Spectra and Integrated yields

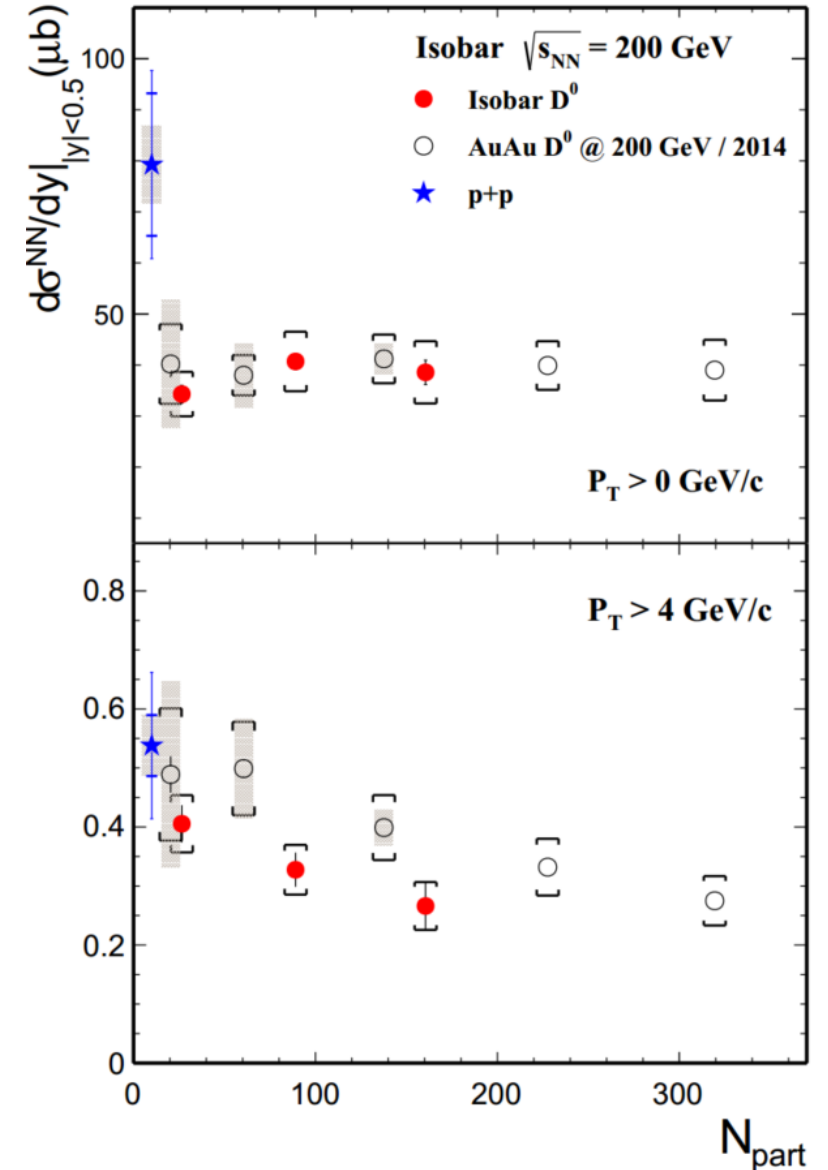
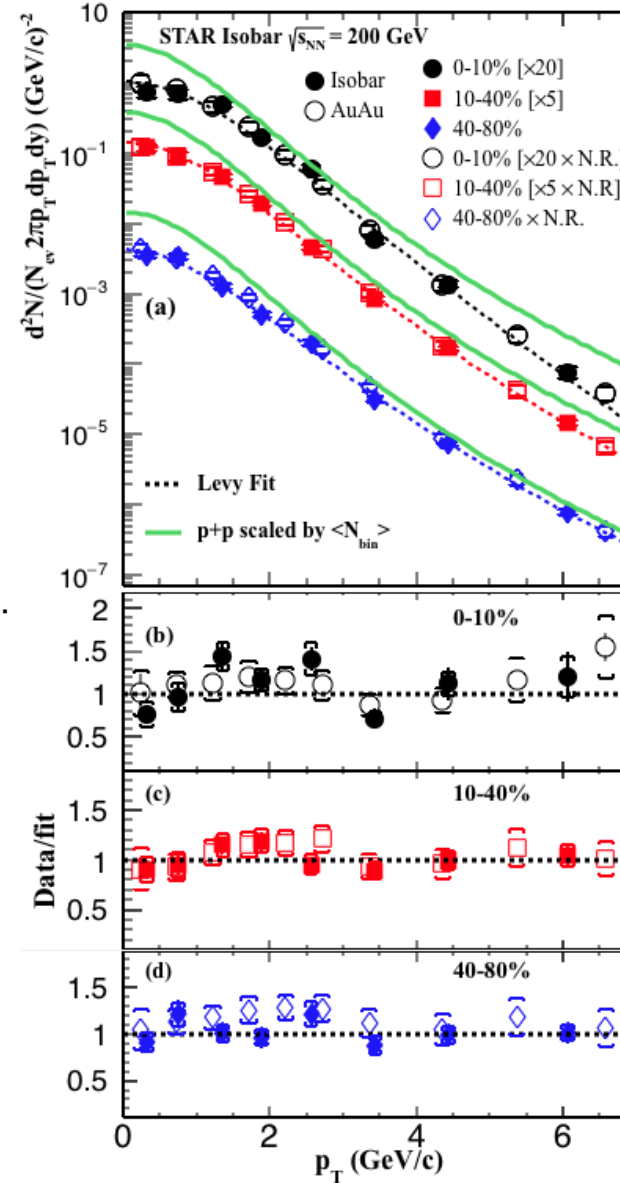
$$E \frac{d^3 N}{d\mathbf{p}^3} = \frac{d^3 N}{p_T dp_T dy d\phi} = \frac{d^2 N}{2\pi p_T dp_T dy}$$

$$\frac{d^2 N}{2\pi p_T dp_T dy} = \frac{\Delta N^{raw} / \epsilon_{D^0}^{tot} / 2}{2\pi p_T \Delta p_T \Delta y \times N_{events} \times B.R.} = \frac{\Delta N_{D^0}^{AA}}{2\pi p_T \Delta p_T \Delta y} = E \frac{d^3 \sigma_{D^0}^{AA}}{d\mathbf{p}^3}$$

- D^0 invariant yields at mid-rapidity ($|y| < 1$) vs. transverse momentum for different centrality classes in isobar (solid) and Au + Au (open) collisions at $\sqrt{s_{NN}} = 200$ GeV.

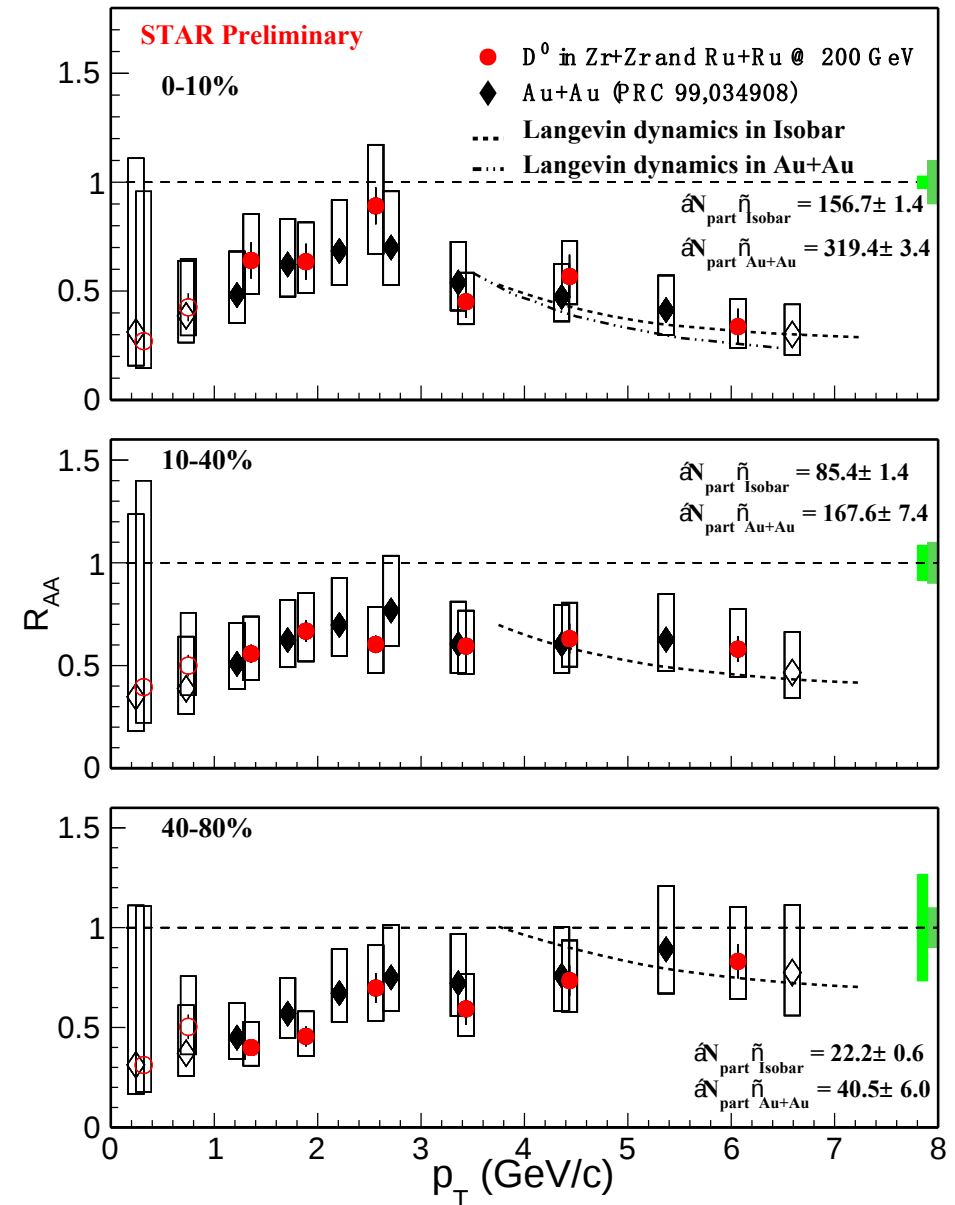
$$\left. \frac{d\sigma_{D^0}^{NN}}{dy} \right|_{y=0} = \left. \frac{dN_{D^0}^{AA}}{dy} \right|_{y=0} \times \frac{\sigma_{inel}^{pp}}{\langle N_{bin} \rangle}$$

- D^0 integrated cross sections per nucleon-nucleon collision in isobar.



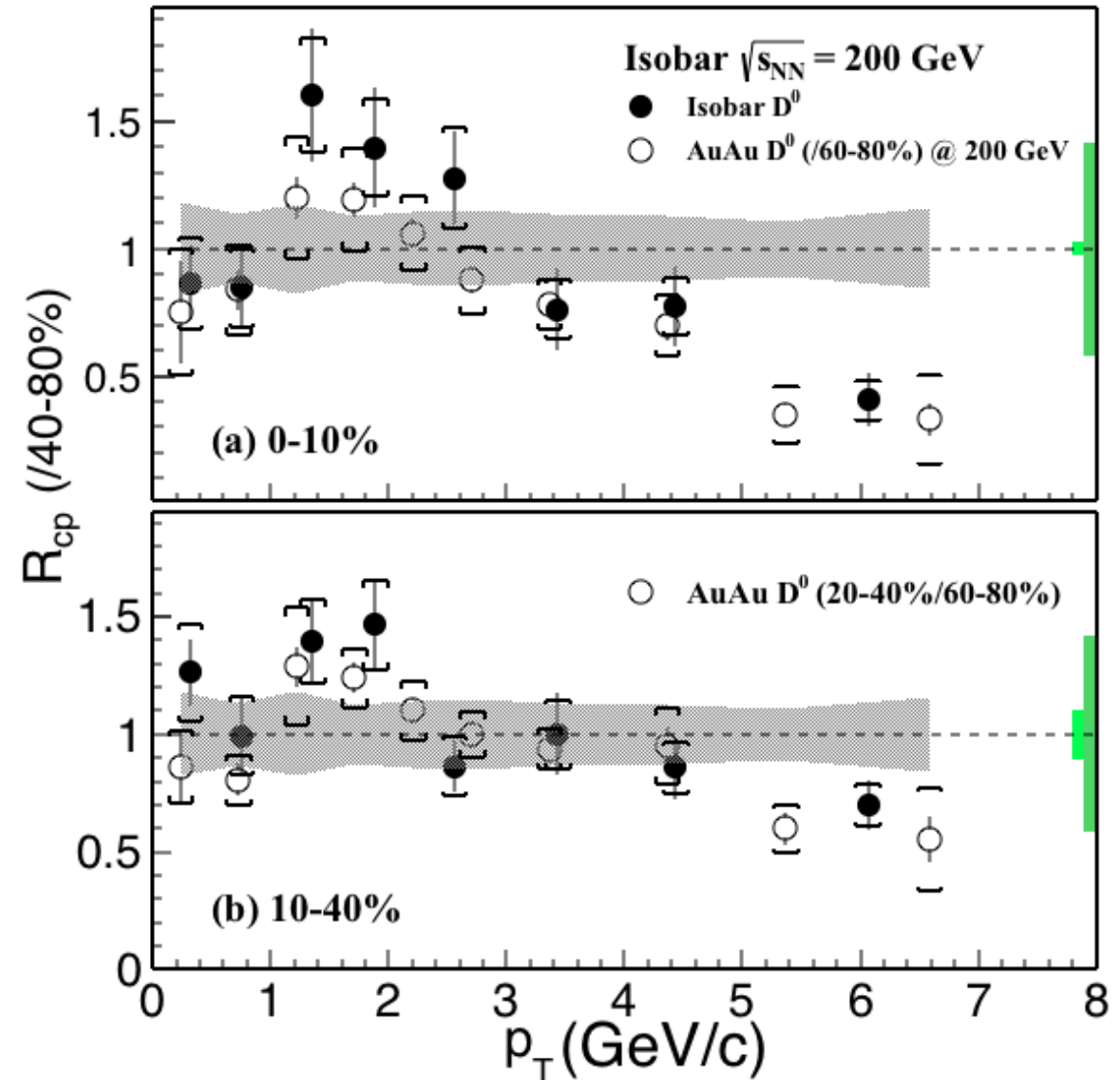
D^0 R_{AA} : Isobar vs. Au+Au

- Significant D^0 R_{AA} suppression observed at high- p_T in central collisions
- Suppression decreases from central to peripheral collisions $\rightarrow$ consistent with decreasing hot medium effects
- Similar suppression between isobar and Au+Au collisions in same centrality
- Theoretical model that includes collisional and radiative energy losses is consistent with observed suppression

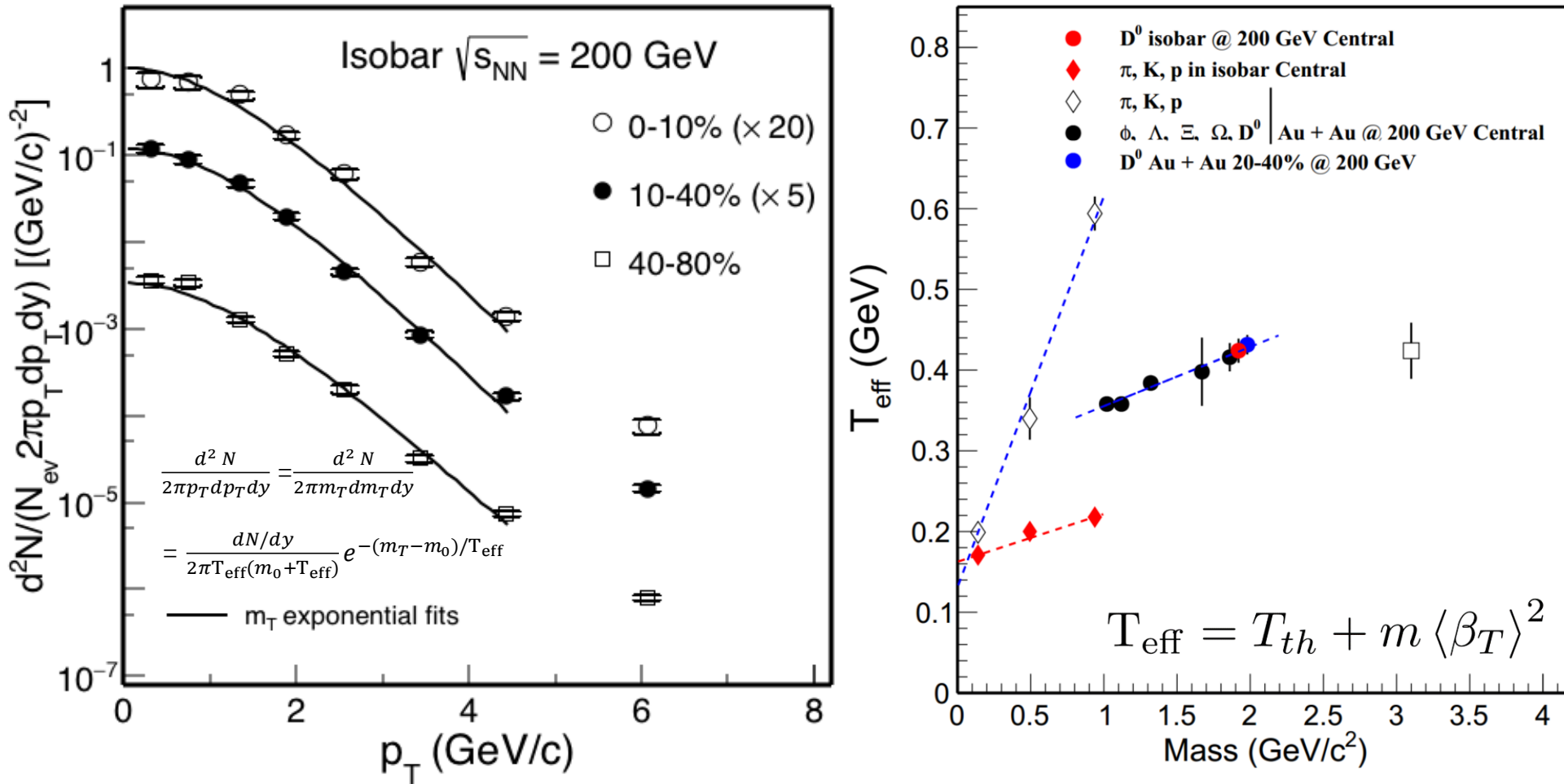


R_{cp} vs. p_T

- R_{cp} with the 40-80% spectrum as the reference for 0-10% and 10-40% centrality in isobar collisions compared to that of Au + Au results.
- Charm quarks suffer significant energy loss in both system evolution



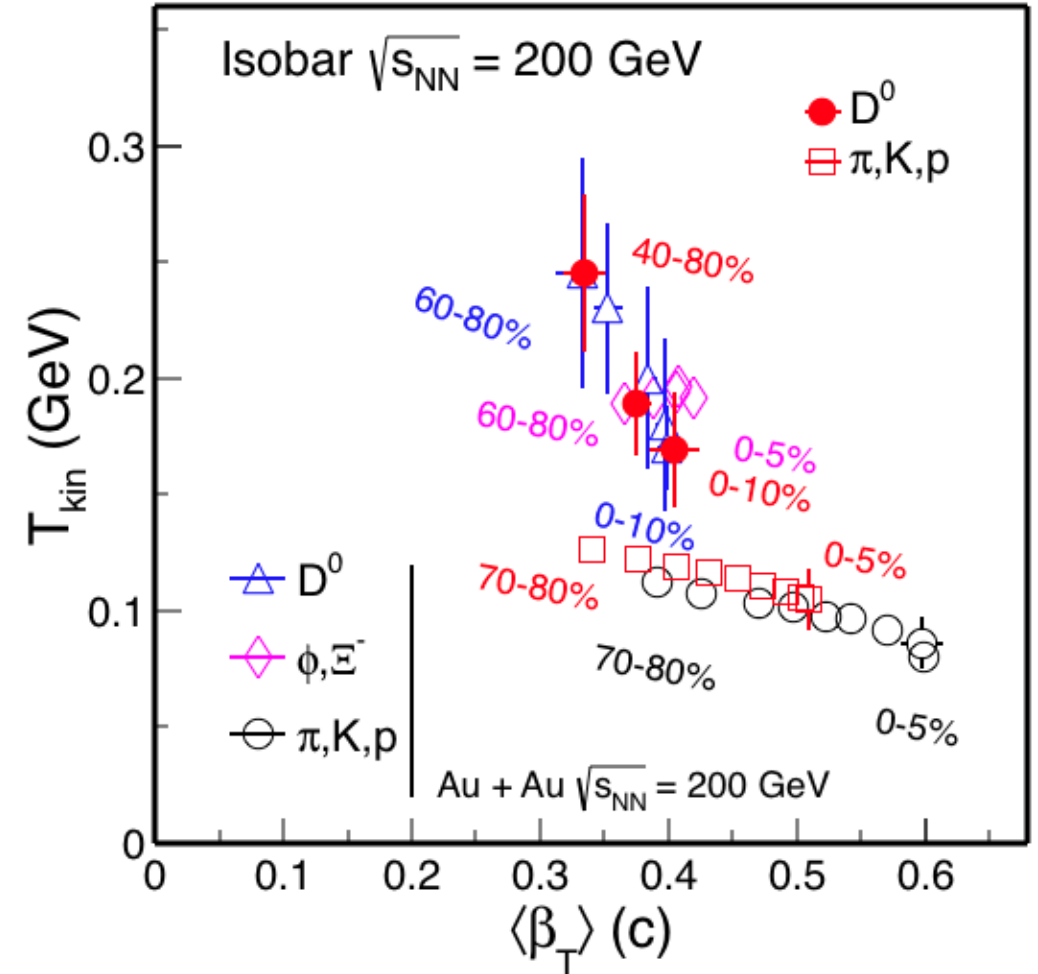
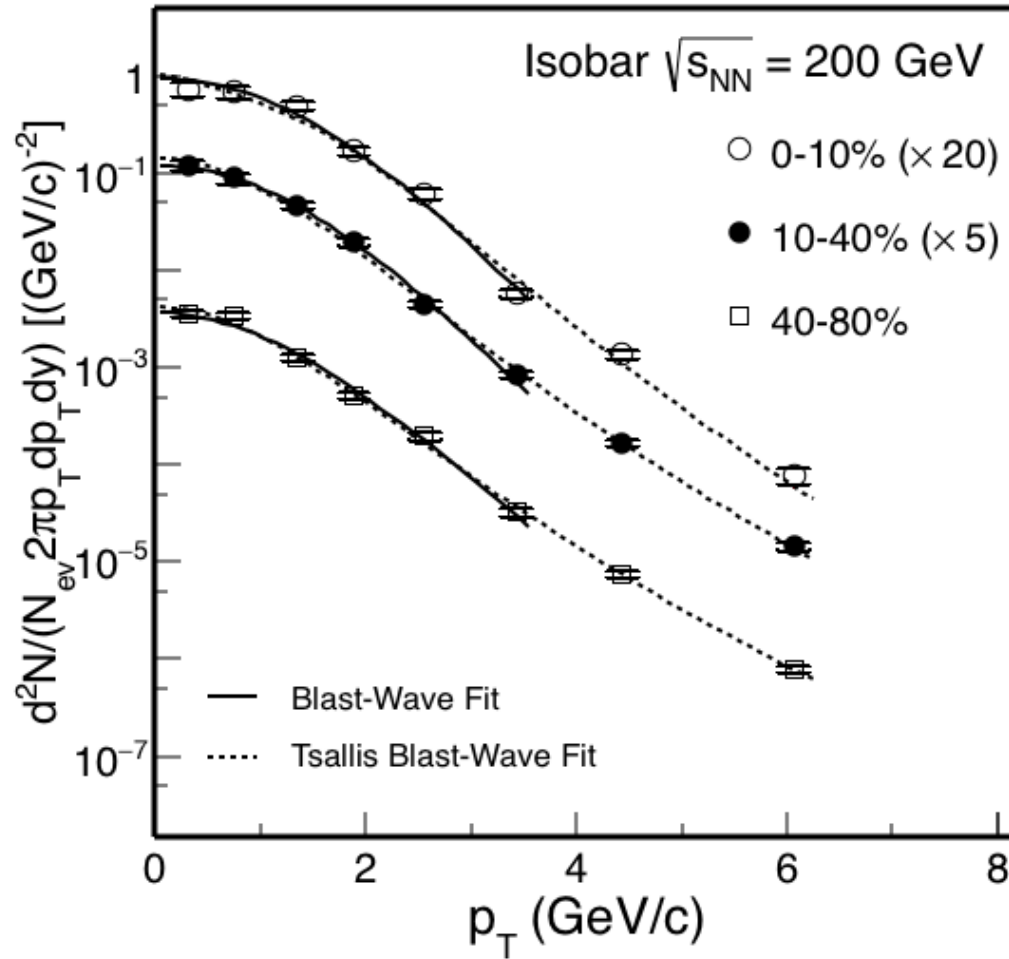
m_T Spectra and Collectivity



system	hadron	T_{eff}
isobar	π	0.176 ± 0.001
	K	0.203 ± 0.005
	p	0.218 ± 0.001
	D^0	0.424 ± 0.015
AuAu	π	0.199 ± 0.002
	K	0.340 ± 0.026
	p	0.594 ± 0.020
	D^0	0.416 ± 0.018
	D^0 (20-40%)	0.431 ± 0.012

- D^0 invariant yield at mid-rapidity ($|y| < 1$) vs. p_T for different centrality bins fitted with m_T distribution.
- T_{eff} for D^0 in central isobar collisions is consistent with that of Au + Au results. The red rhombus from left to right show the fit results of π , K, and proton in isobar collisions respectively.

Blast – Wave Fits



Tsallis blast-wave model fits parameters to D^0 p_T -differential spectra in isobar collisions.

Centrality	$\langle\beta_T\rangle$ (c)	$q - 1$
0-10%	0.282 ± 0.018	0.070 ± 0.007
10-40%	0.207 ± 0.030	0.080 ± 0.007
40-80%	0.189 ± 0.031	0.089 ± 0.005

- D^0 invariant yield at mid-rapidity ($|y| < 1$) vs. p_T for different centrality bins fitted with blast-wave function.
- D^0 freeze out temperature in isobar collisions are consistent with that of in Au + Au collisions for the same centrality. The open red rectangles from left to right shows the result of fitting π , K, and proton simultaneously from peripheral to central collisions.
- The average flow velocity increases with central collision, and $(q - 1)$ is also found to be close to zero.

Summary — D^0 p_T spectra and R_{AA}

- D^0 -meson productions are measured at mid-rapidity ($|y| < 1$) in 200 GeV isobar collisions
- The number of binary collisions scale effect of D^0 production cross section between isobar and Au + Au collisions is observed
- The strong suppression of D^0 nuclear modification factor R_{AA} is observed for $p_T > 3$ GeV/c in central isobar collisions, demonstrating that charm quarks suffer significant energy loss in the QGP.
- No significant systematic dependence of D^0 kinetic freeze-out properties in central collisions between isobar and Au + Au collisions within uncertainties is observed.

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Collective flow in heavy-ion collisions

Acceleration of collective flow

$$u^\nu \partial_\nu u^\rho = \frac{\Delta^{\rho\nu} \partial_\nu P}{e+P} \quad \text{Euler equation}$$

$$\underbrace{\partial_t u^i}_{\text{flow acceleration}} = \underbrace{\frac{\partial^i P}{e+P}}_{\text{inertia of fluid}} \quad \text{Spatial gradient of pressure}$$

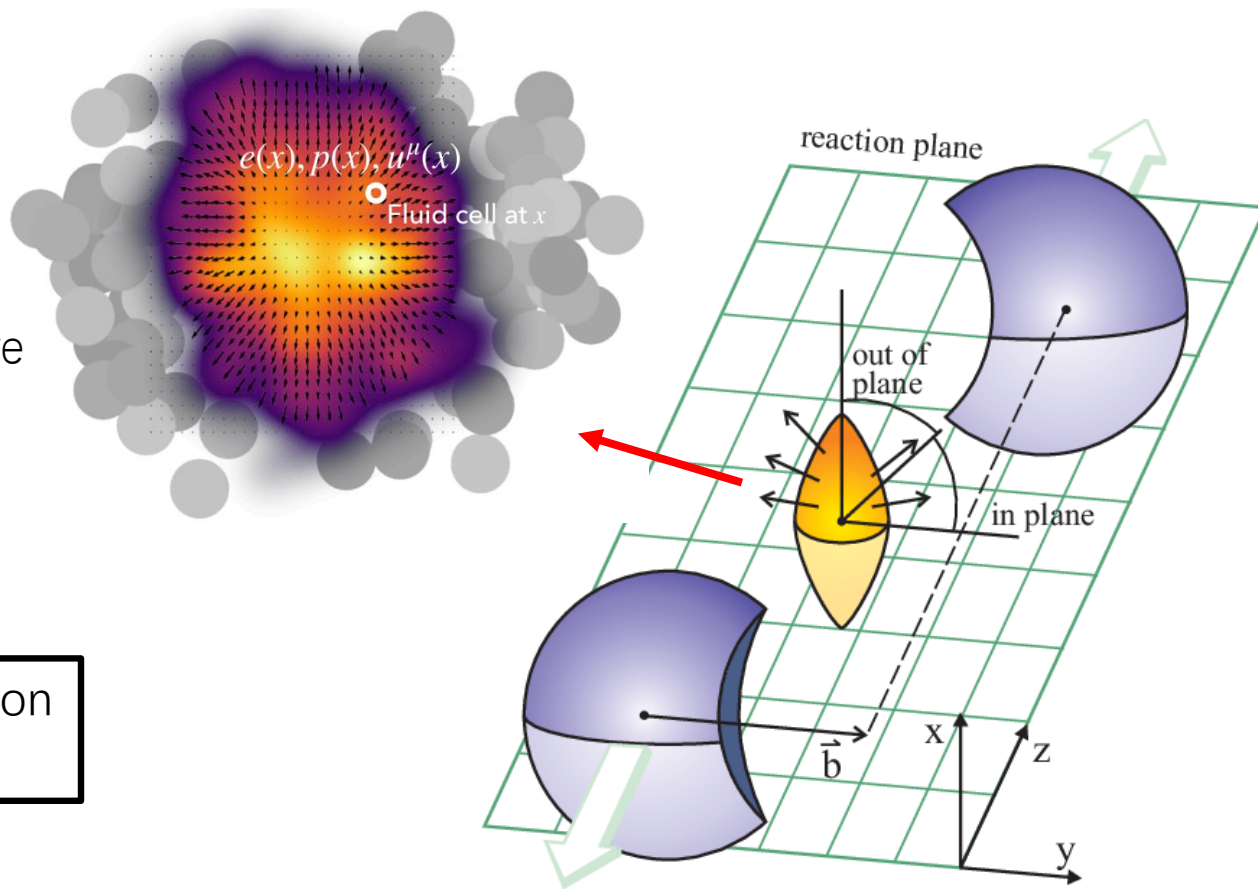
flow acceleration

Pressure gradients drive the hydrodynamic expansion

Anisotropic flow

$$\frac{dN}{d\phi} = \frac{N}{2\pi} \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\phi - \Psi_n)] \right)$$

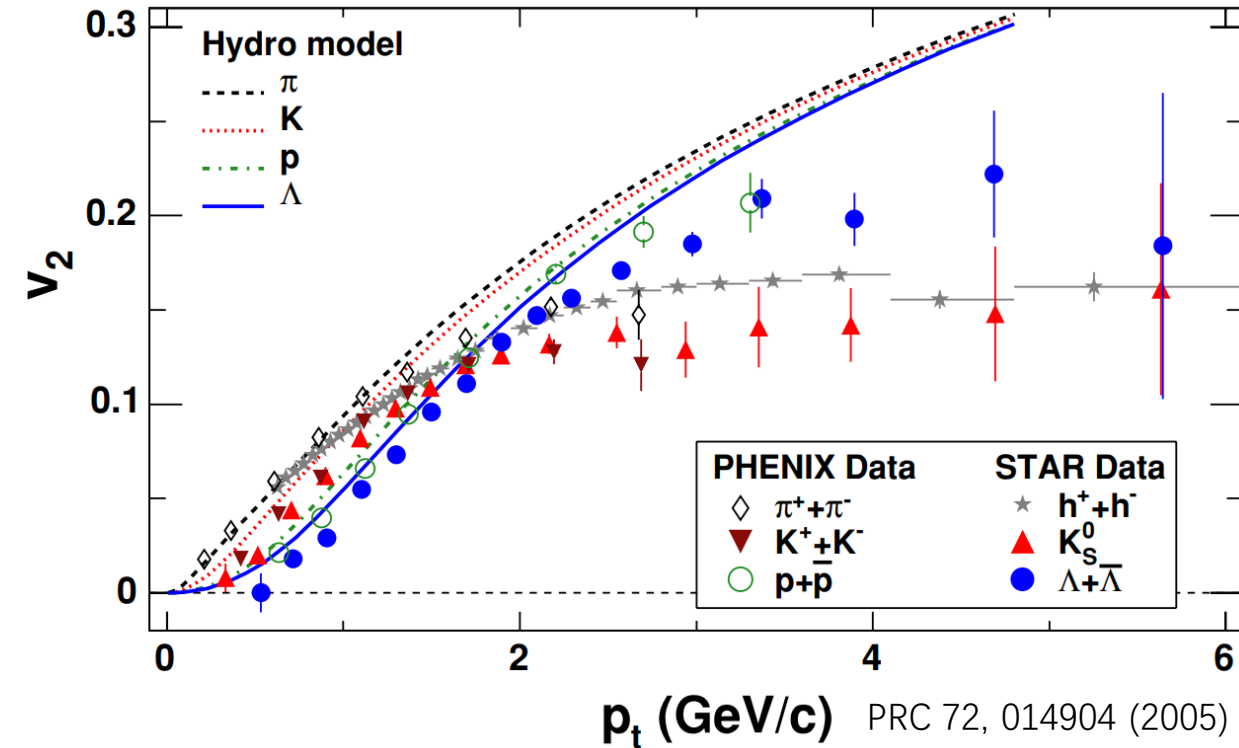
$$v_2 = \langle \cos[2(\phi - \Psi_2)] \rangle$$



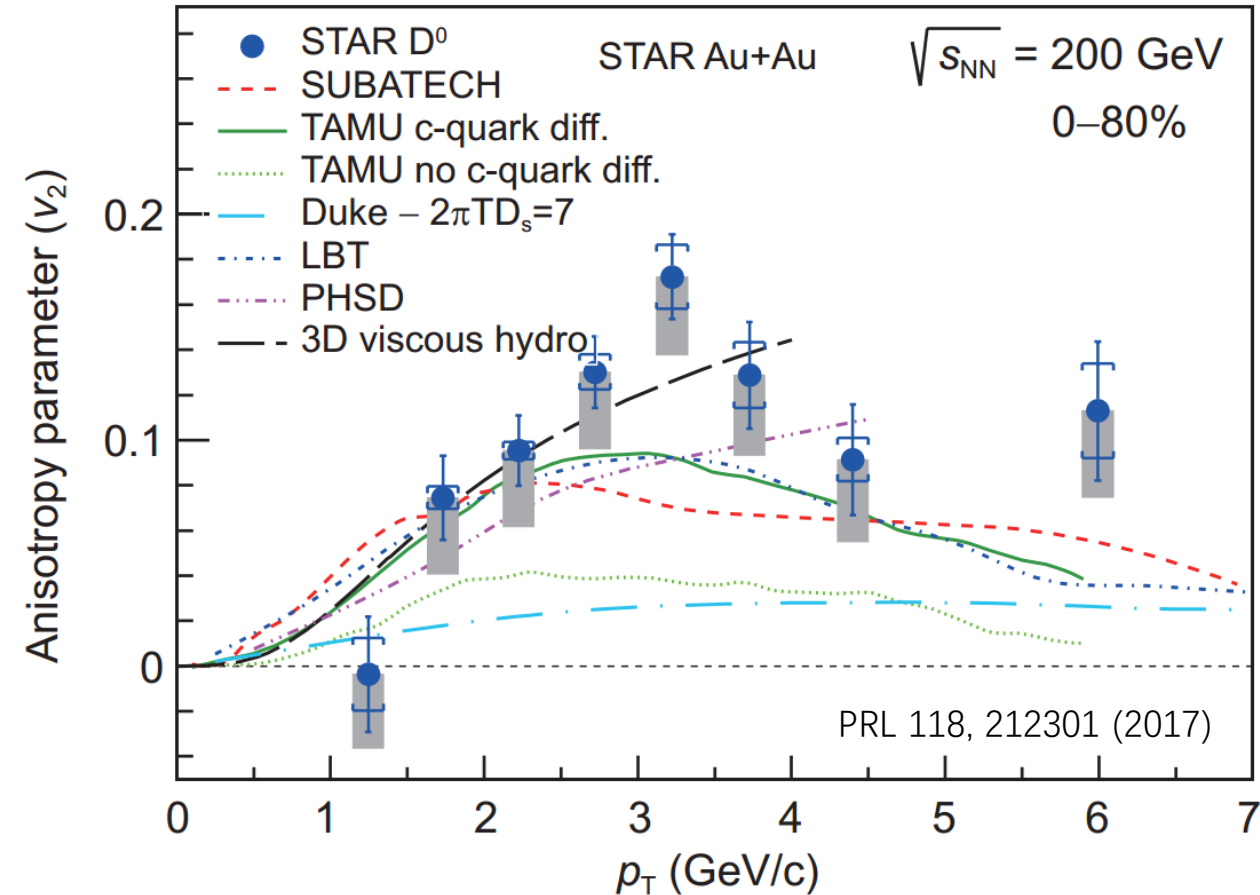
average over azimuthal distribution

v_2 elliptic flow, v_3 triangular flow, ...

Elliptic flow at 200 GeV Au+Au collisions

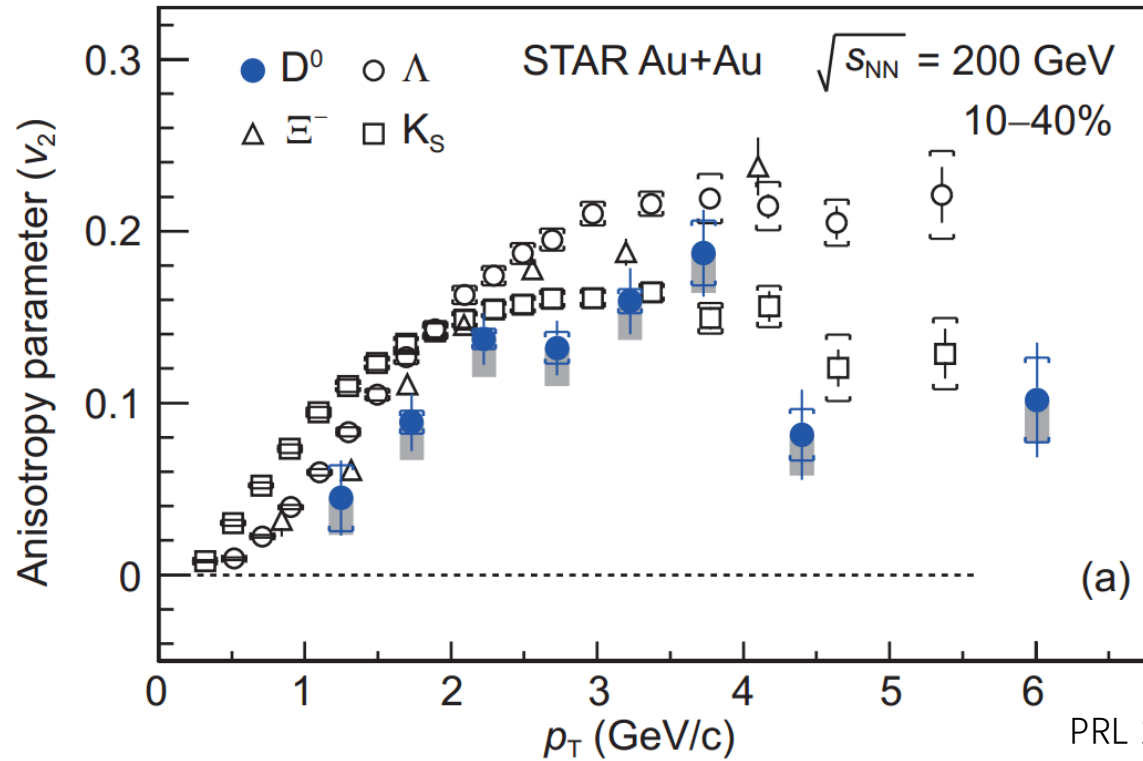


- ⊙ A significant elliptic flow indicate existence of light flavor-medium interactions

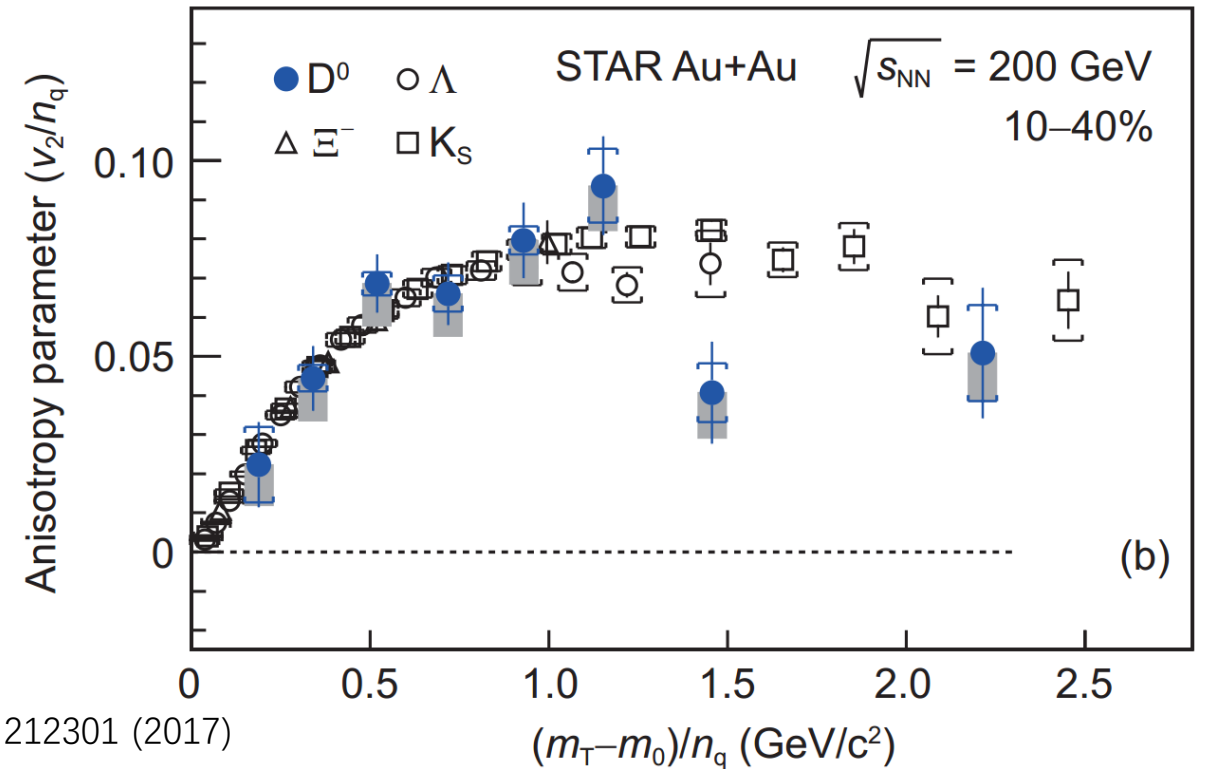


- ⊙ TAMU model with charm quark diffusion can describe data, not vice versa;
- ⊙ A 3D viscous hydro suggests that charm quarks have achieved thermal equilibrium.

D^0 v_2 at 200 GeV Au+Au collisions



PRL 118, 212301 (2017)



● A clear mass ordering for $p_T < 2$ GeV/c;

● Gained significant flow through interactions with the sQGP medium, and may thermalization;

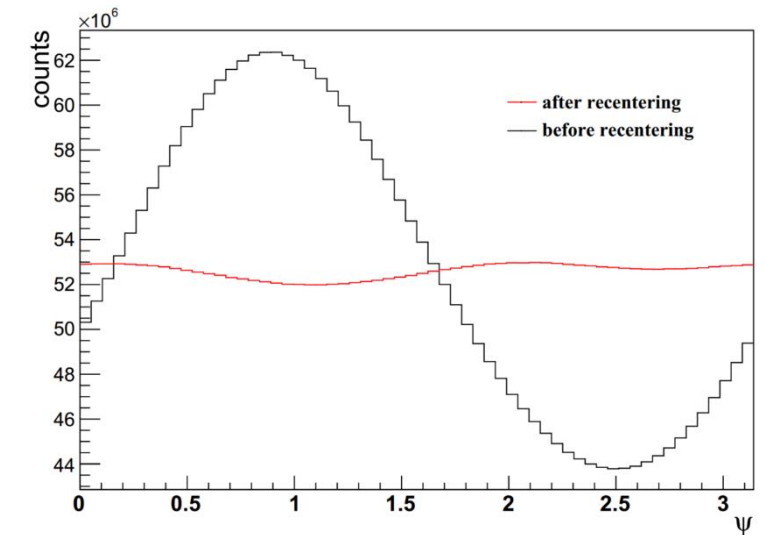
● An experimental evidence for quark deconfinement.

A expected smaller hot nuclei systems in Isobar, so v_2 depends on the system?
Charm quark gets thermalized in Isobar like in Au-Au at 200 GeV?

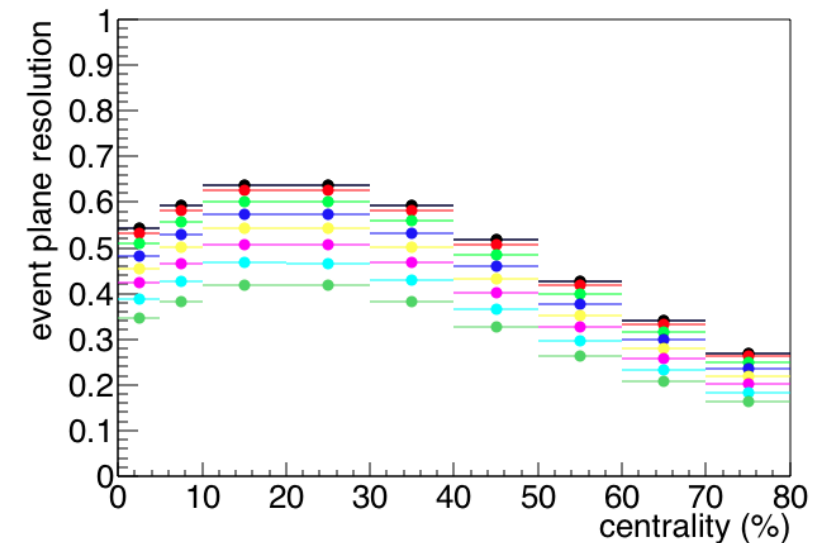
Event plane method to measure $D^0 v_2$

Event Level cuts	
Cuts	Value
$ V_z $	(-35, 25) cm
V_r	< 2. cm
VzDiff	< 3. cm
Badrun_list	
!GoodTrigger	
isBadRun()	
!passnTofMatchRefmultCut()	
0-80%	

Track cuts for EP	
Cuts	Value
p_T (GeV/c)	(0.2, 2.)
$ \eta $	< 1.
TPC Nhits	≥ 15
nHits/HitsMax	[0.52, 1.2]
gDCA	$\leq 2.$ cm
charge	± 1
Primary tracks	

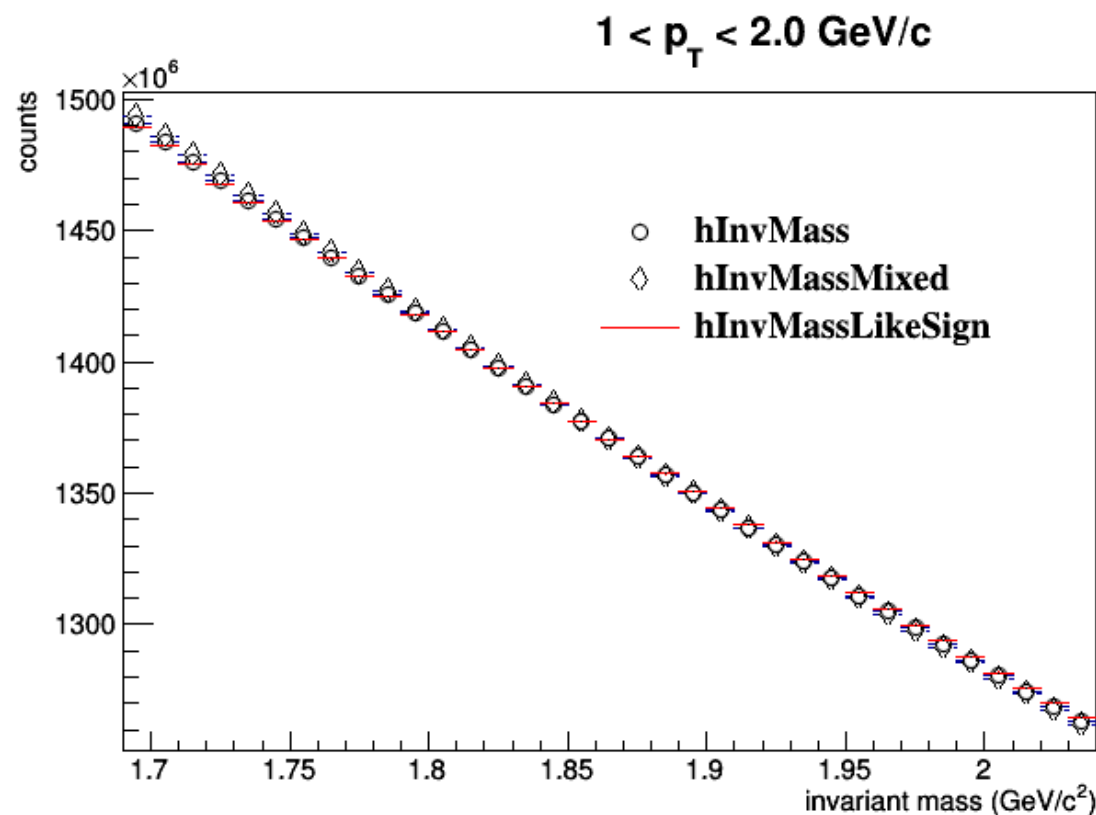


Event plane reconstruction and recentring

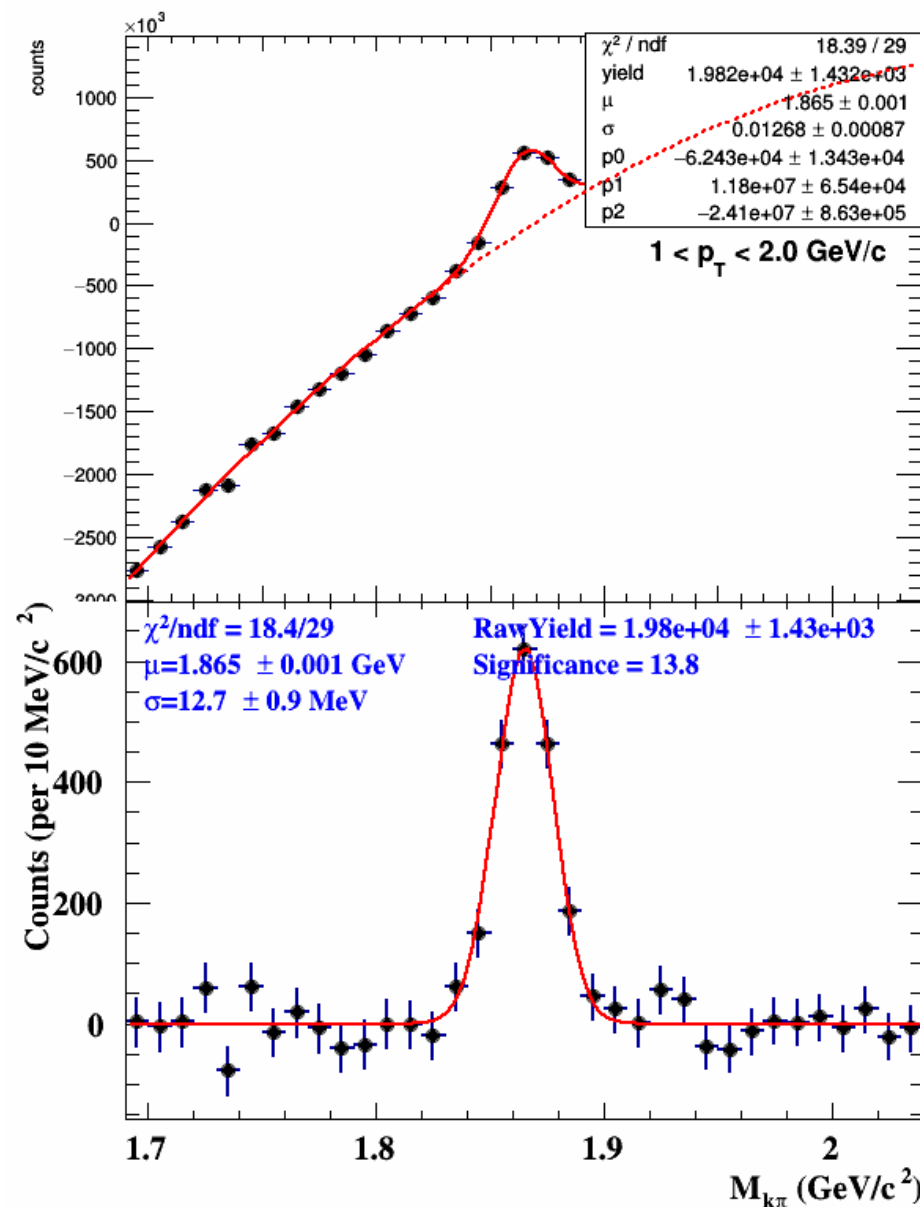


Event plane resolution vs. Centrality

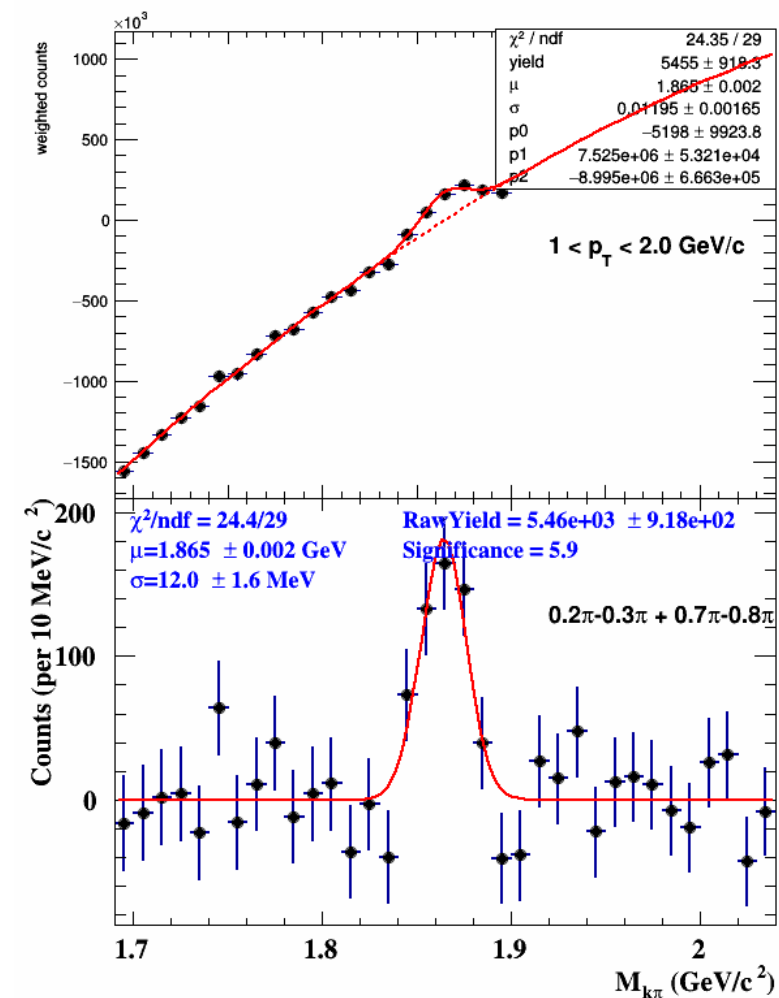
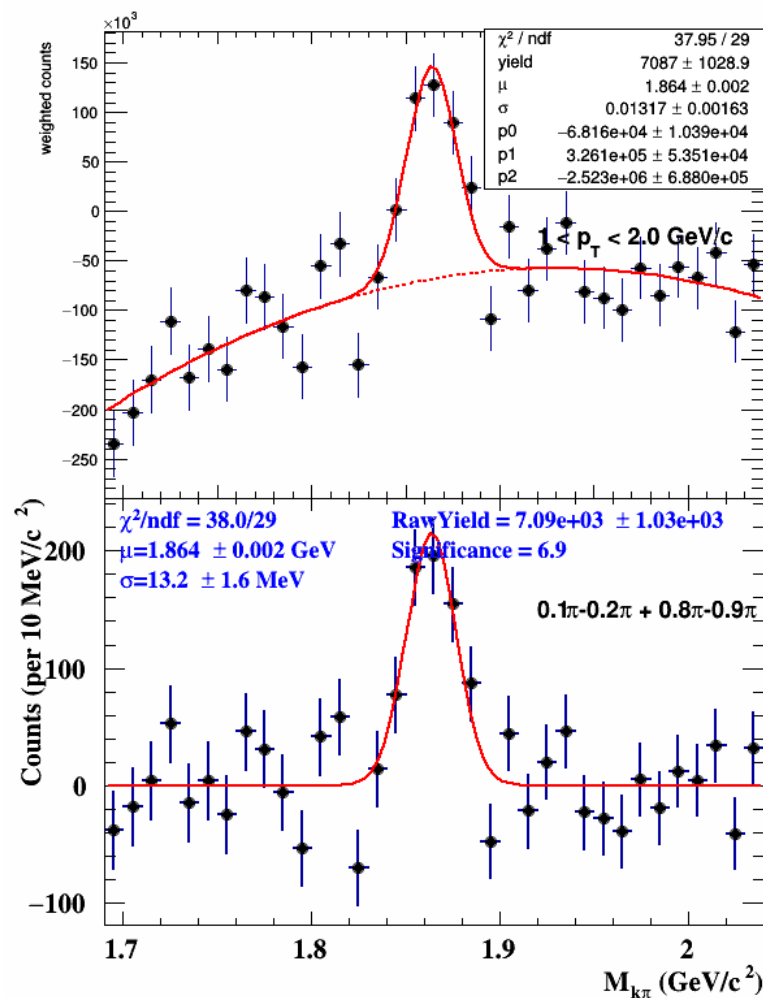
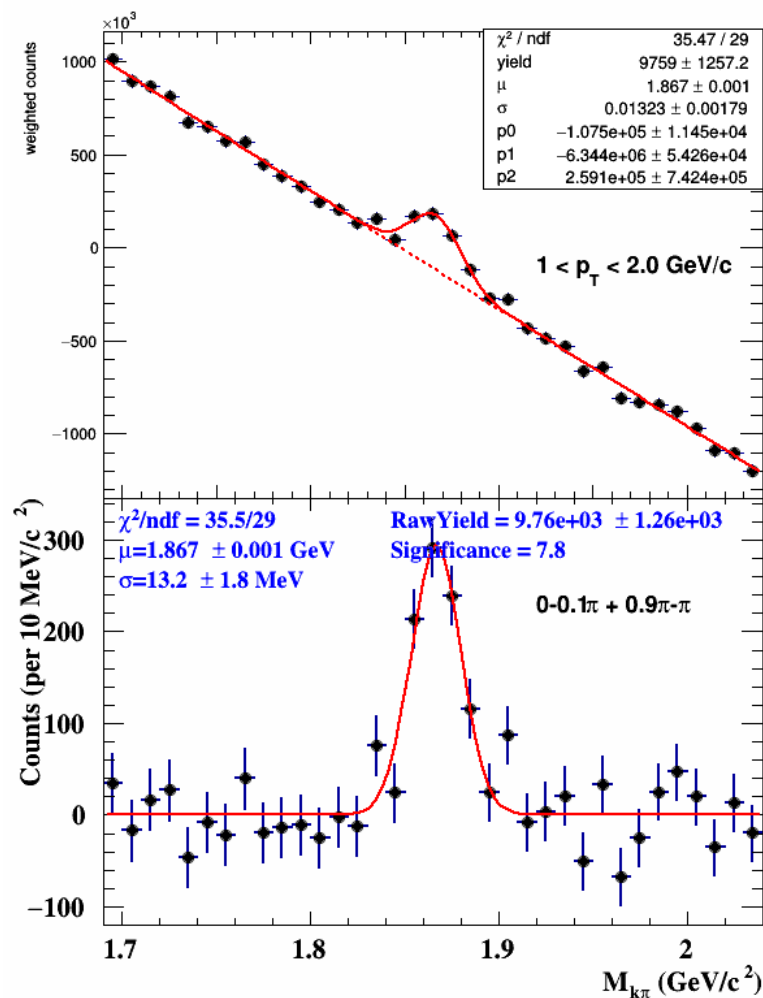
D^0 signal reconstruction



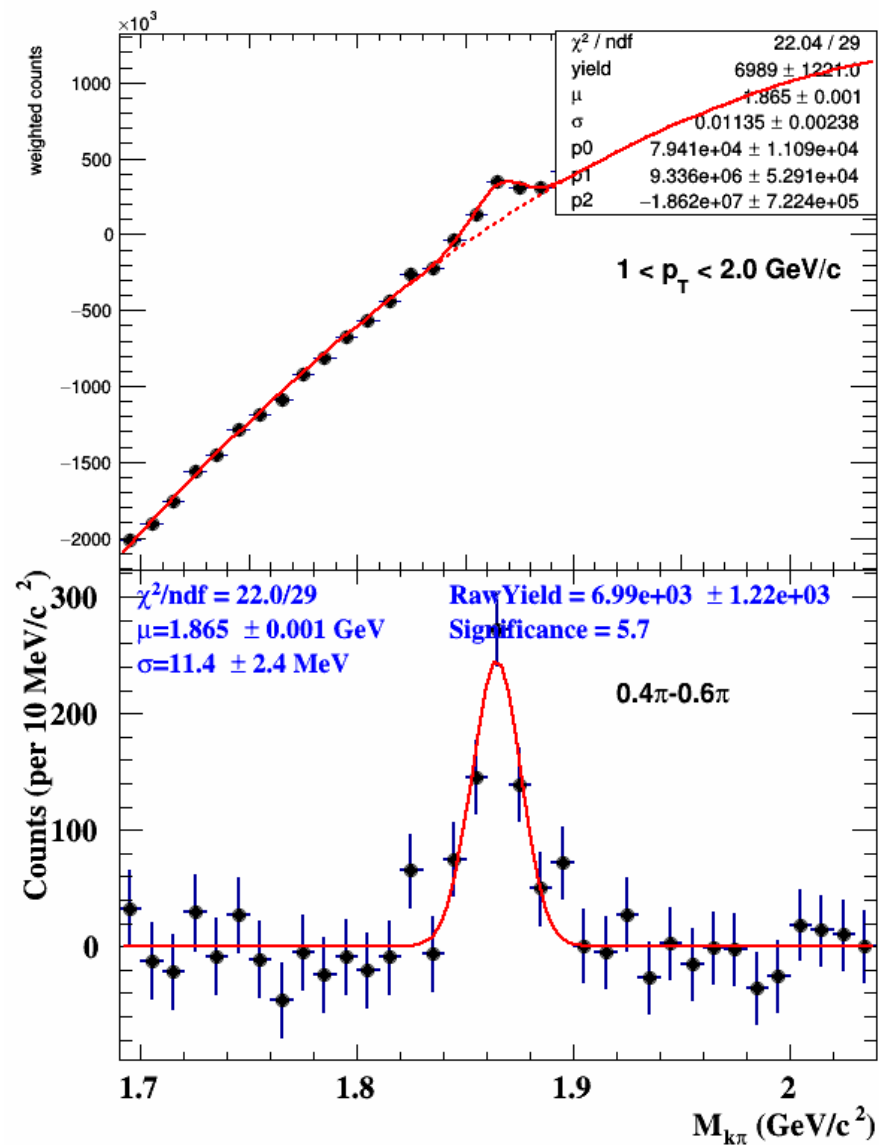
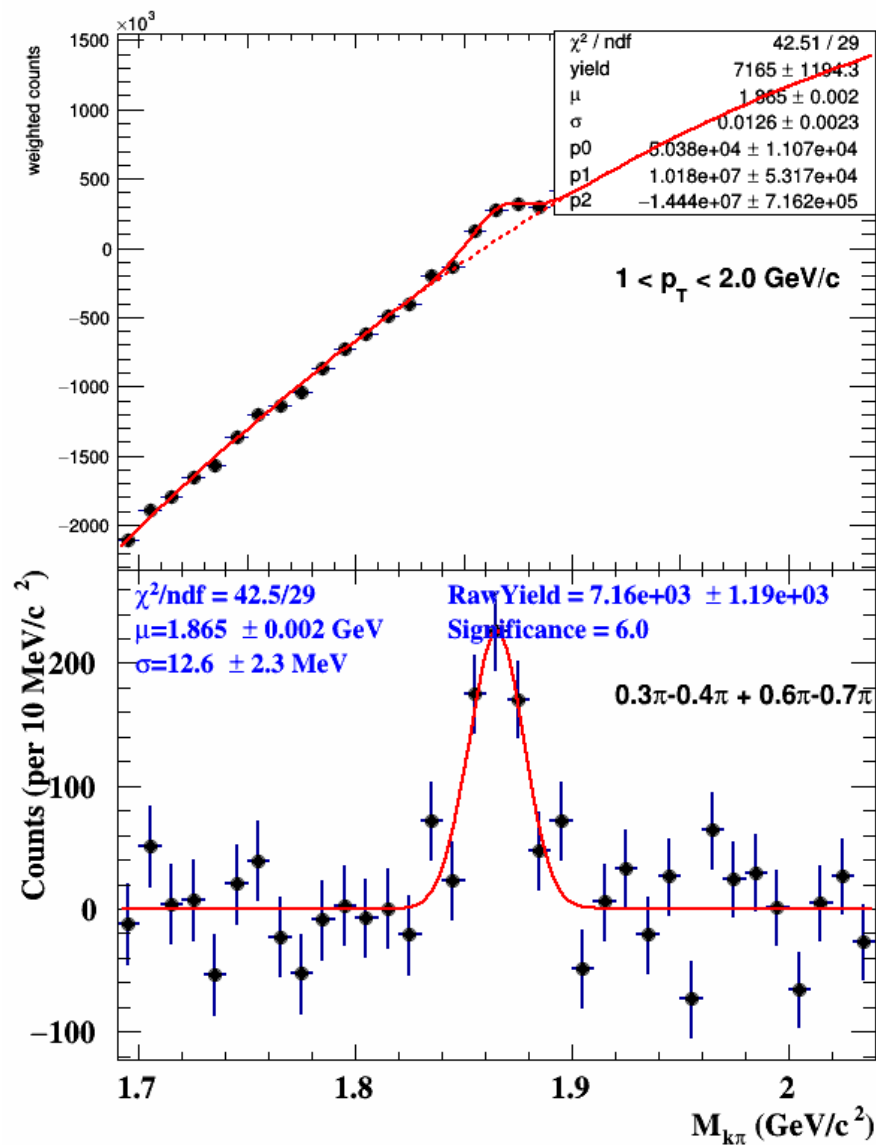
● Total D^0 signal at $1.0 < p_T < 2.0 \text{ GeV/c}$



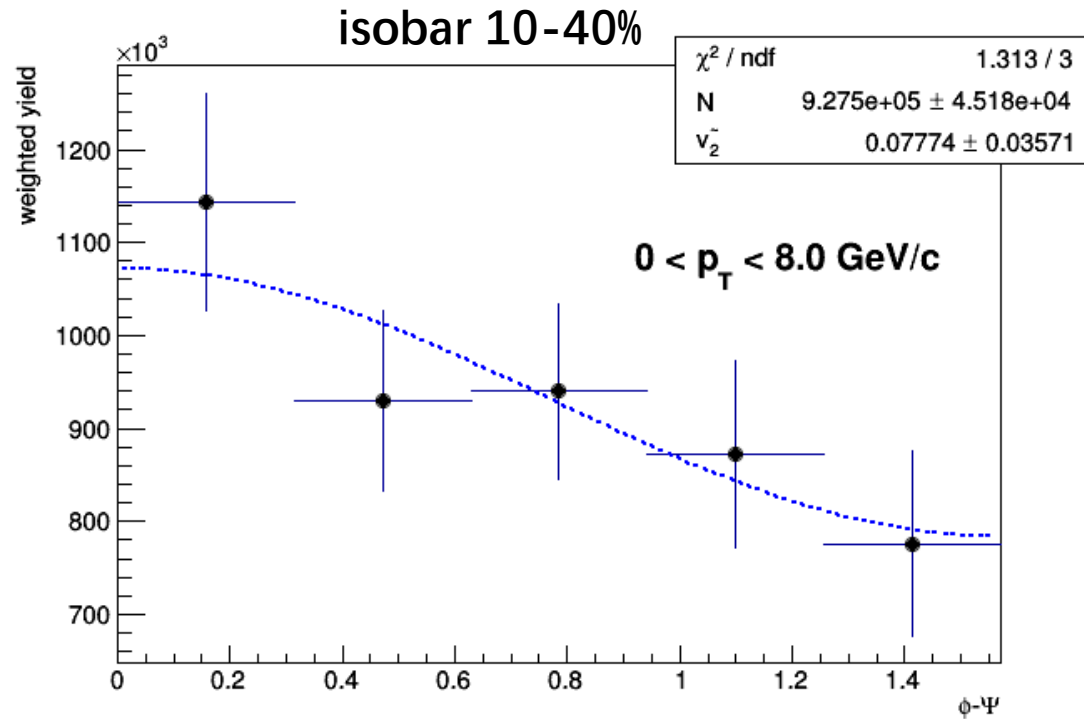
D^0 signal reconstruction



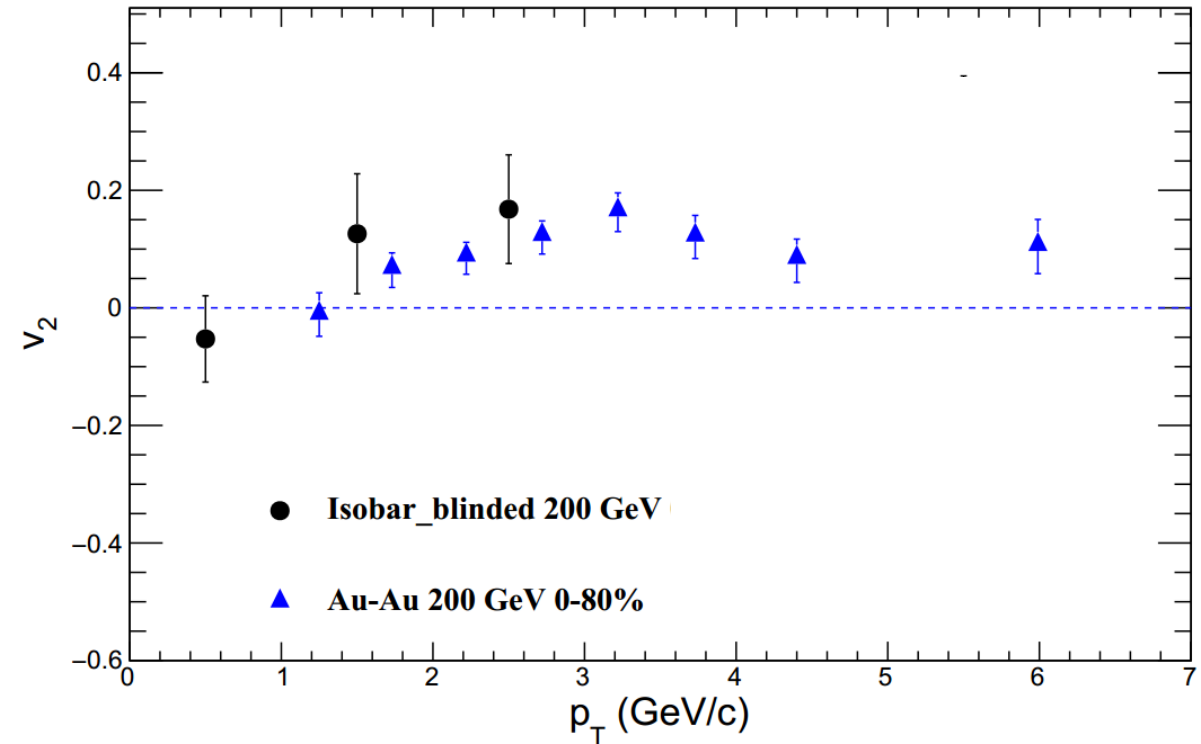
D^0 signal reconstruction



v_2 correction



$$N(1 + 2v_2 \cos(2(\phi - \Psi)))$$



- ⦿ A scale of $\Delta/\sin\Delta = 1.01664$ is done to correct for finite bin width $\Delta = 0.1\pi$;
- ⦿ The $\widetilde{v}_2$ is divided by $\langle 1/R \rangle$, weighted by the D^0 yield, to get physics v_2 .

- ⦿ Non-zero $D^0 v_2$ is observed.

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 - Measurements of D^0 p_T spectra and R_{AA} in isobar collisions at $\sqrt{s_{NN}} = 200$ at RHIC
- Collectivity
 - Measurement of D^0 elliptic flow v_2 at mid-rapidity in isobar collisions
- Summary

Summary

● D^0 p_T spectra and R_{AA} in Zr+Zr & Ru+Ru collisions at STAR

- Significant suppression observed at $p_T > 3$ GeV/c in central collisions, indicating charm quarks strongly interact with medium and lose energy
- At the same centrality, D^0 seems to behave similarly in isobar and Au+Au collisions, while light hadrons do not

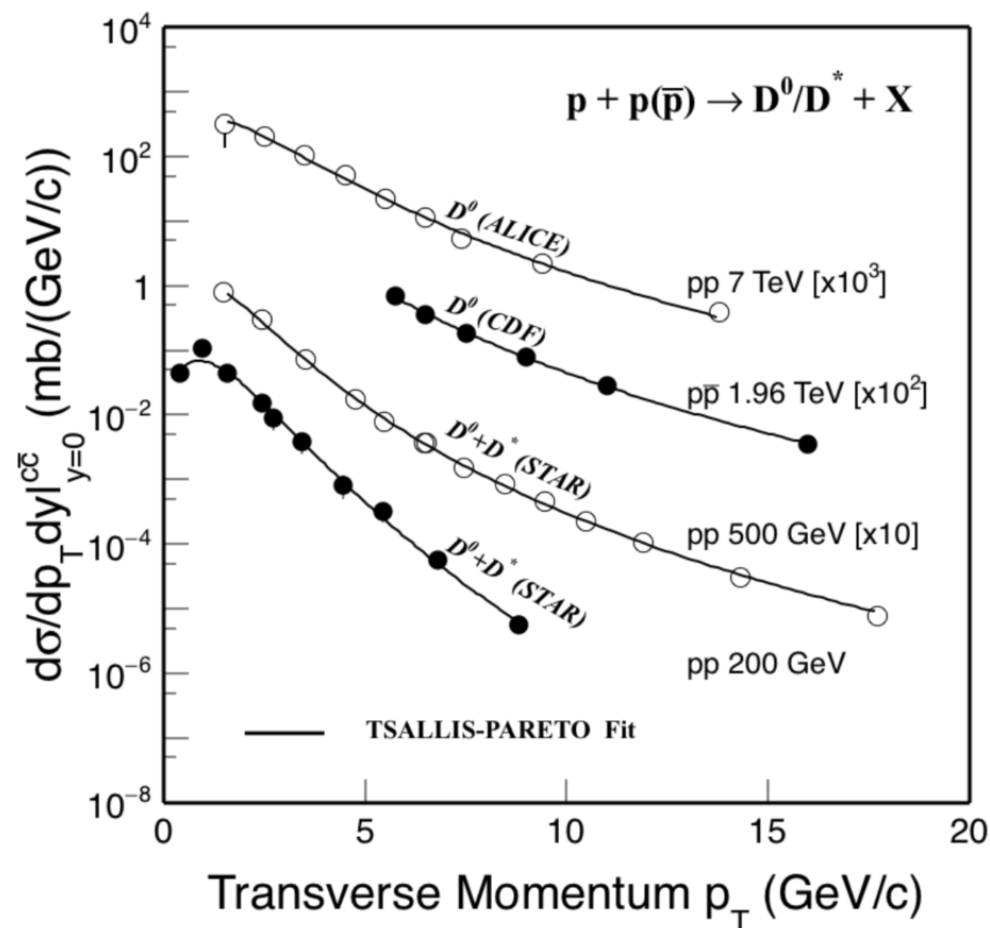
● D^0 elliptic flow v_2 at mid-rapidity in Zr+Zr & Ru+Ru collisions at STAR

- A non-zero v_2 is measured and consistent with Au+Au collisions within the uncertainty.

Thank you!

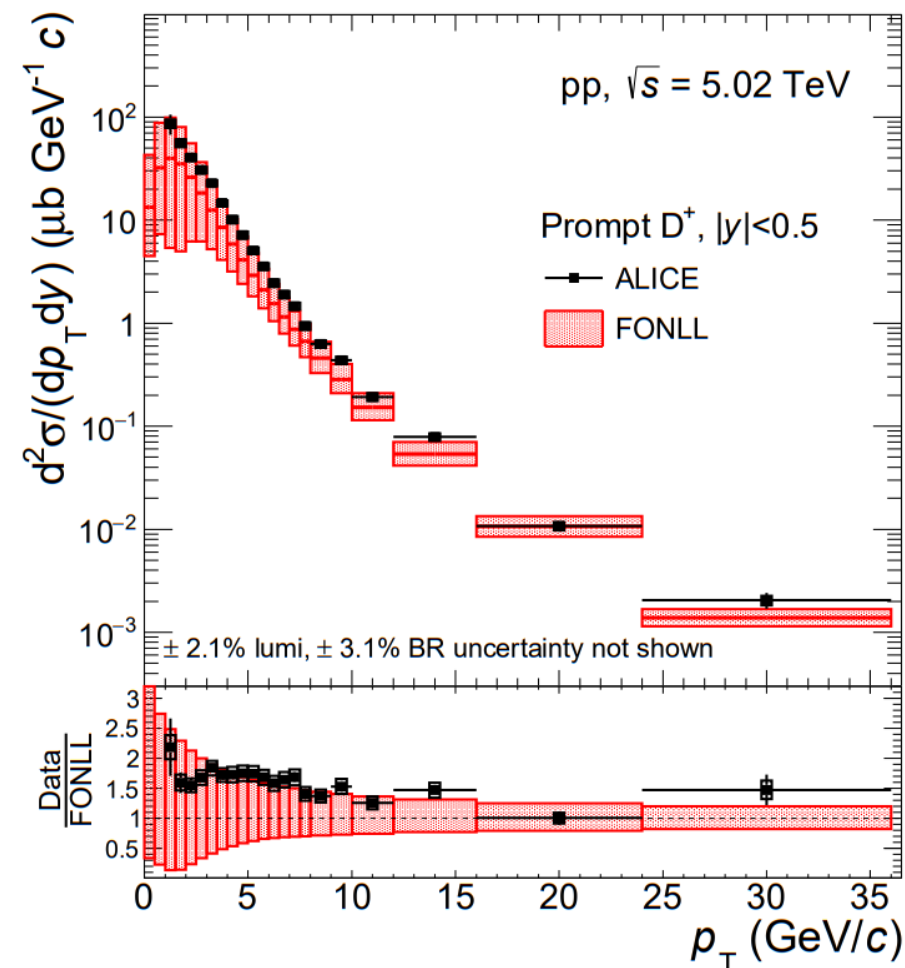
Backup

D^0 production cross section in p+p collisions



- pQCD calculations in p+p collisions, convolution

 - the parton distribution functions (PDFs) of the incoming protons
 - the partonic scattering cross section
 - the fragmentation function

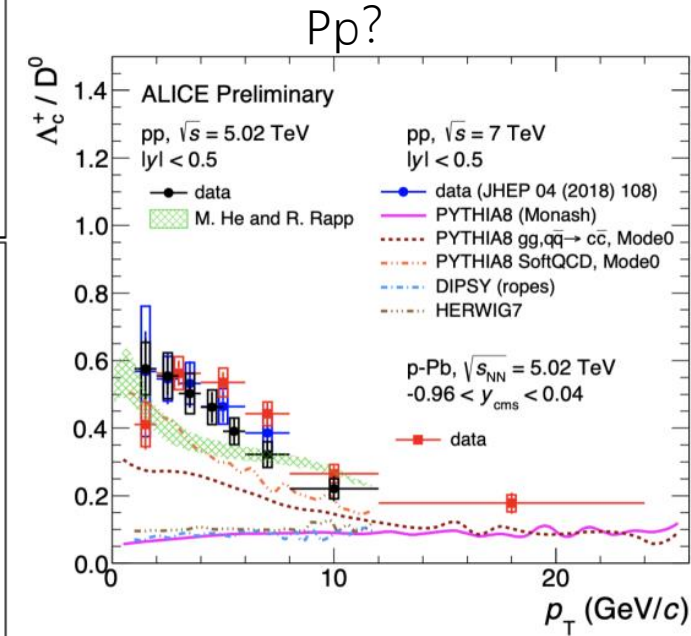
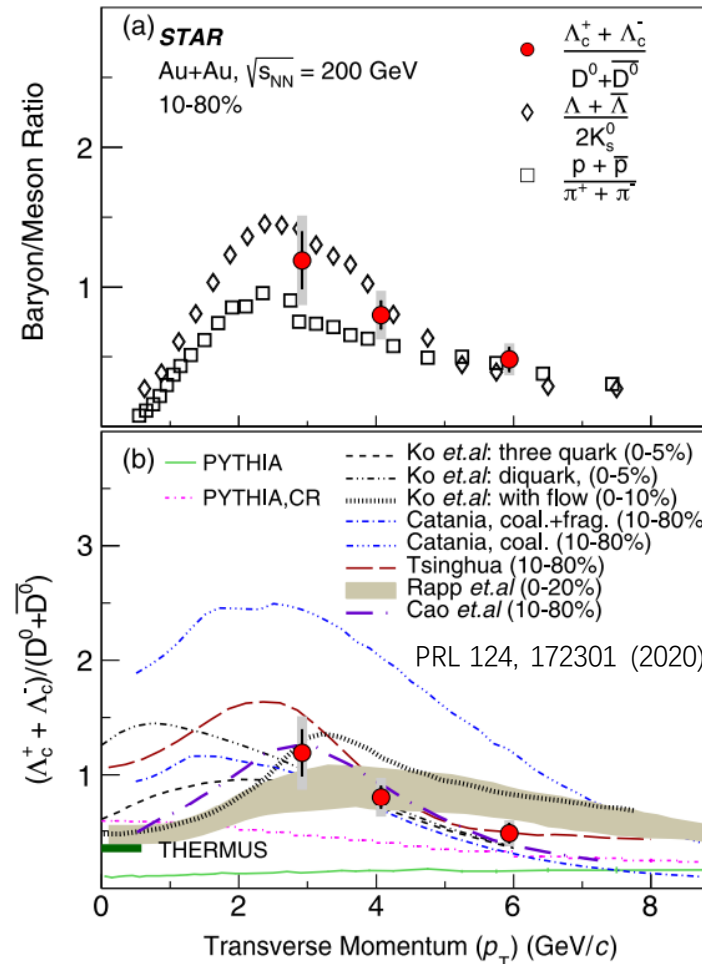
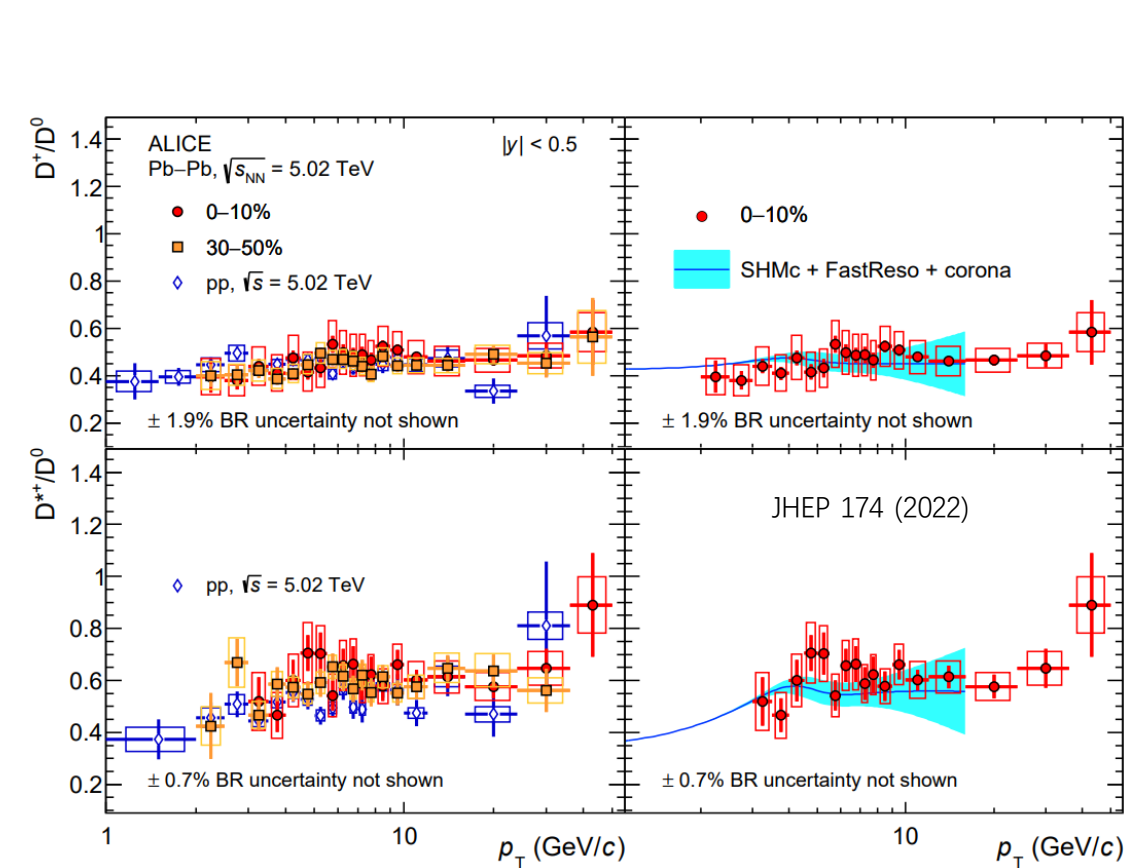


- Measured down to $p_T = 0$ in central collisions

 - provide important constraints to pQCD calculations and to low-x gluon PDFs
 - Represent an essential reference for the study of effects induced by cold and hot strongly-interacting matter

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Particle specie ratios in pp, p-Pb, Au-Au and Pb-Pb



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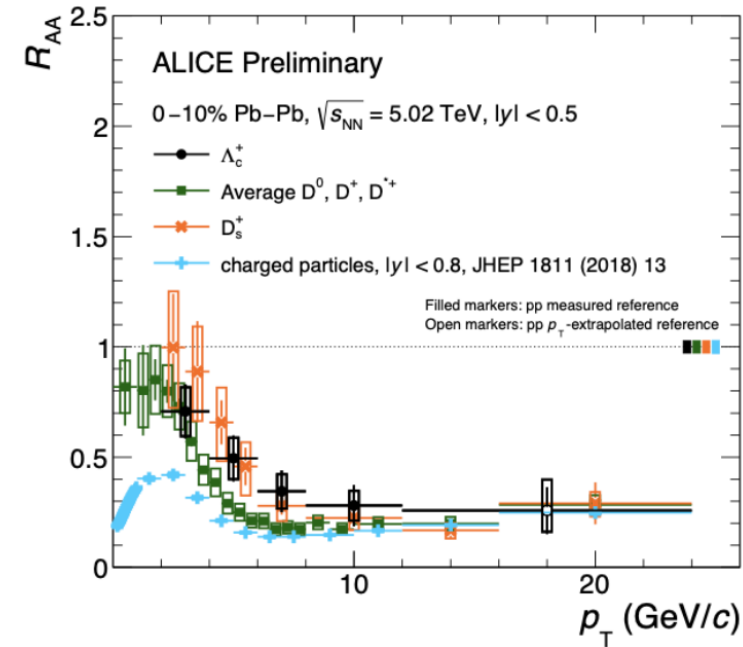
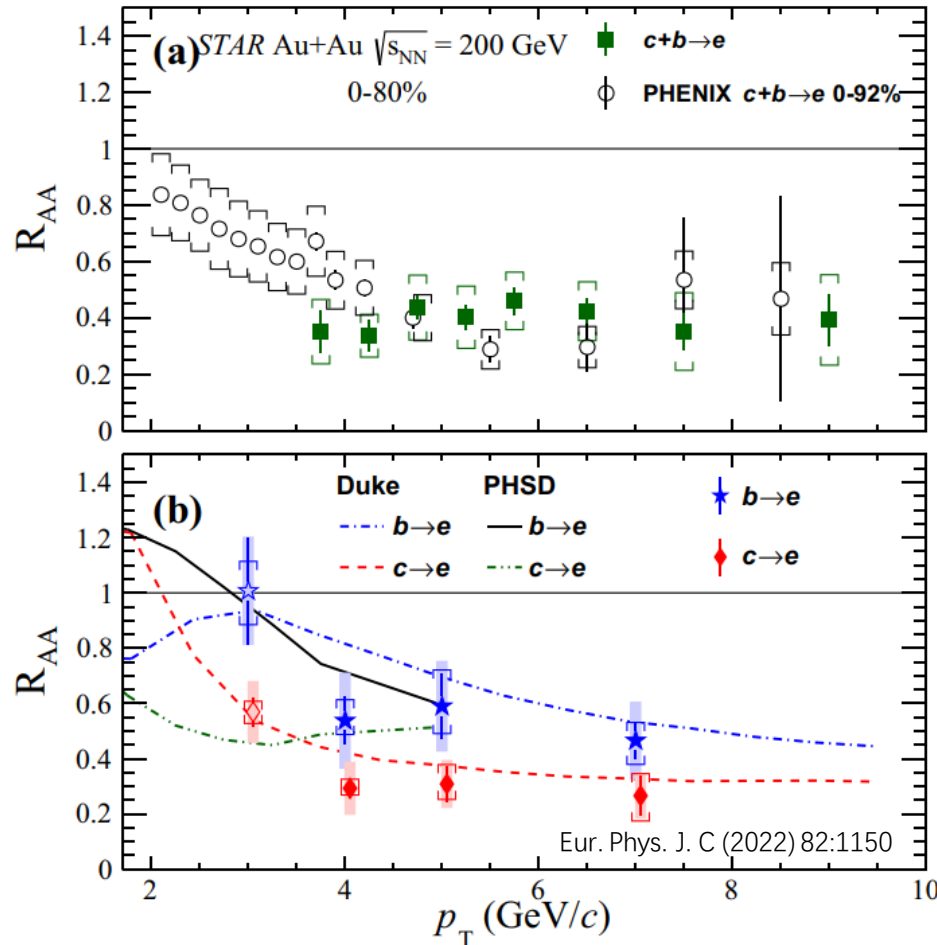
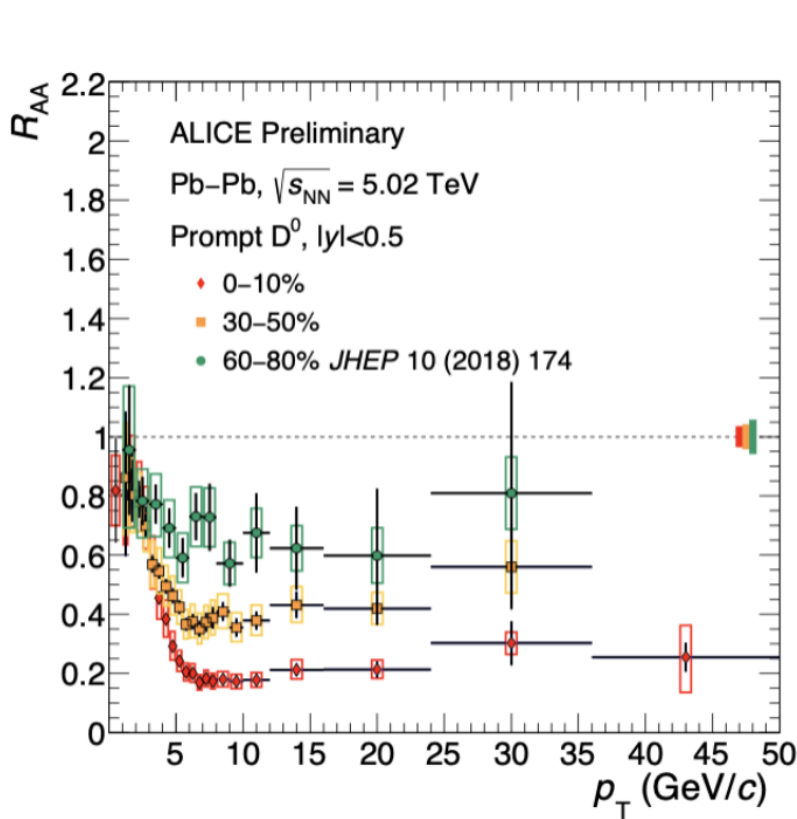
Relative abundances

- fragmentation fractions
- hadronisation scenario (independent on system and energy)

Universality of the fragmentation fractions for charm can not be ruled out in the meson sector.

- Λ_c/D^0 yield ratio is comparable to the baryon-to-meson ratios for light and strange-flavor hadrons
- Baryons contribute significantly to the total charm cross section
- Fragmentation fractions of charm quarks are non-universal

Energy loss mechanisms and Mass ordering



- Measured down to $p_T = 0$ in central collisions
- A large suppression is observed at low-intermediate p_T , and it is larger for the most central class
- Bottom decay electron R_{AA} to be significantly higher than those of charm-decay electrons
- $R_{AA}(\text{gluons}, u, d) < R_{AA}(c) < R_{AA}(b)$

Technical details: Formulas

● Spectra

$$E \frac{d^3 N}{d\mathbf{p}^3} = \frac{d^3 N}{p_T dp_T dy d\phi} = \frac{d^2 N}{2\pi p_T dp_T dy}$$

$$\frac{d^2 N}{2\pi p_T dp_T dy} = \frac{\Delta N^{raw} / \epsilon_{D^0}^{tot} / 2}{2\pi p_T \Delta p_T \Delta y \times N_{events} \times B.R.} = \frac{\Delta N_{D^0}^{AA}}{2\pi p_T \Delta p_T \Delta y} = E \frac{d^3 \sigma_{D^0}^{AA}}{d\mathbf{p}^3}$$

● Nuclear modification factor

$$R_{AA} = \frac{d^2 N_{AA}^{D^0} / dp_T dy}{\langle T_{AA} \rangle d^2 \sigma^{PP} / dp_T dy} = \frac{\sigma_{inel}^{NN} d^2 N_{AA}^{D^0} / dp_T dy}{\langle N_{coll} \rangle d^2 \sigma^{PP} / dp_T dy}$$

$$R_{cp} = \frac{Y(0-10\%) / \langle N_{coll,0-10} \rangle}{Y(40-80\%) / \langle N_{coll,40-80} \rangle} \quad \langle T_{AA} \rangle \times \sigma_{inel}^{NN} = \langle N_{coll} \rangle \quad 42 \text{ mb}$$

● Systematic uncertainties for integrated D^0 cross section

$$\left. \frac{d\sigma_{D^0}^{NN}}{dy} \right|_{y=0} = \left. \frac{dN_{D^0}^{AA}}{dy} \right|_{y=0} \times \frac{\sigma_{inel}^{pp}}{\langle N_{bin} \rangle} \quad 3$$

$$\left. \frac{dN_{D^0}^{AA}}{dy} \right|_{y=0} = \frac{1}{\Delta y \times N_{events} \times B.R.} \frac{(\Delta N^{raw} - D.C.) / \epsilon_{D^0}^{tot} / 2}{2} \quad 2$$

Open Charm Hadron	Constitution
D^0	$c\bar{u}$
D^+	$c\bar{d}$
D_s^+	$c\bar{s}$
Λ_c^+	udc

● Levy func

$$\frac{A(n-1)(n-2)}{nT(nT+m_0(n-2))} \times \left(1 + \frac{m_T - m_0}{nT}\right)^{-n}$$

● m_T Spectra

$$\frac{d^2 N}{2\pi p_T dp_T dy} = \frac{d^2 N}{2\pi m_T dm_T dy} = \frac{dN/dy}{2\pi T_{eff} (m_0 + T_{eff})} e^{-(m_T - m_0)/T_{eff}}$$

● Blast-Wave func

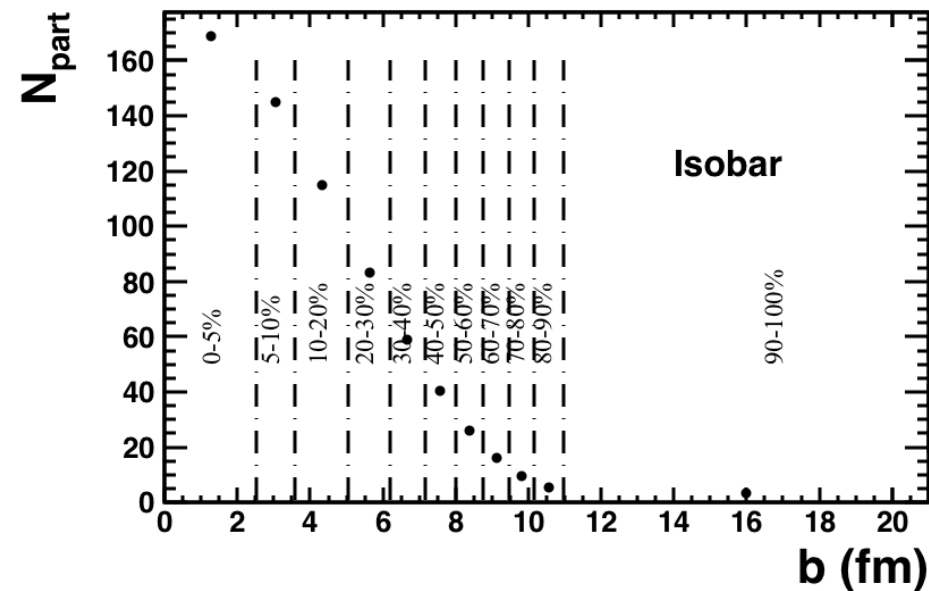
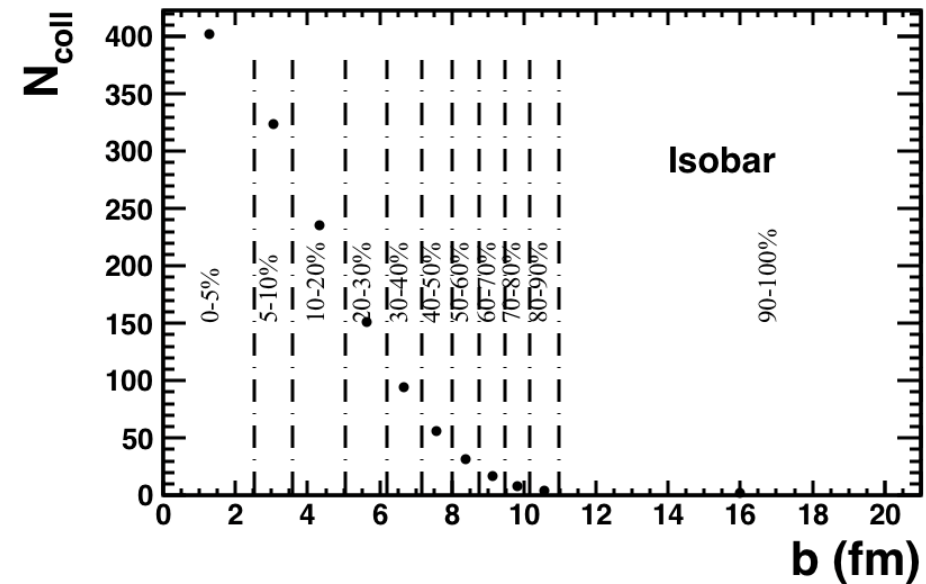
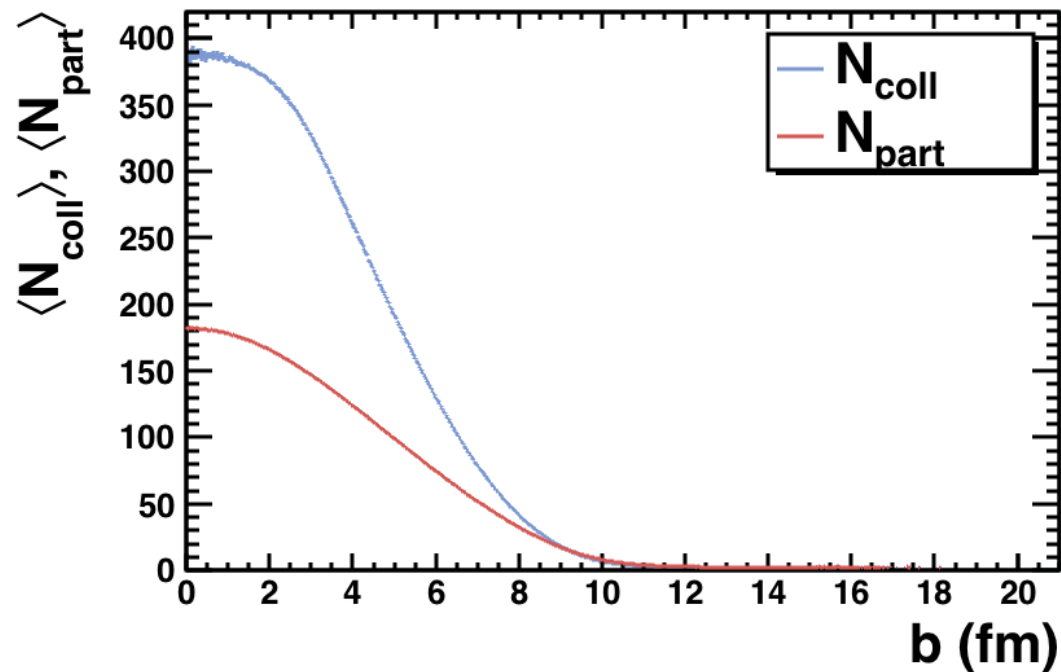
$$E \frac{d^3 N}{dp^3} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh(\rho)}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh(\rho)}{T_{kin}} \right)$$

$$E \frac{d^3 N}{dp^3} \propto m_T \int_{-Y}^{+Y} \cosh(y) dy \int_{-\pi}^{+\pi} d\phi \int_0^R r dr \left(1 + \frac{q-1}{T_{kin}} (m_T \cosh(y) \cosh(\rho) - p_T \sinh(\rho) \cos(\phi)) \right)^{-\frac{1}{q-1}}$$

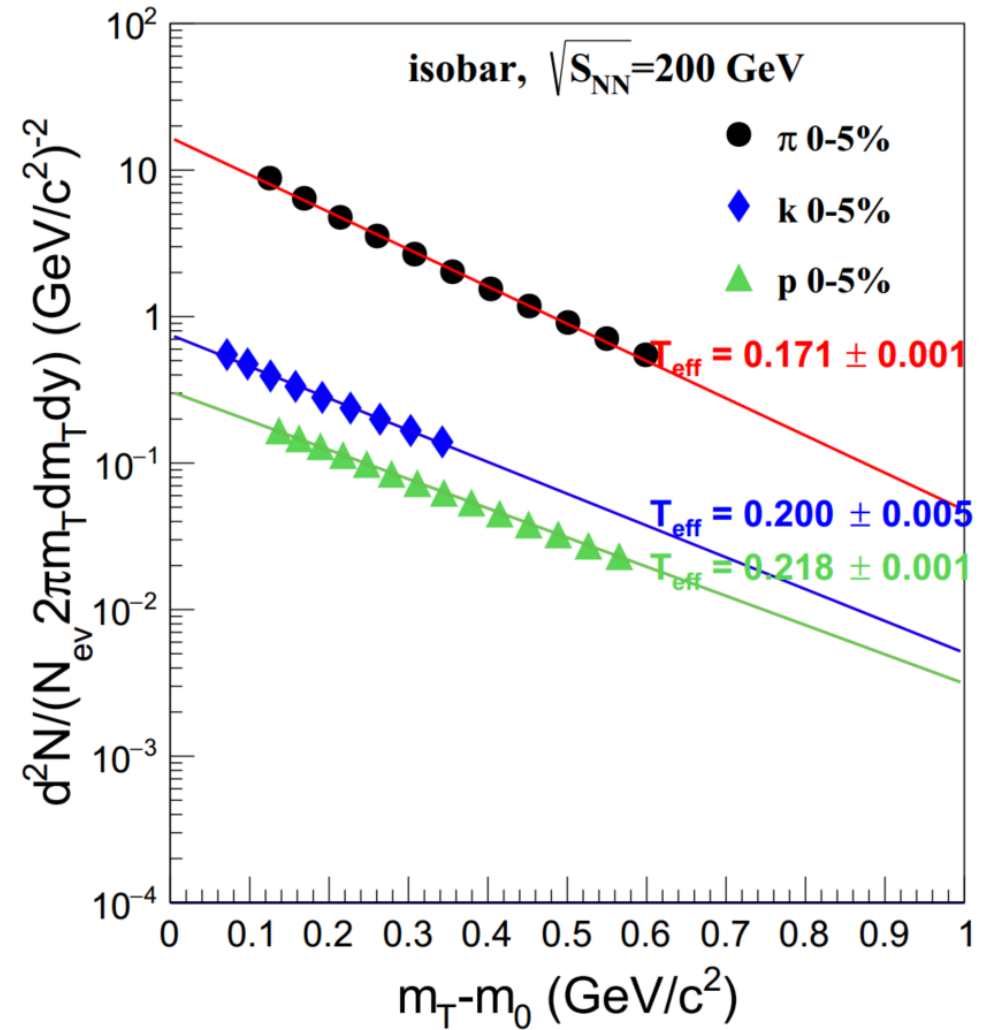
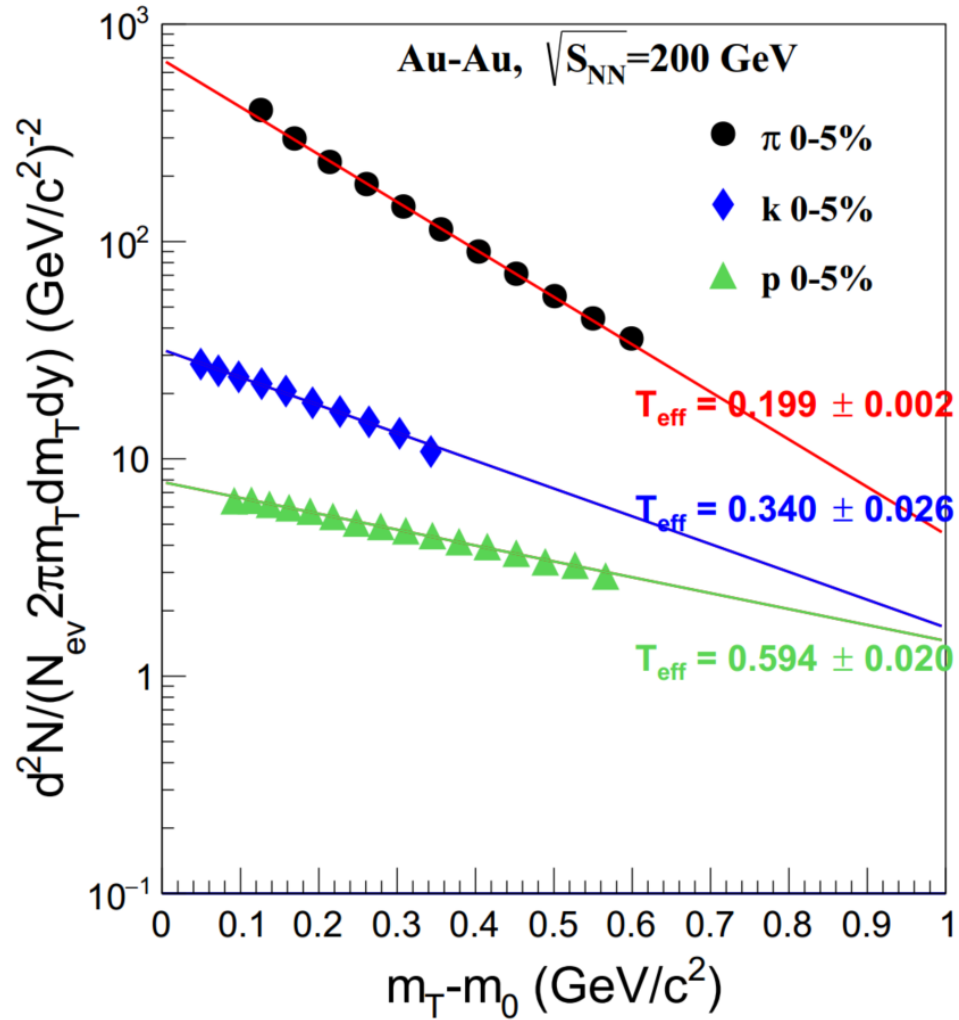
Glauber model simulation

$R=5.02$ fm $a=0.46$ fm

$\sigma_{inel}^{NN} = 42$ mb

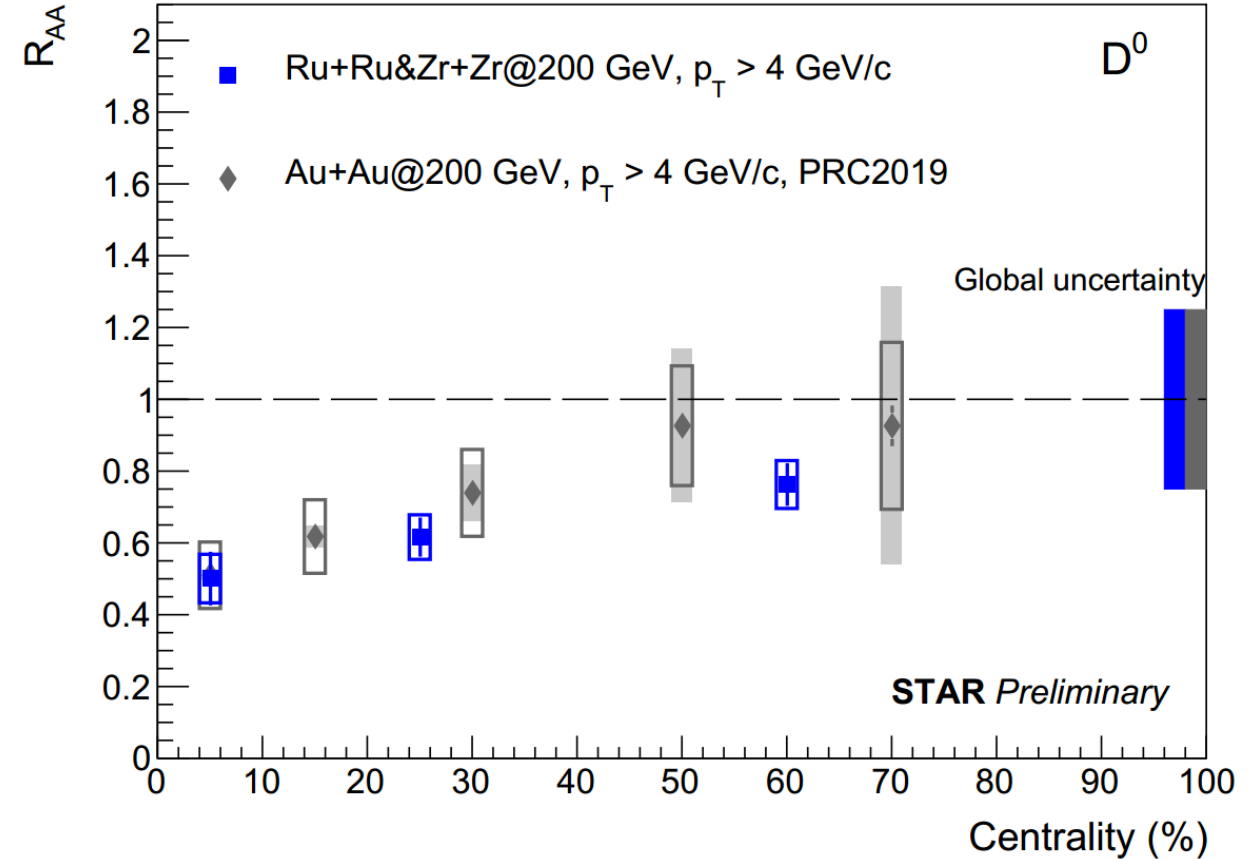
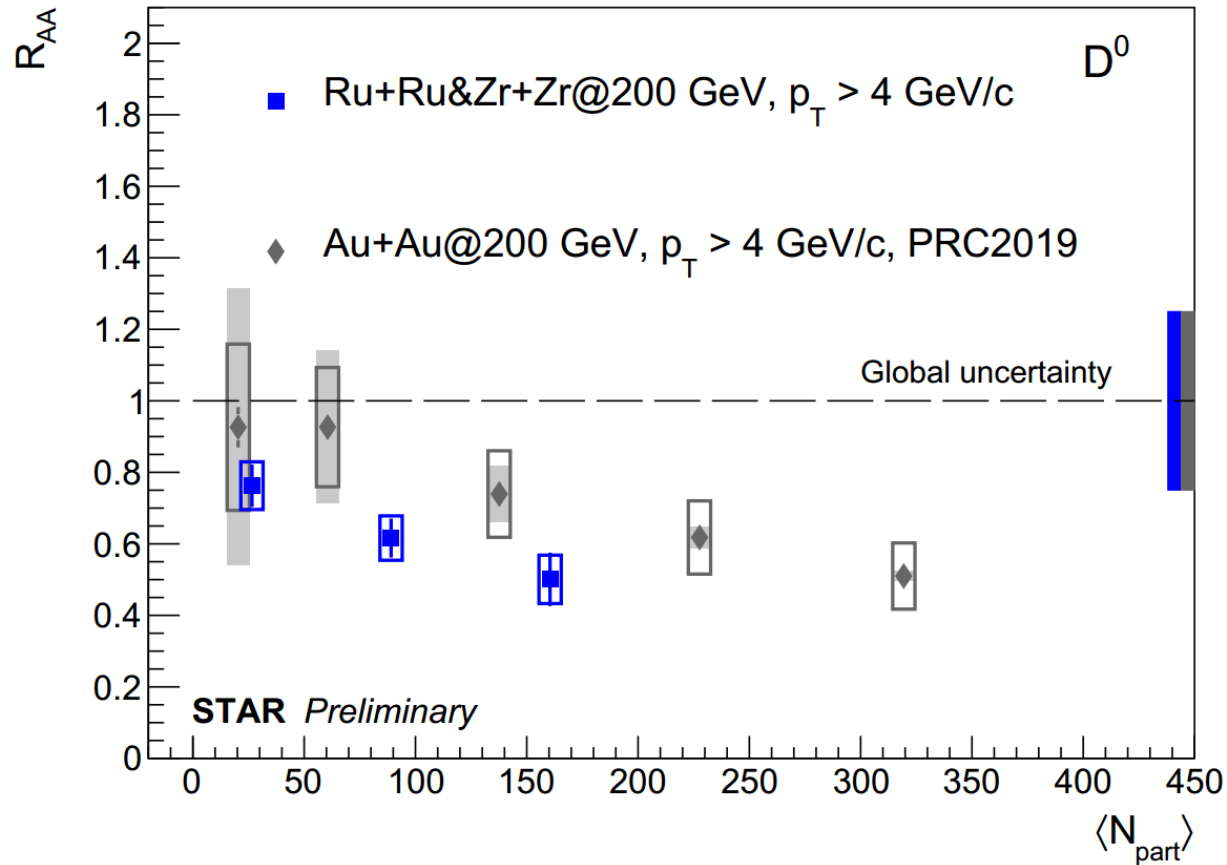


m_T -spectra fit details



⊙ m_T -spectra fit details for Au + Au (left) and isobar (right).

System size dependence of R_{AA}



v_2 measurement

- Event plane construction

$$\Psi_2 = \frac{1}{2} \tan^{-1} \left(\frac{Q_y}{Q_x} \right)$$

$$Q_x = \Sigma (q_x - \langle q_x \rangle)$$

$$q_x = p_T \cos(2\phi)$$

$$Q_y = \Sigma (q_y - \langle q_y \rangle)$$

$$q_y = p_T \sin(2\phi)$$

- v_2 calculation

$$v_2 = \frac{\bar{v}_2}{R} \quad R = \langle \cos(2(\Psi_2 - \Psi_{rp})) \rangle \quad \bar{v}_2 = \langle \cos[2(\phi - \Psi_2)] \rangle$$