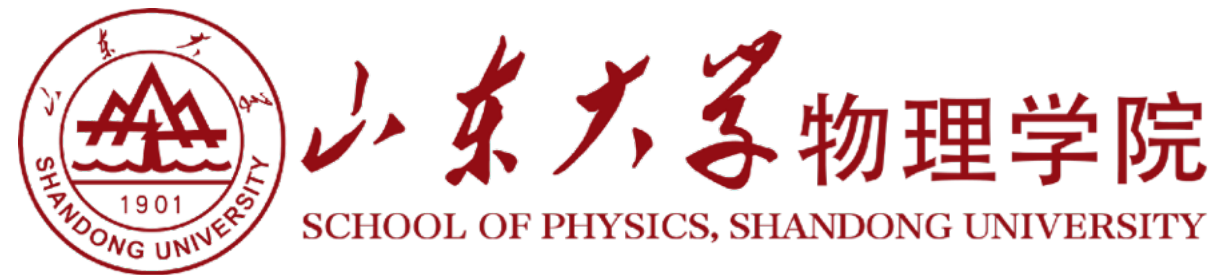


# $\alpha_s$ dependence of DIS EEC in Breit frame

Haitao Li



Based on the work with

Zhi-Hong He, D.Y. Shao in preparation

X. Chen, Zhi-Hong He in preparation

碎裂函数和能量关联研讨会

Institute of Modern Physics, Chinese Academy of Sciences

August 8-10, 2025

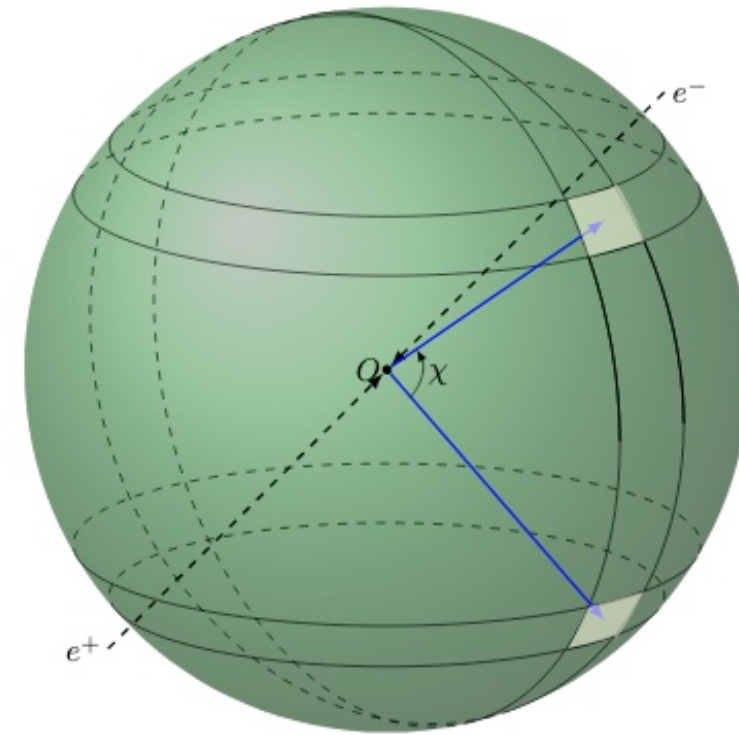
Lanzhou, China

# Outline

- Introduction
- EEC with jets
- EEC and  $\alpha_s$  measurements
- Summary

# Introduction

$e^+e^-$  Collisions



$$\text{EEC} = \sum_{a,b} \int d\sigma_{V \rightarrow a+b+X} \frac{2E_a E_b}{Q^2 \sigma_{\text{tot}}} \delta(\cos(\theta_{ab}) - \cos(\chi))$$

- sum over all the jets for each event
- sum over all the particles for each event

Basham et al 1978  
Moult, Zhu, 2018

Hadronic initial state

observable

$$\text{TEEC} = \sum_{a,b} \int d\sigma_{pp \rightarrow a+b+X} \frac{2E_{T,a} E_{T,b}}{|\sum_i E_{T,i}|^2} \delta(\cos \phi_{ab} - \cos \phi)$$

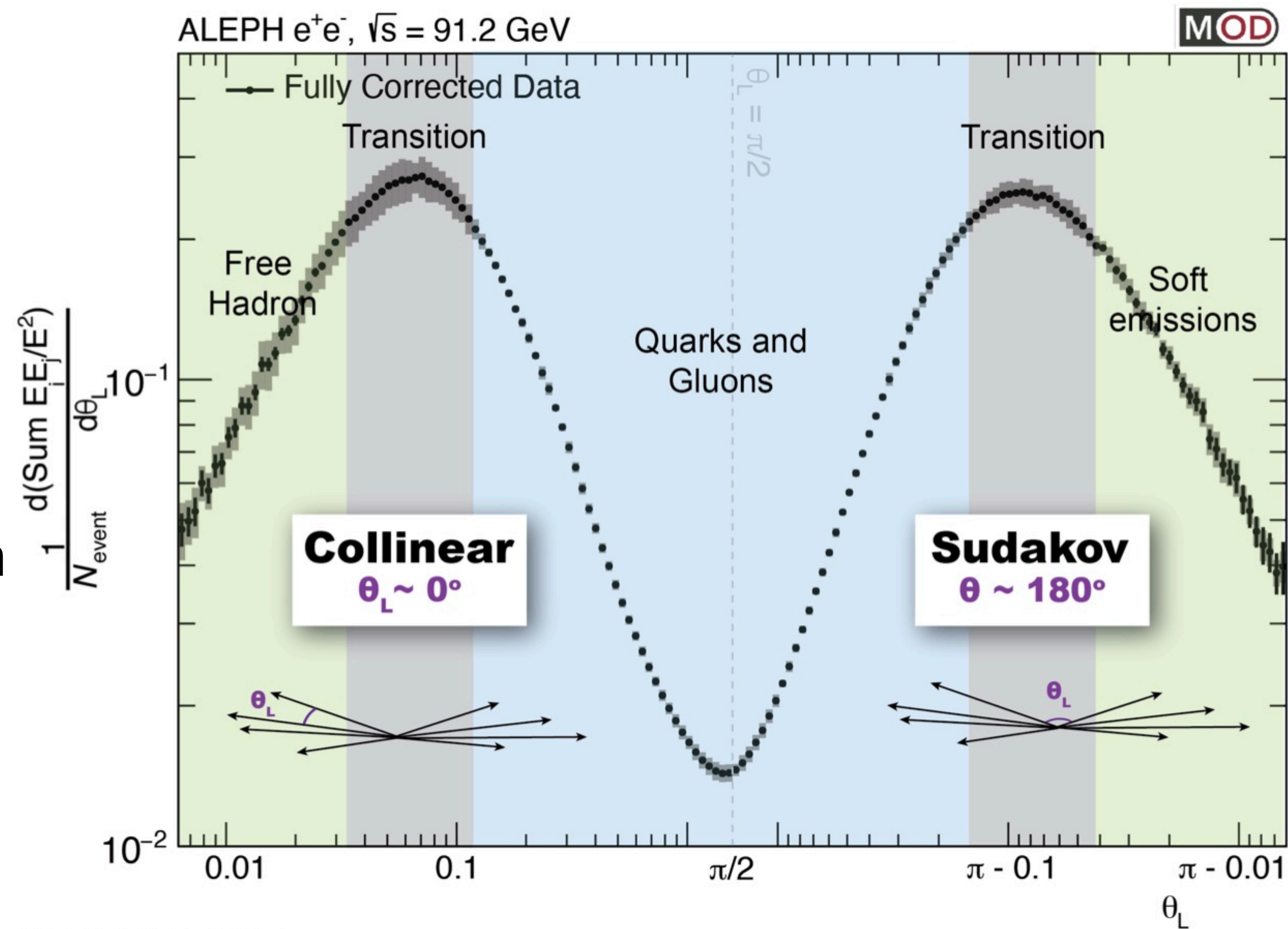
- weighted cross section
- the soft radiation does not contribute directly to the observable at leading power
- soft gluon contributes only via recoil

Ali et al 1984  
Gao, HTL, Moult, Zhu, PRL, 2019 & JHEP



# Introduction

many works recently  
in the collinear region



TMD region

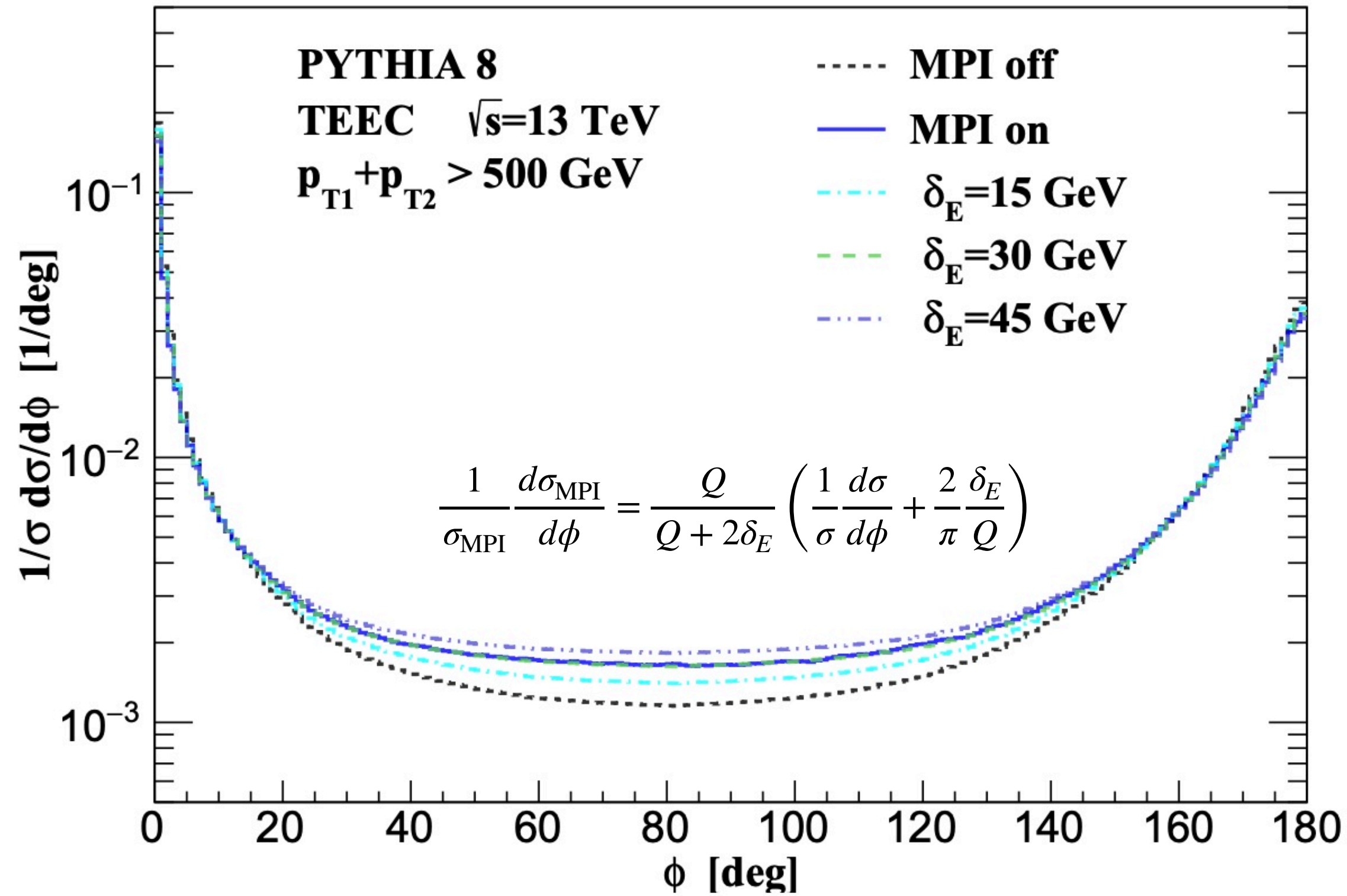
2025 HB - Finalization of Result

arXiv:2505.11828

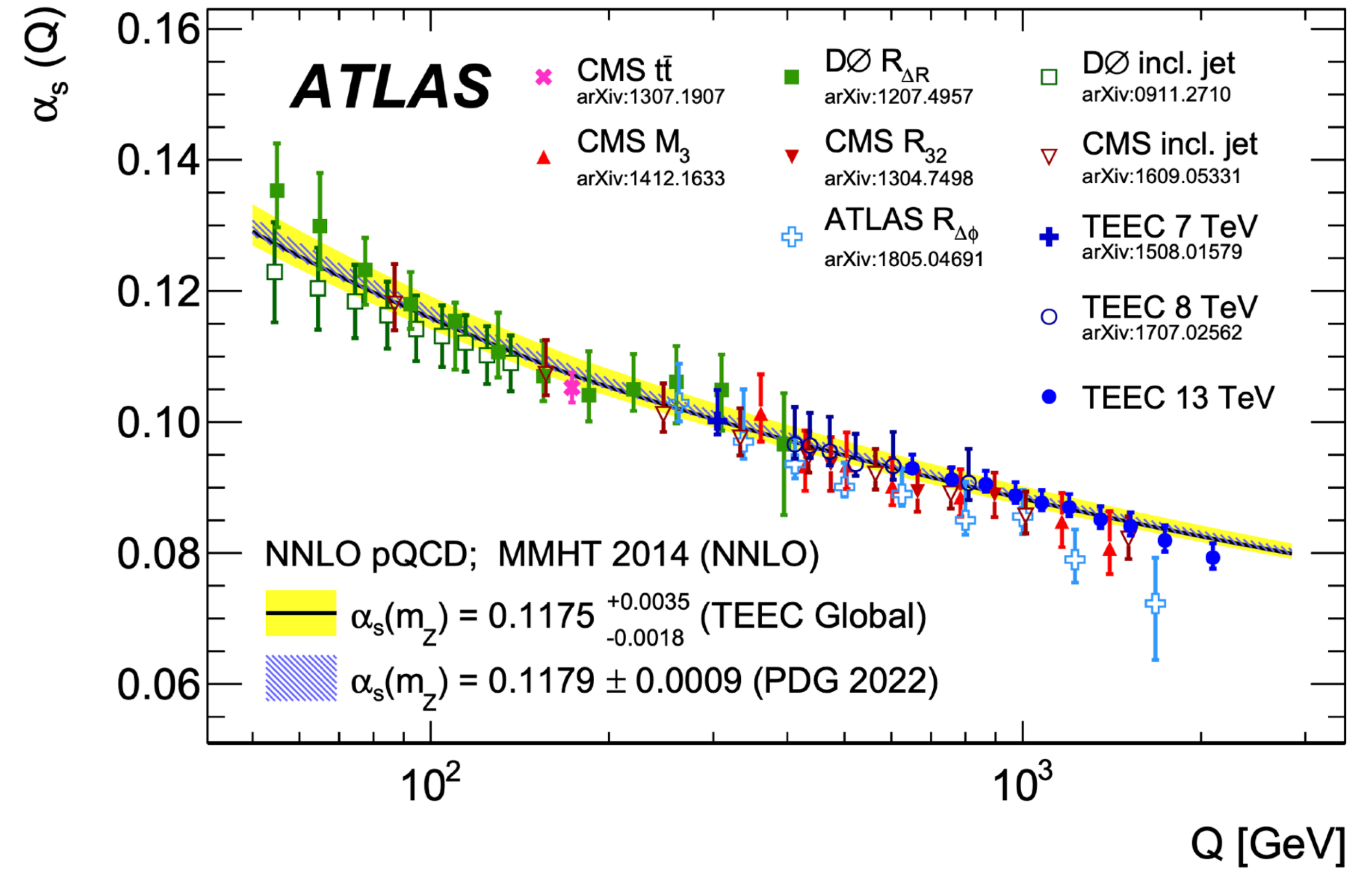


# Introduction

## TEEC in pp



Gao, HTL, Moul, Zhu, PRL, 19 & JHEP 24



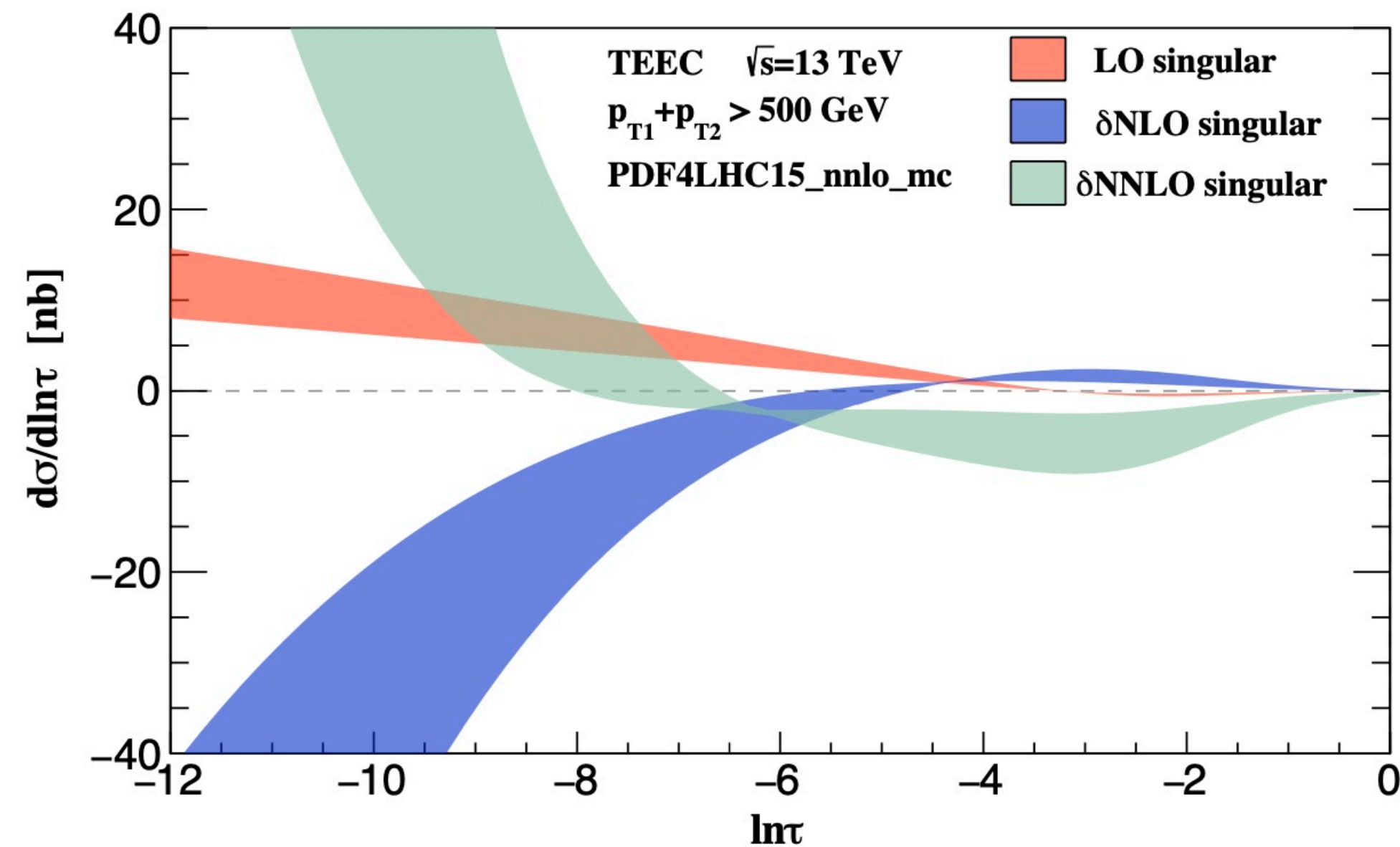
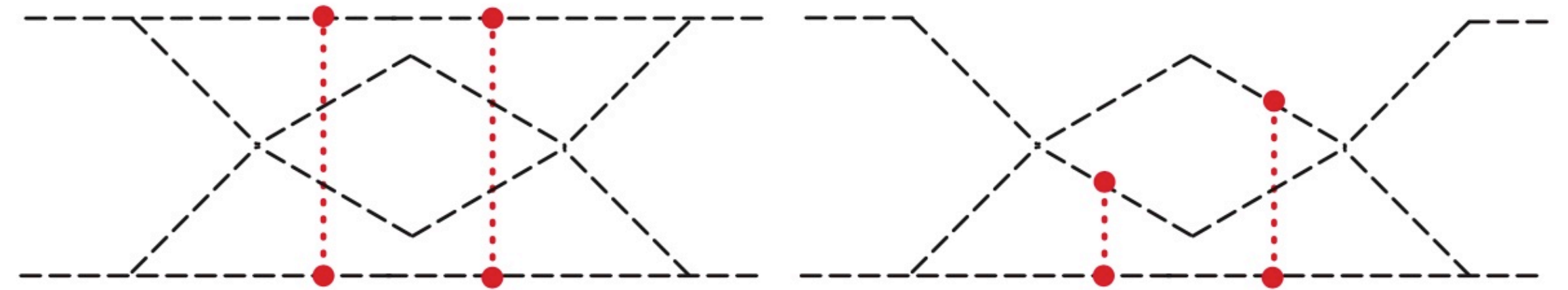
ATLAS, 2023

# Introduction

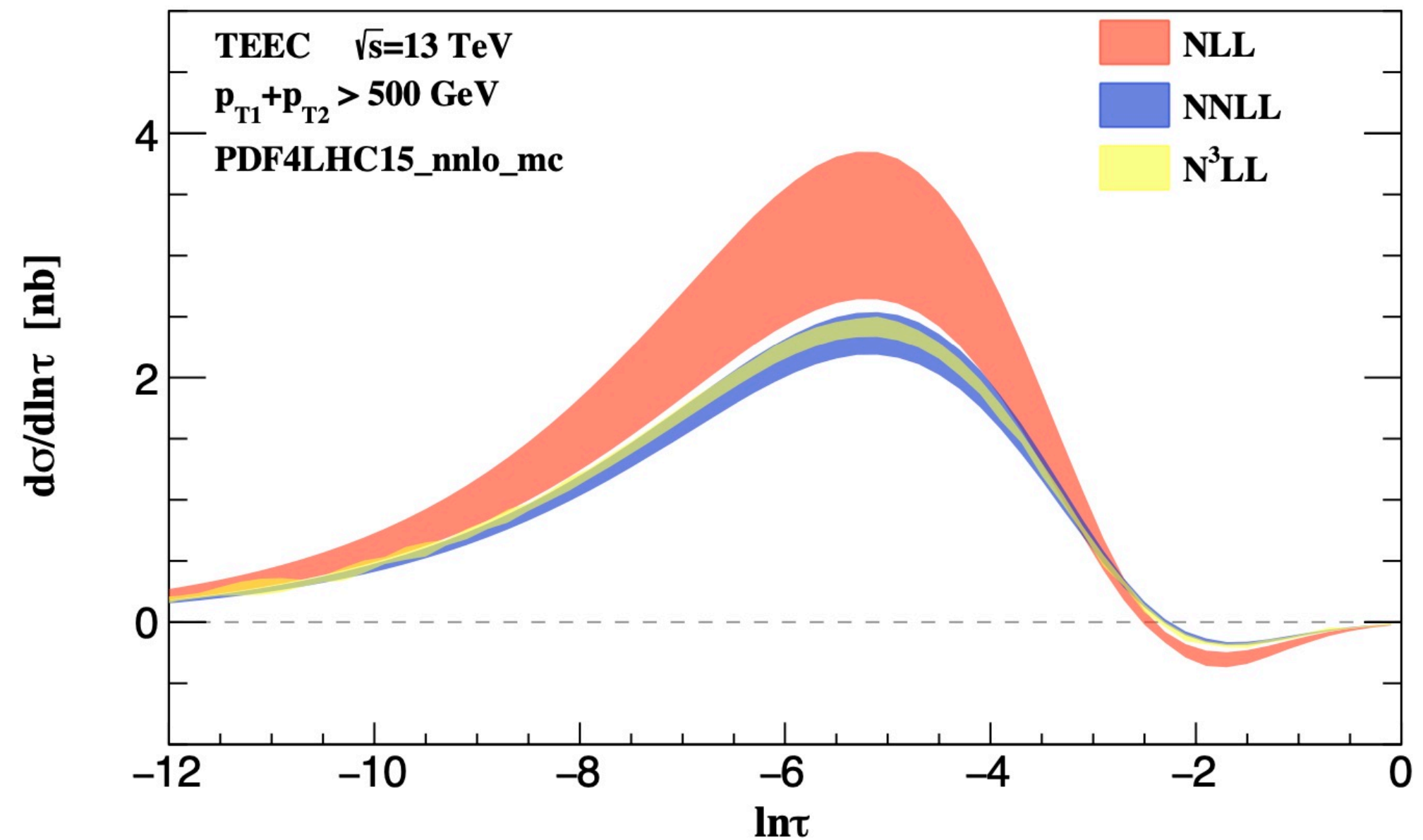
in pp collisions

- ❑ Such as violation in collinear factorization
- ❑ Whether rapidity factorization is still valid
- ❑ RG invariance of the cross section

$$\mathcal{L}^{(0)} = \mathcal{L}_{B_1}^{(0)} + \mathcal{L}_{B_2}^{(0)} + \mathcal{L}_{J_1}^{(0)} + \mathcal{L}_{J_2}^{(0)} + \boxed{\mathcal{L}_G^{(0)}}$$



NNLL accuracy for a hadron collider diet event shape for the first time.

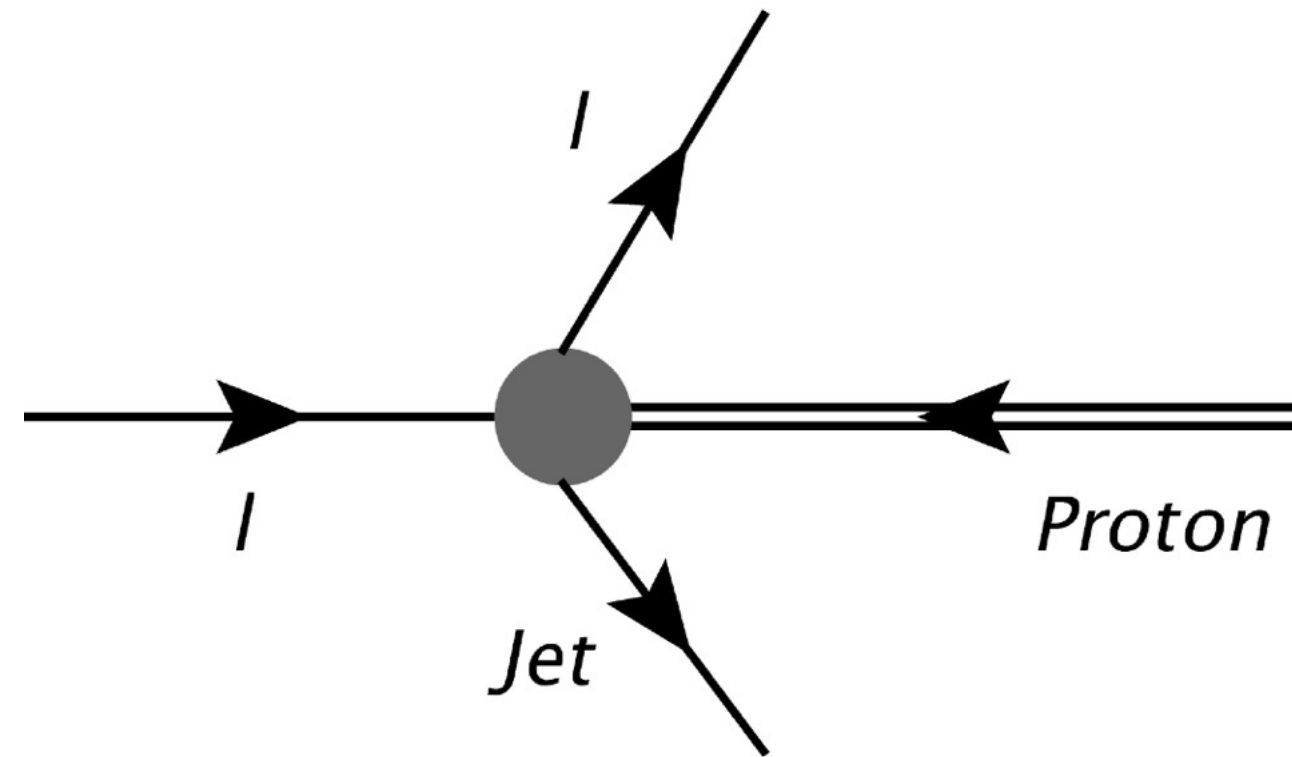


Gao, HTL, Moulton, Zhu, PRL, 19 & JHEP 24



# Introduction

## TEEC in DIS



Definition

$$\text{TEEC} = \sum_a \int d\sigma_{lp \rightarrow l+a+X}$$

sum over  
all hadrons

$$\frac{E_{T,l} E_{T,a}}{E_{T,l} \sum_i E_{T,i}}$$

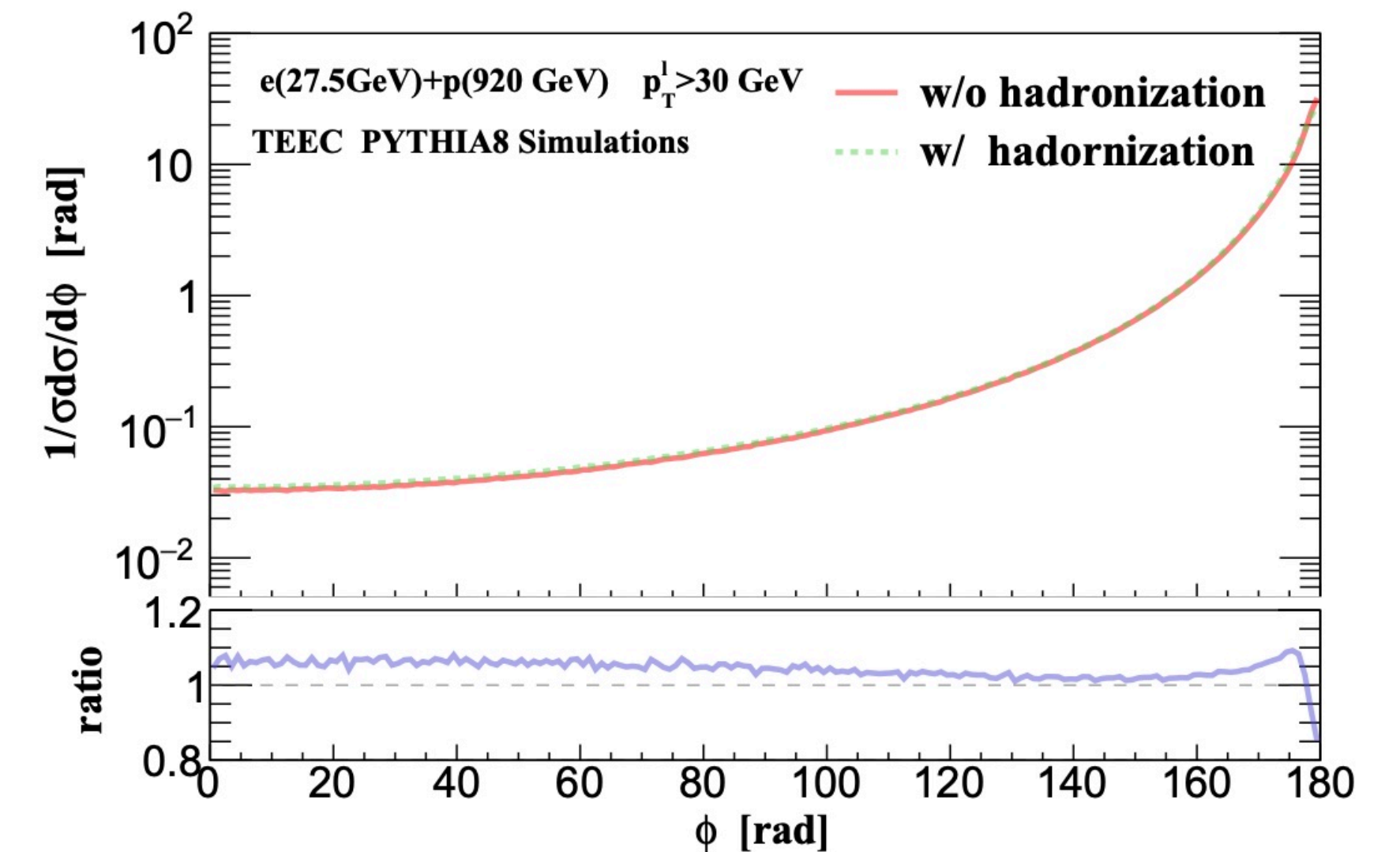
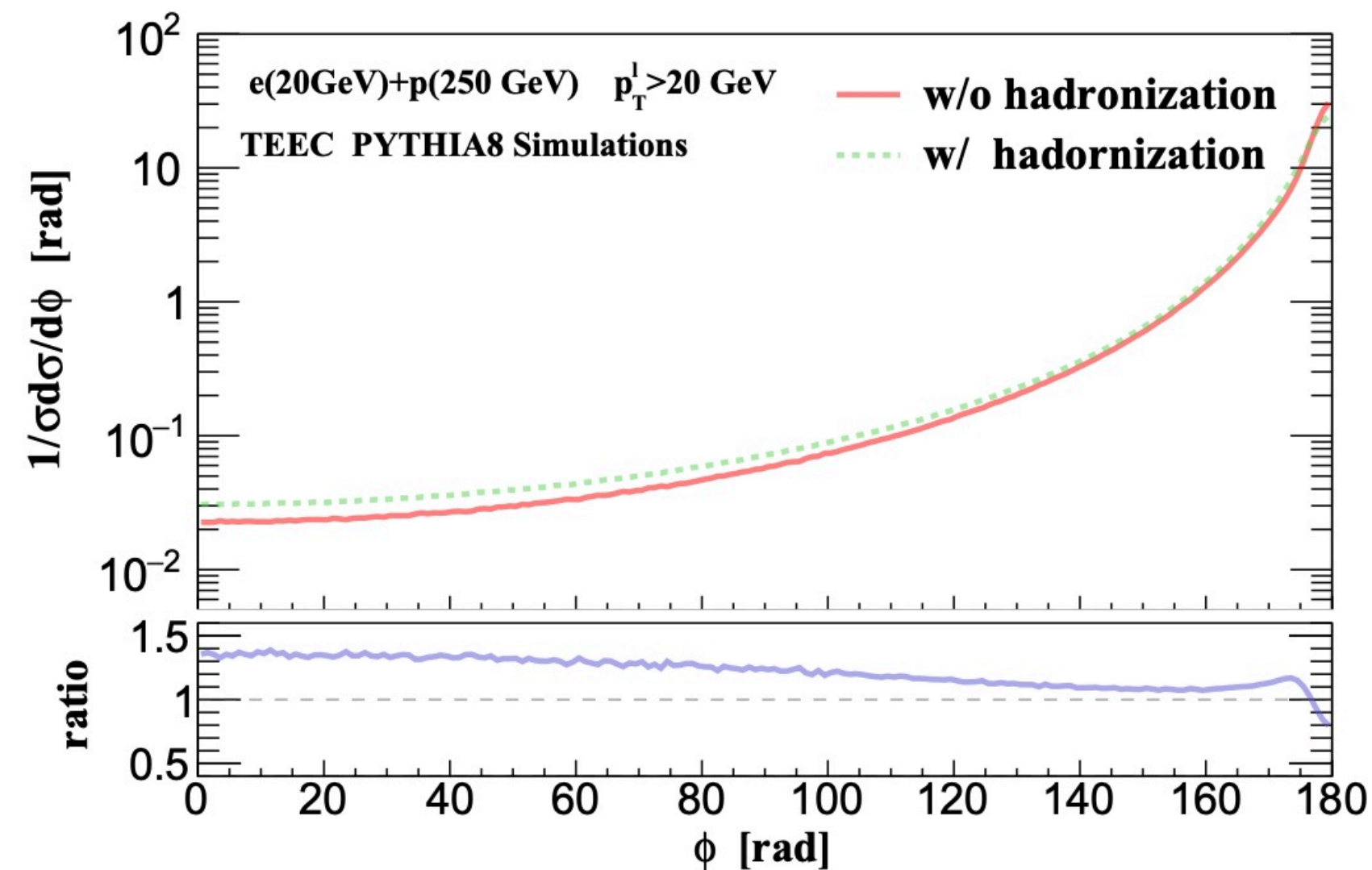
energy weighted

$$\delta(\cos \phi_{la} - \cos \phi)$$

measure azimuthal  
angle correlations

Simulation by Pythia

*HTL, Vitev, Zhu, JHEP, 20*

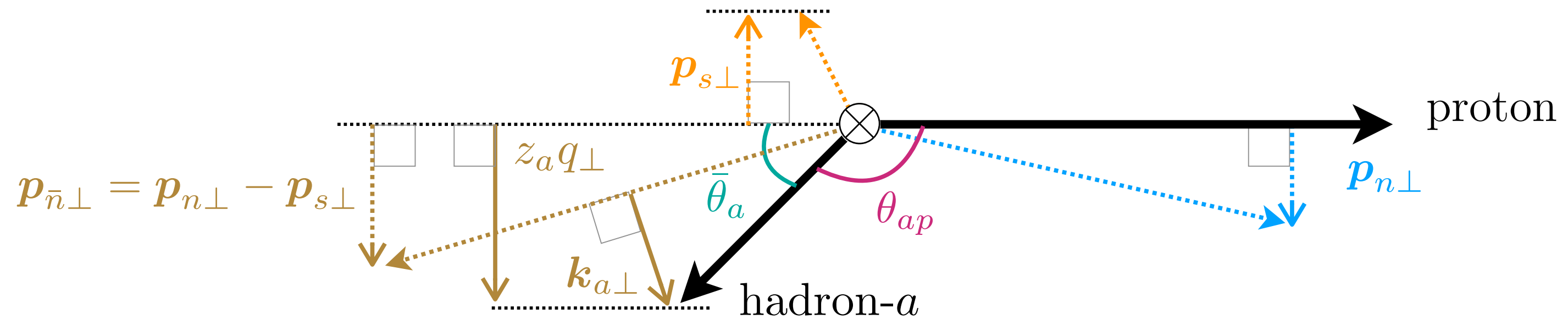


# Introduction

## EEC in DIS

In Breit Frame.

$$\text{proton} + \gamma^* \rightarrow \text{jet/hadron} + X$$



We proposed a new definition of EEC in DIS:

correlation between initial proton and final state hadron

In the back to back limit

$$\text{EEC} = \sum_a \int d\sigma_{lp \rightarrow l+a+X} \left( \frac{p \cdot p_a}{\sum_i p \cdot p_i} \right) \delta(\cos \chi - \cos \theta_{ap})$$

$$\frac{d\sigma_h}{d^2p_{\perp}} = \sum_f \int \frac{d\xi dQ^2}{\xi Q^2} Q_f^2 H(Q, \mu) \int \frac{db}{2\pi} e^{ib_{\perp} \cdot p_{\perp}} f_{f/N}(b, \xi, \mu, \nu) S\left(b, \frac{n_2 \cdot n_4}{2}, \mu, \nu\right) \int dz F_{h/f}(z, b/z, E_4, \mu, \nu)$$



# Introduction

## EEC in DIS

In Breit Frame.

$$\text{proton} + \gamma^* \rightarrow \text{jet/hadron} + X \quad \mathcal{P}$$

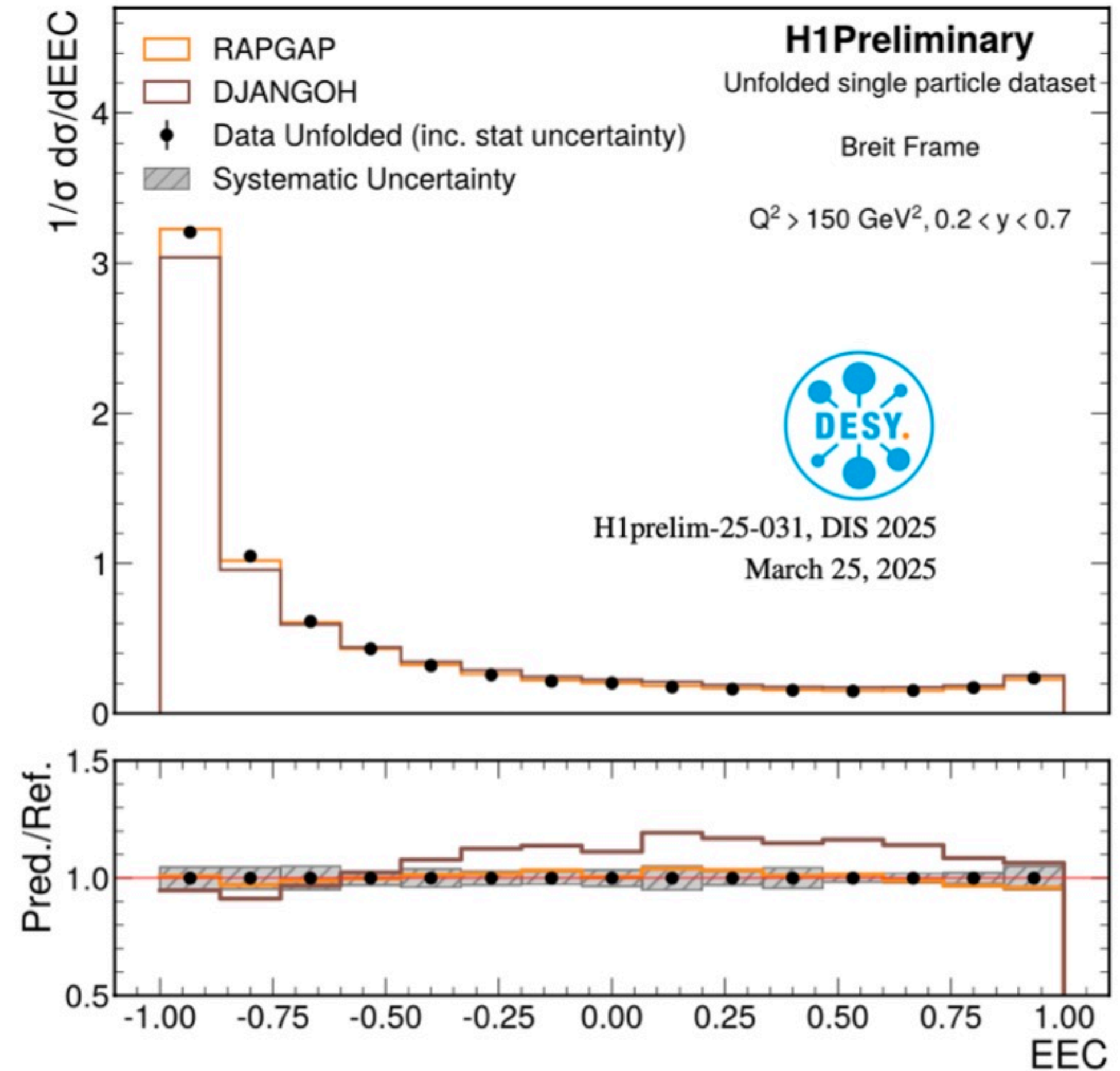
We proposed a new definition of EEC correlation

In the back to back limit

$$\text{EEC} = \sum_a$$

$$\frac{d\sigma_h}{d^2p_\perp} = \sum_f \int \frac{d\xi dQ^2}{\xi Q^2} Q_f^2 H(Q, \mu) \int \frac{db}{2\pi} e^{ib_\perp \cdot p_\perp}$$

HTL, Makris, Vitev, PRD, 2021



→ proton  
 $p_{n\perp}$

)

# Introduction

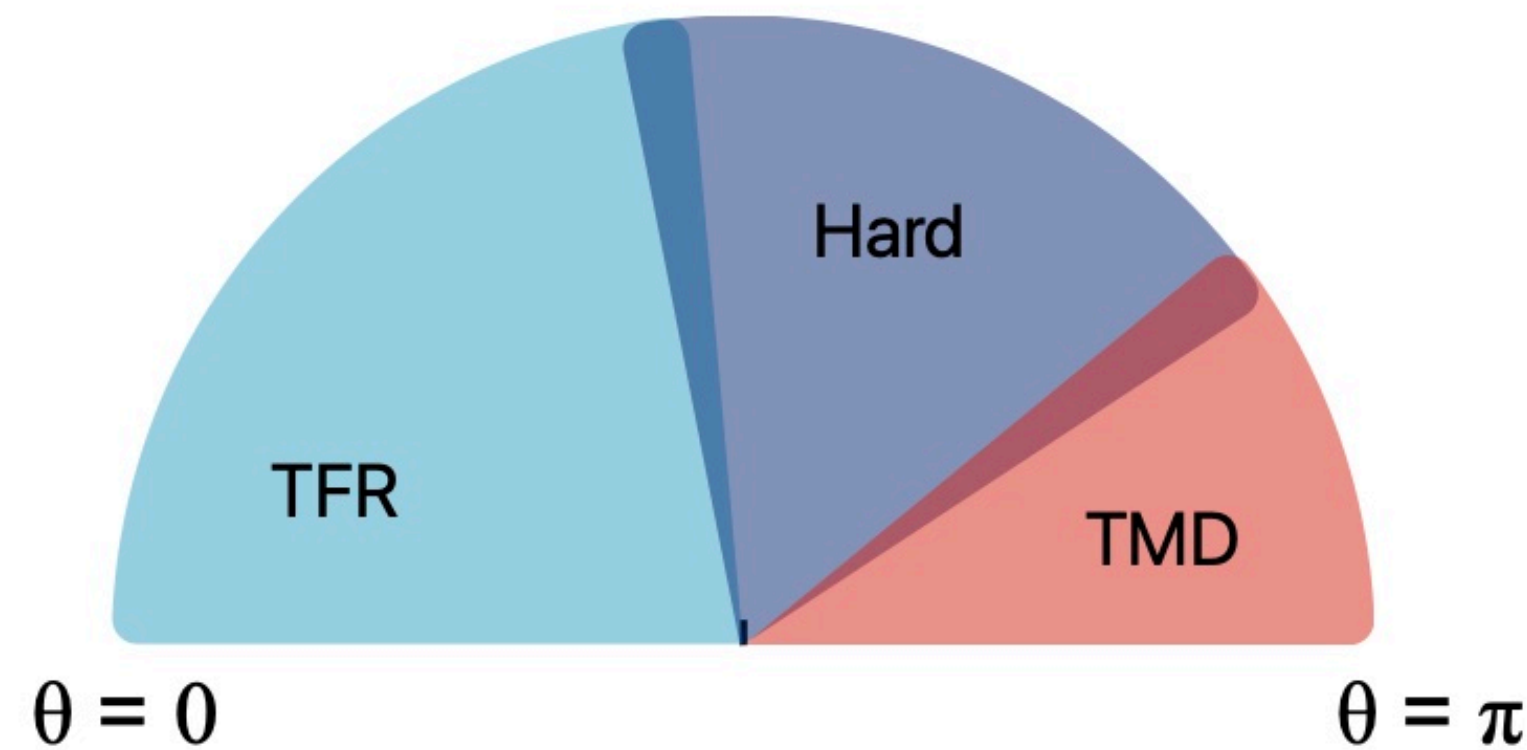
## NEEC in DIS

Nucleon Energy Correlators

$$\Sigma_N(Q^2, \theta^2) = \sum_i \int d\sigma(x_B, Q^2, p_i) x_B^{N-1} \frac{\bar{n} \cdot p_i}{P} \delta(\theta^2 - \theta_i^2)$$

*Liu, Zhu, arXiv:2209.02080*

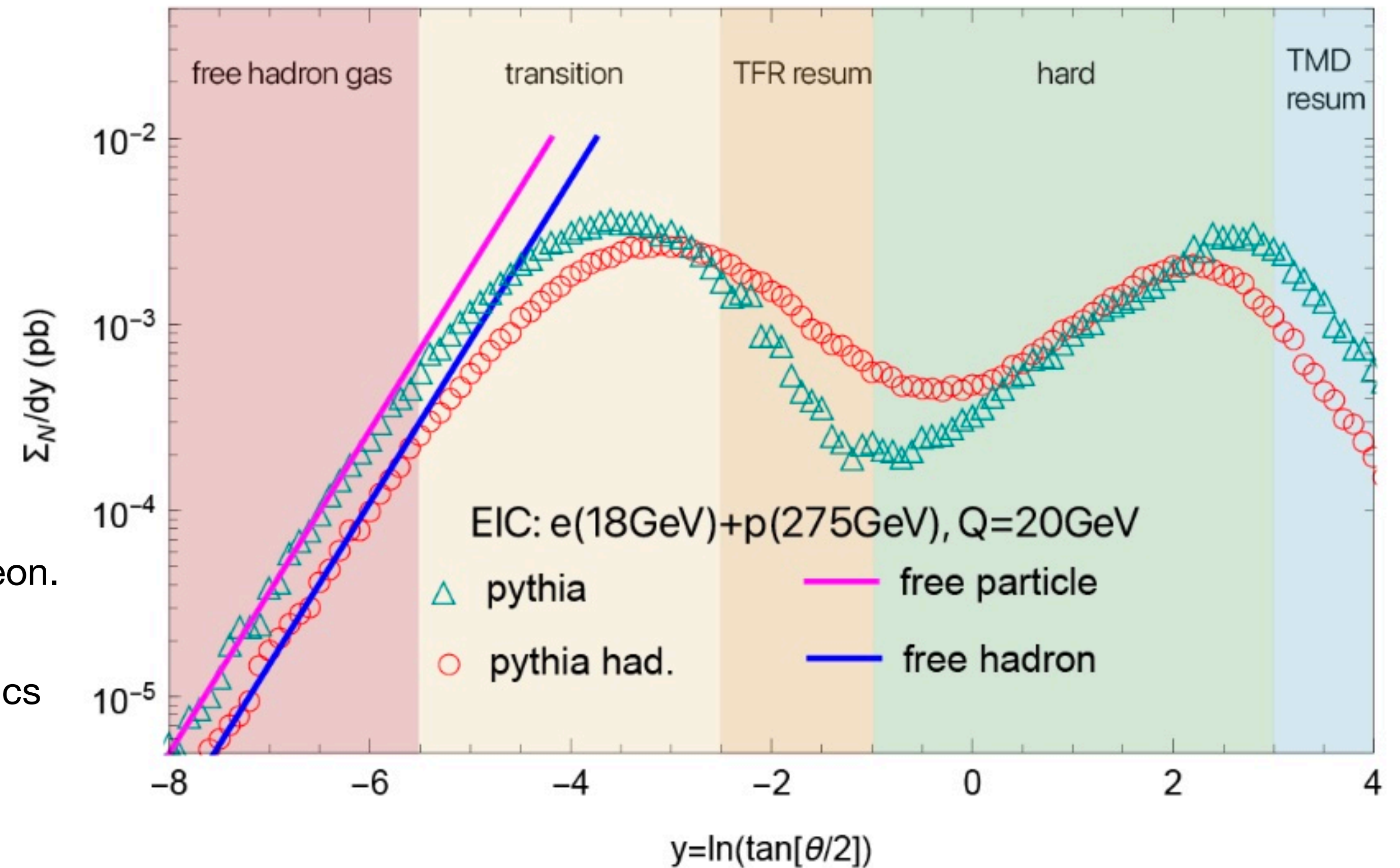
*Cao, Liu, Zhu, arXiv:2303.01530*



TFR: the correlation of the energy flows from the initial nucleon.

Hard: measures the perturbative behavior of QCD

TMD: measures perturbative and nonperturbative TMD physics



*Cao, HTL, Mi, PRD, 24*



# Introduction

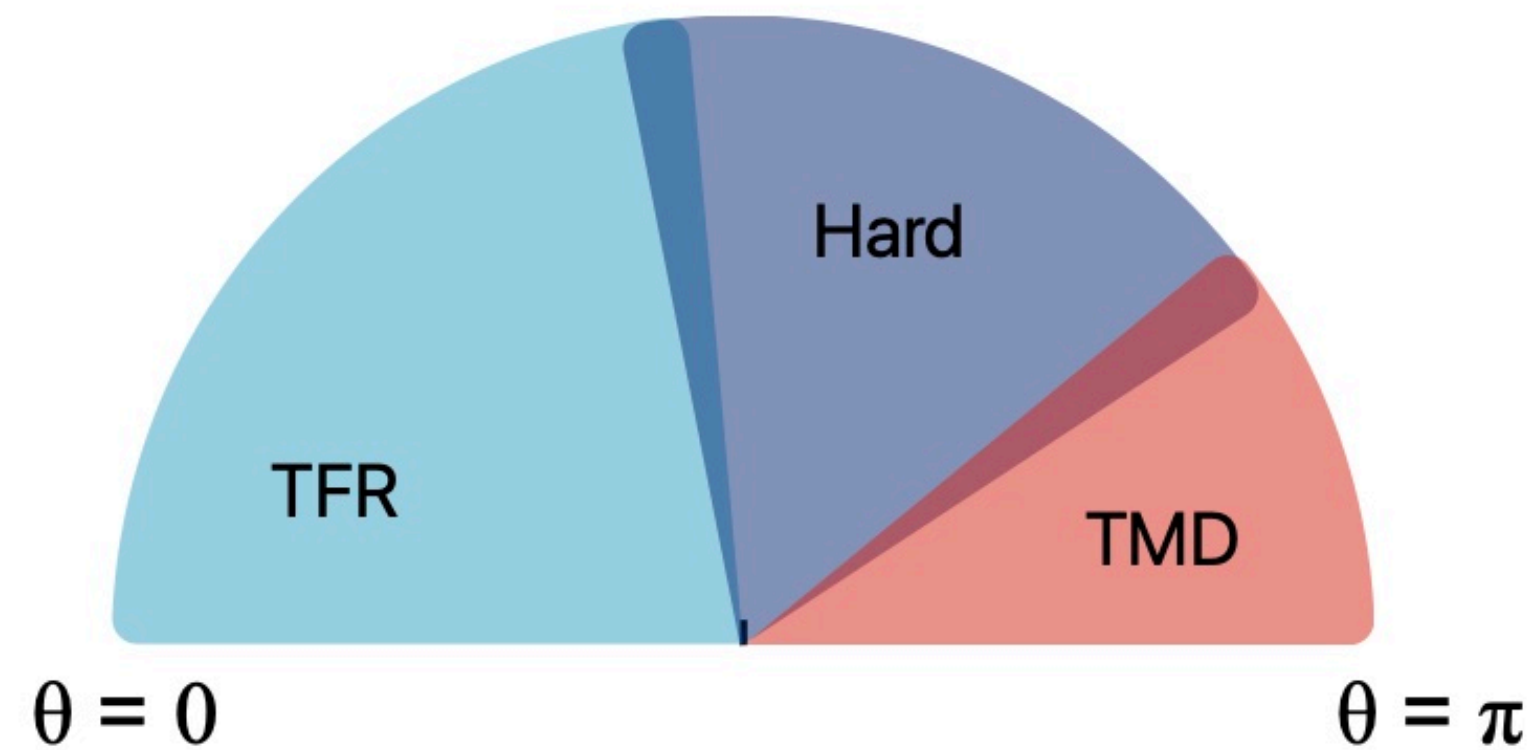
## NEEC in DIS

Nucleon Energy Correlators

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*Liu, Zhu, arXiv:2209.02080*

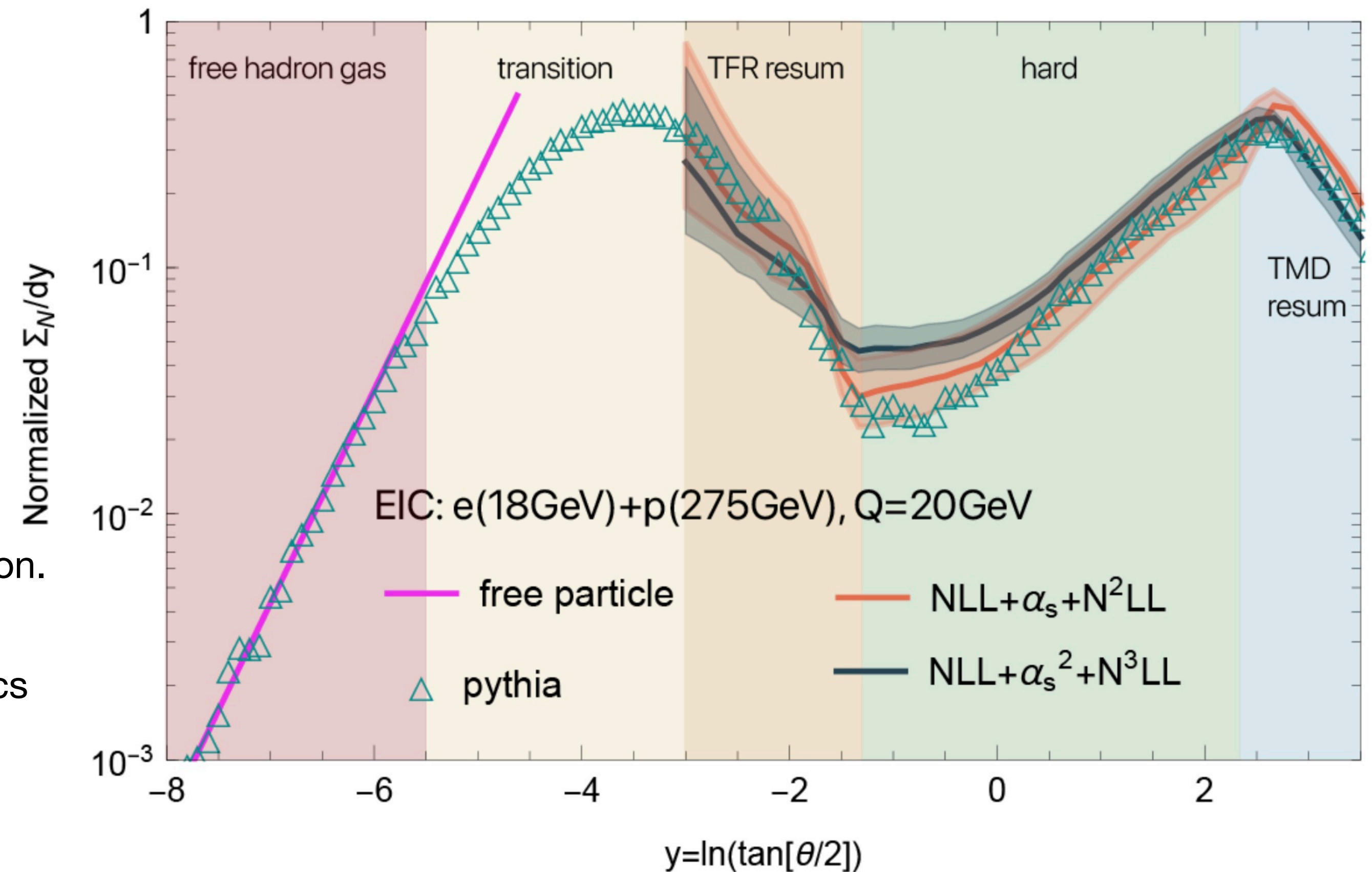
*Cao, Liu, Zhu, arXiv:2303.01530*



TFR: the correlation of the energy flows from the initial nucleon.

Hard: measures the perturbative behavior of QCD

TMD: measures perturbative and nonperturbative TMD physics



*Cao, HTL, Mi, PRD, 24*



# Outline

- Introduction (EEC in back to back limit)
- EEC with jets
- EEC and  $\alpha_s$  measurements
- Summary

# Introduction

$$\text{TEEC} = \sum_{a,b} \int d\sigma_{pp \rightarrow a+b+X} \frac{2E_{T,a}E_{T,b}}{|\sum_i E_{T,i}|^2} \delta(\cos \phi_{ab} - \cos \phi)$$

## Jets as input

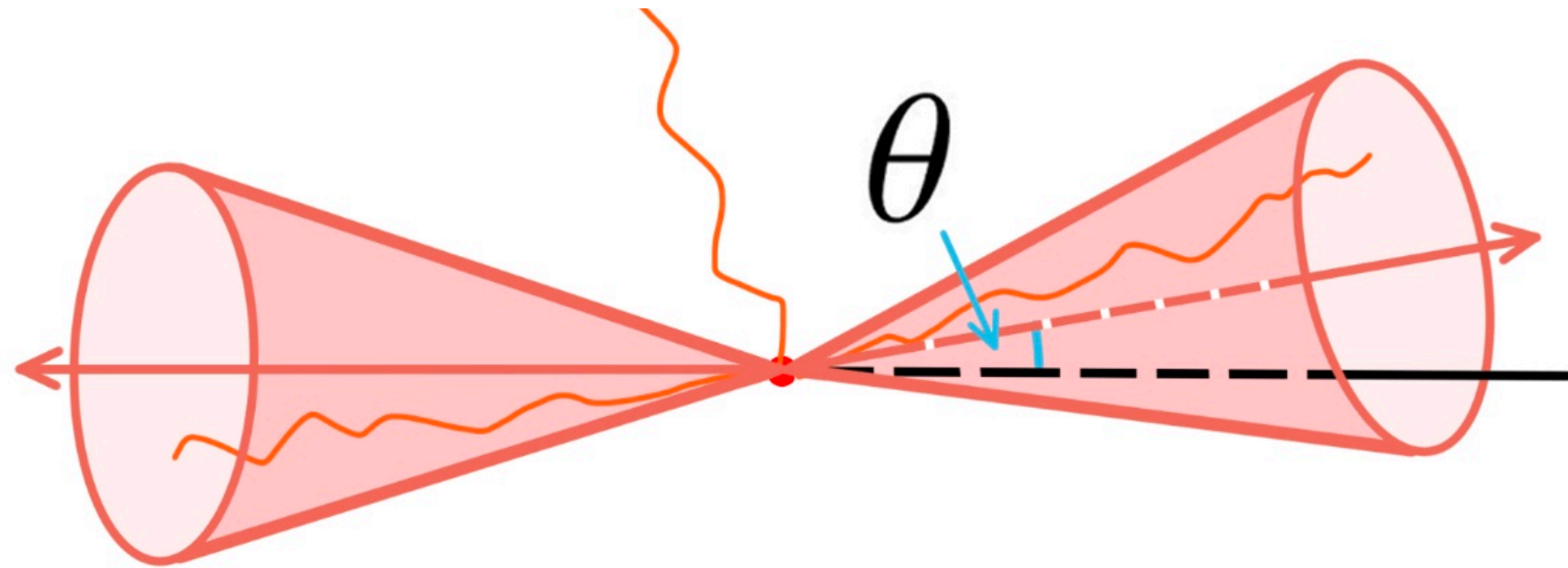
- ❑ The jets are usually defined through an angular resolution parameter R and a transverse momentum cutoff
- ❑ Less contamination from underlying events
- ❑ Introduce new cutoff scale
- ❑ Problem of globalness, i.e. non-global logs
- ❑ Measurable at LHC and EIC

## Particles as input

- ❑ Simpler defined observable
- ❑ relative larger effects from underlying events
- ❑ NO new cutoff scale introduced
- ❑ Global event shapes
- ❑ Measurable at LHC and EIC

# EEC with jets

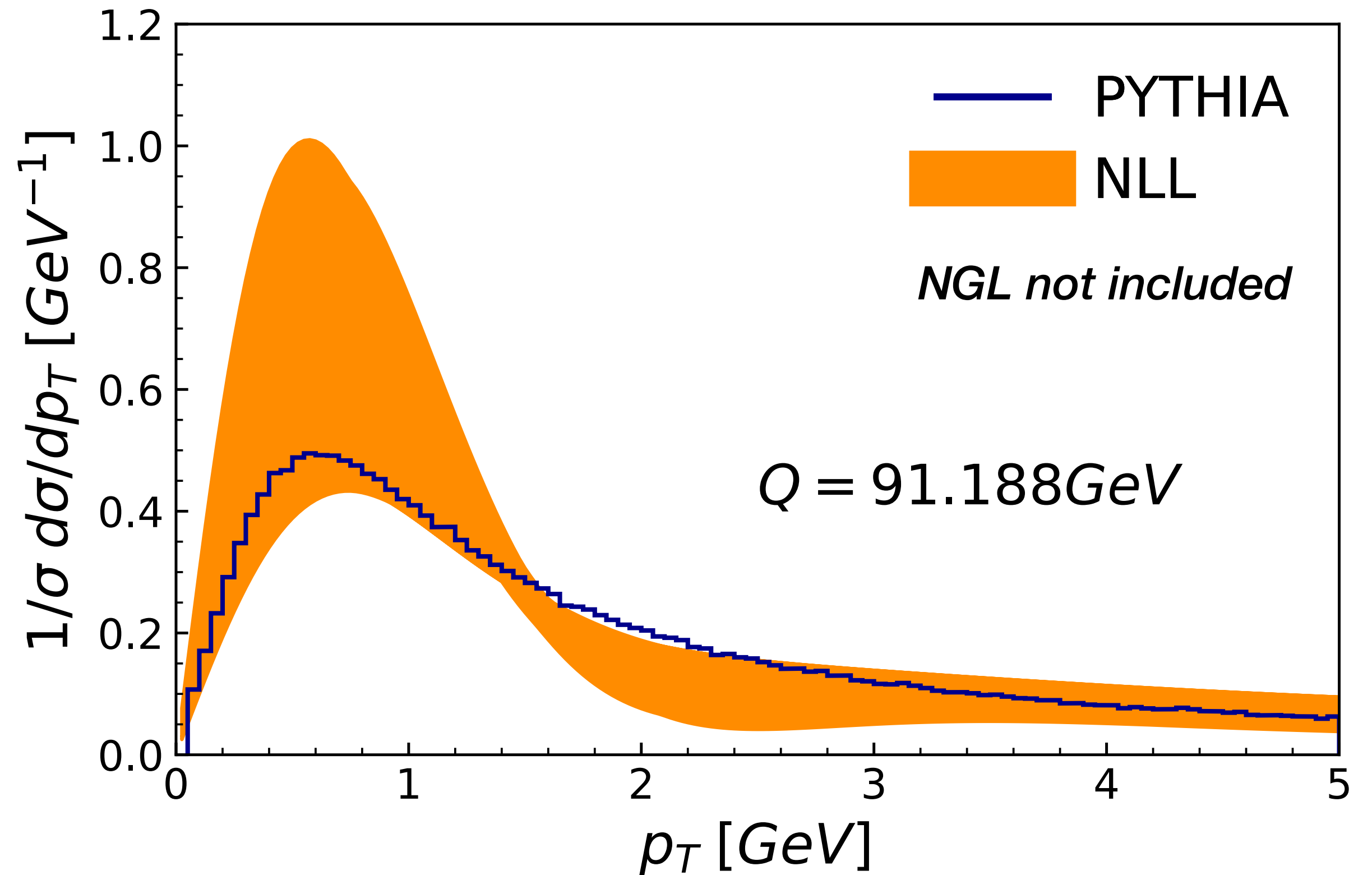
$$\text{EEC} = \sum_{a,b} \int d\sigma_{V \rightarrow a+b+X} \frac{2E_a E_b}{Q^2 \sigma_{\text{tot}}} \delta(\cos(\theta_{ab}) - \cos(\chi))$$



In jet, soft and/or collinear radiations do not change the opening angle of the jets

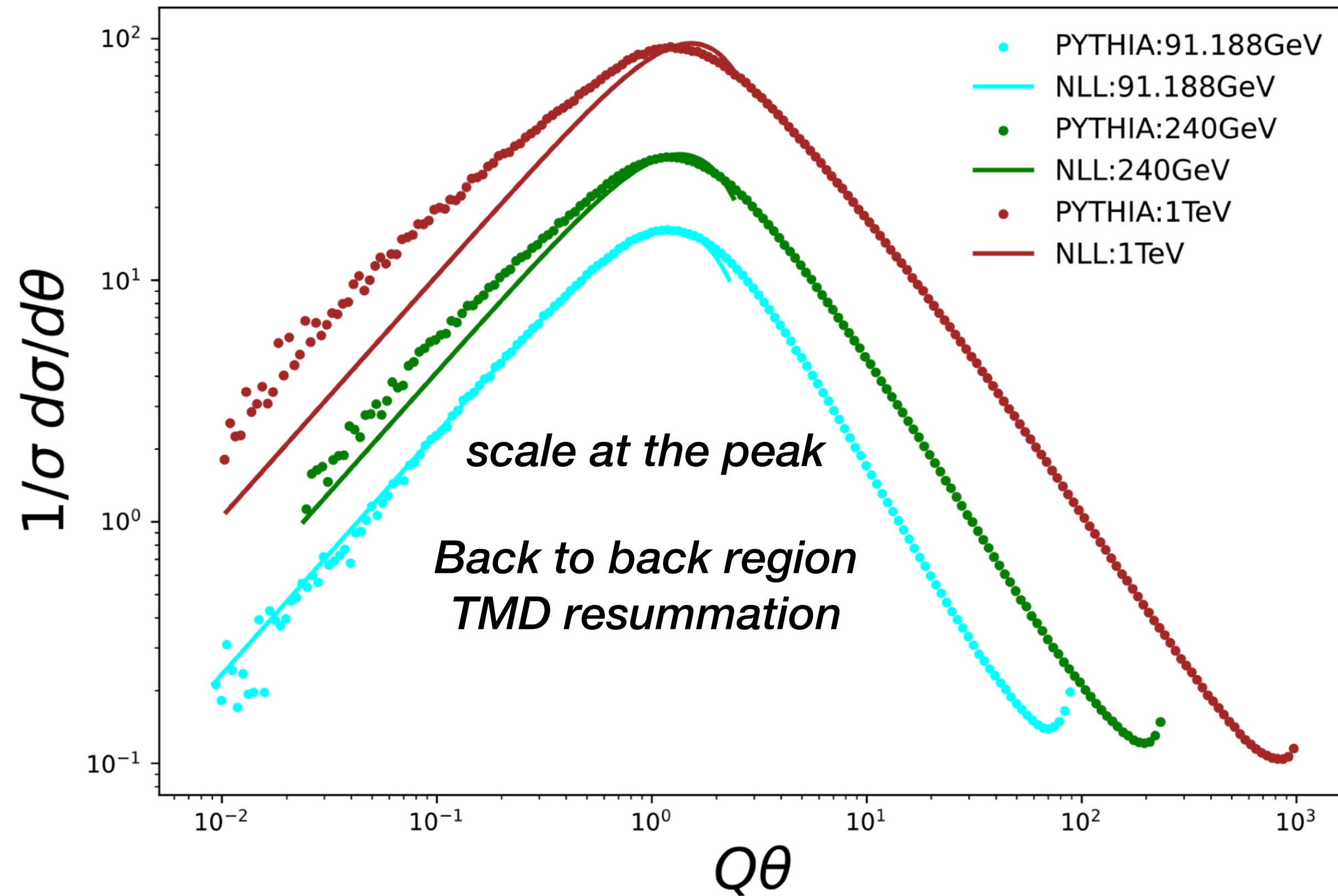
Outside of jet, soft radiation would shift the angle away from back to back region

$$\frac{1}{\sigma_0} \frac{d\sigma}{d^2\mathbf{p}_\perp} = H(Q^2, \mu) J_c(Q^2, \delta, \mu) J_{\bar{c}}(Q^2, \delta, \mu) \times \int \frac{d\vec{b}}{(2\pi)^2} e^{i\vec{b} \cdot \vec{p}_\perp} S(Q^2, \vec{b}, \mu, \nu) S_{sc}(\delta Q^2, \vec{b}, \mu, \nu) S_{s\bar{c}}(\delta Q^2, \vec{b}, \mu, \nu)$$

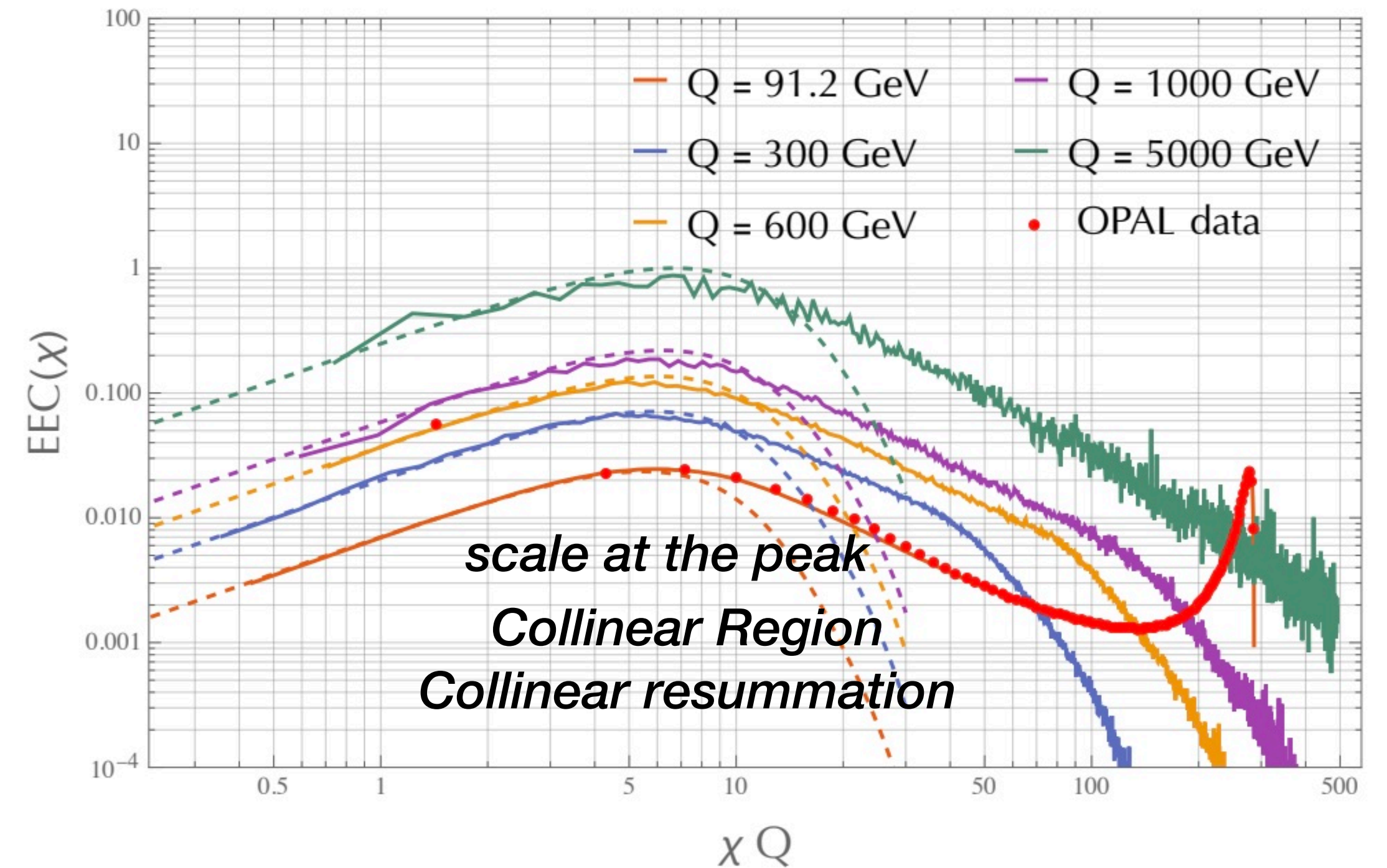




# EEC with jets



He, HTL, Shao, in preparation



$$c^3 N E_J^2 \chi \int db b J_0(c E_J \chi b) e^{-\frac{g_1}{c^2} b^2} \times e^{-\frac{g_2}{2} \ln\left(\frac{b}{b_*}\right) \ln \frac{c E_J}{\mu_0}}$$

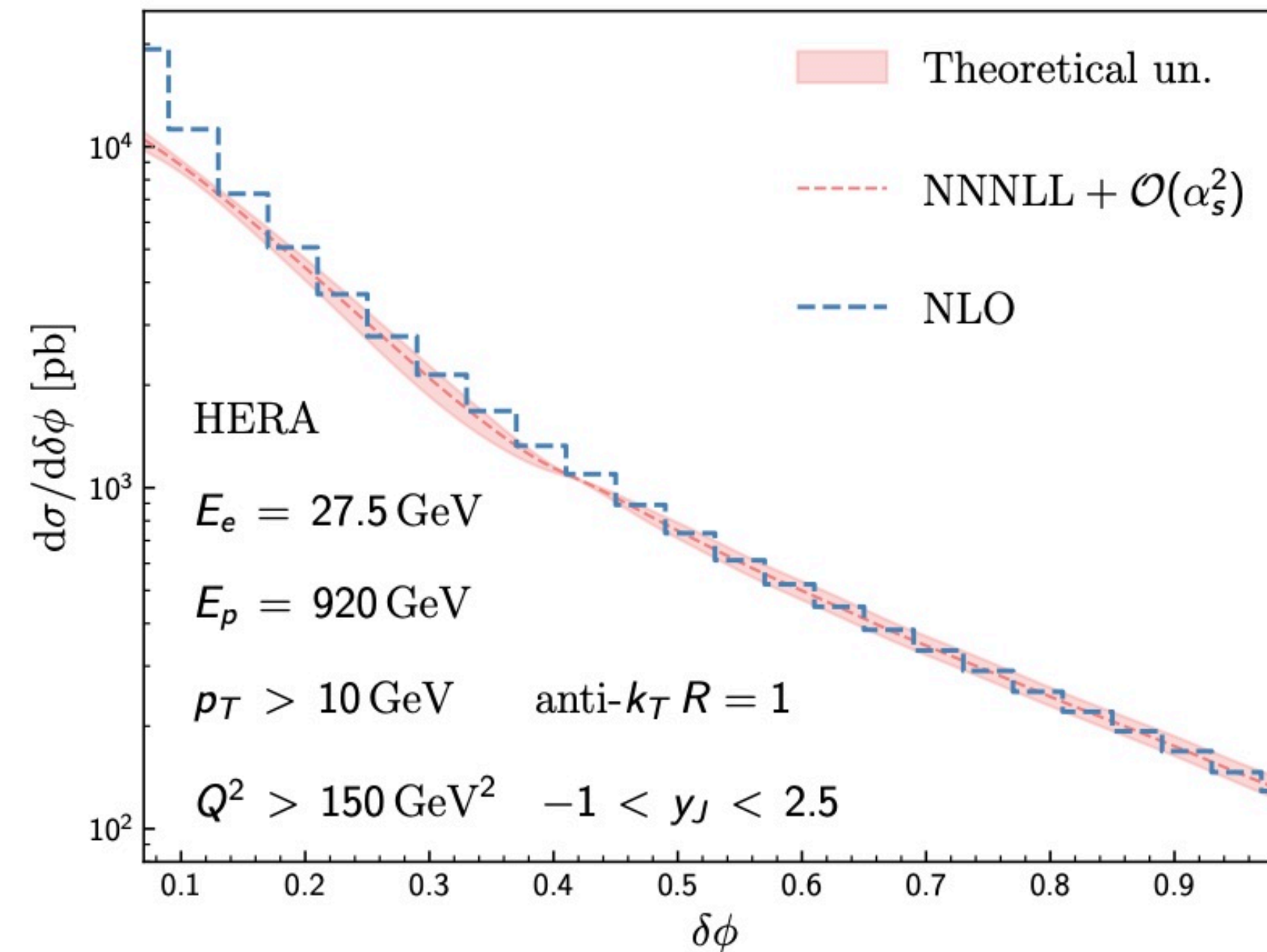
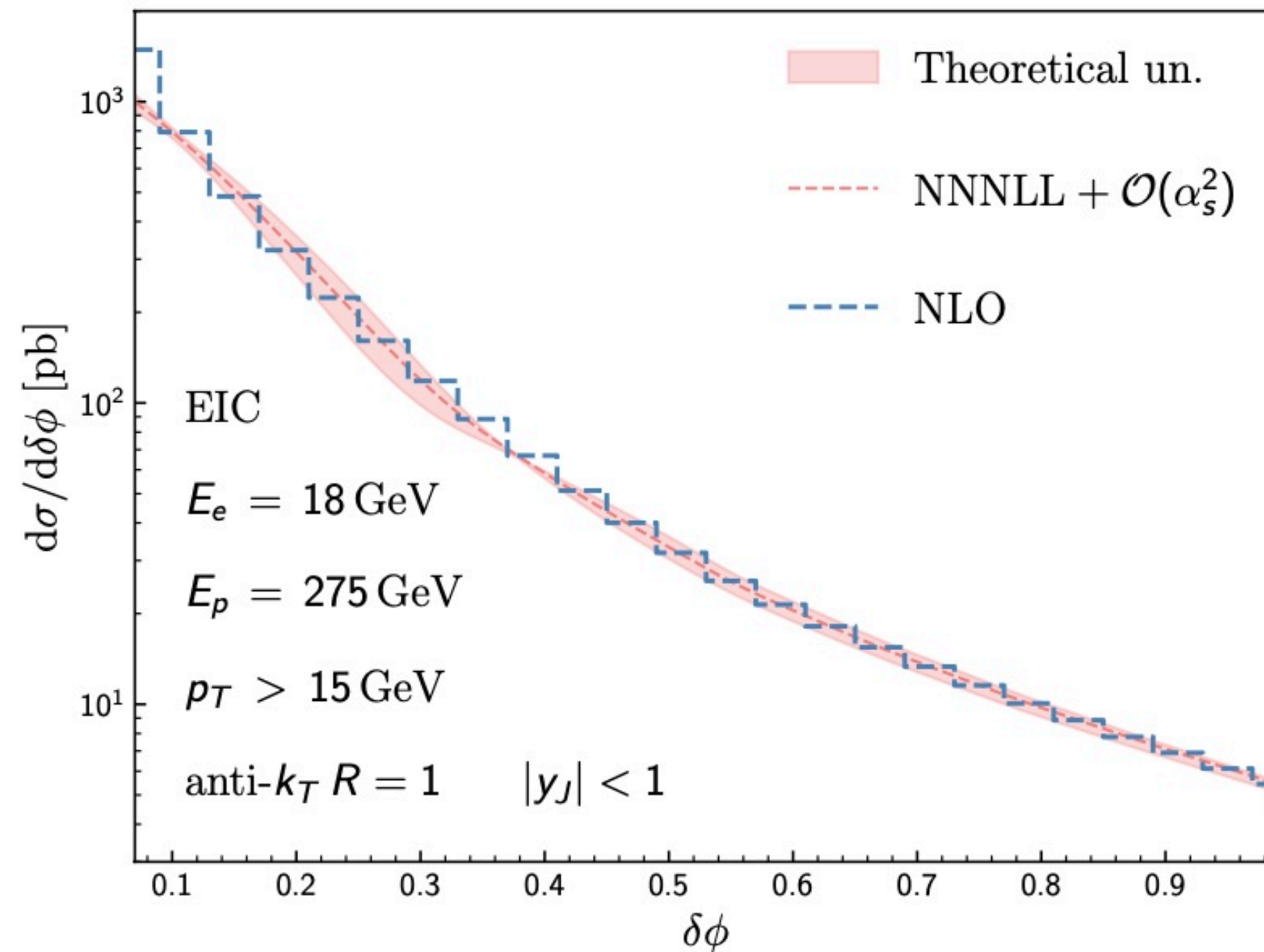
Liu, Vogelsang, Yuan, Zhu, PRL 2025

# EEC with jets

## WTA jet axis in ee and DIS

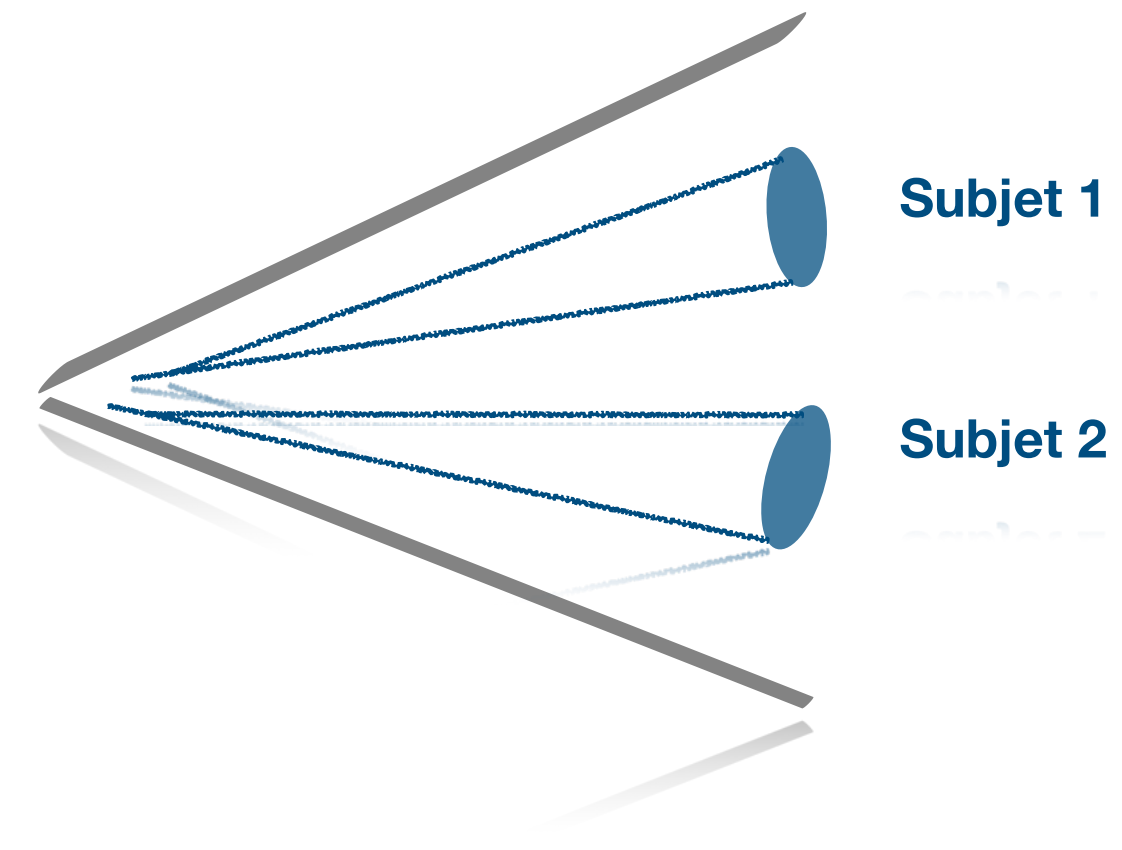
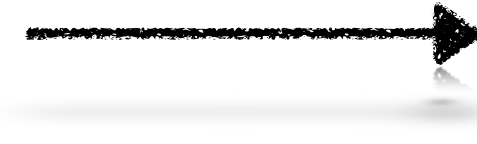
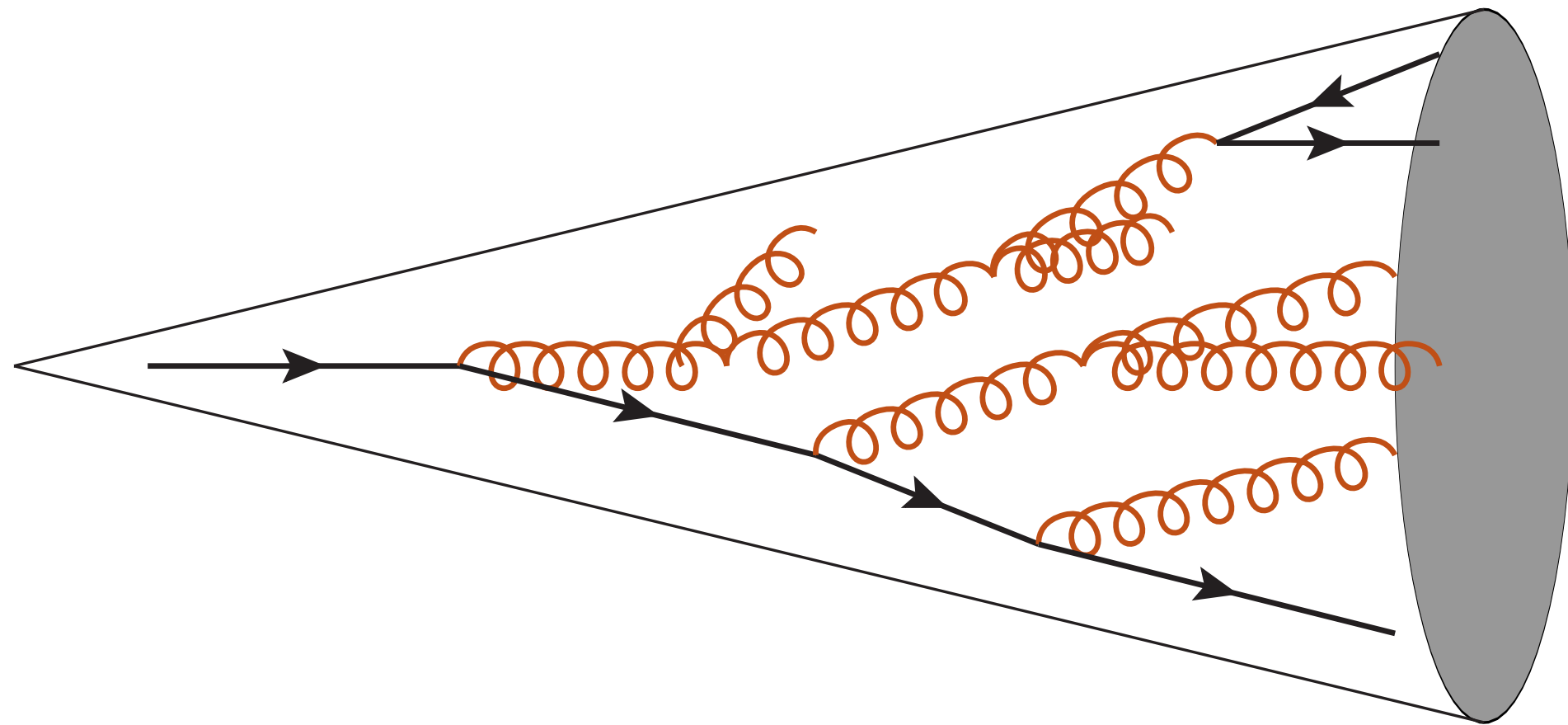
Jets are defined by the anti- $k_T$  clustering algorithm and the winner-take-all recombination scheme.

$$\frac{d\sigma}{d^2p_T dy_J d\lambda_x} = \sigma_0 H(Q, \mu) \int_{-\infty}^{+\infty} \frac{db_x}{2\pi} e^{ib_x \lambda_x} \sum_q e_q^2 \mathcal{B}_{q/p}(x_{bj}, b_x, \mu, \zeta_B/\nu^2) \\ \times \mathcal{J}_q(b_x, \mu, \zeta_J/\nu^2) \mathcal{S}(b_x, n \cdot n_J, \mu, \nu),$$





# EEC with jets



Original jet with radius

Undo last stage of C/A clustering

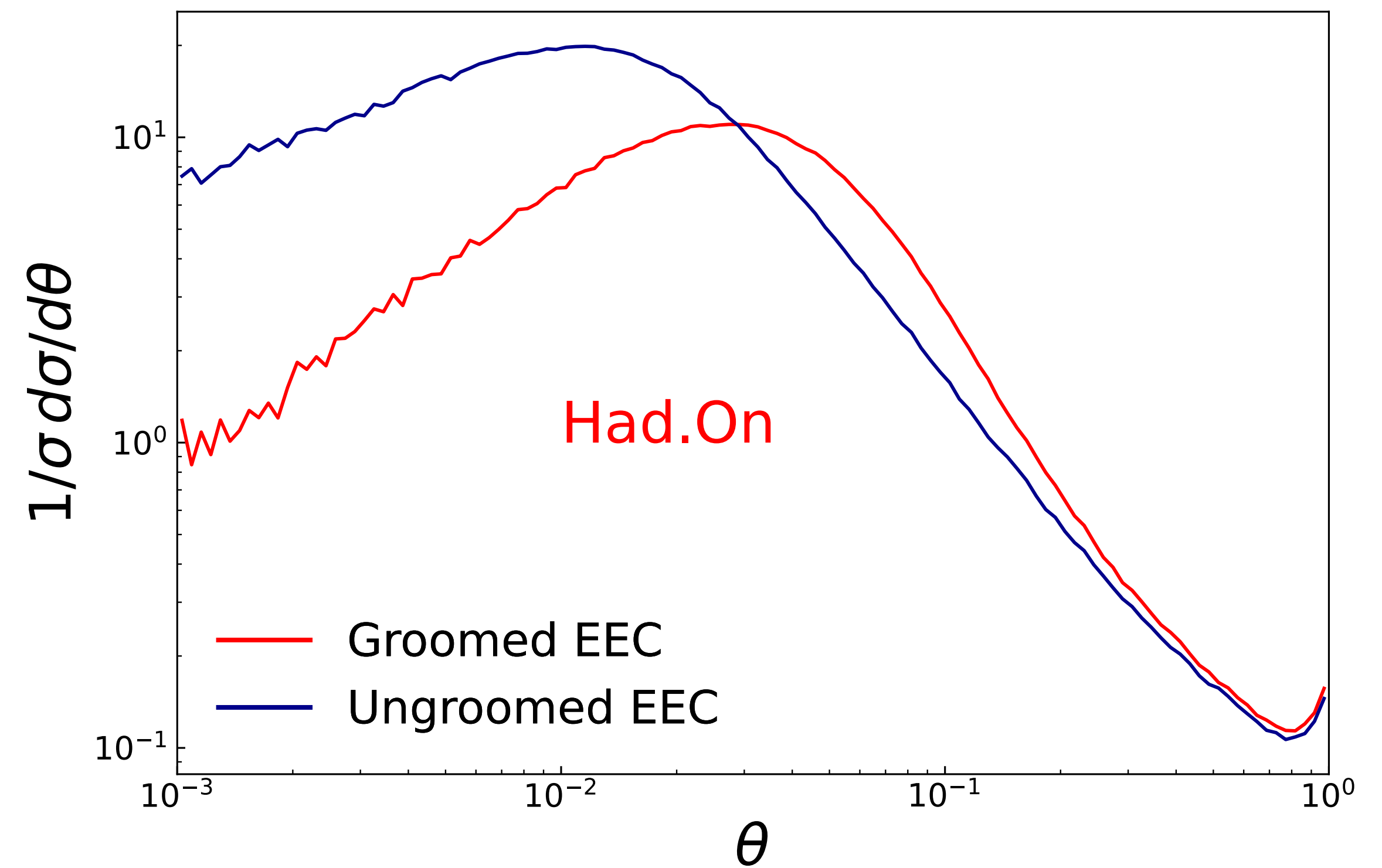
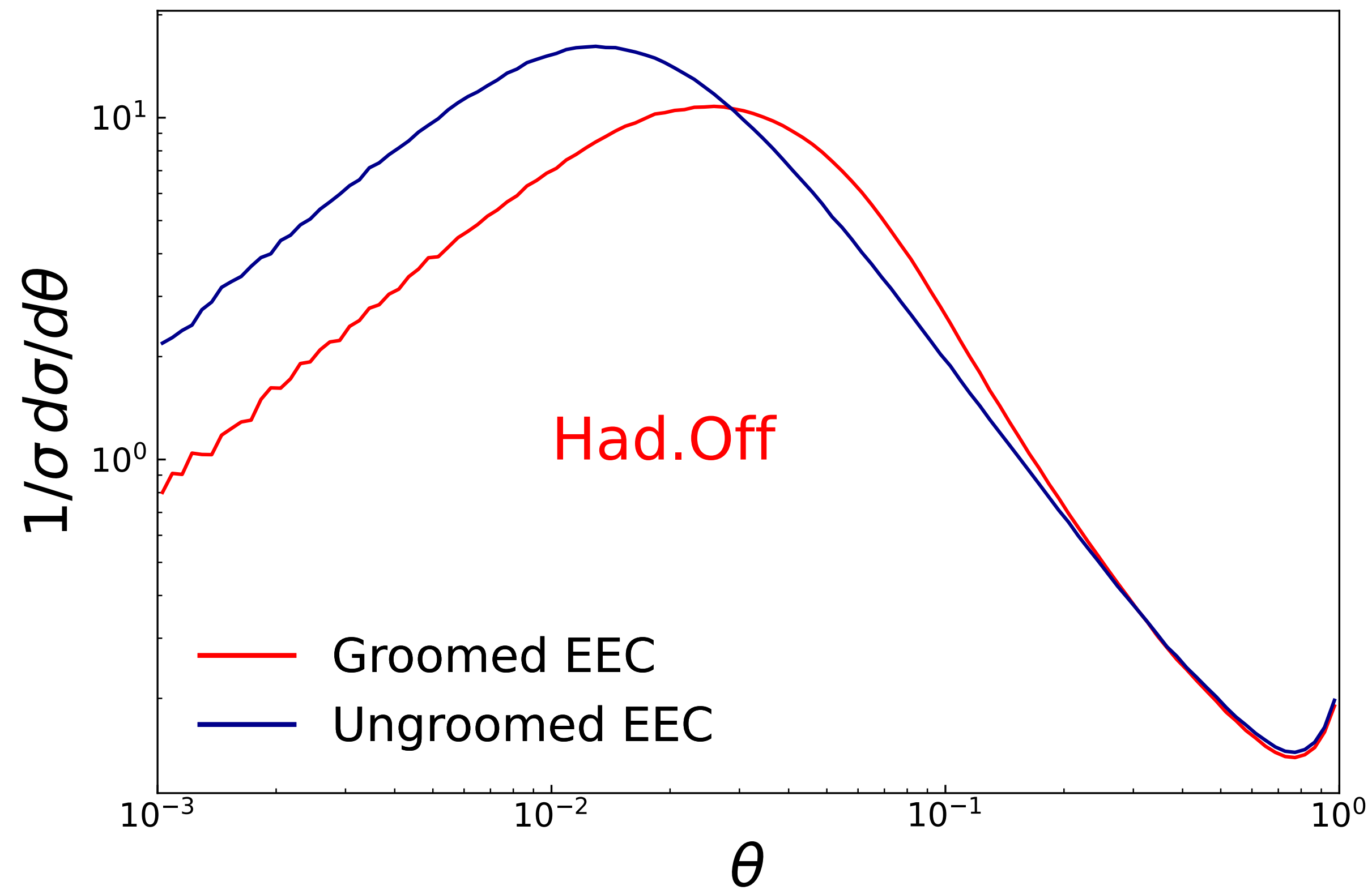
Define 
$$z_g = \frac{\min(p_{T1}, p_{T2})}{p_{TG1} + p_{T2}}$$

If  $z_g < z_{\text{cut}} \left( \frac{\Delta R_{12}}{R_0} \right)^\beta$  redefine j to be the harder one,  
else we have the two-prong subjects



# EEC with jets

Pythia simulation in back to back region

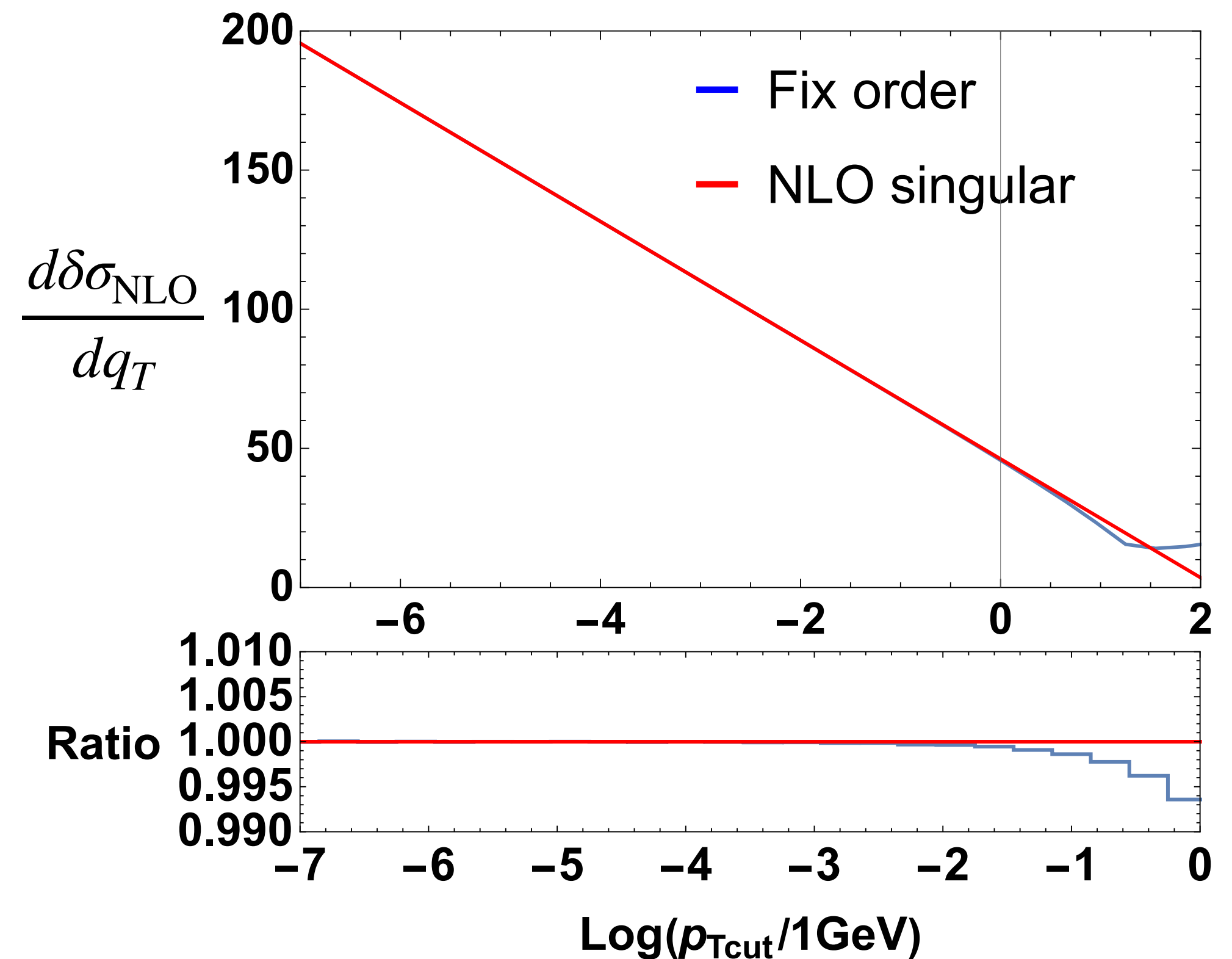
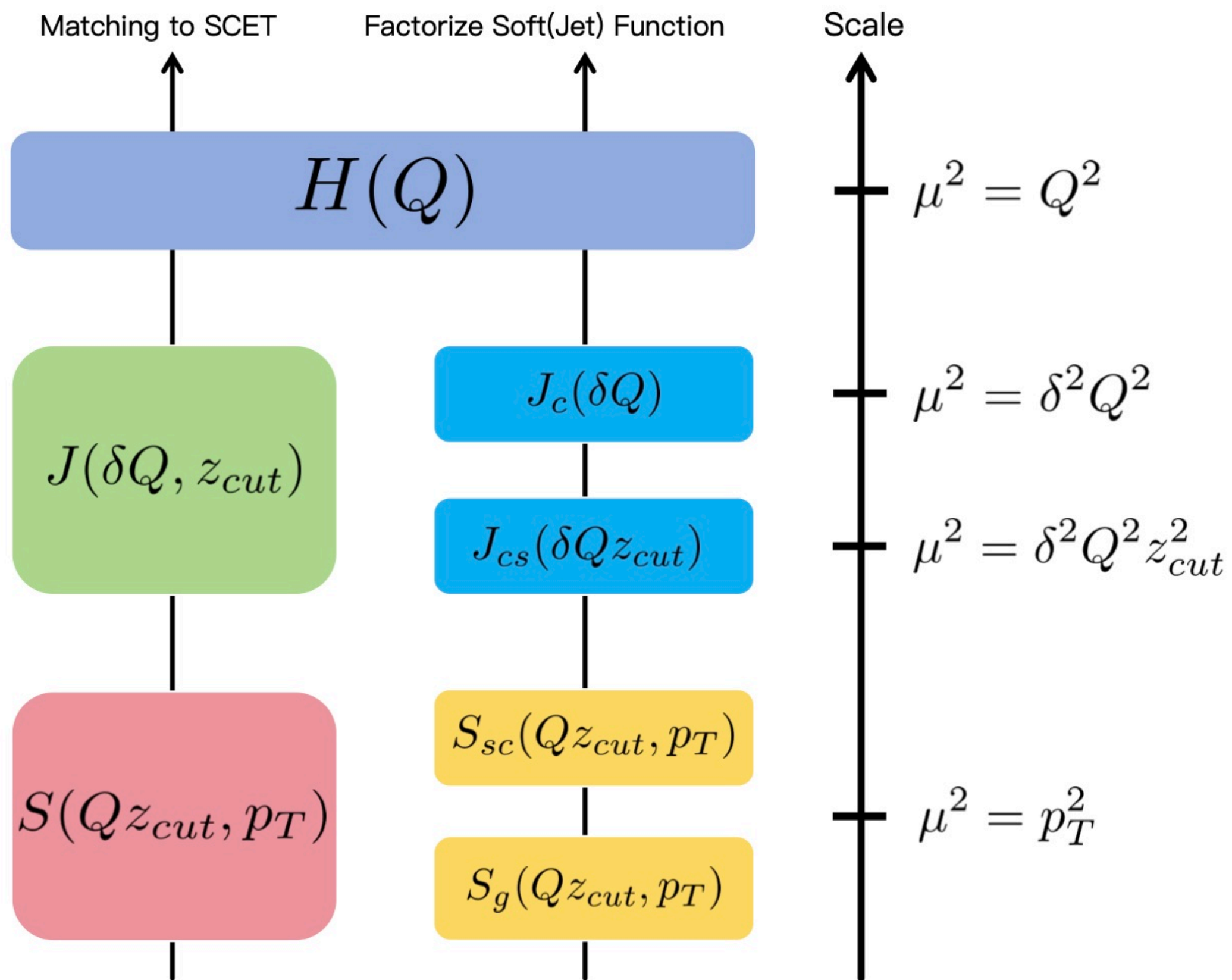


# EEC with jets

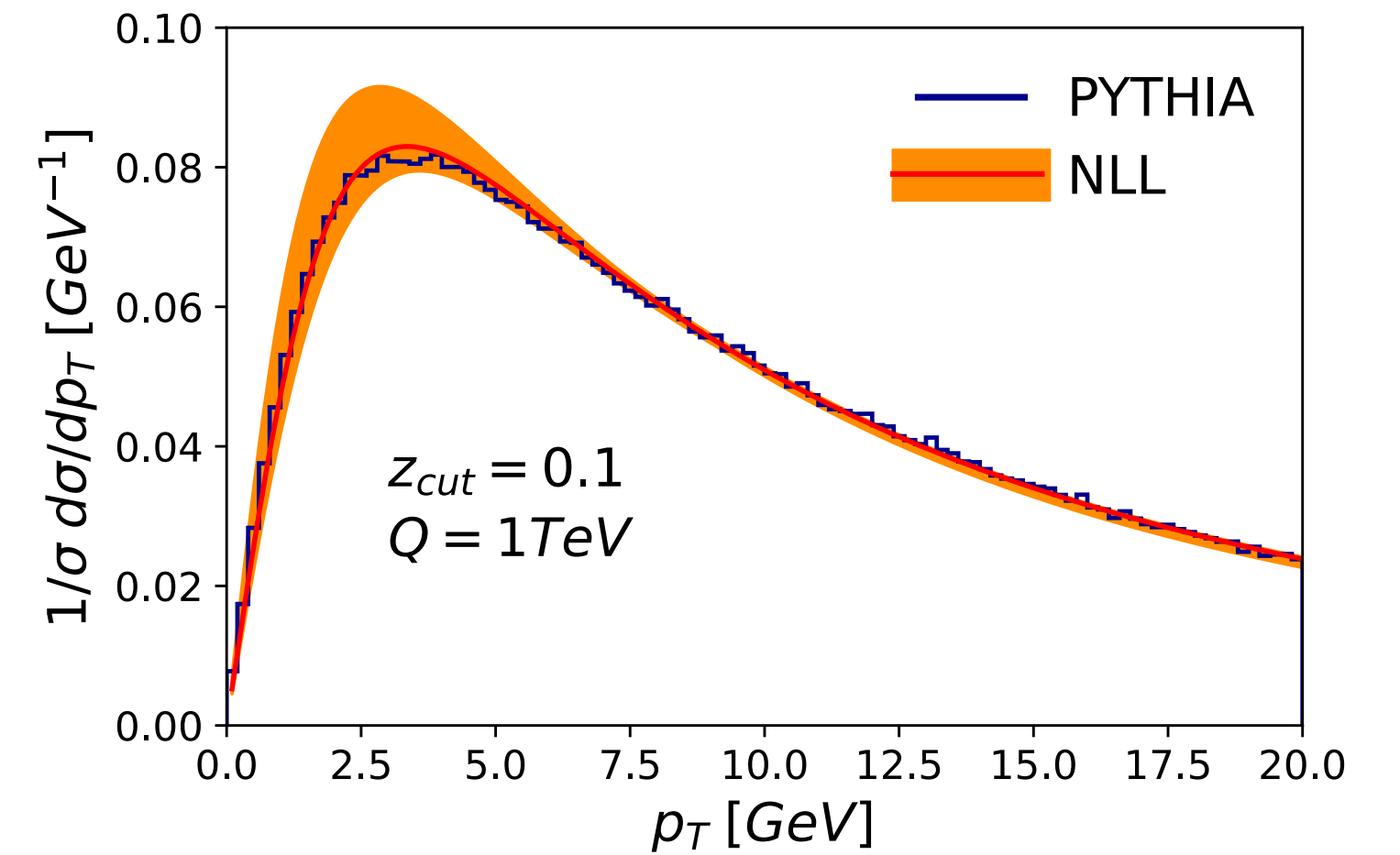
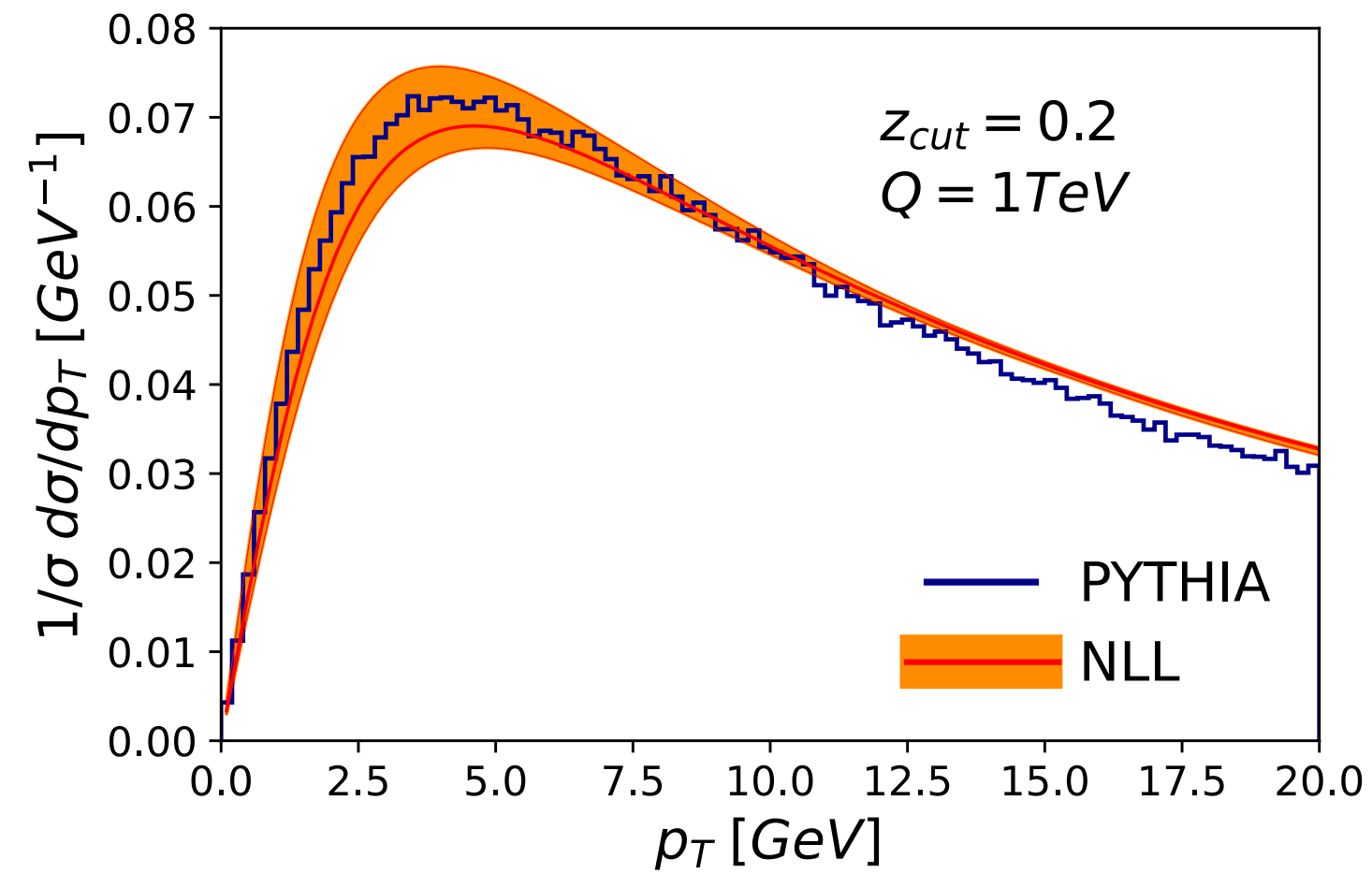
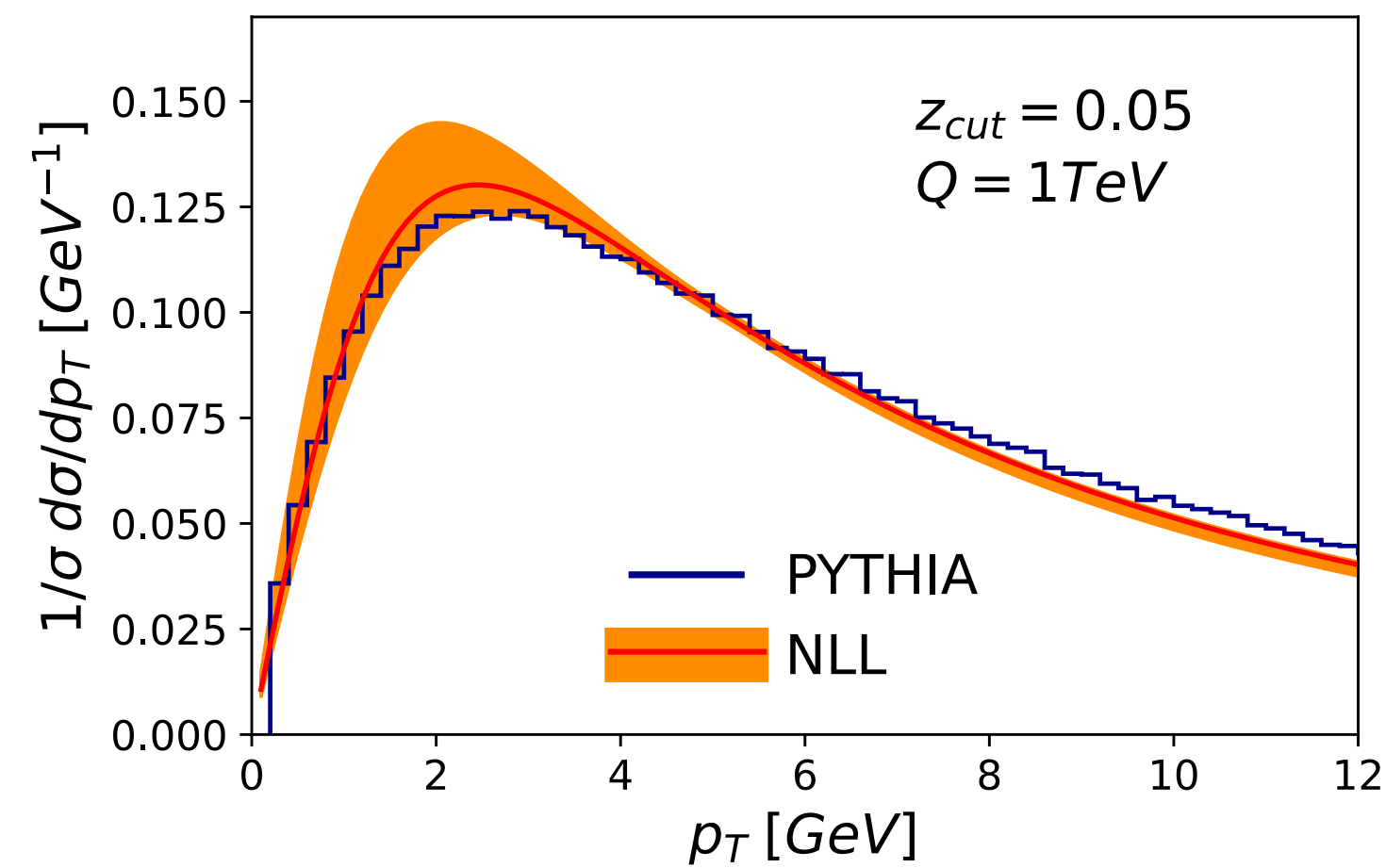
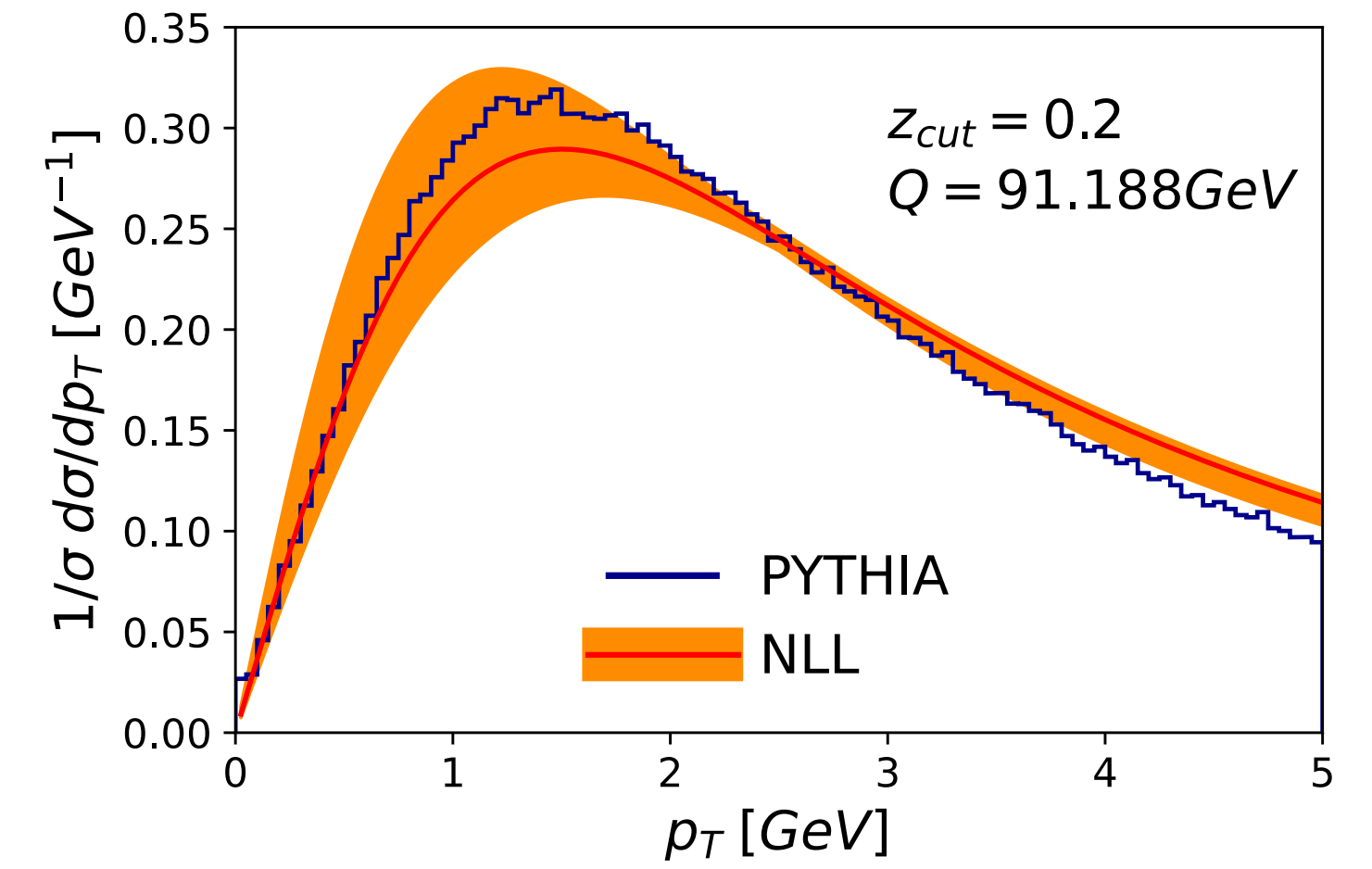
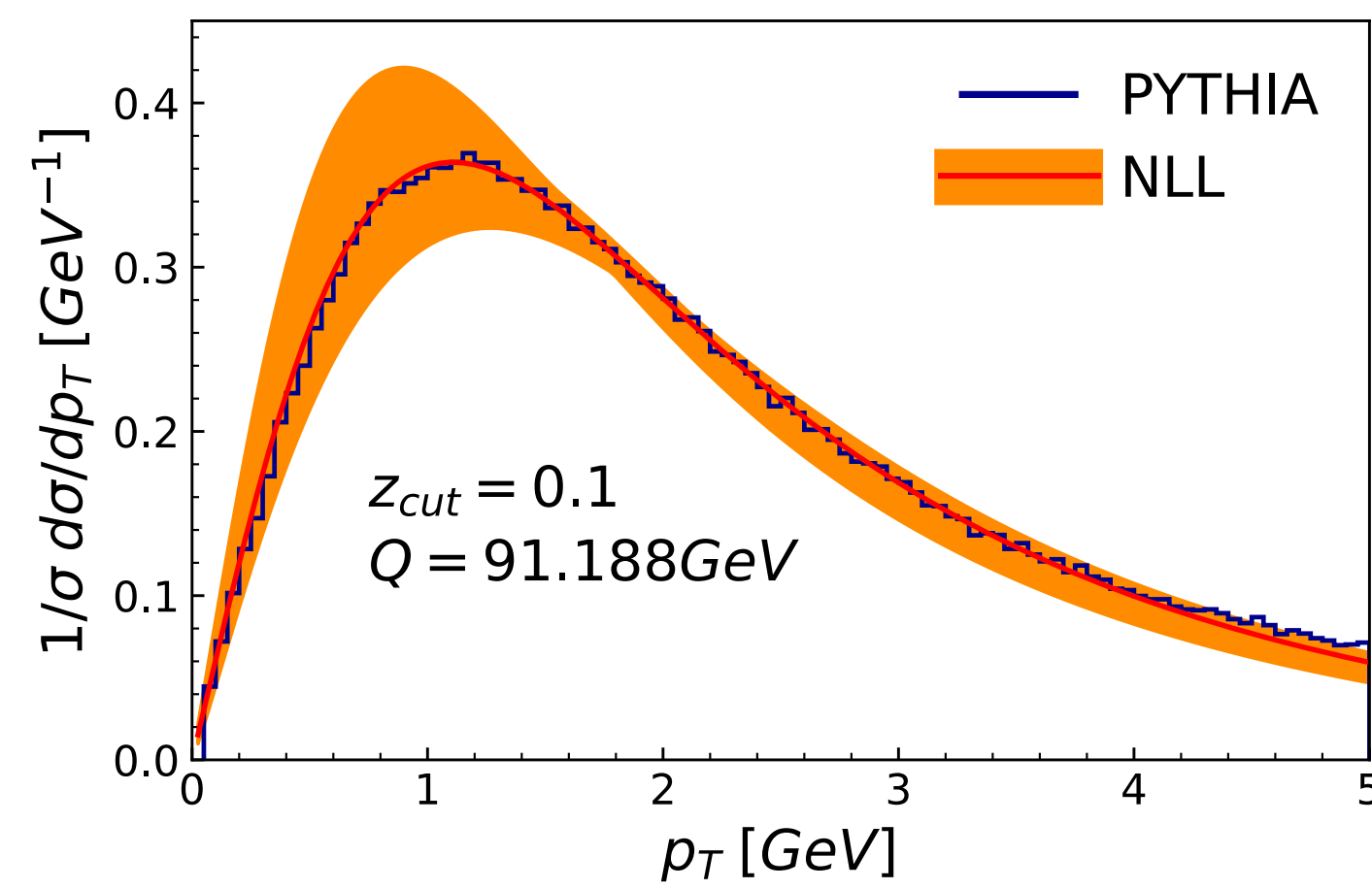
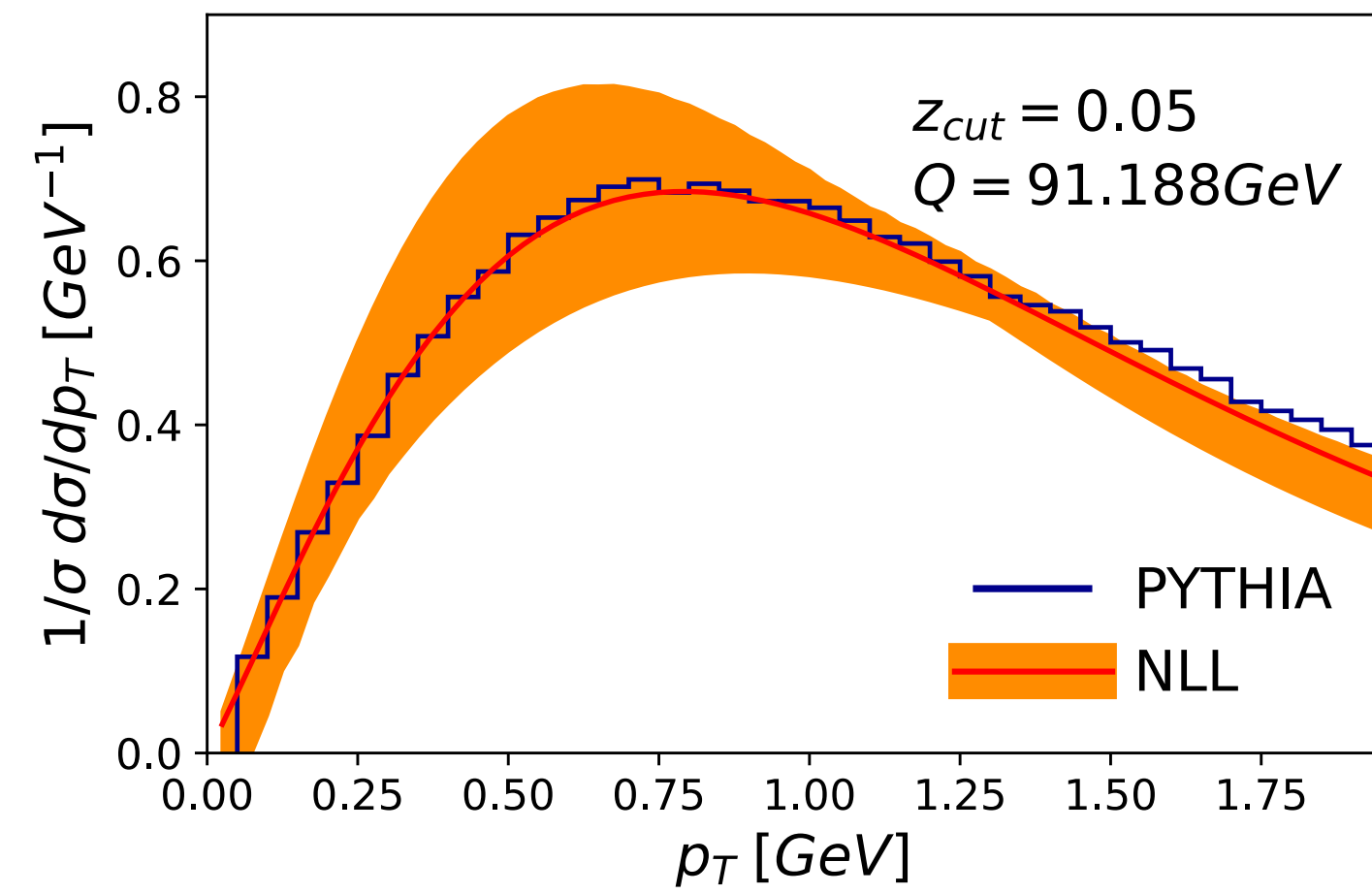
## Groomed jets in back to back limits

$$\frac{1}{\sigma_0} \frac{d\sigma}{d^2\mathbf{p}_\perp} = H(Q^2, \mu) J_c(Q^2, \delta, \mu) J_{\bar{c}}(Q^2, \delta, \mu) J_{cs}(Q^2, \delta, z_{\text{cut}}, \mu) J_{\bar{c}s}(Q^2, \delta, z_{\text{cut}}, \mu)$$

$$\times \int \frac{d\vec{b}}{(2\pi)^2} e^{i\vec{b} \cdot \vec{p}_\perp} S(Q^2, \vec{b}, \mu, \nu) S_{sc}(Q^2, \vec{b}, z_{\text{cut}}, \mu, \nu) S_{s\bar{c}}(Q^2, \vec{b}, z_{\text{cut}}, \mu, \nu)$$



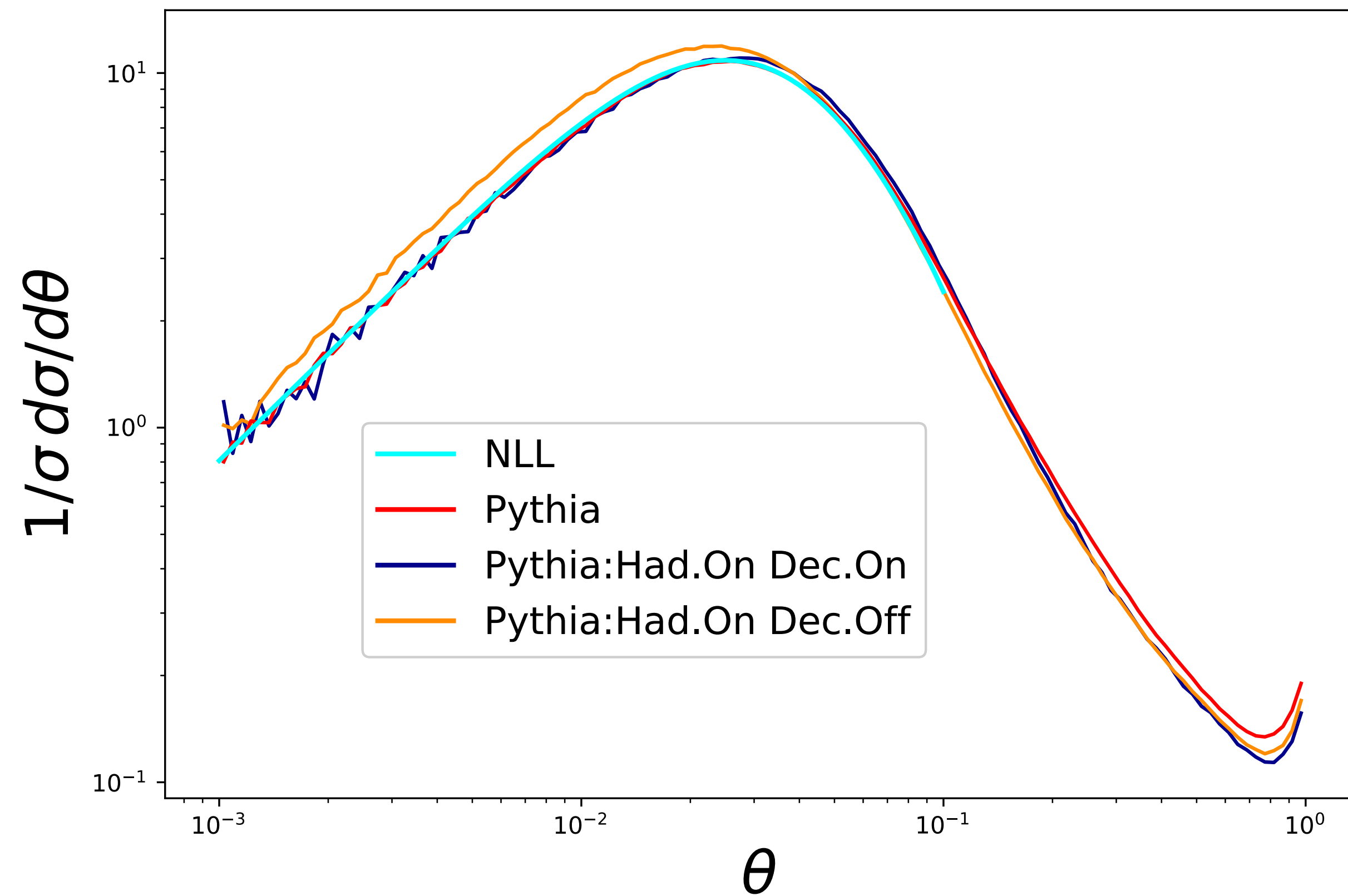
# EEC with jets



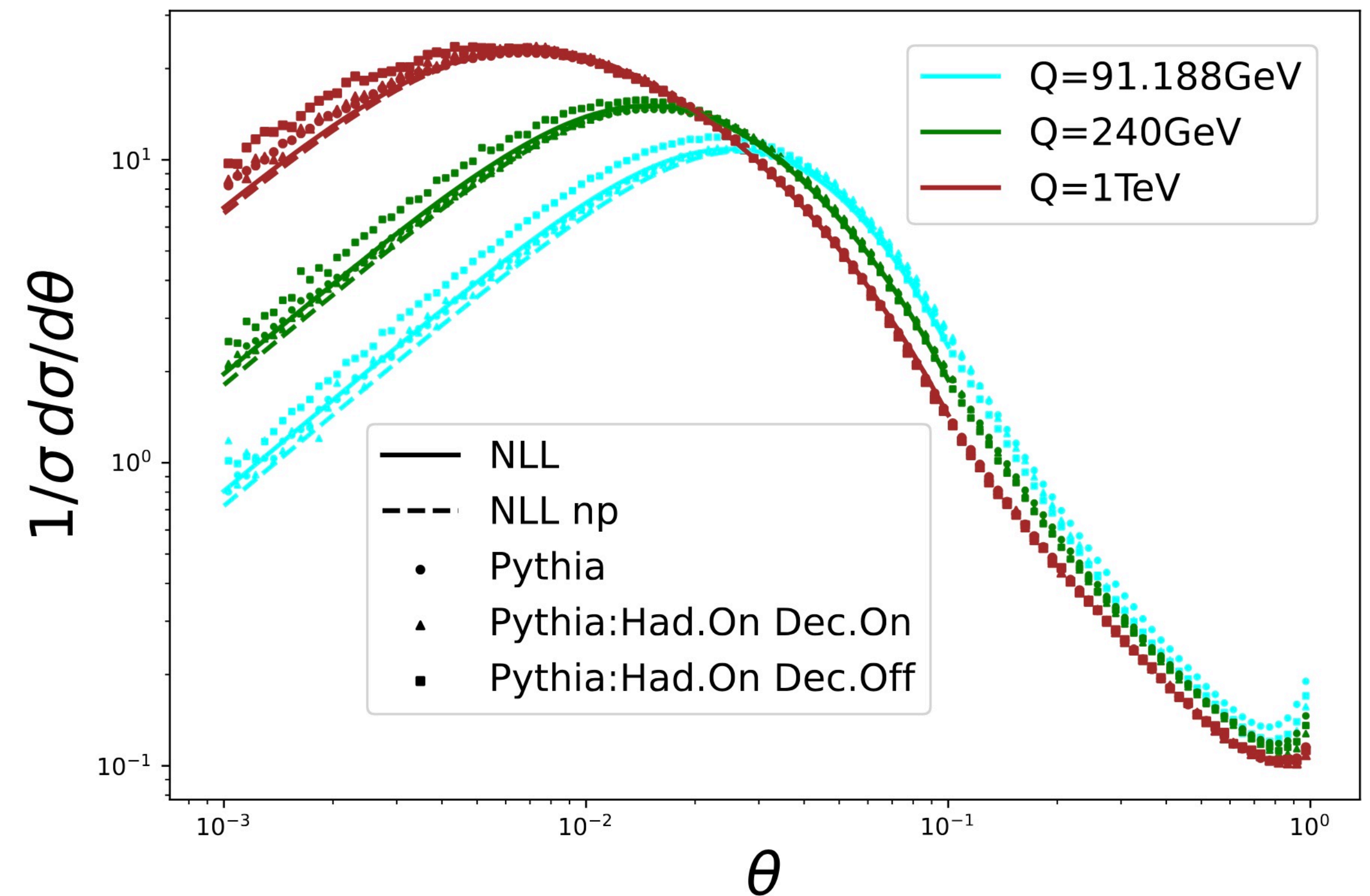


# EEC with jets

EEC with groomed jet



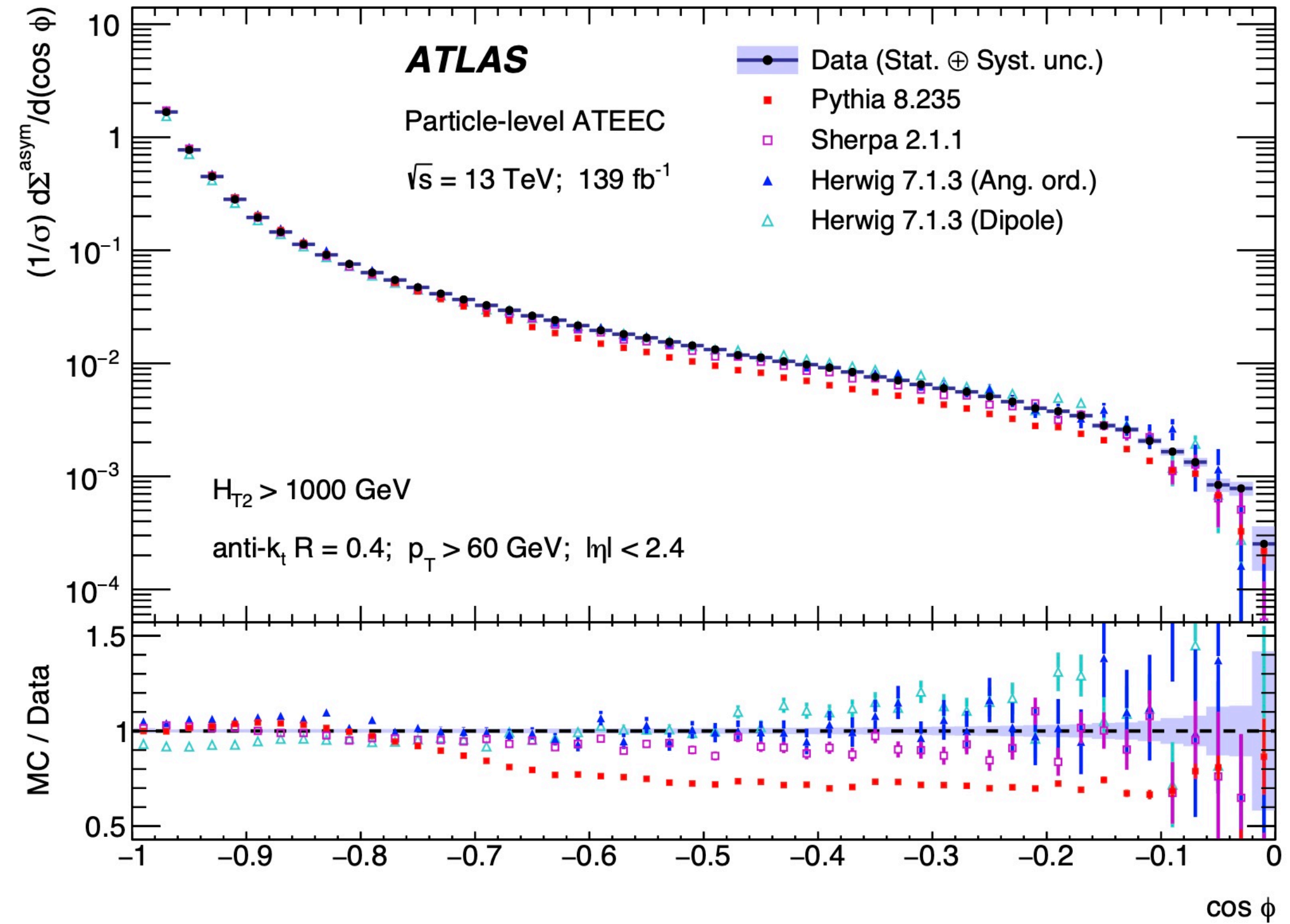
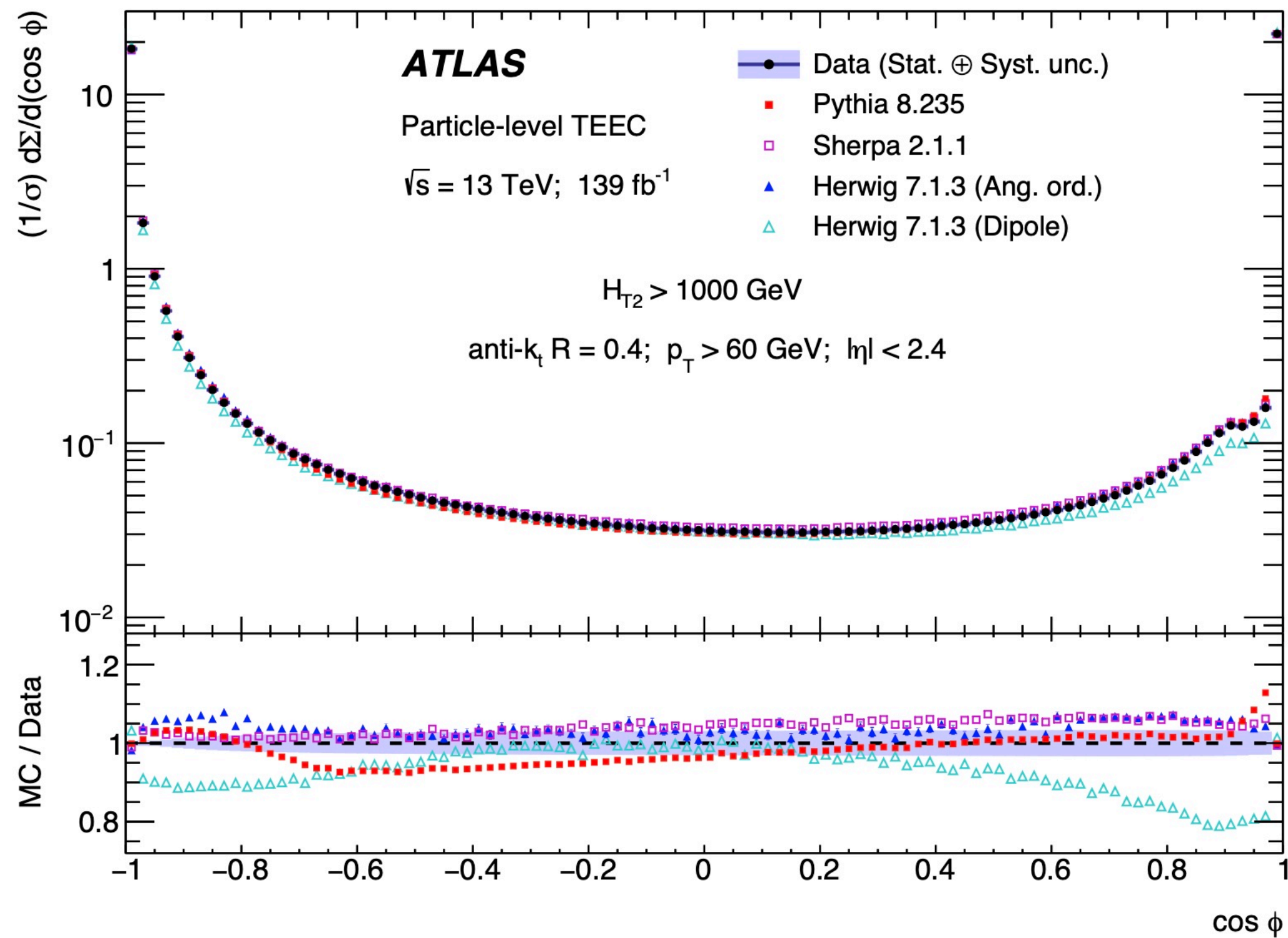
*He, HTL, Shao, in preparation*



# Outline

- Introduction (EEC in back to back limit)
- EEC with jets
- EEC and  $\alpha_s$  measurements
- Summary

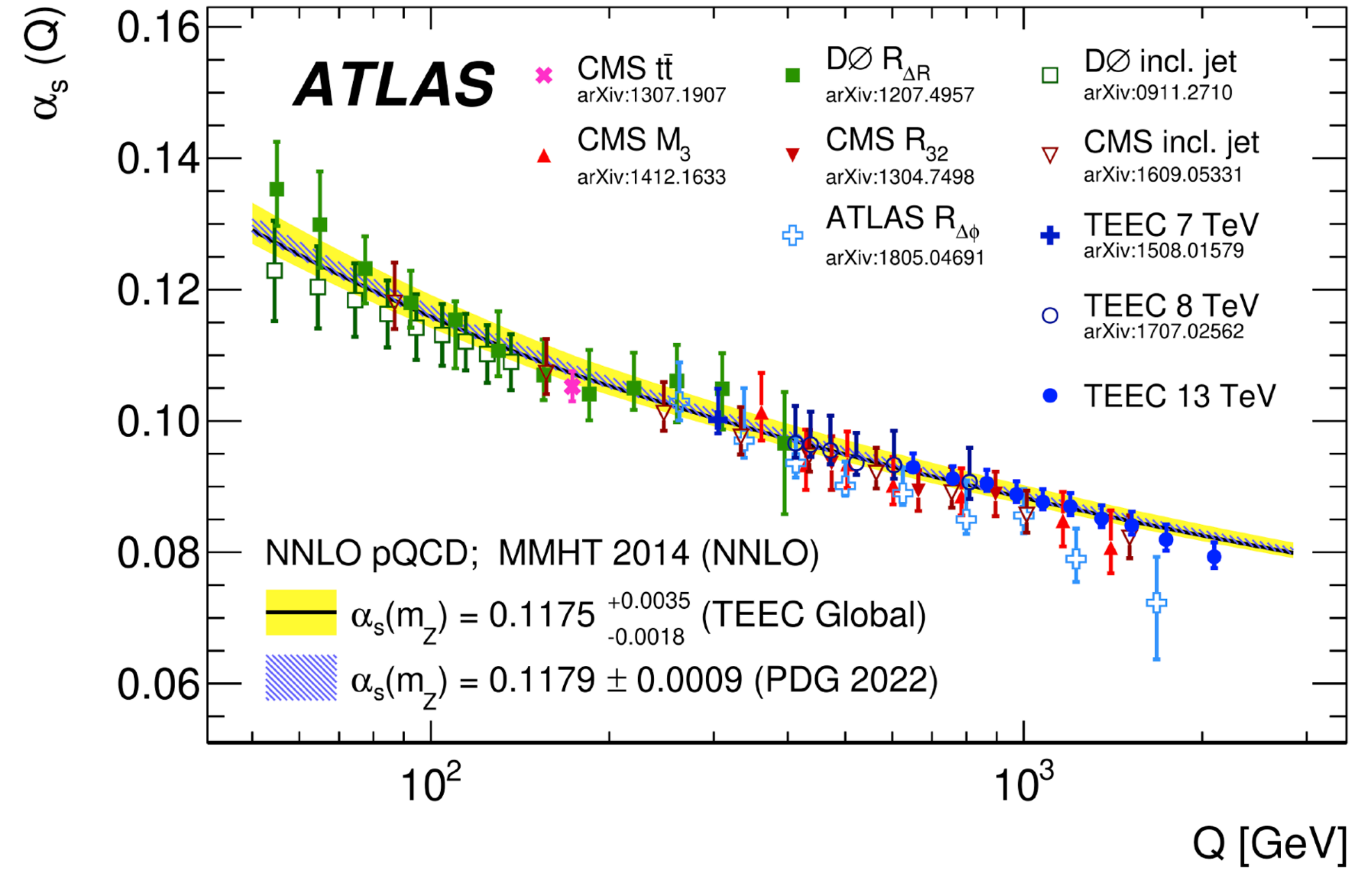
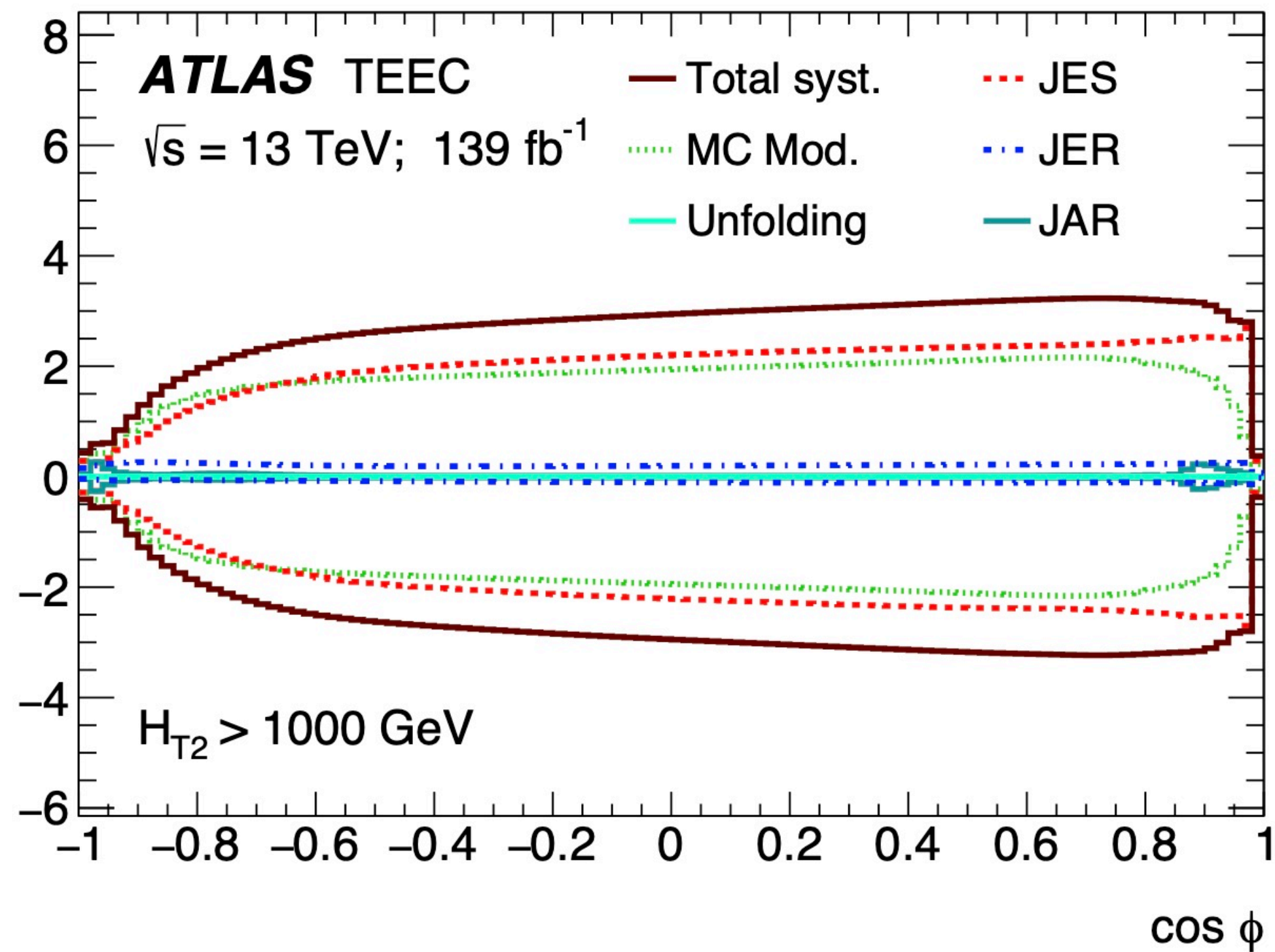
# EEC and $\alpha_s$ measurement





# EEC and $\alpha_s$ measurement

Systematic uncertainty [%]



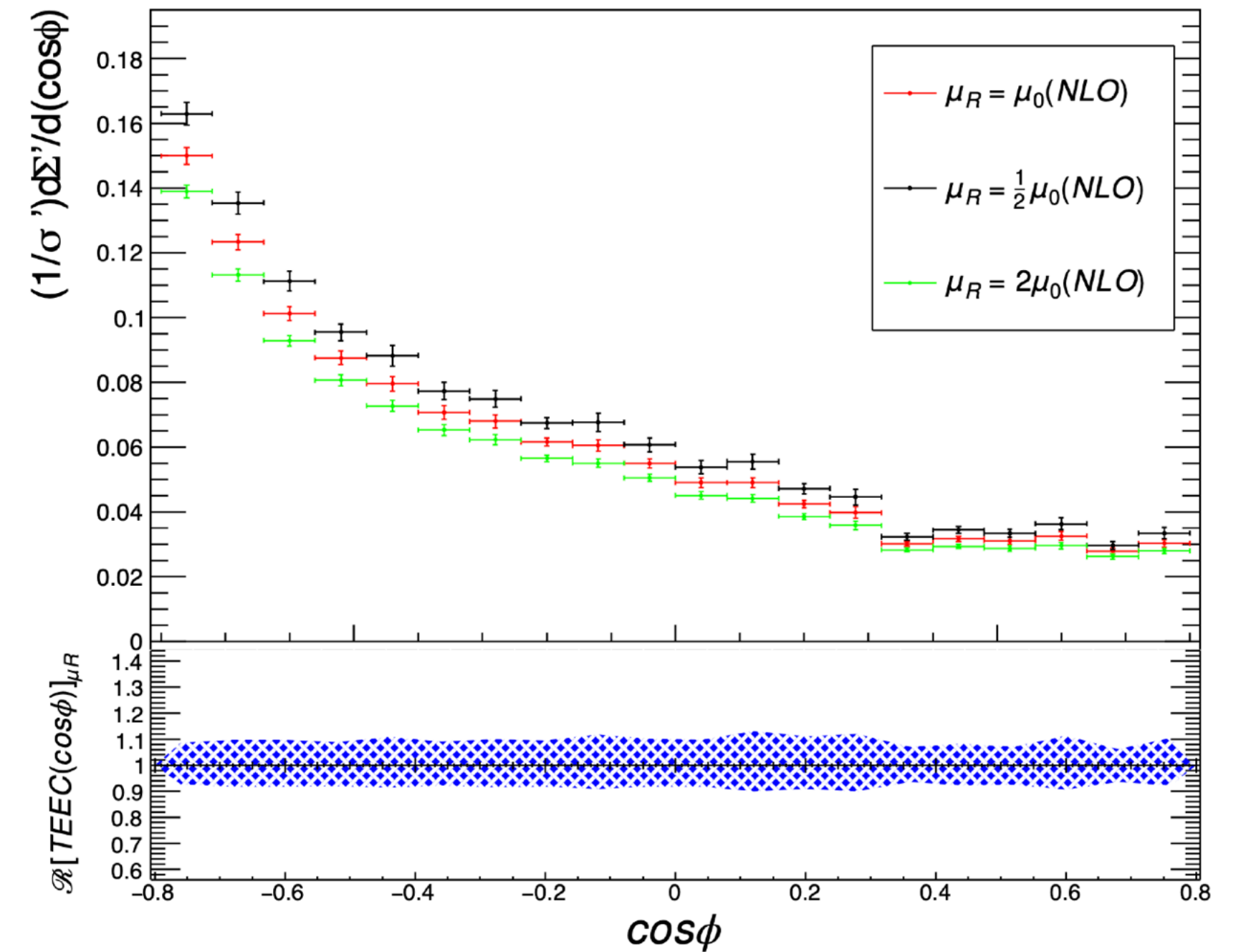
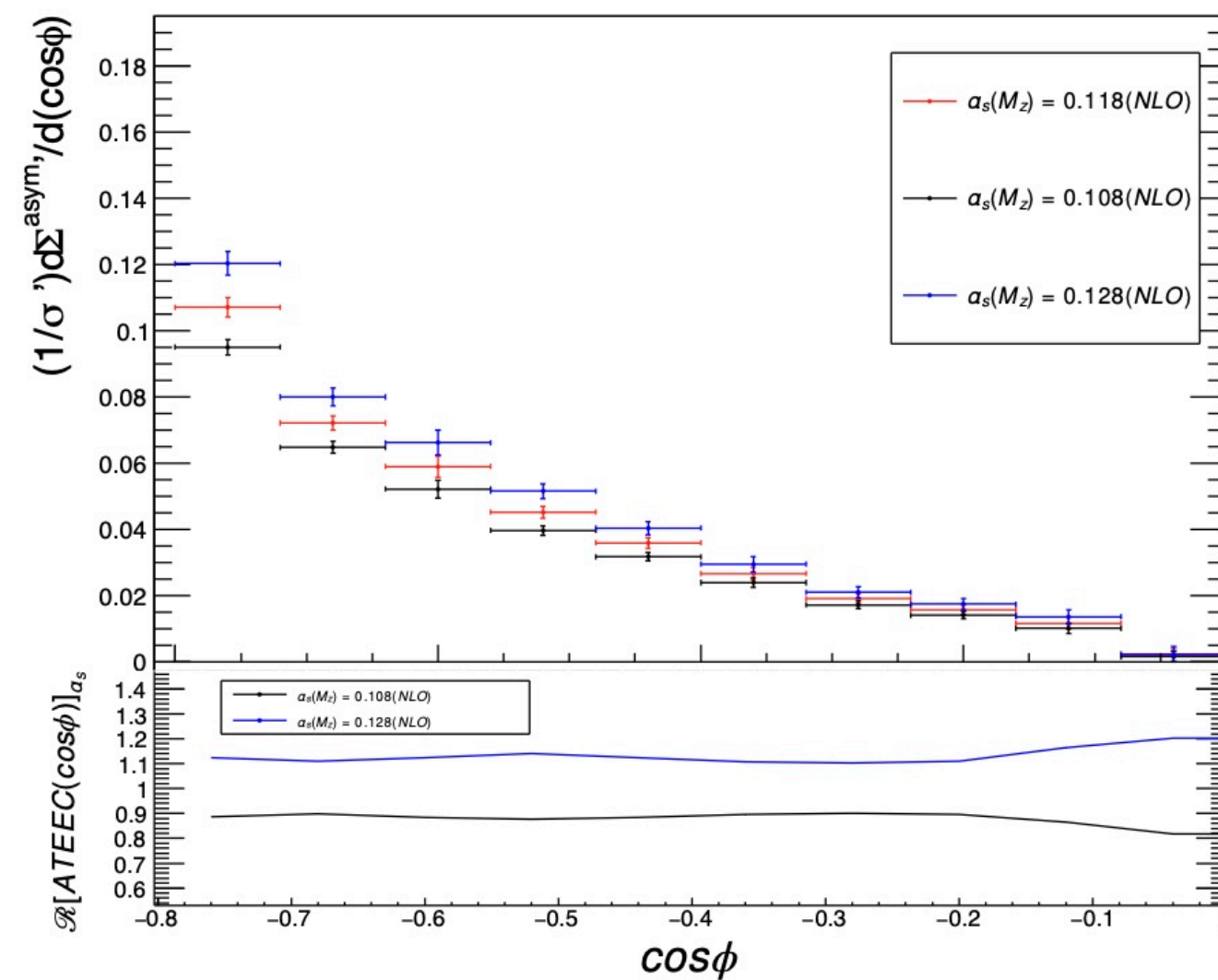
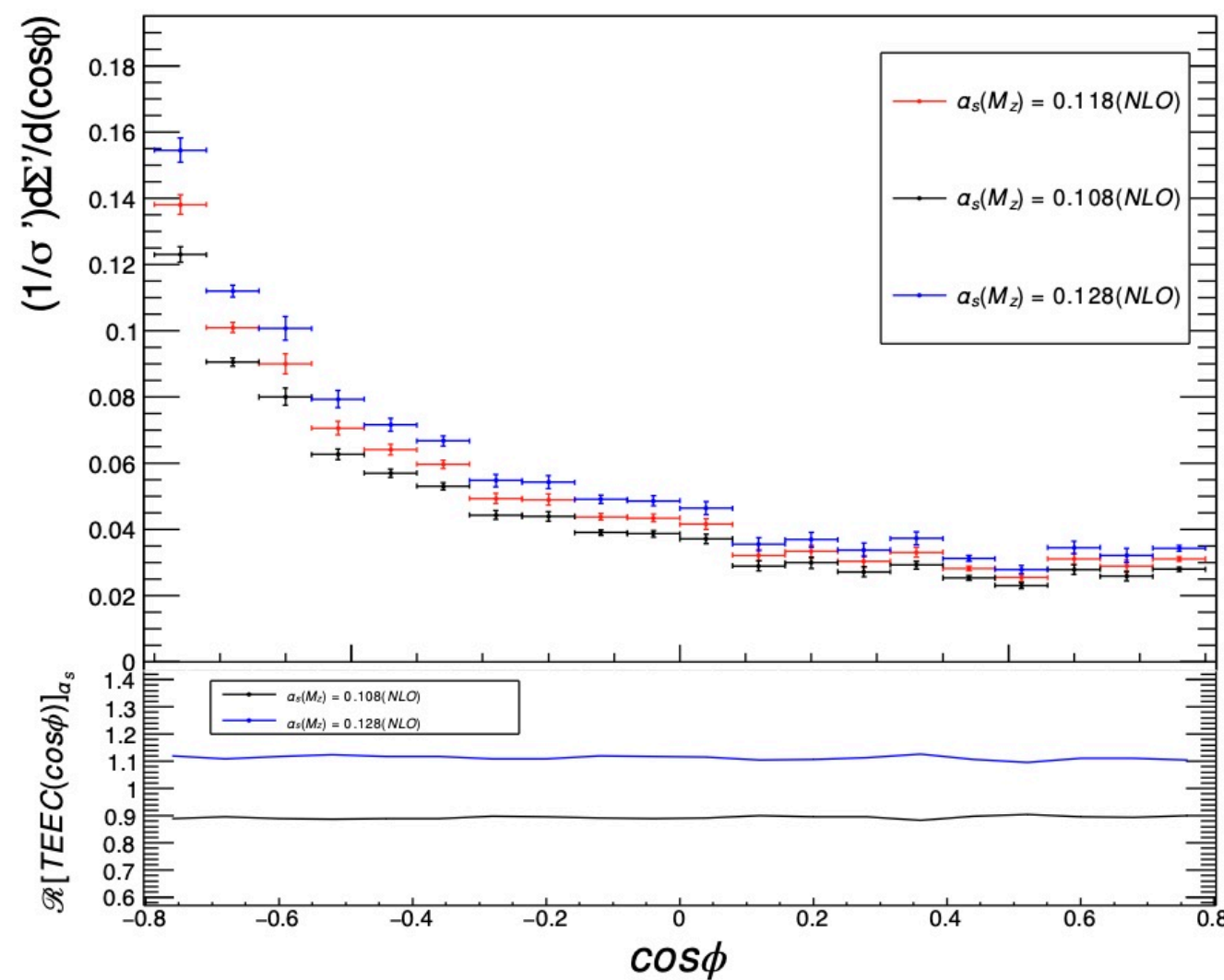
From TEEC  $\alpha_s(m_Z) = 0.1185 \pm 0.0009 \text{ (exp.)}^{+0.0025}_{-0.0012} \text{ (theo.)}$ .

From ATEEC  $\alpha_s(m_Z) = 0.1175 \pm 0.0006 \text{ (exp.)}^{+0.0034}_{-0.0017} \text{ (theo.)}$

# EEC and $\alpha_s$ measurement

$$\frac{1}{\sigma'} \frac{d\Sigma'}{d\cos\phi} \equiv \frac{\sum_{a,b} \int dE_T d\cos\phi_{ab} \frac{d\sigma_{\gamma p \rightarrow a+b+X}}{dE_T d\cos\phi_{ab}} \frac{2E_{T,a}E_{T,b}}{|\sum_i E_{T,i}|^2} \delta(\cos\phi_{ab} - \cos\phi)}{\int dE_T d\sigma_{\gamma p \rightarrow a+b+X}/dE_T}$$

$$\frac{1}{\sigma'} \frac{d\Sigma'^{asym}}{d\cos\phi} \equiv \frac{1}{\sigma'} \frac{d\Sigma'}{d\cos\phi} \Big|_{\phi} - \frac{1}{\sigma'} \frac{d\Sigma'}{d\cos\phi} \Big|_{\pi-\phi}.$$



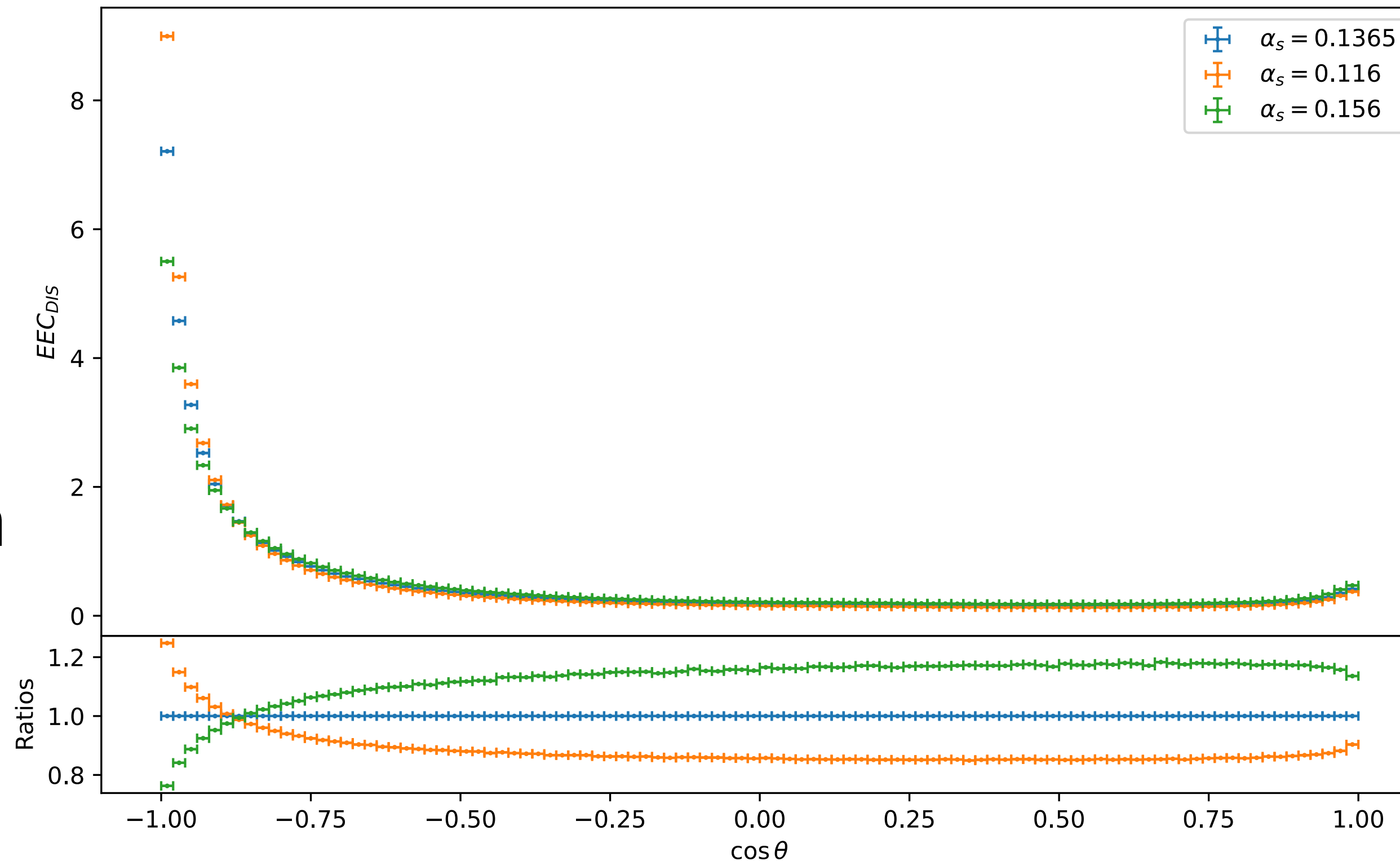
scale uncertainties

Ali, Li, Wang, Xing, 2020



# EEC and $\alpha_s$ measurement

Back to Back  
Resummation region



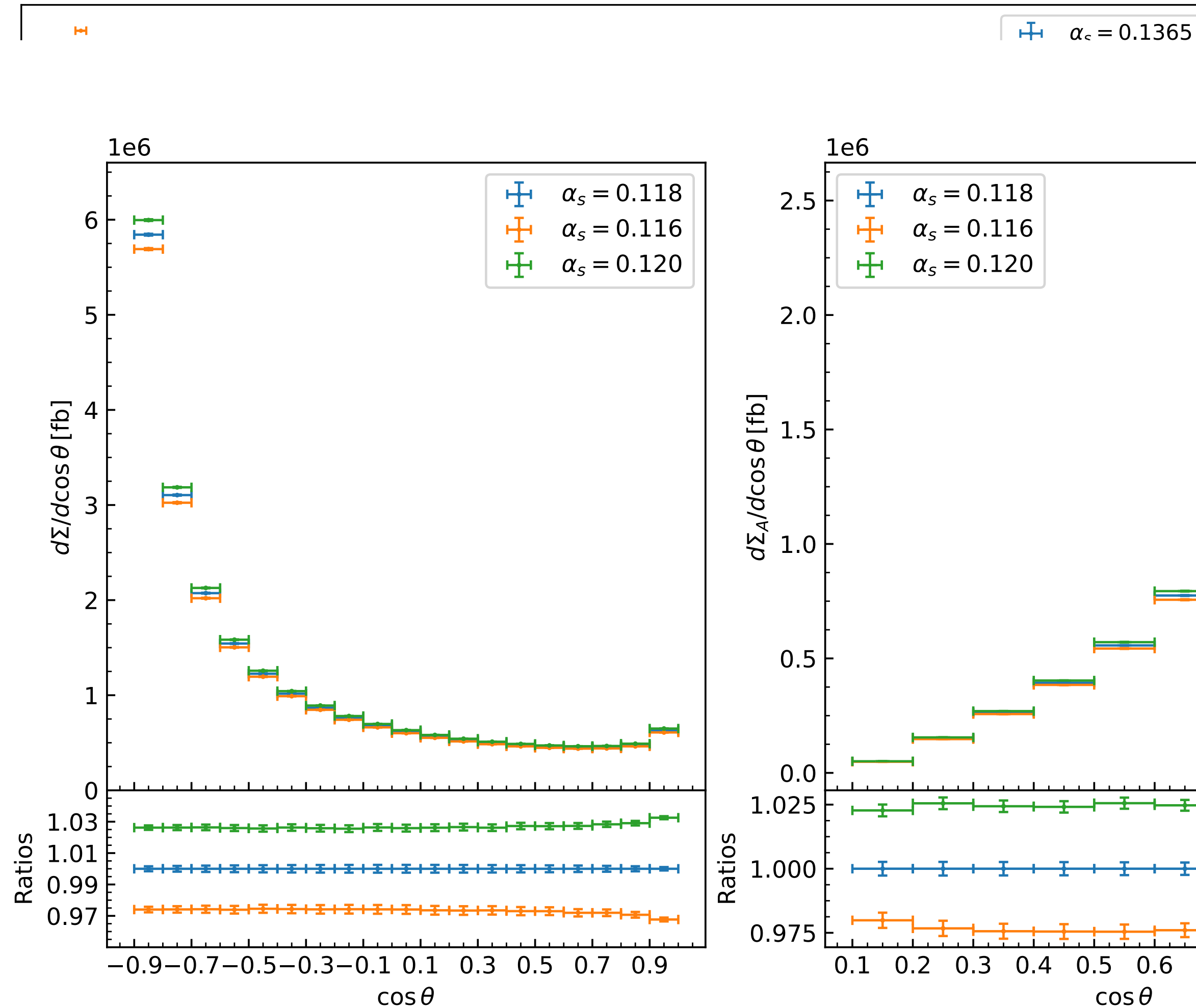
Fixed order  
dominated region

From Pythia Simulation, vary the default  $\alpha_s$  by  $\pm 15\%$



# EEC and $\alpha_s$ measurement

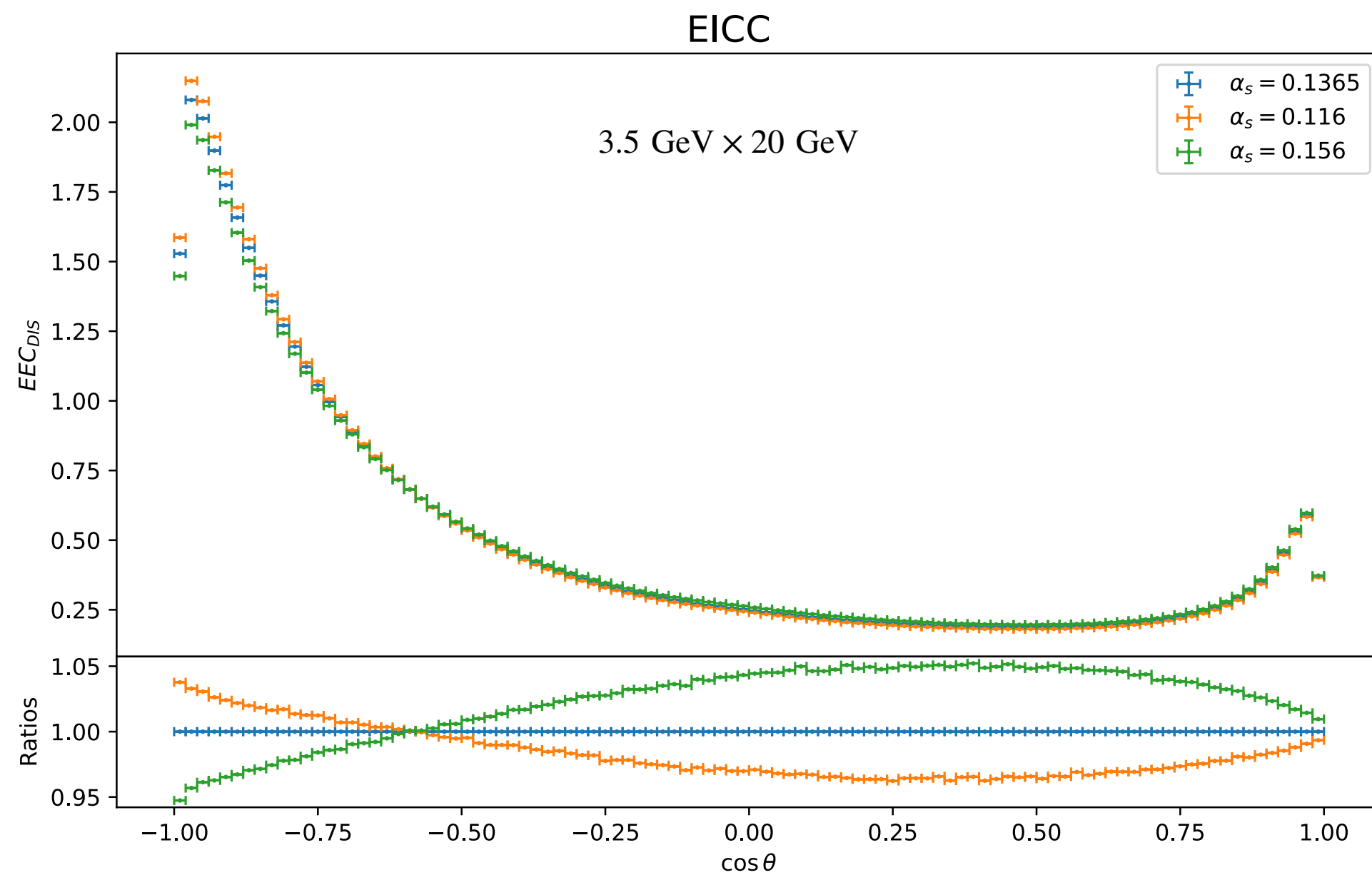
Back to Back  
Resummation regio



order  
1 region

Chen, He, HTL, in preparati

# Summary



- Showed EEC/TEEC in the pp and ep collisions as event shape
- Discussed EEC with jets and groomed jets which shows interesting perturbative and non perturbative properties
- Presented EEC distributions and  $\alpha_s$  measurements

# Thank You!