

Prospects of τ EDM Measurements at STCF

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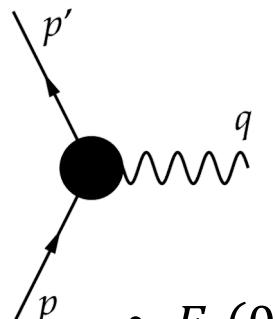
2025.07.03 @ 2025年超级陶粲装置研讨会

Xulei Sun, Y. Wu, Xiaorong Zhou; arXiv: 2411.19469



Introduction

- Electric Dipole Moment
 - Current Strongest CPV test
 - **Matter-antimatter asymmetry**
- Electromagnetic form factors


$$= eQ_f \bar{u}(p') \left[\gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2m_\ell} F_2(q^2) + \frac{\sigma^{\mu\nu}q_\nu\gamma_5}{2m_\ell} F_3(q^2) + \left(\gamma^\mu - \frac{2m_\ell q^\mu}{q^2} \right) \gamma^5 F_4(q^2) \right] u(p)$$

- $F_1(0) = 1, \quad F_2(0) = a_f = \left(\frac{g-2}{2}\right)_f, \quad F_3(0) = \frac{2m_f}{eQ_f} d_f, \quad F_4(0)$

Total Charge

Anomalous Magnetic Moment

EDM

Anapole Moment

- BSM Contributions
 - Lepton Universality
 - Flavor Structure

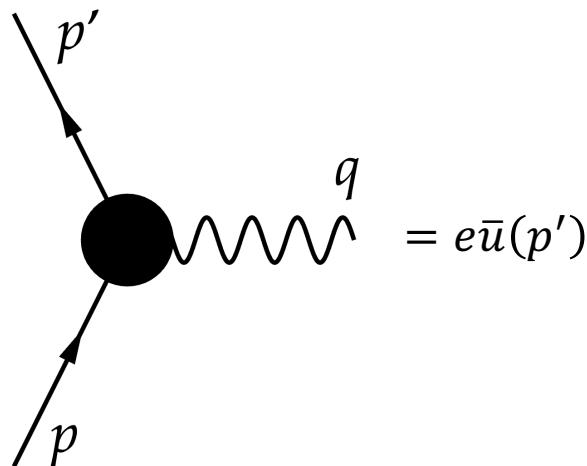
$$\frac{\delta d_\ell}{\delta d_e} \sim \frac{m_\ell}{m_e}$$
$$\delta a_\ell \sim \frac{m_\ell^2}{\Lambda^2}$$

$$\frac{m_\tau}{m_e} \sim 3460$$
$$\left(\frac{m_\tau}{m_\mu}\right)^2 \sim 280$$

Mod.Phys.Lett.A22(2007)159
Phys.Lett.B255(1991)611
Phys.Lett.B395(1997)369
Rev.Mod.Phys.87(2015)531

Introduction

- Electromagnetic form factors


$$= e\bar{u}(p') \left[\gamma^\mu F_1(q^2) + \frac{i\sigma^{\mu\nu}q_\nu}{2m_\ell} F_2(q^2) + \frac{\sigma^{\mu\nu}q_\nu\gamma_5}{2m_\ell} F_3(q^2) + \left(\gamma^\mu - \frac{2m_\ell q^\mu}{q^2} \right) \gamma^5 F_4(q^2) \right] u(p)$$

- $F_2(q^2), F_3(q^2)$: Measured at corresponding energy scale
 - Can develop imaginary part

$$\frac{e}{2m_\ell} \Re(F_3(q^2)) \rightarrow \Re d_\ell \quad \frac{e}{2m_\ell} \Im(F_3(q^2)) \rightarrow \Im d_\ell$$

$$\Re(F_3(q^2)) = \Re(F_3^{SM}(q^2)) + \Re(F_3^{BSM}(q^2)) \approx \Re(F_3^{BSM}(q^2))$$

Current Limits for d_τ

• Current Results

Citation: S. Navas *et al.* (Particle Data Group), Phys. Rev. D **110**, 030001 (2024) and 2025 update

τ ELECTRIC DIPOLE MOMENT (d_τ)

A nonzero value is forbidden by both T invariance and P invariance.

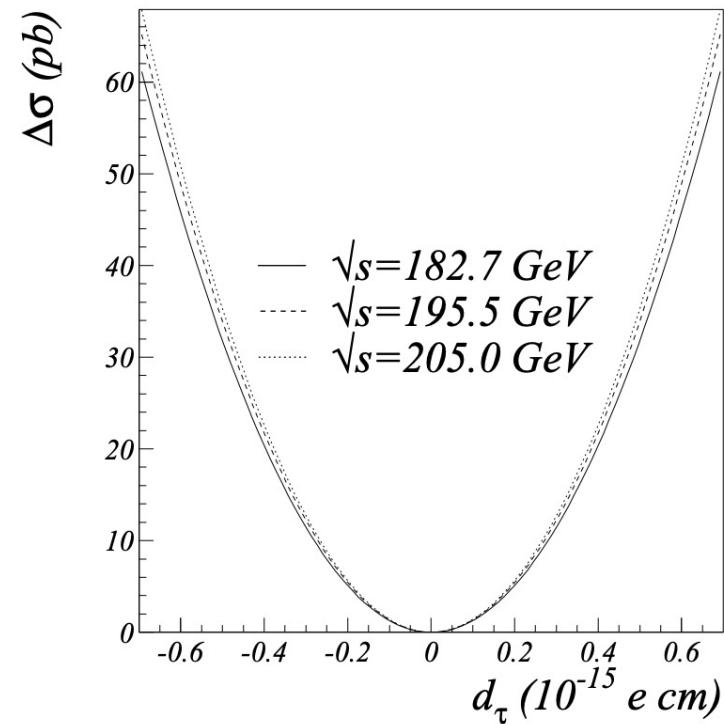
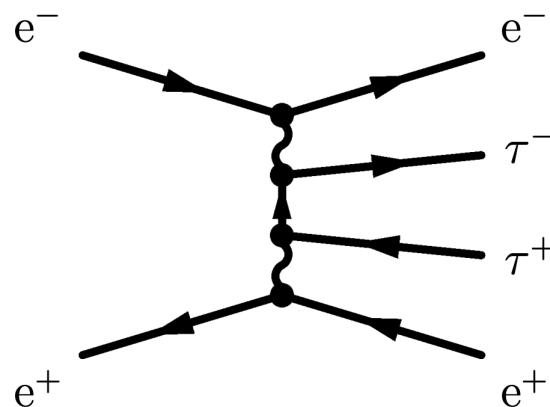
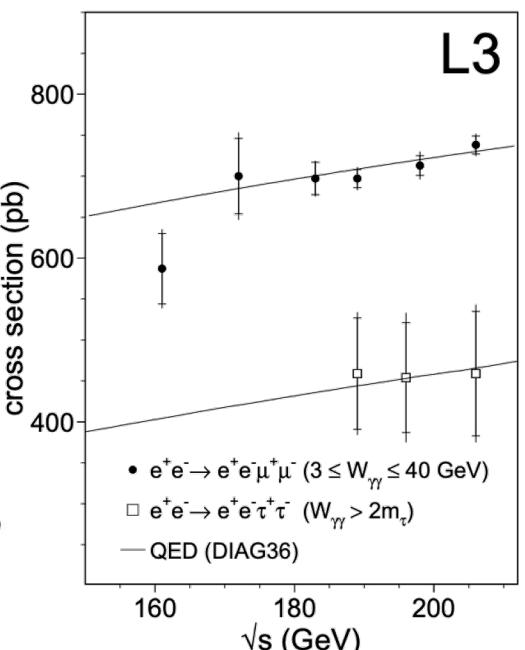
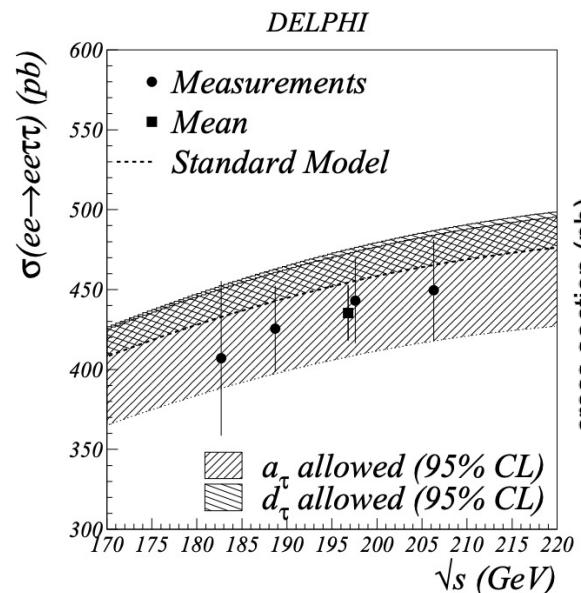
The q^2 dependence is expected to be small providing no thresholds are nearby.

$\text{Re}(d_\tau)$						
VALUE (10^{-16} ecm)	CL%	DOCUMENT ID	TECN	COMMENT		
- 0.185 to 0.061	95	1 INAMI	22	BELL	$E_{\text{cm}}^{\text{ee}} = 10.6 \text{ GeV}$	
• • • We do not use the following data for averages, fits, limits, etc. • • •						
< 2.3	90	2 GROZIN	09A	RVUE	From e EDM limit	
< 3.7	95	3 ABDALLAH	04K	DLPH	$e^+ e^- \rightarrow e^+ e^- \tau^+ \tau^-$ at LEP2	
< 11.4	95	4 ACHARD	04G	L3	$e^+ e^- \rightarrow e^+ e^- \tau^+ \tau^-$ at LEP2	
- 0.22 to 0.45	95	5 INAMI	03	BELL	$E_{\text{cm}}^{\text{ee}} = 10.6 \text{ GeV}$	
< 4.6	95	6 ALBRECHT	00	ARG	$E_{\text{cm}}^{\text{ee}} = 10.4 \text{ GeV}$	
> -3.1 and < 3.1	95	ACCIARRI	98E	L3	1991–1995 LEP runs	
> -3.8 and < 3.6	95	ACKERSTAFF	98N	OPAL	1990–1995 LEP runs	
< 0.11	95	8,9 ESCRIBANO	97	RVUE	$Z \rightarrow \tau^+ \tau^-$ at LEP	
< 0.5	95	10 ESCRIBANO	93	RVUE	$Z \rightarrow \tau^+ \tau^-$ at LEP	
< 7	90	GRIFOLS	91	RVUE	$Z \rightarrow \tau \tau \gamma$ at LEP	
< 1.6	90	DELAGUILA	90	RVUE	$e^+ e^- \rightarrow \tau^+ \tau^-$ $E_{\text{cm}}^{\text{ee}} = 35 \text{ GeV}$	

Current Limits for d_τ

- Direct Measurement
 - Production cross section
 - LEP
 - $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$

Phys.Rev.D 53 (1996) 1181
Eur.Phys.J.C 35 (2004) 159 - DELPHI
Phys.Lett.B 585 (2005) 53 – L3

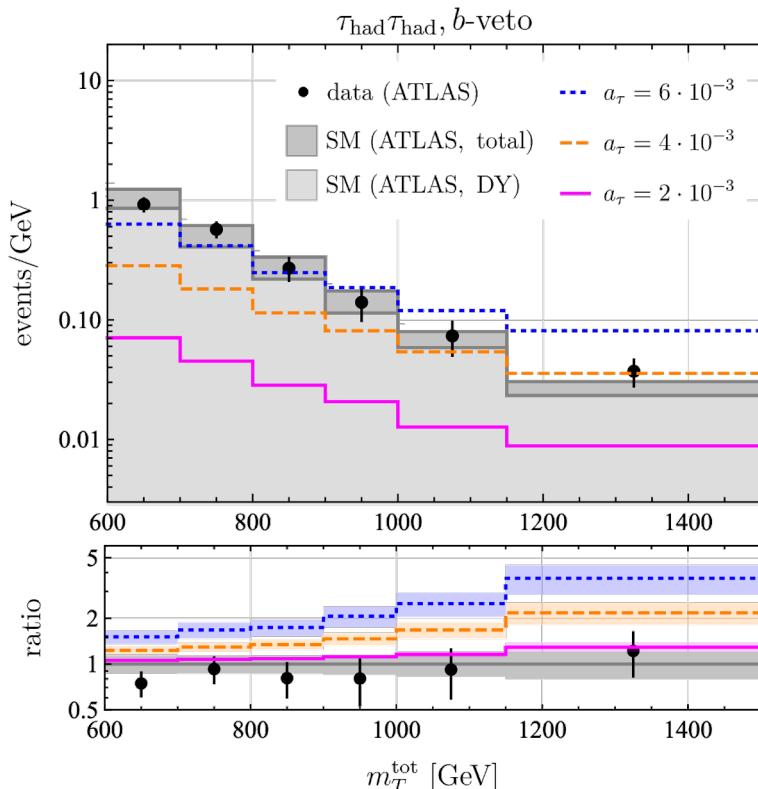


L3 $|d_\tau| \leq 1.14 \times 10^{-15} e \cdot cm$

DELPHI $|d_\tau| \leq 3.7 \times 10^{-16} e \cdot cm$

Current Limits for d_τ

- Direct Measurement
 - Production cross section
 - LHC
 - $q\bar{q} \rightarrow \tau^+\tau^-$
 - Contribution from d_τ/a_τ enhance at high energies



SciPost Phys. 16 (2024) 2, 048

Based on ATLAS measurements: *Phys.Rev.Lett. 125 (2020) 051801*

$$|d_\tau| < 1.0 \times 10^{-17} e \cdot \text{cm}$$

Comparable with Belle measurements

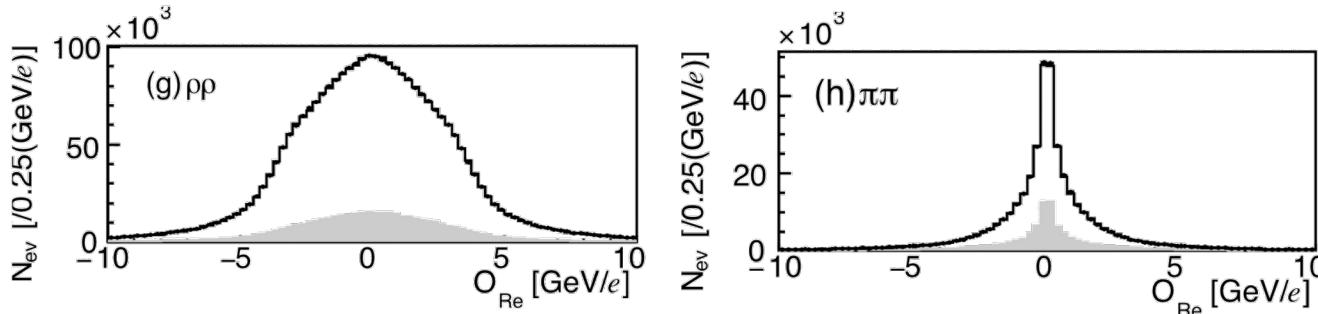
Current Limits for d_τ

- Direct Measurement - Belle
 - Spin correlation at low energy lepton collider
 - $e^+e^- \rightarrow \gamma^* \rightarrow \tau^+\tau^-$
 - 8 different channels:
 - $(ev\bar{v})(\mu\nu\bar{\nu}), (ev\bar{v})(\pi\nu), (\mu\nu\bar{v})(\pi\nu), (ev\bar{v})(\rho\nu), (\mu\nu\bar{v})(\rho\nu), (\pi\nu)(\rho\nu), (\rho\nu)(\rho\bar{\nu}), (\pi\nu)(\pi\bar{\nu})$
 - Matrix Element

JHEP 04 (2022) 110
Phys.Lett.B 551 (2003) 16

$$\chi_{\text{prod}} = \chi_{\text{SM}} + \text{Re}(d_\tau)\chi_{\text{Re}} + \text{Im}(d_\tau)\chi_{\text{Im}} + |d_\tau|^2\chi_{d^2},$$

- Optimal Observable



$$\Re(d_\tau) \in [-1.85, 0.61] \times 10^{-17} e \cdot \text{cm}$$

$$\Im(d_\tau) \in [-1.03, 0.23] \times 10^{-17} e \cdot \text{cm}$$

Current Limits for d_τ

- Indirect Constraints from d_e
 - d_τ contributes to d_e through 3-loop diagram

A. Grozin, I. Khriplovich, A. Rudenko; Nucl.Phys.B 821 (2009) 285; Phys.Atom.Nucl. 72 (2009) 1203

$$d_e \approx \left(\left(\frac{15}{4} \zeta(3) - \frac{31}{12} \right) \frac{m_e}{m_\tau} \left(\frac{\alpha}{\pi} \right)^3 \right) d_\tau \approx 6.9 \times 10^{-12} d_\tau$$

- Current Measurement (Prospects) of d_e

$$|d_e| < 4.1 \times 10^{-30} e \cdot cm \text{ @ 90% CL}$$

Roussy et al. Science 381 (2023) 6653, 46

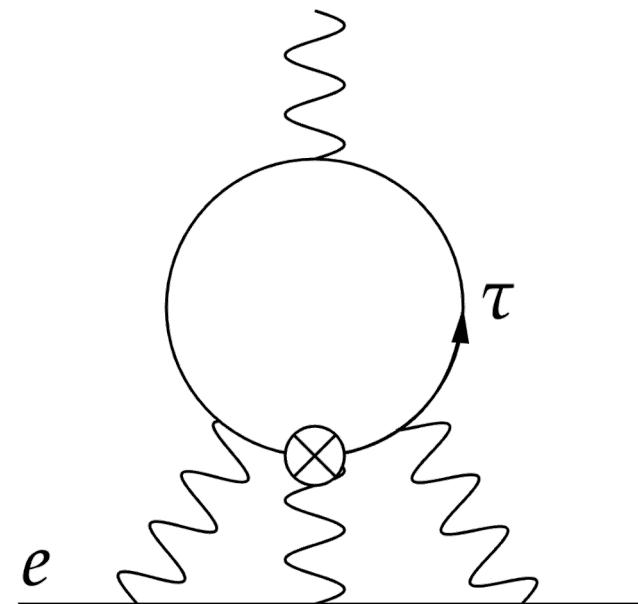
$$|d_e| \lesssim 4 \times 10^{-31} e \cdot cm$$

Next generation of ACME



$$|d_\tau| < 5.9 \times 10^{-19} e \cdot cm$$

$$|d_\tau| < 5.80 \times 10^{-20} e \cdot cm$$



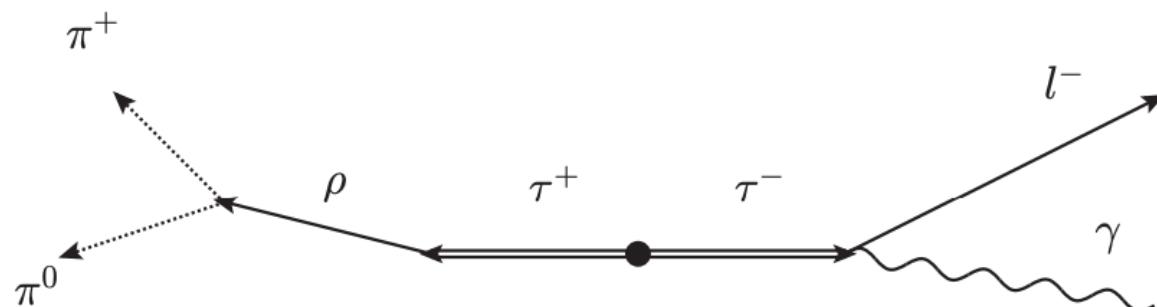
Model Dependent

Current Studies

- Radiative Decay

JHEP 03 (2016) 140

- $\tau^\pm \rightarrow \ell^\pm \nu_\tau \nu_\ell \gamma$



	Prospects				Current Limits	
	Belle (ρ)	Belle II (ρ)	Belle (full)	Belle II (full)	DELPHI [2]	Belle [52]
\tilde{a}_τ	0.16	0.023	0.085	0.012	0.017	—
$(m_\tau/e) \tilde{d}_\tau$	0.15	0.021	0.080	0.011	—	0.0015

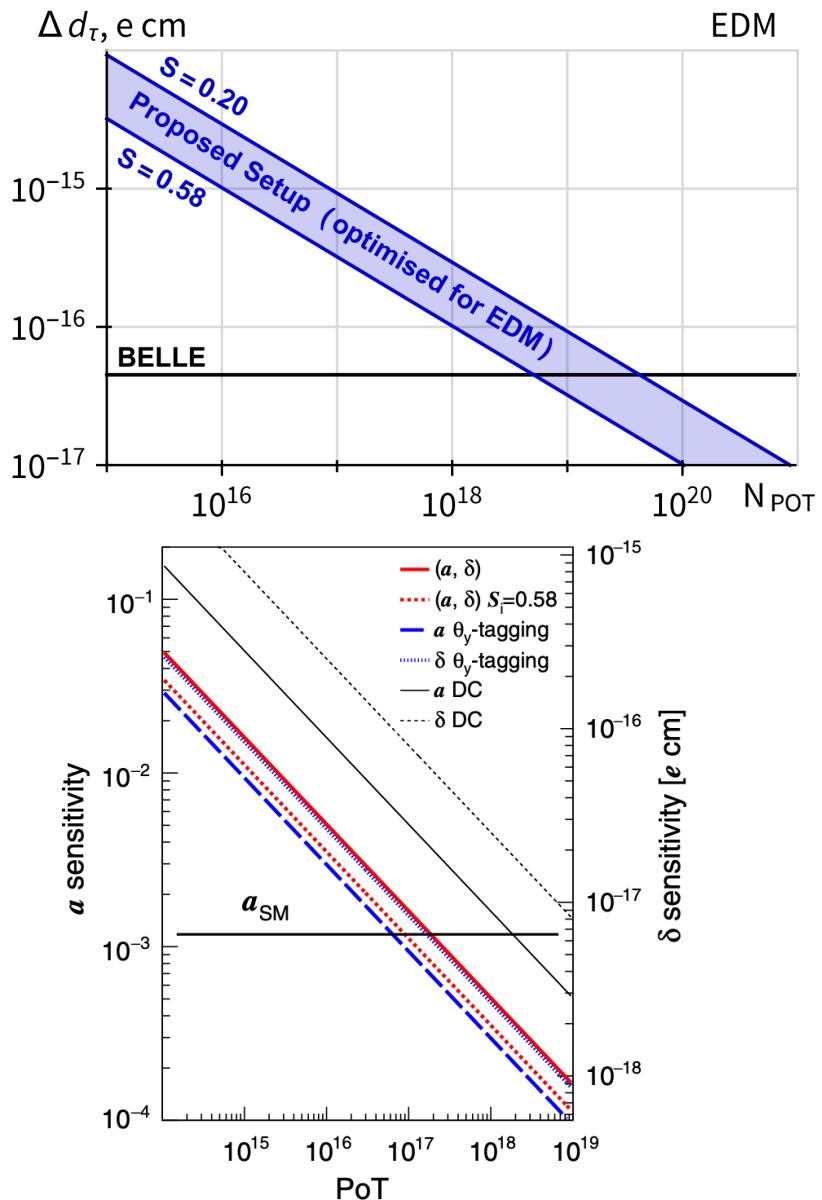
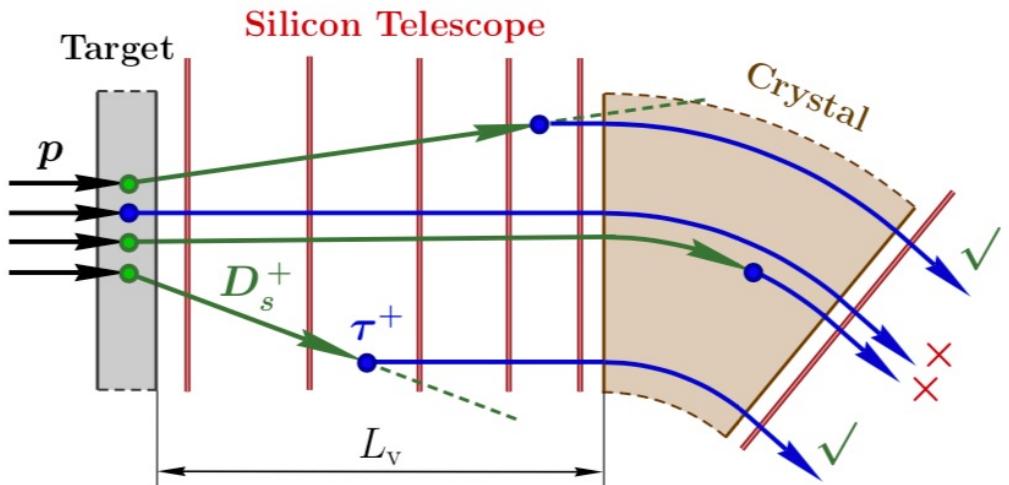
Current Studies

- Bent Crystal

JHEP 03 (2019) 156

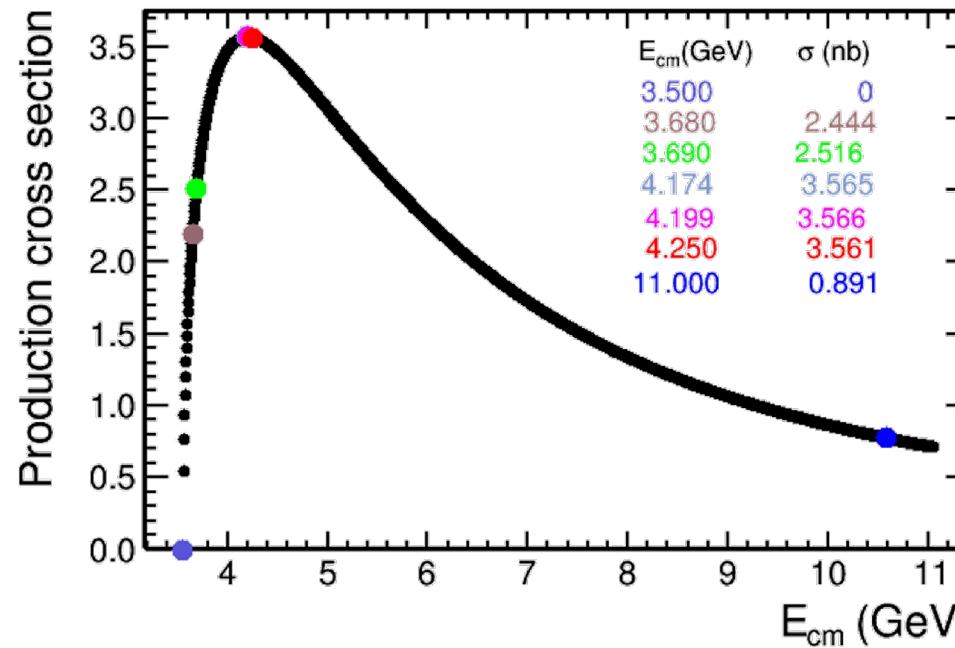
Phys. Rev. Lett 123 (2019) 011801

$$pp \rightarrow D_s^+ X, \quad D_s^+ \rightarrow \tau^+ \nu_\tau, \quad \tau^+ \rightarrow \pi^+ \pi^+ \pi^- \bar{\nu}_\tau.$$



Super Tau-Charm Facility (STCF)

- Large amount of tau-pair $e^+e^- \rightarrow \gamma^* \rightarrow \tau^+\tau^-$
 - 3.5×10^9 tau pair per year @ $\sqrt{s} = 4.26$ GeV
 - 1.7×10^9 @ $\sqrt{s} = 7.0$ GeV



- STCF can significantly improve the sensitivity for tau-physics

τ EDM at STCF

- Cross Section
 - Summing over decay phase space



$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \beta}{4s} \left((2 - \beta^2 \sin^2 \theta) (|F_1|^2 - \gamma^2 |F_2|^2) + 2(1 + \gamma^2) |F_2|^2 + \beta^2 \gamma^2 \sin^2 \theta |F_3|^2 + 4\Re(F_1 F_2^*) \right)$$

- $F_1(0) = 1, \quad F_2(0) = a_f = \left(\frac{g-2}{2}\right)_f, \quad F_3(0) = \frac{2m_f}{eQ_f} d_f$
- $\beta = \sqrt{1 - 4m_\tau^2/s}, \quad \gamma = 1/\sqrt{1 - \beta^2} = \sqrt{s}/(2m_\tau)$
- Weak Sensitivity on d_τ
- Decay phase space matters – Spin Correlations

τ EDM at STCF

- Squared spin density matrix

$$e^+(\vec{p})e^-(-\vec{p}) \rightarrow \tau^+(\vec{k}, \vec{s}_+) \tau^-(-\vec{k}, \vec{s}_-)$$

$$|\mathcal{M}|^2 = |\mathcal{M}_{SM}|^2 + \Re(d_\tau)v|\mathcal{M}_{Re}|^2 + \Im(d_\tau)v|\mathcal{M}_{Im}|^2 + |d_\tau|^2v^2|\mathcal{M}_{d^2}|^2$$

$$|\mathcal{M}_{Re}|^2 = \frac{4e^3}{k_0 v} |\vec{k}| \left[- \left(m_\tau + (k_0 - m_\tau) (\hat{k} \cdot \hat{p})^2 \right) (\vec{s}_+ \times \vec{s}_-) \cdot \hat{k} + k_0 (\hat{k} \cdot \hat{p}) (\vec{s}_+ \times \vec{s}_-) \cdot \hat{p} \right]$$

CP-odd

- Key factor: the momentum \vec{k} and spin \vec{s} of the τ leptons

The spin of τ

- For different τ decay modes:

$$\tau \rightarrow \ell \nu_\ell \nu_\tau \quad \vec{s}_\pm = \frac{4c_\pm - m_\tau^2 - 3m_\ell^2}{3m_\tau^2 c_\pm - 4c_\pm^2 - 2m_\ell^2 m_\tau^2 + 3c_\pm m_\ell^2} \left(\pm m_\tau \vec{p}_\ell - \frac{c_\pm + E_\ell m_\tau}{k_0 + m_\tau} \vec{k} \right) \quad c_\pm = k_0 E_\ell \mp \vec{k} \cdot \vec{p}_\ell$$

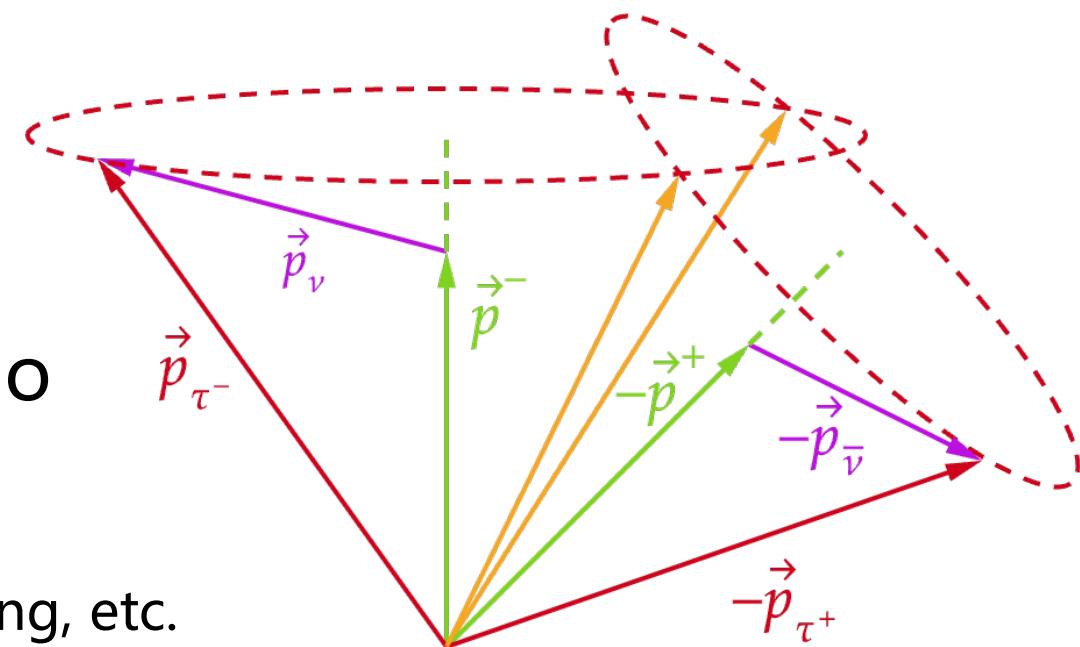
$$\tau \rightarrow \pi \nu_\tau \quad \vec{s}_\pm = \frac{2}{m_\tau^2 - m_\pi^2} \left(\pm m_\tau \vec{p}_\pi + \frac{m_\tau^2 + m_\pi^2 + 2m_\tau E_\pi}{2(E_\tau + m_\tau)} \vec{k} \right)$$

$$\tau \rightarrow \pi \pi^0 \nu_\tau \quad \vec{s}_\pm = \mp \frac{1}{\vec{k}_\pm \cdot \vec{H}_\pm - m_\tau^2 (p_{\pi^\pm} - p_{\pi^0})^2} \left(\mp H_0^\pm \vec{k} + m_\tau \vec{H}^\pm + \frac{\vec{k}(\vec{k} \cdot \vec{H}^\pm)}{E_\tau + m_\tau} \right)$$
$$H_\pm^\mu = 2 \left(\vec{k}_\pm \cdot (p_{\pi^\pm} - p_{\pi^0}) \right) (p_{\pi^\pm} - p_{\pi^0})^\mu + (p_{\pi^\pm} - p_{\pi^0})^2 (p_{\pi^\pm} + p_{\pi^0})^\mu$$

- Reconstruction of the τ momentum
 - Decay products – Neutrinos
 - Two-fold ambiguity (Even worse for leptonic decay)

τ Momentum

- Reconstruction (semi-leptonic)
 - The visible part (π) can be measured (and well paired)
 - \vec{p}^- and \vec{p}^+
 - $2 \rightarrow 2$ system with known E_{cm} and m_τ
 - $|\vec{p}_\nu|$ and $|\vec{p}_{\bar{\nu}}|$
 - $\vec{p}_\nu \cdot \vec{p}^-$ and $\vec{p}_{\bar{\nu}} \cdot \vec{p}^+$
 - \vec{p}_{τ^-} and \vec{p}_{τ^+} are back-to-back
 - Two solutions (orange arrows)
 - One solution/No solution
- More information about neutrino
 - Decay vertex detection
 - Global fitting
 - Remove ambiguity, improve the pairing, etc.



Xin Chen, Y. Wu; JHEP10(2019)089, 1803.00501

Xin Chen, Y. Wu; Eur.Phys.J. C77(2017)697, 1703.04855

Xin Chen, Y. Wu; Phys.Lett.B790(2019)332, 1708.02882

τ EDM at STCF

- Optimal Observable (OO)

$$|\mathcal{M}|^2 = |\mathcal{M}_{SM}|^2 + \Re(d_\tau)v|\mathcal{M}_{Re}|^2$$

$$\mathcal{O}_{Re} = \frac{|\mathcal{M}_{Re}|^2}{|\mathcal{M}_{SM}|^2}$$

$$\langle \mathcal{O}_{Re} \rangle = \int \mathcal{O}_{Re} |\mathcal{M}|^2 d\Pi = a_{Re} \cdot \Re(d_\tau) + b_{Re}$$

- Procedure:
 - Reconstruct the final state momentum
 - τ momentum and spin
 - Matrix elements – Optimal observable - EDM

Event Reconstruction

- Signal:

$$e^+ e^- \rightarrow \tau^+ \tau^- (\tau^\pm \rightarrow \pi^\pm \pi^0 \nu_\tau, \pi^0 \rightarrow \gamma\gamma)$$

$$(B(\tau \rightarrow \pi\pi^0\nu) \approx 25\%)^2 \approx 6.2\%$$

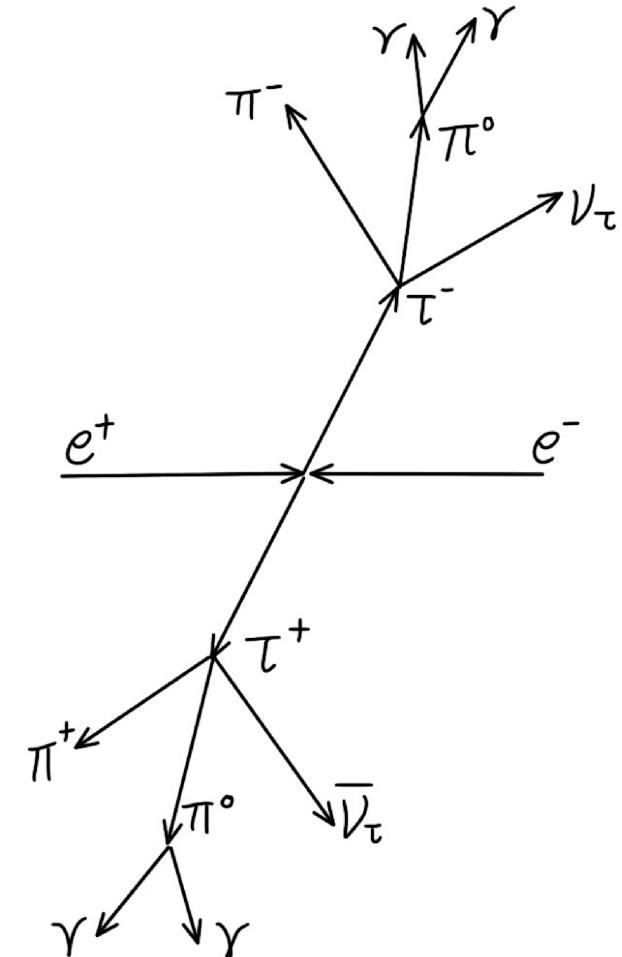
- Backgrounds:
 - Contamination from other decays

$$\tau^\pm \xrightarrow{17.8\%} \nu_\tau e^\pm \nu_e$$

$$\tau^\pm \xrightarrow{17.4\%} \nu_\tau \mu^\pm \nu_\mu$$

$$\tau^\pm \xrightarrow{10.8\%} \nu_\tau \pi^\pm$$

$$\tau^\pm \xrightarrow{9.3\%} \nu_\tau \pi^\pm \pi^0 \pi^0$$



Event Reconstruction

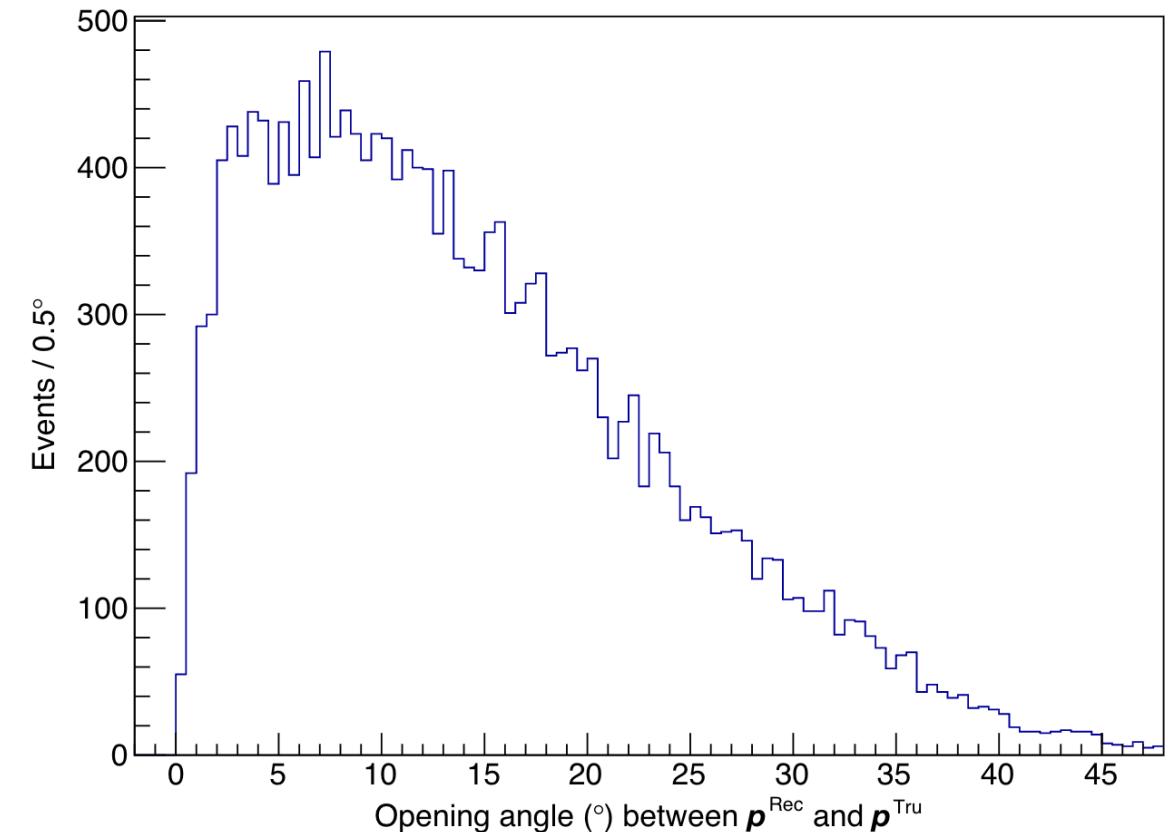
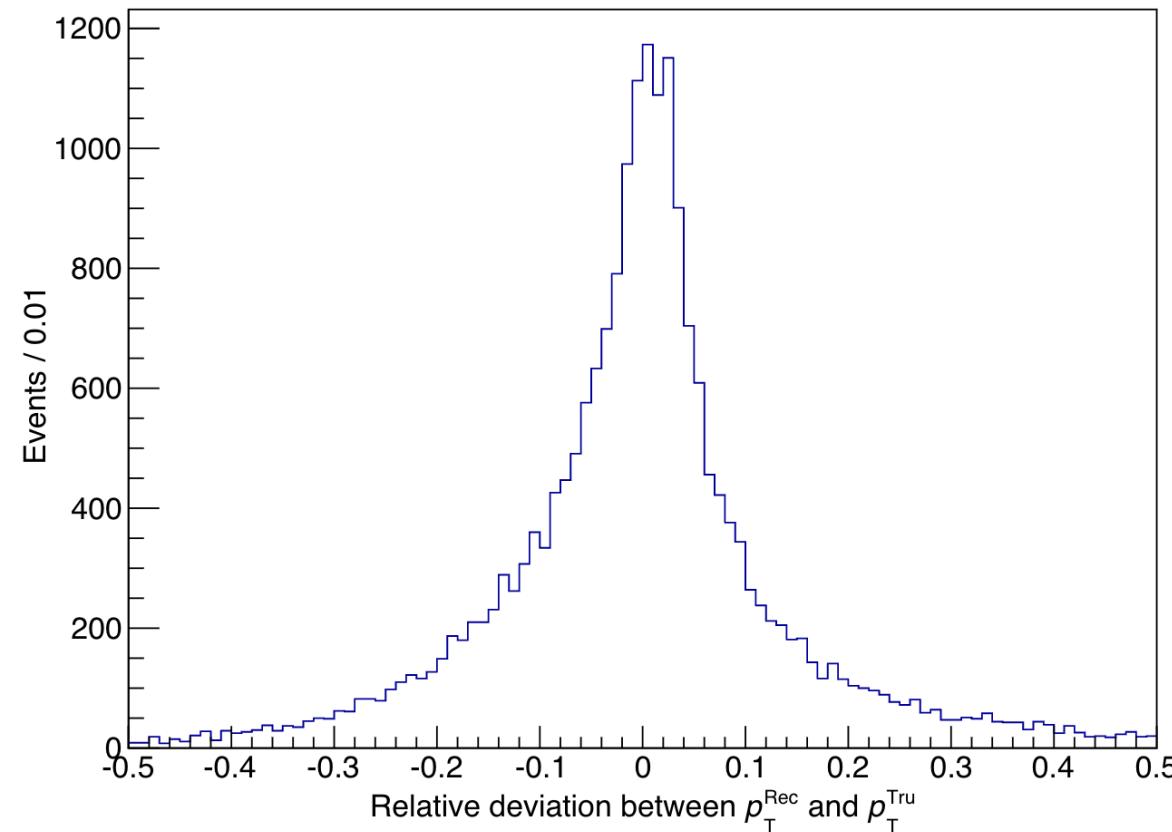
- Selections

No.	Step	Percentage of previous step		Signal purity
		Inclusive events	Signal events	
0	Total events	-	-	6.2%
1	Number of charged tracks = 2, total charge = 0	58.3%	76.5%	8.1%
2	Number of photons = 4	7.2%	23.7%	26.7%
3	Number of $\pi^+ = 1$, Number of $\pi^- = 1$	81.8%	92.1%	30.0%
4	Passed the particle pairing	25.2%	52.3%	62.5%
5	Passed event-level machine learning selection	57.9%	73.4%	79.3%
6	Passed the τ momentum reconstruction	97.0%	97.7%	80.0%
		Overall	0.49%	6.3%

Main background: multi- π^0 processes, mainly $\tau^\pm \rightarrow \pi^\pm \pi^0 \pi^0 \nu_\tau$ (about 14%)

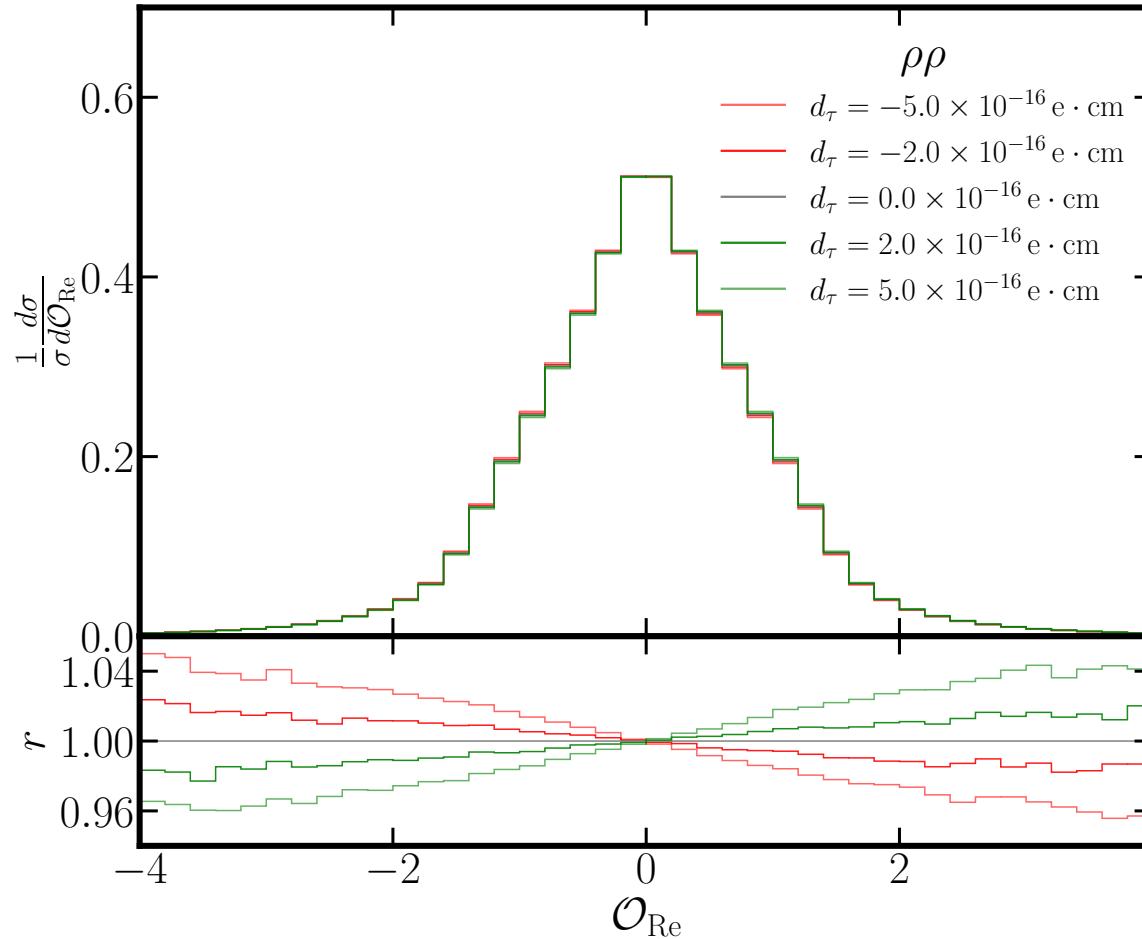
Tau Reconstruction Results

- Constructed vs. Truth level information



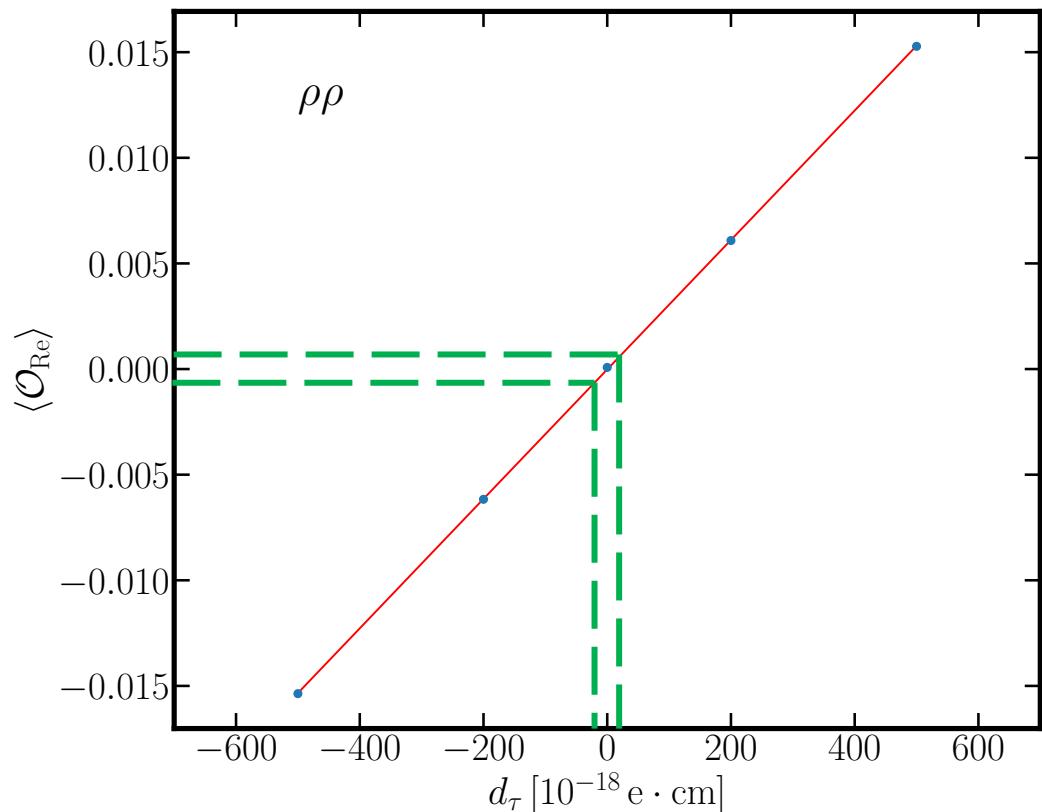
The distribution of OO

- Distribution of OO for $\rho\rho$ channel
 - Average over the two solutions of τ momentum



The distribution of $\langle \mathcal{O}_{Re} \rangle$

- $\langle \mathcal{O}_{Re} \rangle$ vs. $\Re(d_\tau)$ in $\rho\rho$ channel



$$\langle \mathcal{O}_{Re} \rangle = a_{Re} \cdot \Re(d_\tau) + b_{Re}$$

Uncertainties of $\langle \mathcal{O}_{Re} \rangle$: 1.09×10^{-4}

$$a_{Re} \approx 3.06 \times 10^{-5} / (10^{-18} e \cdot cm)$$

$\pm 1.43 \times 10^{-7}$

$$b_{Re} \approx -1.85 \times 10^{-5}$$

$\pm 4.87 \times 10^{-5}$

$$|\Re(d_\tau)| < 3.89 \times 10^{-18} e \cdot cm$$

3.55×10^{-18}

If neutrino can be reconstructed by using impact parameters
Another factor of 4 improvement can be achieved

$$|\Re(d_\tau)| \lesssim 9.73 \times 10^{-19} e \cdot cm$$

8.88×10^{-19}

Depends on the resolution of impact parameters

Summary and Outlook

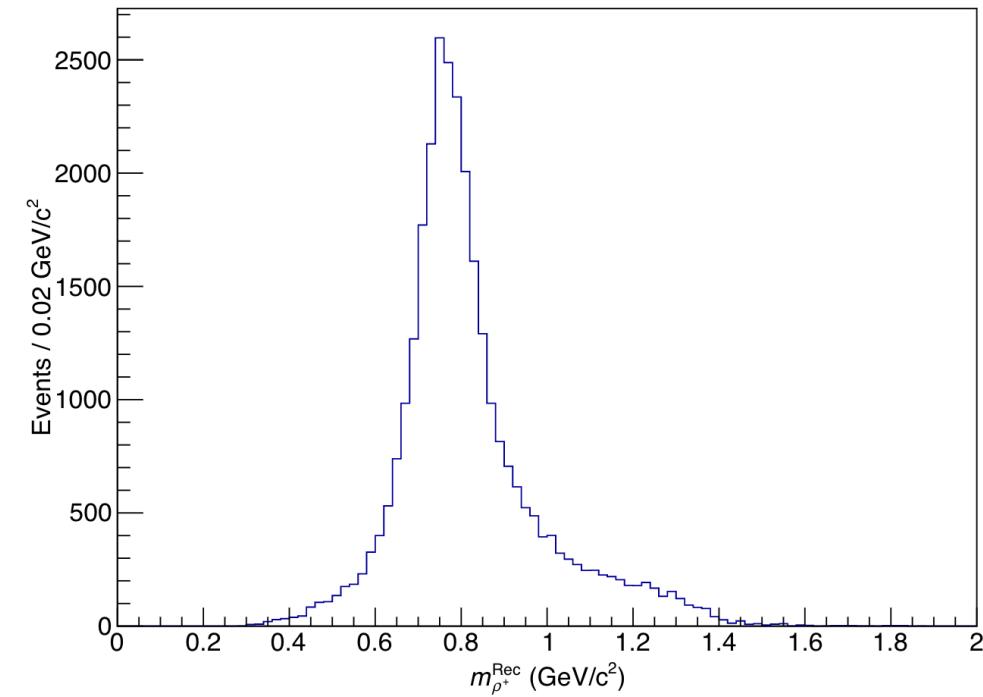
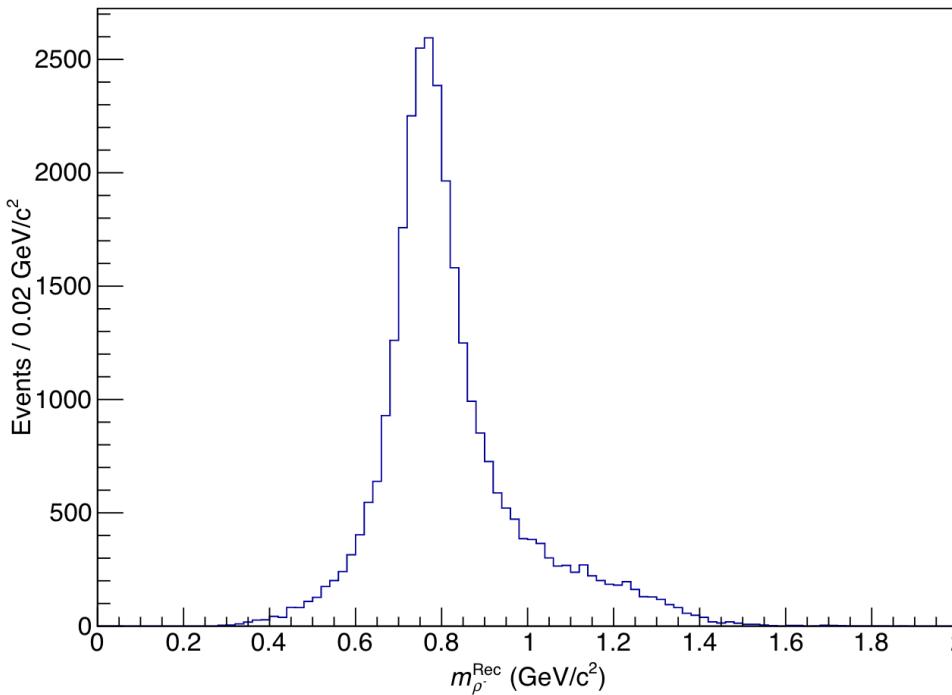
- STCF produce tremendous τ pair events
 - 3.5×10^9 per year
 - From $\rho\rho$ channel: $|\Re(d_\tau)| < 3.89 \times 10^{-18} e \cdot cm$
- Extension to other channels:
 - $\pi\pi, \pi\rho, \dots$
- The polarization of the beams
- Reconstruction of the neutrino - vertex detector
 - Removing the ambiguity
 - Factor of 4 improvement - [How good do we need for the vertex detection?](#)
- Application to $(g - 2)_\tau$

Thanks for your attention!

Backups

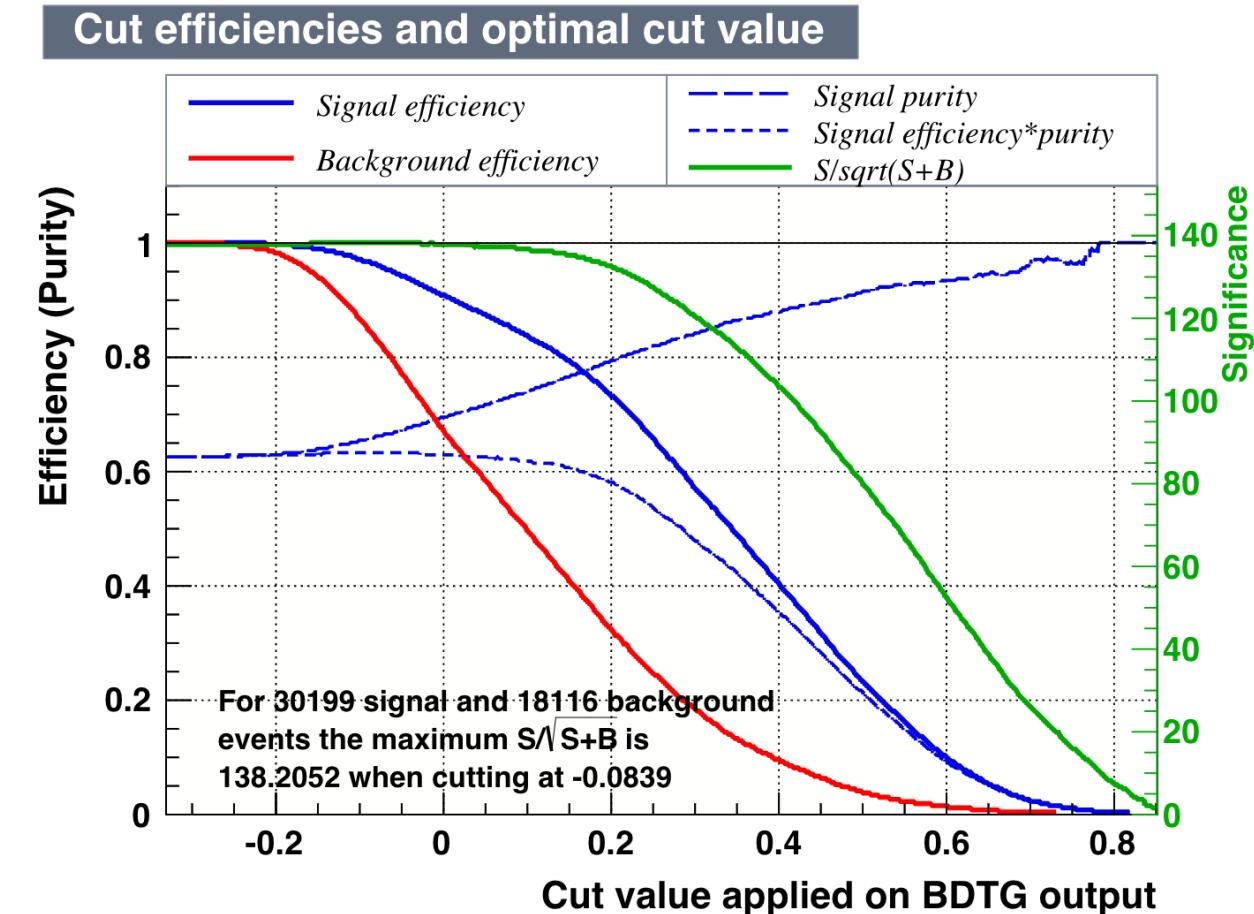
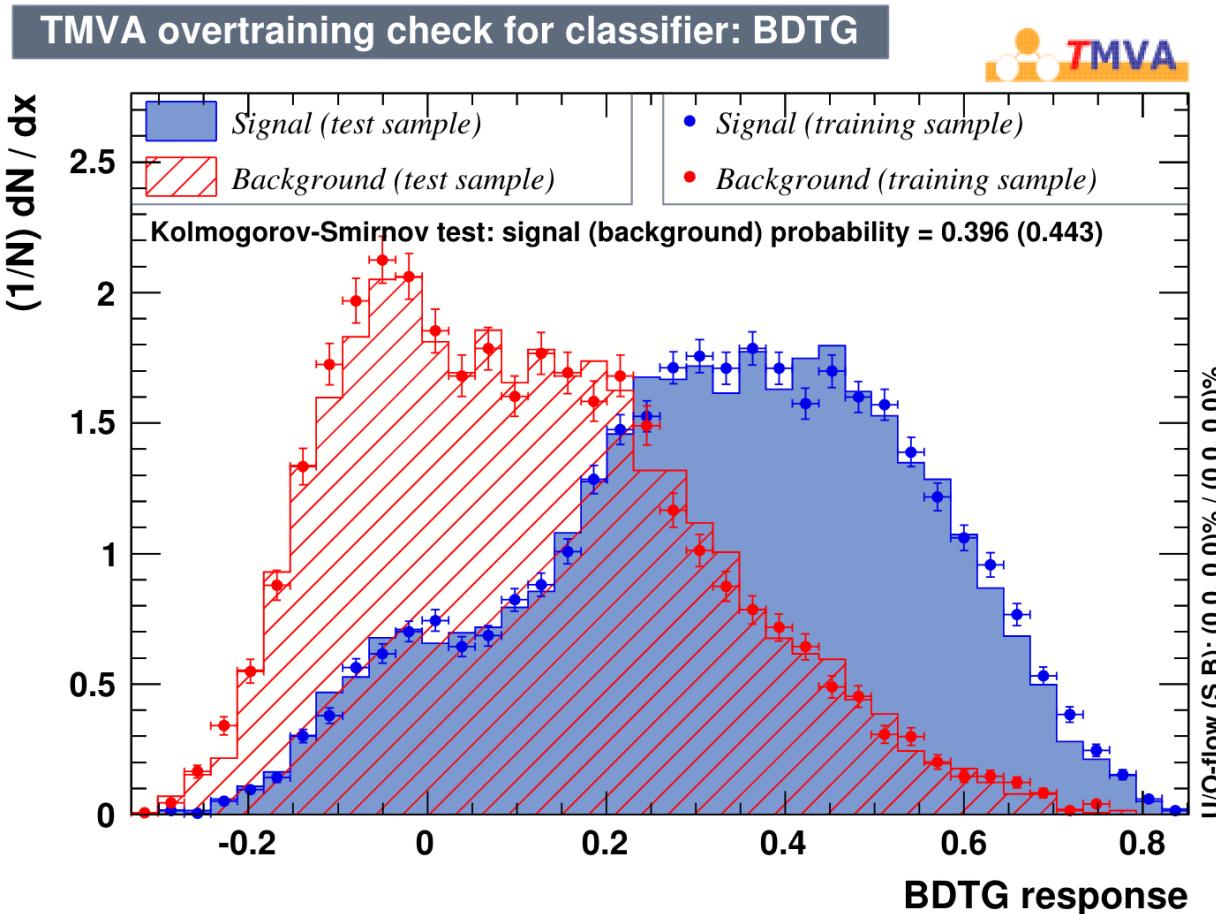
Particle Pairing

- Pairing the photons with the corresponding charged pions
 - Using kinematic constraints
 - Energy-momentum conservation
 - Mass constraints of the intermediate states



Event-level ML Selection

- BDTG
 - Inputs: momenta of π^\pm and four photons

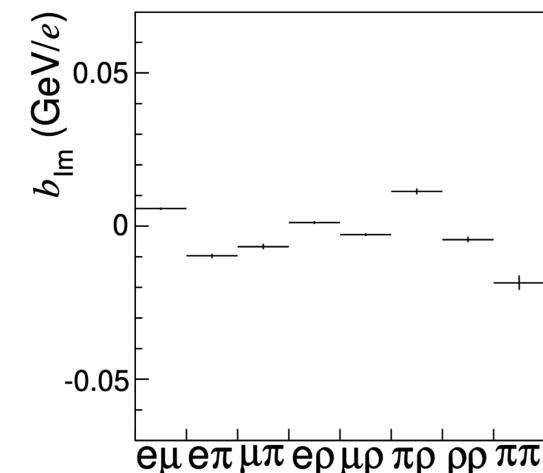
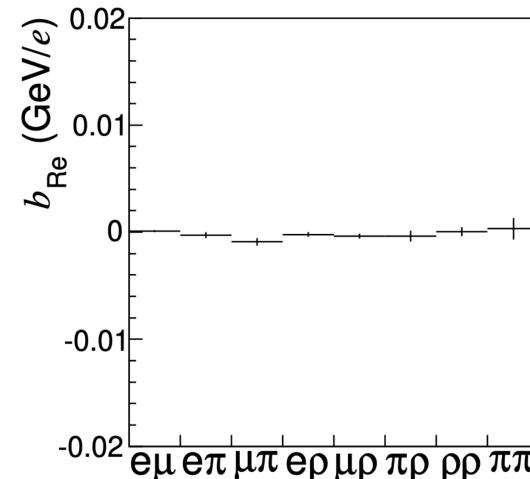
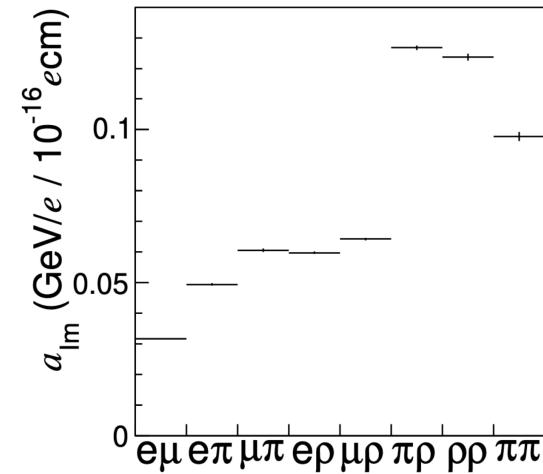
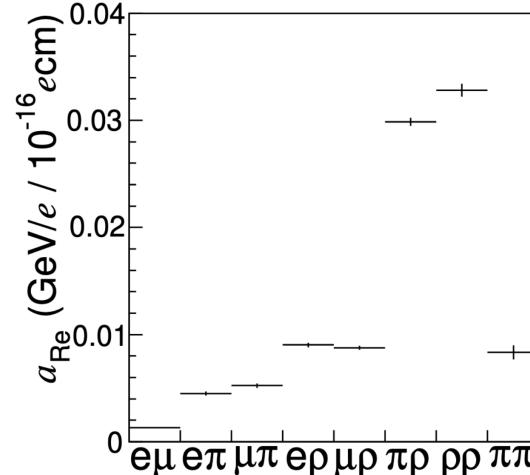


Different Decay Channels

- Comparison among different channels

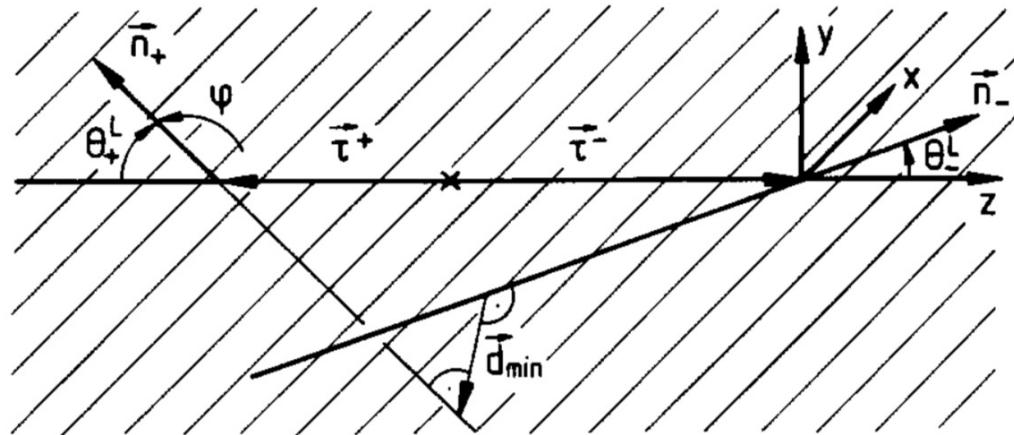
τ^\pm Decay Channels:

- $\tau^\pm \rightarrow \pi^\pm \nu$ (~10.8%)
- $\tau^\pm \rightarrow \pi^\pm \pi^0 \nu (\rho^\pm)$ (~25.4%)
- $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp \nu (a^\pm)$ (~9.3%)
- $\tau^\pm \rightarrow \ell^\pm \nu_\ell \nu$ (~17% $\times 2$)



Impact Parameters

- The ambiguity can be resolved by using the impact parameters



Phys. Lett. B 313 (1993) 458

- Advancement in silicon trackers
 - Much better resolutions

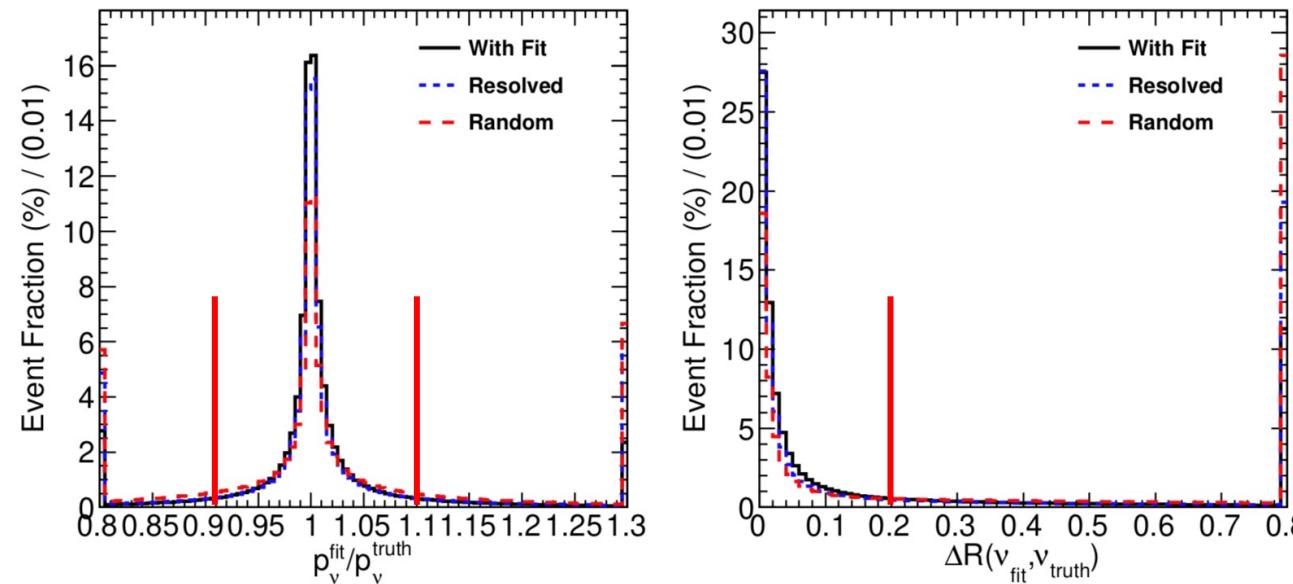
$$\sigma = a \oplus b/(p_T \sin^{1/2} \theta)$$

In mm	a	b/GeV
d_0	0.015	0.007
Belle z_0	0.020	0.010

Belle II TDR - 1011.0352

Performance

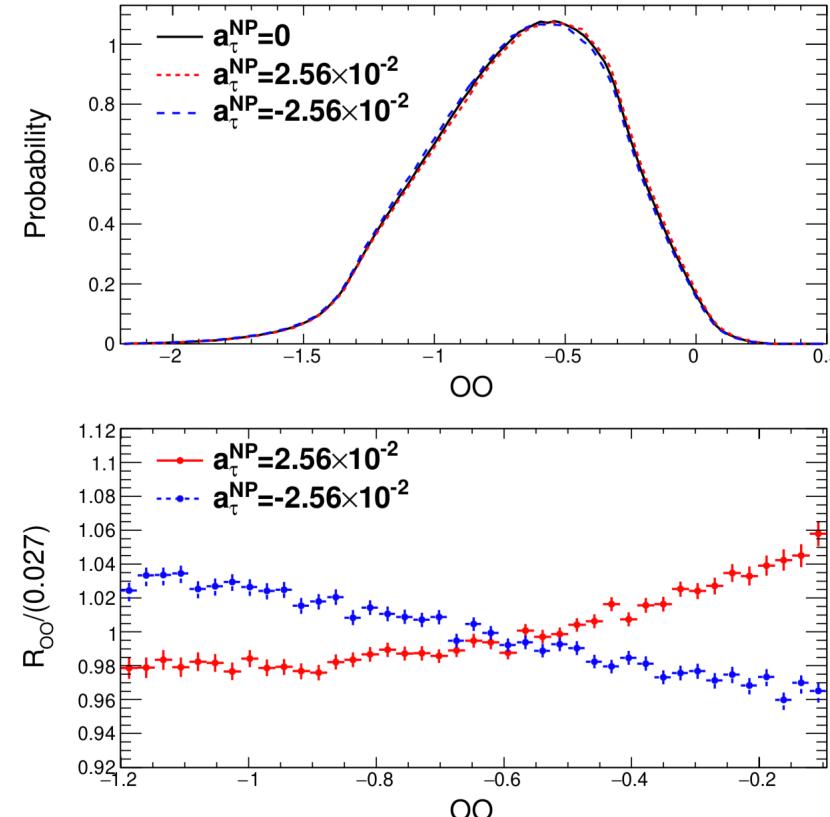
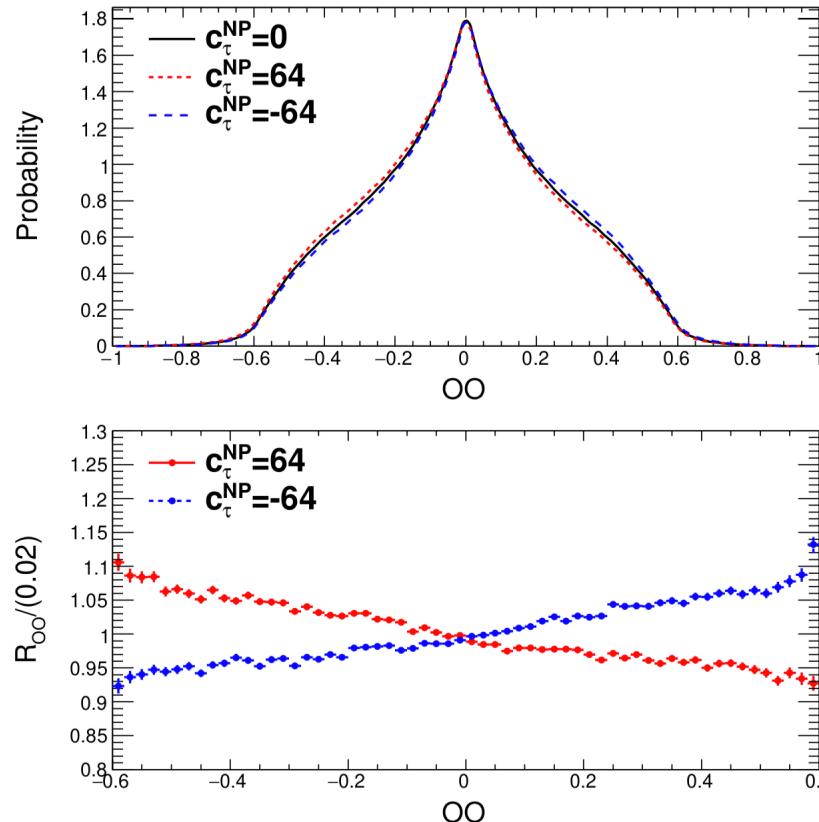
- The tau/neutrino momentum can be reconstructed by minimizing χ^2 event by event
- Compared with two other cases:
 - **Random**: Choose one of the two solutions randomly (PLB 313, 458)
 - **Resolved**: no fitting, but using the impact parameter to resolve the two-fold ambiguity



- Fractions of good reconstructed neutrinos: 42% → 59% → 65%

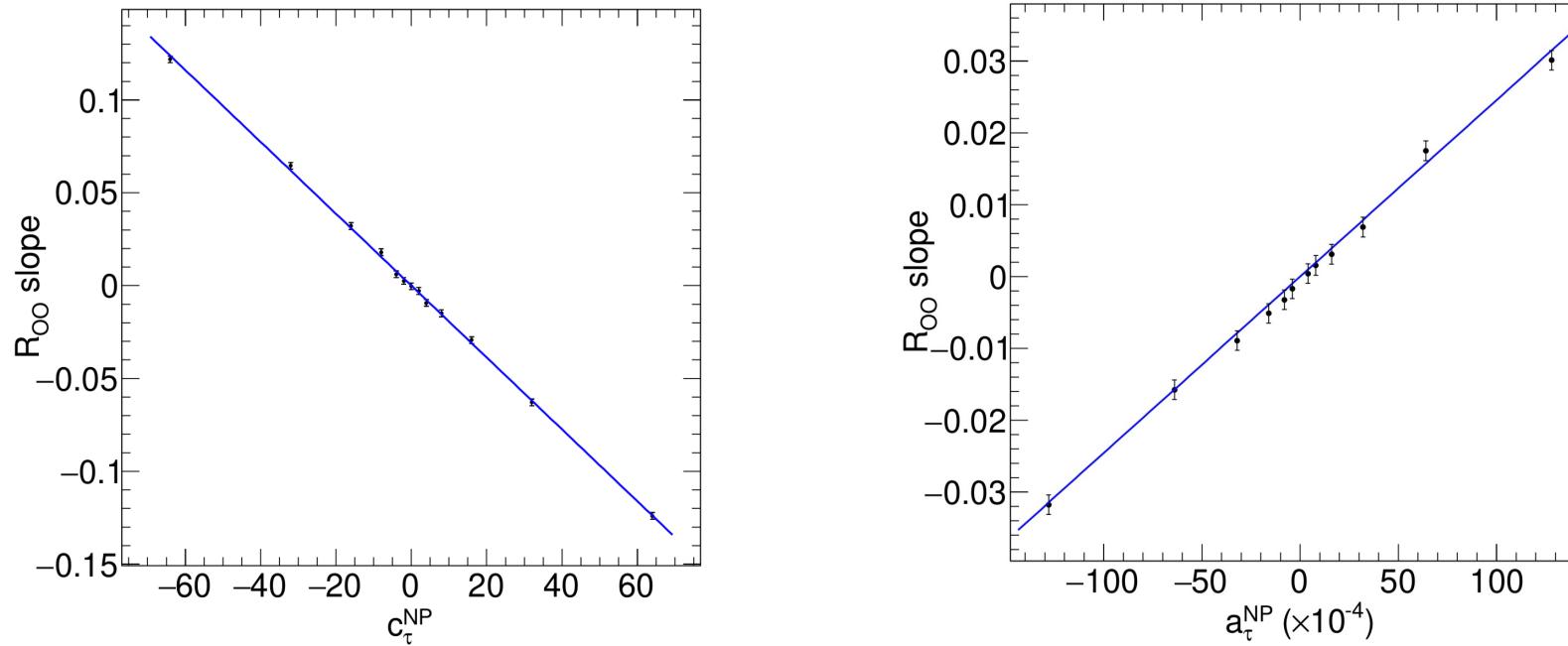
Sensitivities: OO distributions

with neutrino reconstructed



- For small c_τ and a_τ , the ratio of the distribution of different value can be parameterized as
 - $R_{OO} = 1 + b(OO - x_0)$

Sensitivities with neutrino reconstructed



\mathcal{L}	1 ab^{-1}	10 ab^{-1}	50 ab^{-1}
$ d_{\tau}^{\text{NP}} \text{ (e}\cdot\text{cm)}$	1.44×10^{-18}	4.56×10^{-19}	2.04×10^{-19}
$ a_{\tau}^{\text{NP}} $	1.24×10^{-4}	3.92×10^{-5}	1.75×10^{-5}

$$a_{\tau}^{\text{SM}} = 0.00117721(5)$$

Comparison with Belle's method is also performed, factor of 4 better

Xin Chen, Y. Wu; JHEP10(2019)089, 1803.00501
Xin Chen, Y. Wu; Eur.Phys.J. C77(2017)697, 1703.04855
Xin Chen, Y. Wu; Phys.Lett.B790(2019)332, 1708.02882