



Neutrino CP Violation

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August 27, 2023
FIND CPV @ USTC

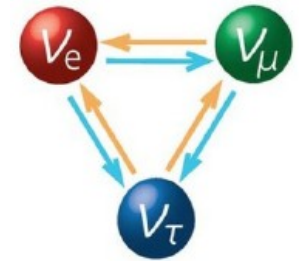
李政道研究所

Tsung-Dao Lee Institute

- **CP Violation in the Neutrino Sector**
- Dirac CP Phase & New Physics
- Accelerator + μ DAR for better CP measurement
- Majorana CP Phase
- Neutrino Electromagnetic Properties

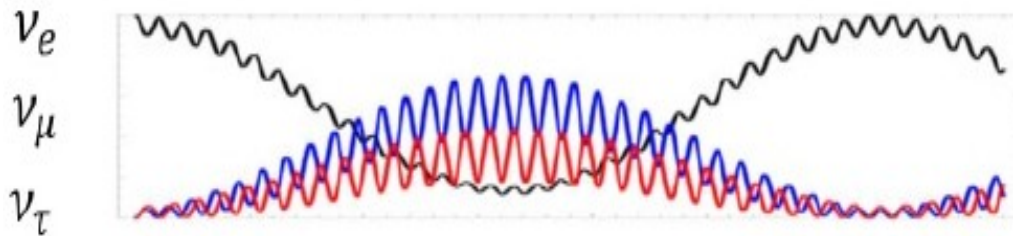
● PMNS Matrix

$$U_{\text{PMNS}} = \mathcal{P} \begin{pmatrix} c_s c_r & s_s c_r & s_r e^{-i\delta_D} \\ -s_s c_a - c_s s_a s_r e^{i\delta_D} & +c_s c_a - s_s s_a s_r e^{i\delta_D} & s_a c_r \\ +s_s s_a - c_s c_a s_r e^{i\delta_D} & -c_s s_a - s_s c_a s_r e^{i\delta_D} & c_a c_r \end{pmatrix} \mathcal{Q}$$



with $\mathcal{P} \equiv \text{diag}(e^{i\phi_1}, e^{i\phi_2}, e^{i\phi_3})$ & $\mathcal{Q} \equiv \text{diag}(e^{i\delta_{M1}}, 1, e^{i\delta_{M3}})$
 [(s, a, r) \equiv (12, 23, 13) for (solar, atmospheric, reactor) angles]

● Oscillation



Mass Splitting + Mixing



Neutrino Oscillation

$$P_{\alpha\beta}|_{\alpha \neq \beta} \equiv |A_{\alpha\beta}|^2 = \sin^2 2\theta \sin^2 \left(\delta m^2 \frac{L}{4E} \right)$$

Amplitude Frequency

Current Status of ν Oscillation

		Normal Ordering (best fit)		Inverted Ordering ($\Delta\chi^2 = 7.0$)	
		bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range
with SK atmospheric data	$\sin^2 \theta_{12}$	$0.304^{+0.012}_{-0.012}$	0.269 \rightarrow 0.343	$0.304^{+0.013}_{-0.012}$	0.269 \rightarrow 0.343
	$\theta_{12}/^\circ$	$33.45^{+0.77}_{-0.75}$	31.27 \rightarrow 35.87	$33.45^{+0.78}_{-0.75}$	31.27 \rightarrow 35.87
	$\sin^2 \theta_{23}$	$0.450^{+0.019}_{-0.016}$	0.408 \rightarrow 0.603	$0.570^{+0.016}_{-0.022}$	0.410 \rightarrow 0.613
	$\theta_{23}/^\circ$	$42.1^{+1.1}_{-0.9}$	39.7 \rightarrow 50.9	$49.0^{+0.9}_{-1.3}$	39.8 \rightarrow 51.6
	$\sin^2 \theta_{13}$	$0.02246^{+0.00062}_{-0.00062}$	0.02060 \rightarrow 0.02435	$0.02241^{+0.00074}_{-0.00062}$	0.02055 \rightarrow 0.02457
	$\theta_{13}/^\circ$	$8.62^{+0.12}_{-0.12}$	8.25 \rightarrow 8.98	$8.61^{+0.14}_{-0.12}$	8.24 \rightarrow 9.02
	$\delta_{CP}/^\circ$	230^{+36}_{-25}	144 \rightarrow 350	278^{+22}_{-30}	194 \rightarrow 345
	$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.42^{+0.21}_{-0.20}$	6.82 \rightarrow 8.04	$7.42^{+0.21}_{-0.20}$	6.82 \rightarrow 8.04
	$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.510^{+0.027}_{-0.027}$	+2.430 \rightarrow +2.593	$-2.490^{+0.026}_{-0.028}$	-2.574 \rightarrow -2.410

Gonzalez-Garcia, Maltoni & Schwetz [2111.03086]

Daya Bay heralded a new era of precision measurement in 2012!

Kam-Biu Luk @ Neutrino 2022

- Heavy neutrinos (N)

$$\bar{\nu} M_D \mathcal{N} + h.c. + \bar{\mathcal{N}} M_N \mathcal{N} = \begin{pmatrix} \bar{\nu} & \bar{\mathcal{N}} \end{pmatrix} \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \begin{pmatrix} \nu \\ \mathcal{N} \end{pmatrix}$$

required by **Grand Unification Theory (大统一理论)**

- Seeaw Mechanism

The diagonalization of the full mass matrix

$$\begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \Rightarrow M_\nu = -M_D \frac{1}{M_N} M_D^T$$

$$M_D \sim O(100) \text{ GeV}, \quad M_N \sim O(10^{15}) \text{ GeV} \quad \rightarrow \quad M_\nu \sim O(0.01) \text{ eV}$$

Light neutrino mass M_ν is suppressed by the heavy ones



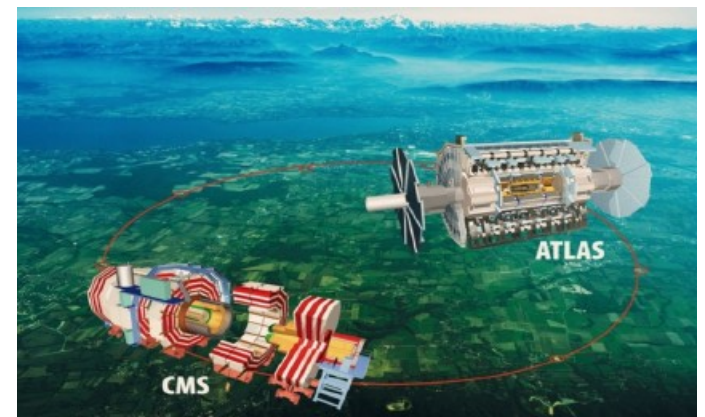
柳田勉、何小刚：
大家一起来玩跷跷板！

Neutrino vs Higgs

- **Higgs boson** \Rightarrow electroweak symmetry breaking & mass. $\sim O(100)\text{GeV}$
- **Chiral symmetry breaking** \Rightarrow majority of mass.
- The world seems not affected by the tiny neutrino mass?
 - Neutrino mass \Rightarrow Mixing
 - 3 Neutrino \Rightarrow possible **CP violation**
 - CP violation \Rightarrow **Leptogenesis**
 - \Rightarrow **Matter-Antimatter Asymmetry**
 - There is something left in the Universe.
 - **EW Baryogenesis** is not enough.



Daya Bay @ **March 8, 2012**



LHC @ **July 4, 2012**

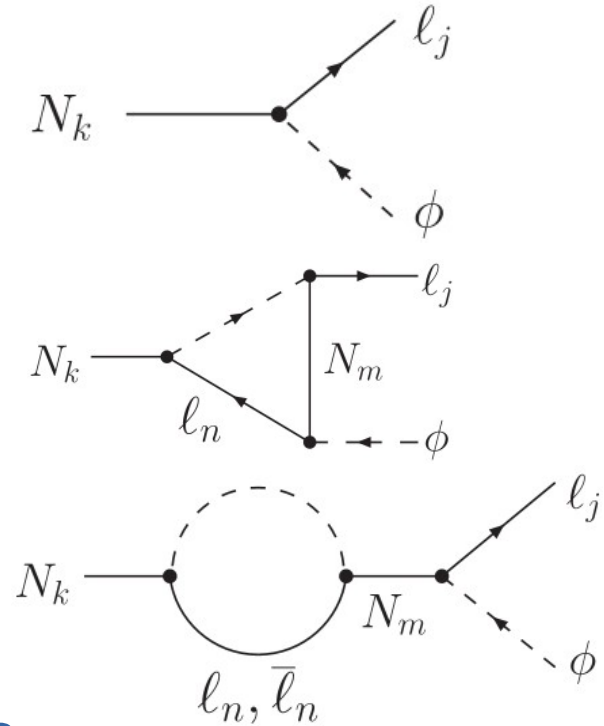
Leptogenesis

With decreasing temperature, heavy N decays to light SM particles.

$$N_k \rightarrow \begin{cases} \ell_j + \bar{\phi} \\ \bar{\ell}_j + \phi \end{cases}$$



Matter-Antimatter Asymmetry



- Interference between tree & loop diagrams

$$\Gamma = \Gamma_{\text{tree}} + \Gamma_{\text{loop}}(+\delta_D, +\delta_M)$$

$$\bar{\Gamma} = \Gamma_{\text{tree}} + \Gamma_{\text{loop}}(-\delta_D, -\delta_M)$$

The matter-antimatter asymmetry needs Dirac/Majorana CP phases.



Yanagida



Fukugita

- CP Violation in the Neutrino Sector
- **Dirac CP Phase & New Physics**
- Accelerator + μ DAR for better CP measurement
- Majorana CP Phase
- Neutrino Electromagnetic Properties

Current Status of CP Measurement

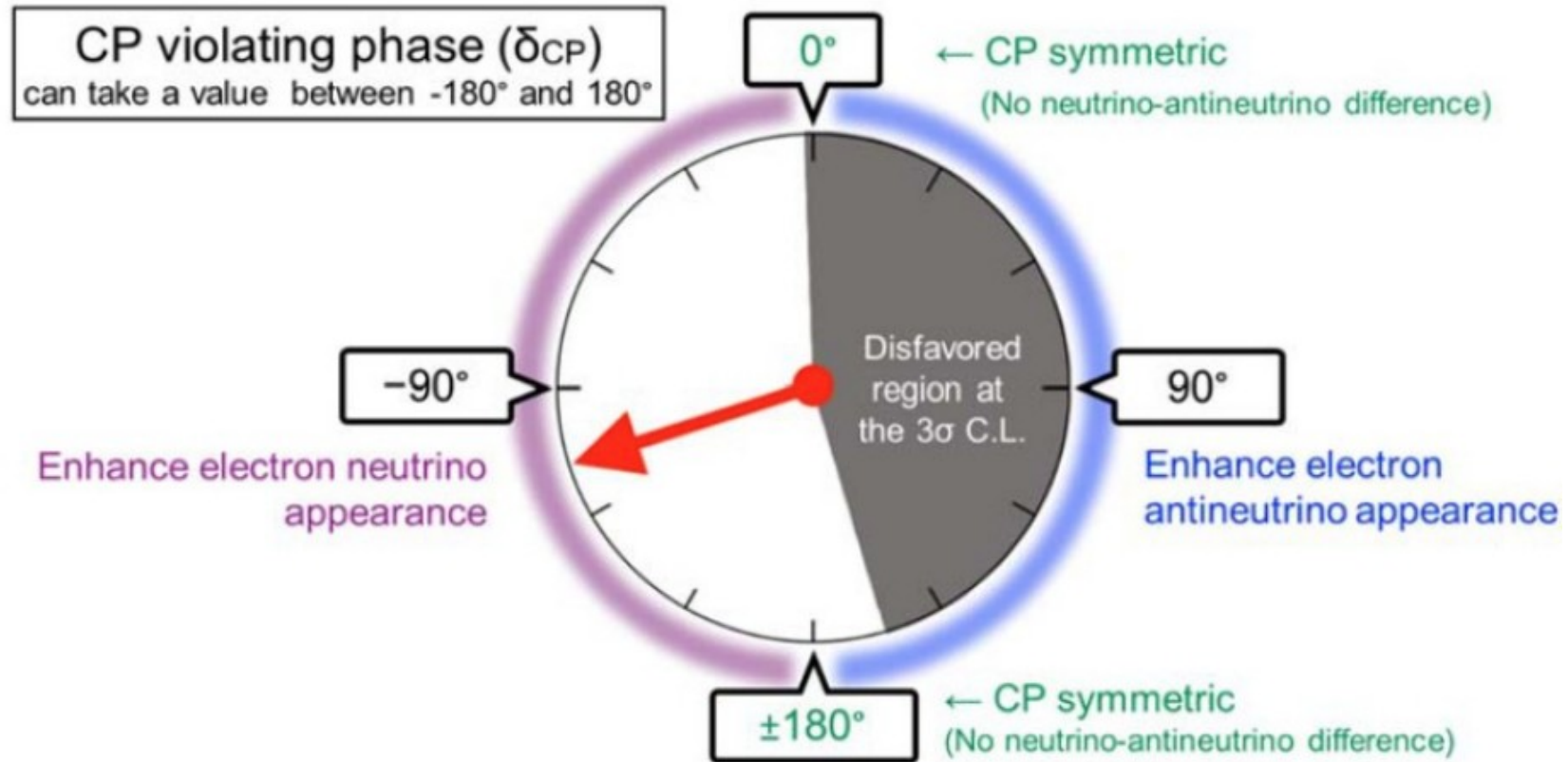


Fig.1 The arrow indicates the value most compatible with the data. The gray region is disfavored at 99.7% (3σ) confidence level. Nearly half of the possible values are excluded.

Nature vol. 580, pages 339-344(2020)

https://www.kek.jp/en/newsroom/attic/PR20200416_T2K_E.pdf

PHYSICAL REVIEW LETTERS

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Feasibility of Using High-Energy Neutrinos to Study the Weak Interactions

M. Schwartz
Phys. Rev. Lett. **4**, 306 – Published 15 March 1960

Theoretical Discussions on Possible High-Energy Neutrino Experiments

T. D. Lee and C. N. Yang
Phys. Rev. Lett. **4**, 307 – Published 15 March 1960

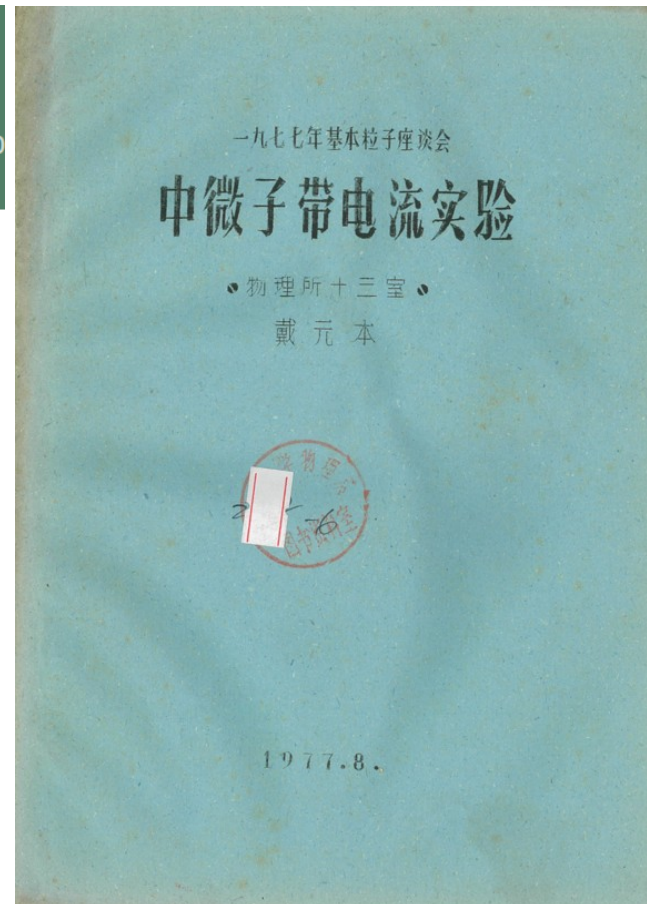
历史印记——60年前高能加速器中微子实验原理诞生

🕒 2021-01-04 👁 50

作者：葛韶锋

这项60年前的工作，今天还在为粒子物理研究指引着前进的方向

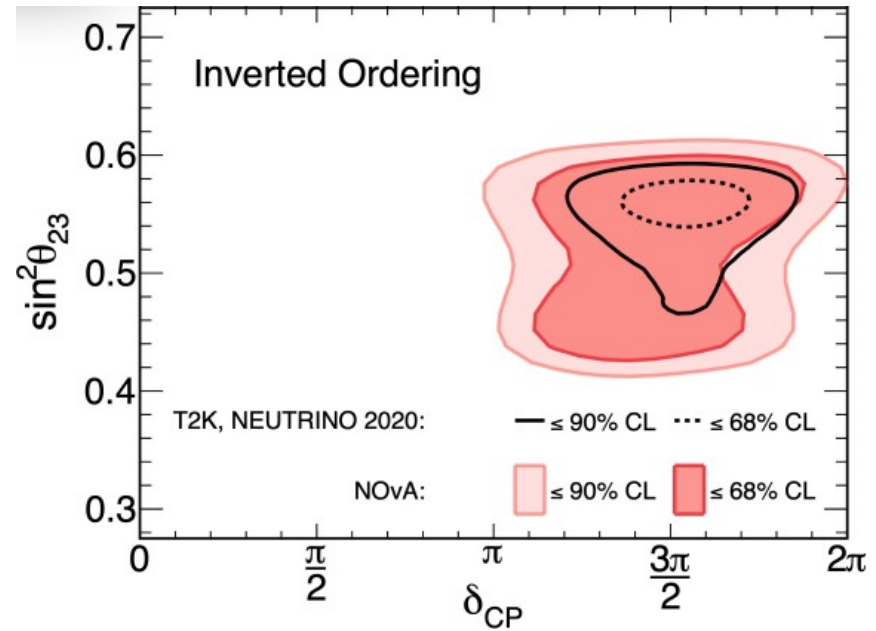
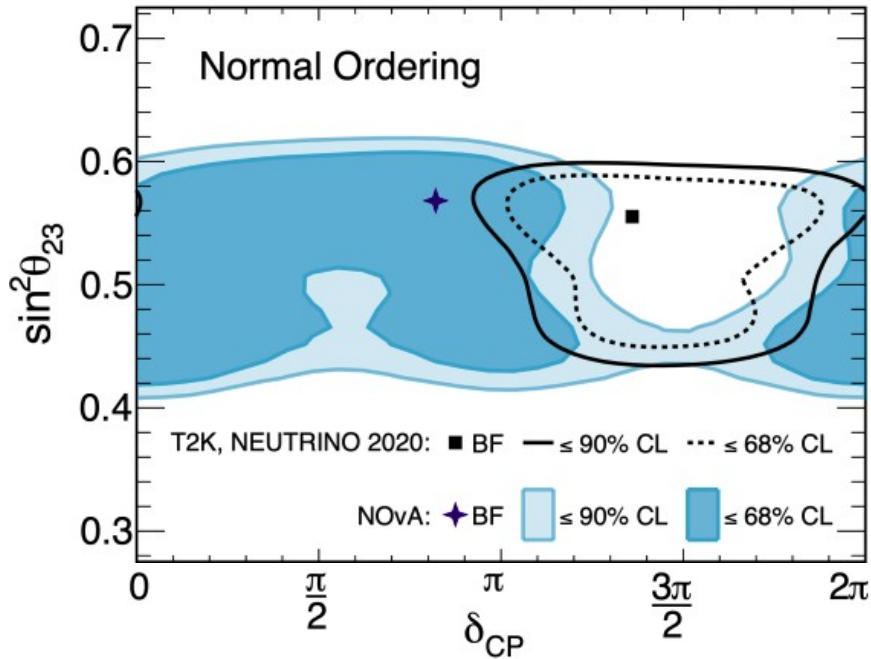
中微子是粒子物理标准模型中非常独特而重要的基本粒子。中微子只参与弱相互作用，能够穿透几光年的铅板，是非常难以探测的鬼魅粒子。然而，恰恰是中微子率先给出了超出标准模型的第一个新物理现象——中微子振荡——这已经得到大量实验的证实，并于2015年获诺贝尔物理学奖，成为指引超出标准模型新物理研究的重要线索。在上世纪五六十年代粒子物理发展的黄金期，施瓦兹、李政道与杨振宁提出的高能加速器中微子实验原理，是一个非常革命性的思想，一直在有力地推动着包括中微子在内的粒子物理不断向前发展。



何小刚，李政道先生和现代中微子物理，《现代物理知识杂志》2021

<https://tdli.sjtu.edu.cn/CN/customize/436?columnId=35>

Tension between T2K & NOvA



NOvA [2108.08219]

Denton, Gehrlein & Pestes [2008.01110]

Rahaman, Razzaque & Sankar [2201.03250]

Chatterjee & Palazzo [2005.10338]

- Non-Standard Interactions
- Non-Unitary Mixing
- Lorentz Violation
- Sterile Neutrinos

● Non-Standard Interactions

$$\mathcal{H} \equiv \frac{1}{2\mathbf{E}_\nu} U \begin{pmatrix} 0 & & \\ & \Delta m_s^2 & \\ & & \Delta m_a^2 \end{pmatrix} U^\dagger + V_{cc} \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}$$

Du, Li, Tang, Vihonen & Yu [2011.14292, 2106.15800]

Tang & Zhang [1705.09500]

Du, Li, Tang, Vihonen & Yu [2106.15800]

Scalar NSI: SFG & Parke [1812.08376] Smirnov & Xu [1909.07505]

● Non-Unitary Mixing

$$N = N^{NP} U = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ |\alpha_{21}| e^{i\phi} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

SFG, Pasquini, Tortola & Valle [1605.01670]

Tang, Zhang & Li [1708.04909]

Hu, Ling, Tang & Wang [2008.09730]

● Sterile neutrinos

Chatterjee & Palazzo [2005.10338]

● Lorentz violation

Rahaman, Razzaque & Sankar [2201.03250]

● Dark NSI

Berlin [1608.01307]

Liao, Marfatia & Whisnant [1803.01773]

Zhao [1701.02735]

Chao, Hu, Jiang & Jin [2009.14703]

Brdar, Kopp, Liu, Prass & Wang [1705.09455]

SFG, Murayama [1904.02518]

Non-Standard Interactions (NSI)

$$\mathcal{H} \equiv \frac{1}{2E_\nu} U \begin{pmatrix} 0 & & \\ & \Delta m_s^2 & \\ & & \Delta m_a^2 \end{pmatrix} U^\dagger + V_{cc} \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}$$

- Standard Interaction – V_{cc} (also V_{nc})

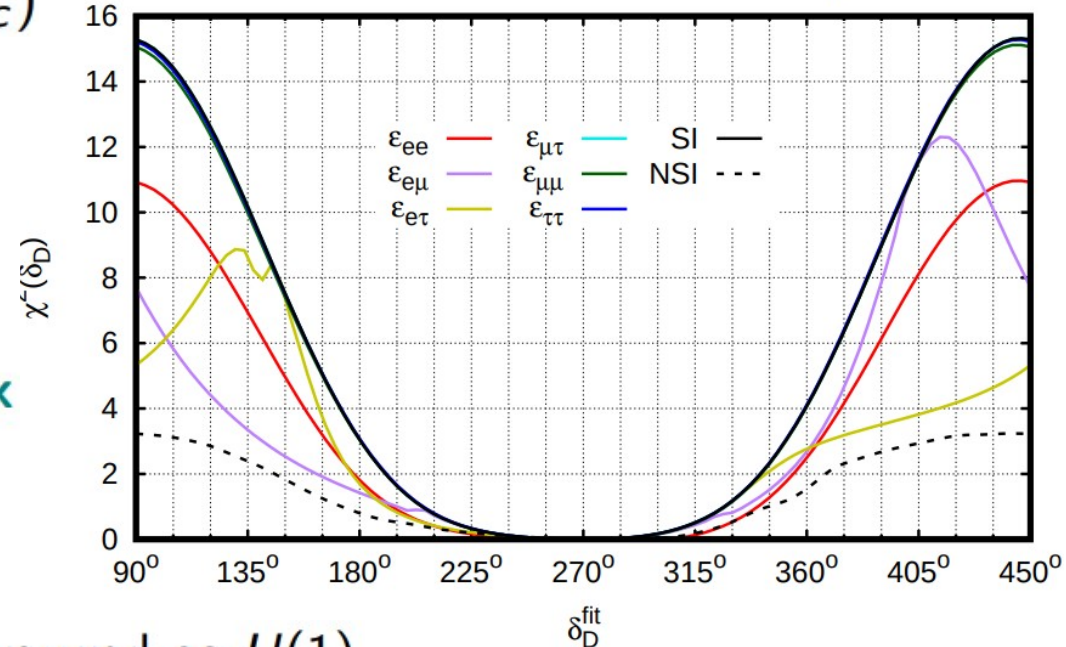
- Non-Standard Interaction – $\epsilon_{\alpha\beta}$

- Diagonal $\epsilon_{\alpha\alpha}$ are real
- Off-diagonal $\epsilon_{\alpha\neq\beta}$ are complex
- Both can fake CP

- Z' in LMA-Dark model with $L_\mu - L_\tau$ gauged as $U(1)$

- $M_{Z'} \sim \mathcal{O}(10)\text{MeV}$
- $g_{Z'} \sim 10^{-5}$

The effect of NSI on the CP sensitivity at T2K [$\delta_D^{\text{true}} = -90^\circ$]

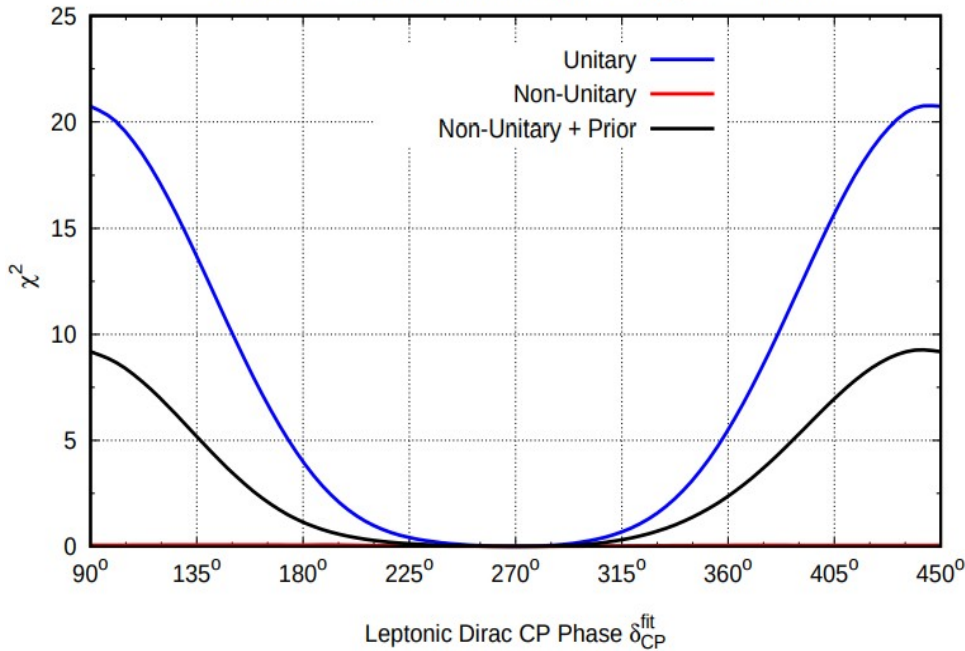


Non-Unitarity Mixing (NUM)

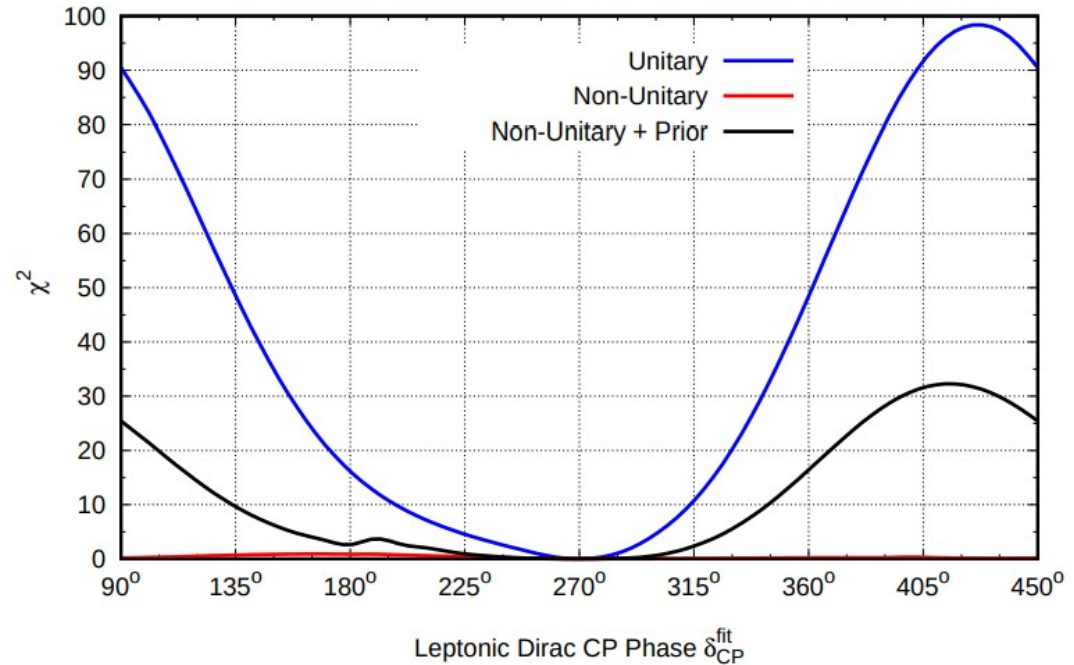
$$N = N^{NP} U = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ |\alpha_{21}| e^{i\phi} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

$$P_{\mu e}^{NP} = \alpha_{11}^2 \left\{ \alpha_{22}^2 \left[c_a^2 |S'_{12}|^2 + s_a^2 |S'_{13}|^2 + 2c_a s_a (\cos \delta_D \mathbb{R} - \sin \delta_D \mathbb{I})(S'_{12} S'_{13}^*) \right] + |\alpha_{21}|^2 P_{ee} \right. \\ \left. + 2\alpha_{22} |\alpha_{21}| \left[c_a (c_\phi \mathbb{R} - s_\phi \mathbb{I})(S'_{11} S'_{12}^*) + s_a (c_{\phi+\delta_D} \mathbb{R} - s_{\phi+\delta_D} \mathbb{I})(S'_{11} S'_{13}^*) \right] \right\} .$$

The effect of including non-unitarity at T2K [$\delta_{CP}^{\text{true}} = -90^\circ$, NH]



The effect of including non-unitarity at T2HK [$\delta_{CP}^{\text{true}} = -90^\circ$, NH]



SFG, Pasquini, Tortola & Valle [1605.01670]

- Heavy Neutrinos

$$\bar{\nu} M_D \mathcal{N} + h.c. + \bar{\mathcal{N}} M_N \mathcal{N} = \begin{pmatrix} \bar{\nu} & \bar{\mathcal{N}} \end{pmatrix} \begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \begin{pmatrix} \nu \\ \mathcal{N} \end{pmatrix}$$

The diagonalization of the full mass matrix

$$\begin{pmatrix} 0 & M_D \\ M_D^T & M_N \end{pmatrix} \Rightarrow M_\nu = -M_D \frac{1}{M_N} M_D^T$$



$$N = N^{NP} U = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ |\alpha_{21}| e^{i\phi} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

NUM is not exotic at all!

- CP Violation in the Neutrino Sector
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- **Accelerator + μ DAR for better CP measurement**
- Majorana CP Phase
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$$P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e} - P_{\nu_\mu \rightarrow \nu_e} = \alpha\pi \sin(2\theta_s) \sin(2\theta_r) \sin(2\theta_a) \cos\theta_r \sin\delta_D$$

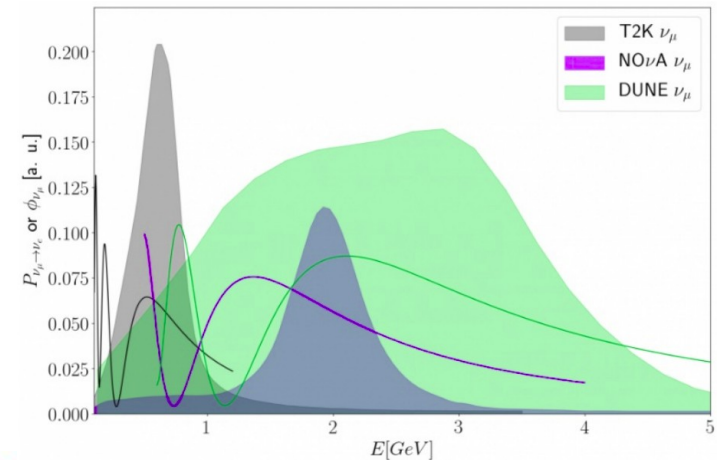
@ 1st oscillation peak

● Efficiency:

- Proton accelerators produce ν more efficiently than $\bar{\nu}$ ($\sigma_\nu > \sigma_{\bar{\nu}}$).
- The $\bar{\nu}$ mode needs more beam time [$\mathbf{T_{\bar{\nu}} : T_\nu = 2 : 1}$].
- Undercut statistics \Rightarrow Difficult to reduce the uncertainty.

● Degeneracy:

- Only $\sin\delta_D$ appears in $P_{\nu_\mu \rightarrow \nu_e}$ & $P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e}$.
- Cannot distinguish δ_D from $\pi - \delta_D$.

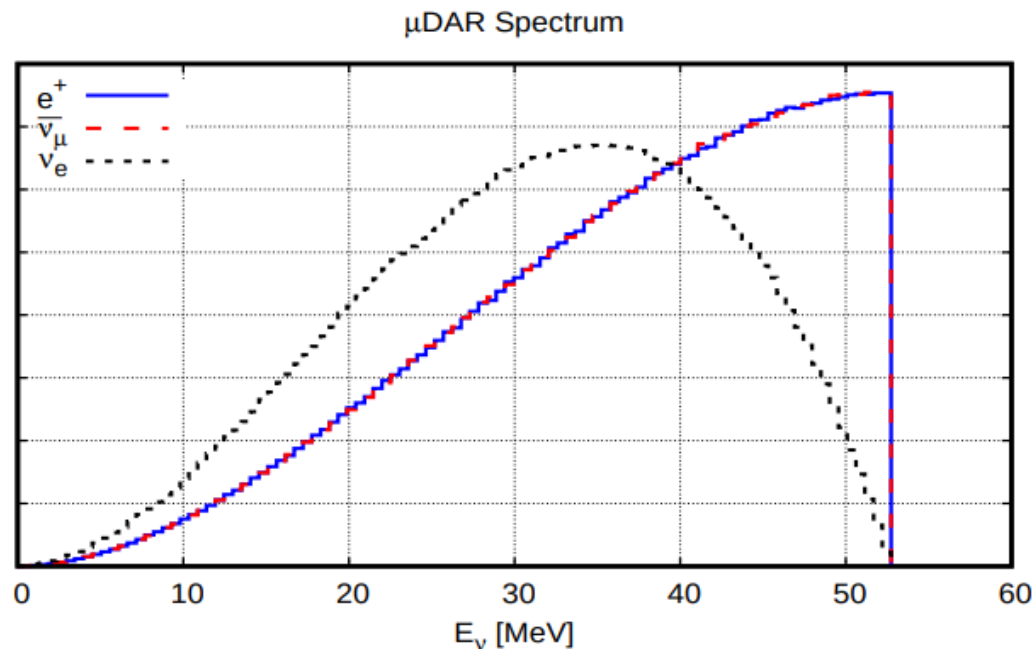


- **CP Uncertainty** $\frac{\partial P_{\mu e}}{\partial \delta_D} \propto \cos\delta_D \Rightarrow \Delta(\delta_D) \propto \mathbf{1 / \cos\delta_D}$

SFG [1704.08518, PoS NuFact2019 (2020) 108]

Solution: Muon Decay at Rest

- A cyclotron produces 800 MeV proton beam @ fixed target.
- Produce π^\pm which stops &
 - π^- is absorbed,
 - π^+ decays @ rest: $\pi^+ \rightarrow \mu^+ + \nu_\mu$.
- μ^+ stops & decays @ rest: $\mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e$



Well predicted spectrum!

- $\bar{\nu}_\mu$ travel in all directions, oscillating as they go.
- A detector measures the $\bar{\nu}_e$ from $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ oscillation.

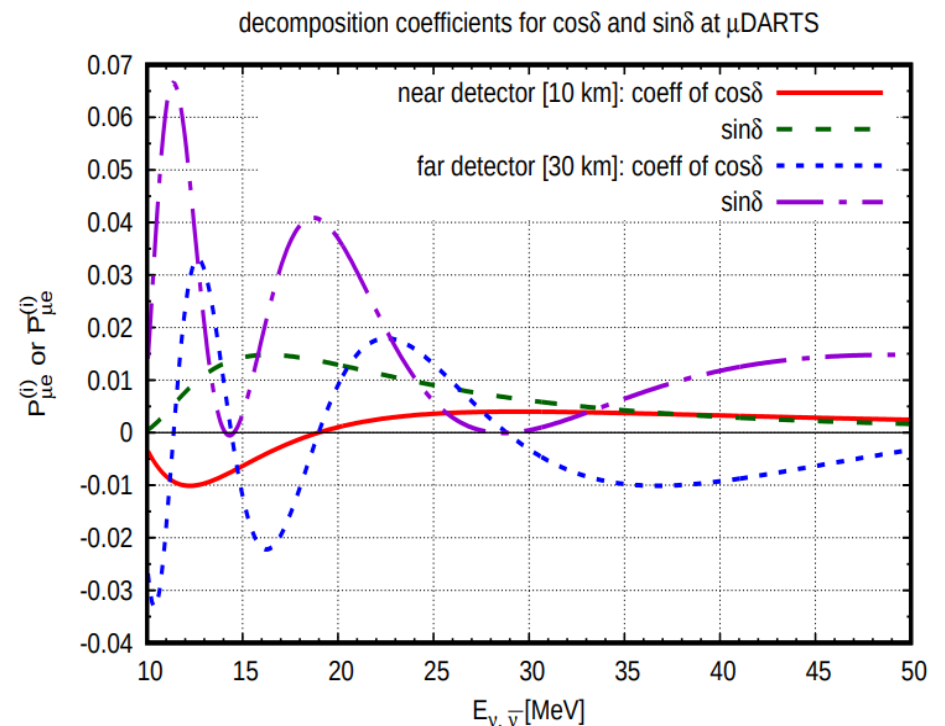
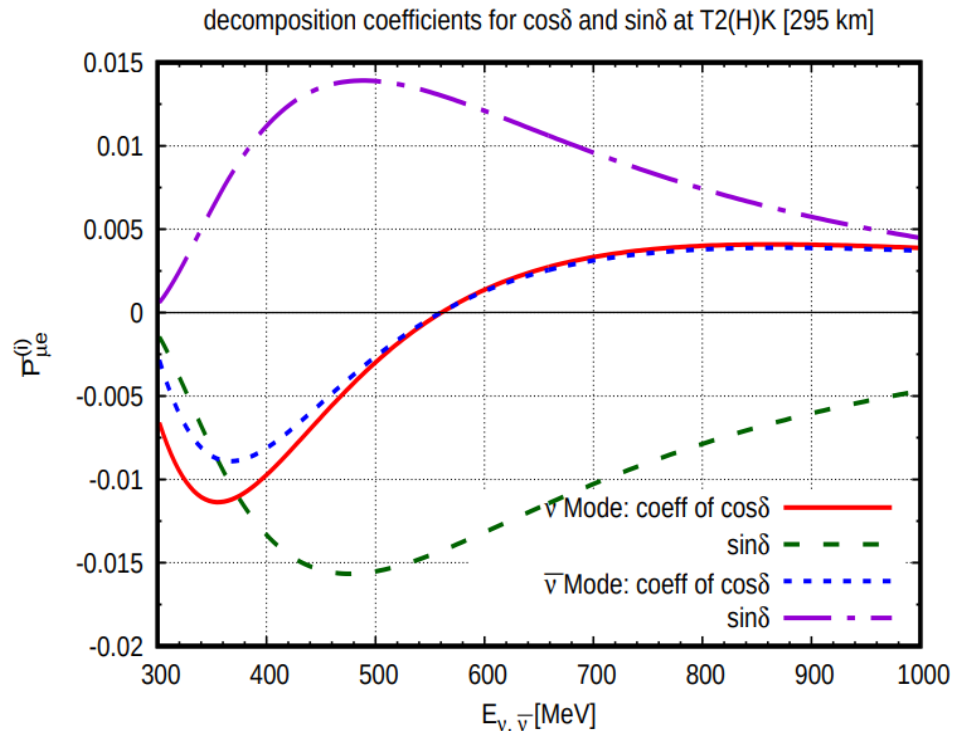
μ DAR & Accelerator ν

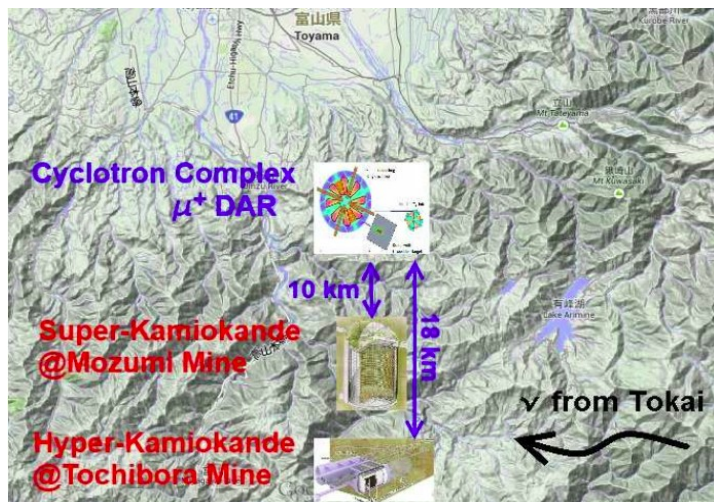
Combining $\nu_\mu \rightarrow \nu_e$ @ accelerator [narrow peak @ 550 MeV] & $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ @ μ DAR [wide peak \sim 45 MeV] solves the 2 problems:

- **Efficiency:**

- $\bar{\nu}$ @ high intensity, μ DAR is plentiful enough.
- Accelerator Exps can devote all run time to the ν mode. With same run time, the statistical uncertainty drops by $\sqrt{3}$.

- **Degeneracy:** (decomposition in propagation basis [1309.3176])





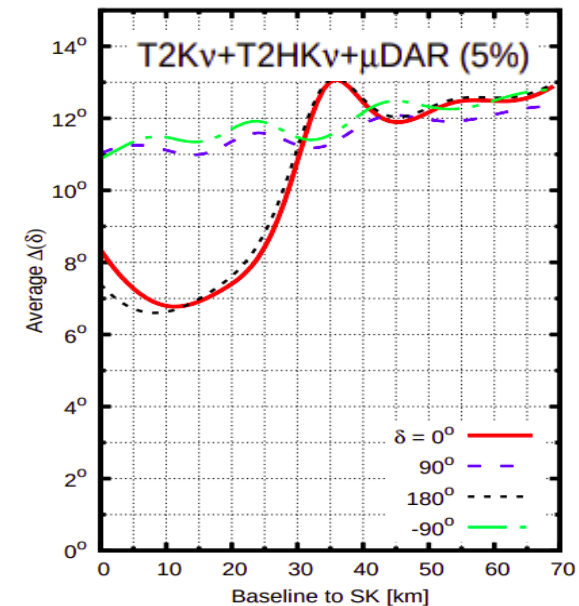
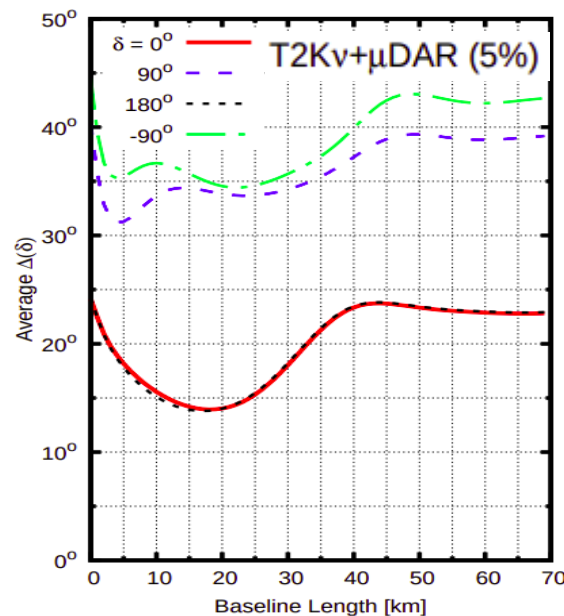
1 μ DAR source + 2 detectors

Advantages

- Full (**100%**) duty factor!
- **Lower** intensity: $\sim 9\text{mA}$ [$\sim 4\times$ lower than DAE δ ALUS]
- Not far beyond the current state-of-art technology of cyclotron [2.2mA @ Paul Scherrer Institute]
- MUCH **cheaper** & technically **easier**.
 - Only one cyclotron.
 - Lower intensity.

Disadvantage: A second detector!

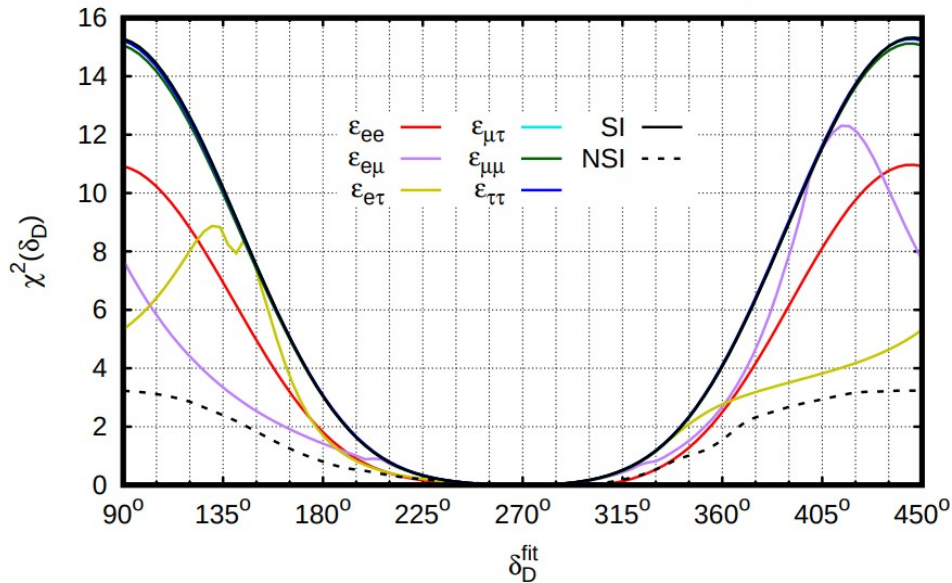
Harnik, Kelly & Machado [1911.05088]



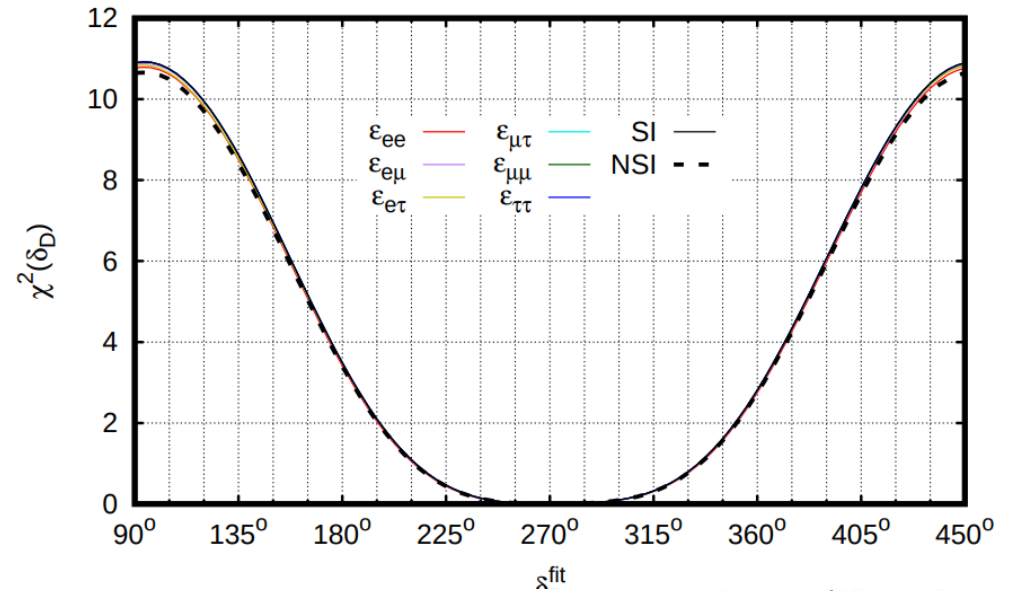
Evslin, SFG, Hagiwara [1506.05023]
SFG, Pasquini, Tortola, Valle [1605.01670]
SFG, Smirnov [1607.08513]

Guarantee CP against NSI

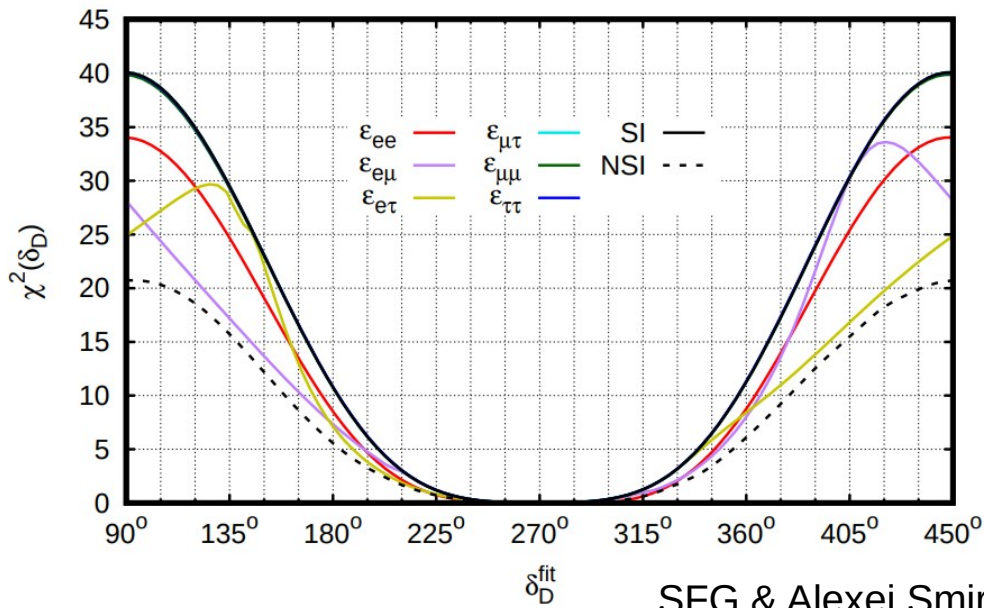
The effect of NSI on the CP sensitivity at T2K [$\delta_D^{\text{true}} = -90^\circ$]



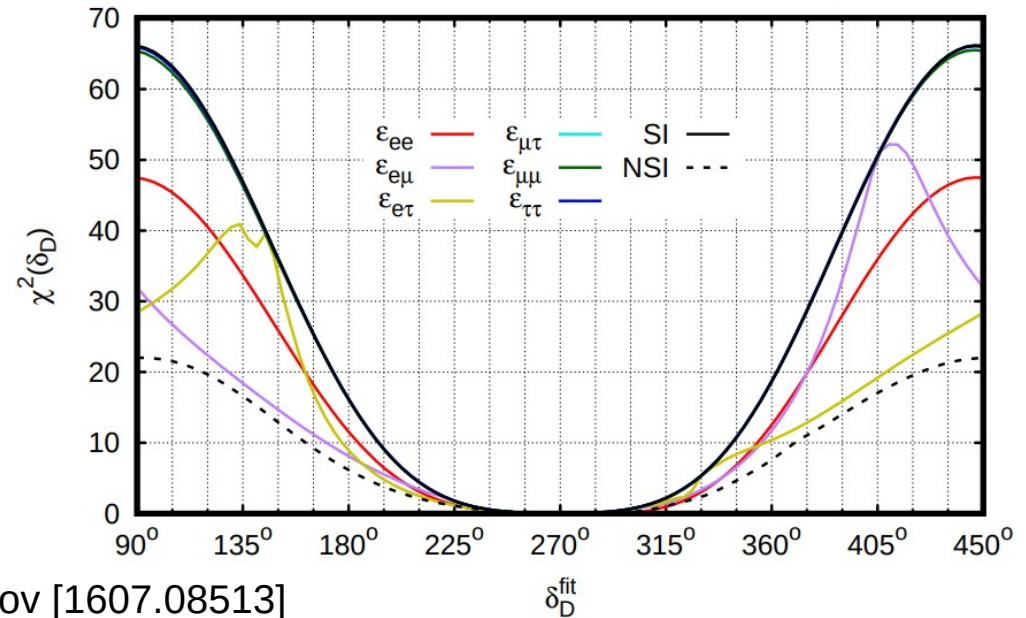
The effect of NSI on the CP sensitivity at μ SK [$\delta_D^{\text{true}} = -90^\circ$]



The effect of NSI on the CP sensitivity at T2K+ μ SK [$\delta_D^{\text{true}} = -90^\circ$]



The effect of NSI on the CP sensitivity at ν T2K+ μ SK [$\delta_D^{\text{true}} = -90^\circ$]

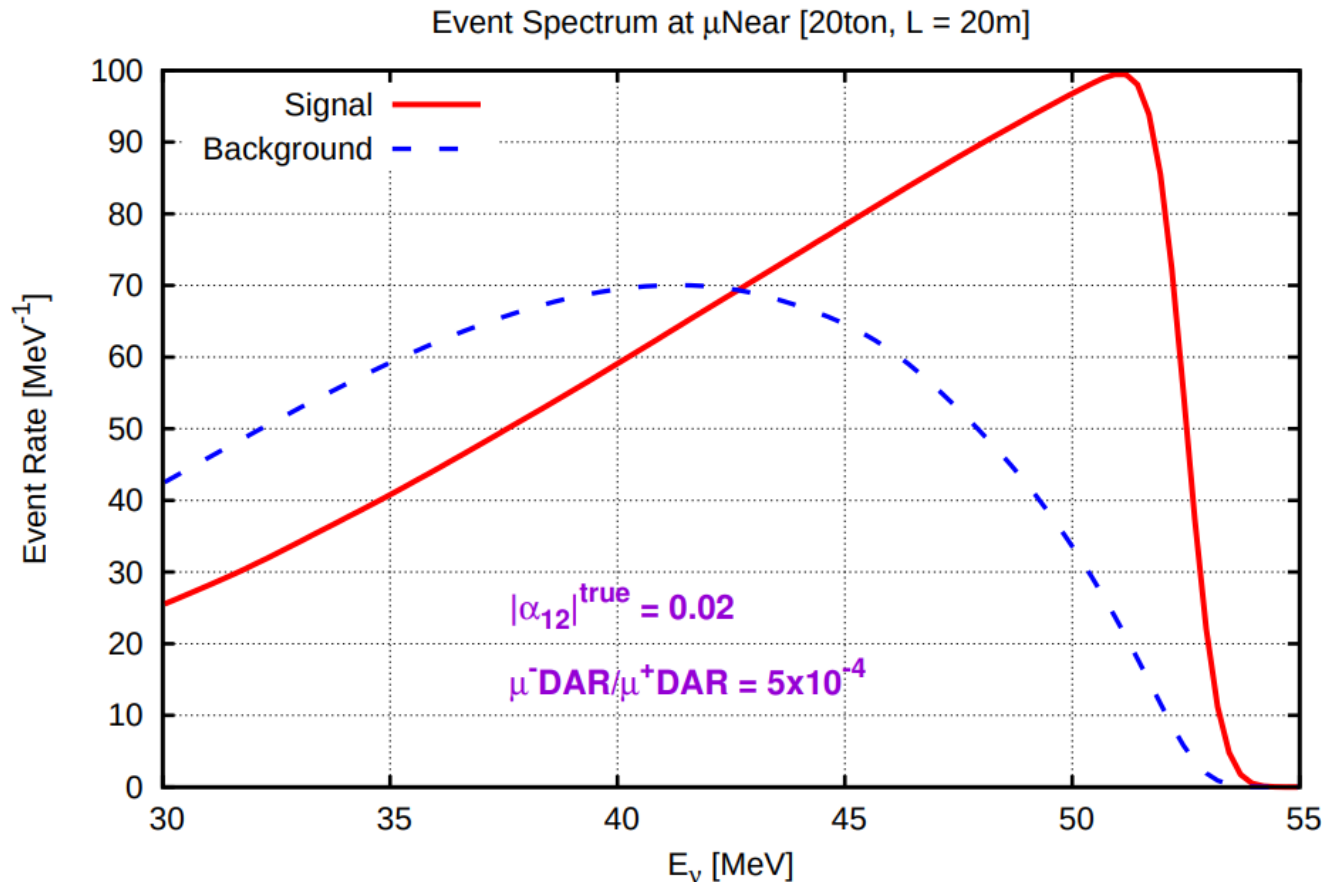


SFG & Alexei Smirnov [1607.08513]

NUM Zero Distance Effect

$$N = N^{NP} U = \begin{pmatrix} \alpha_{11} & 0 & 0 \\ |\alpha_{21}| e^{i\phi} & \alpha_{22} & 0 \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} U$$

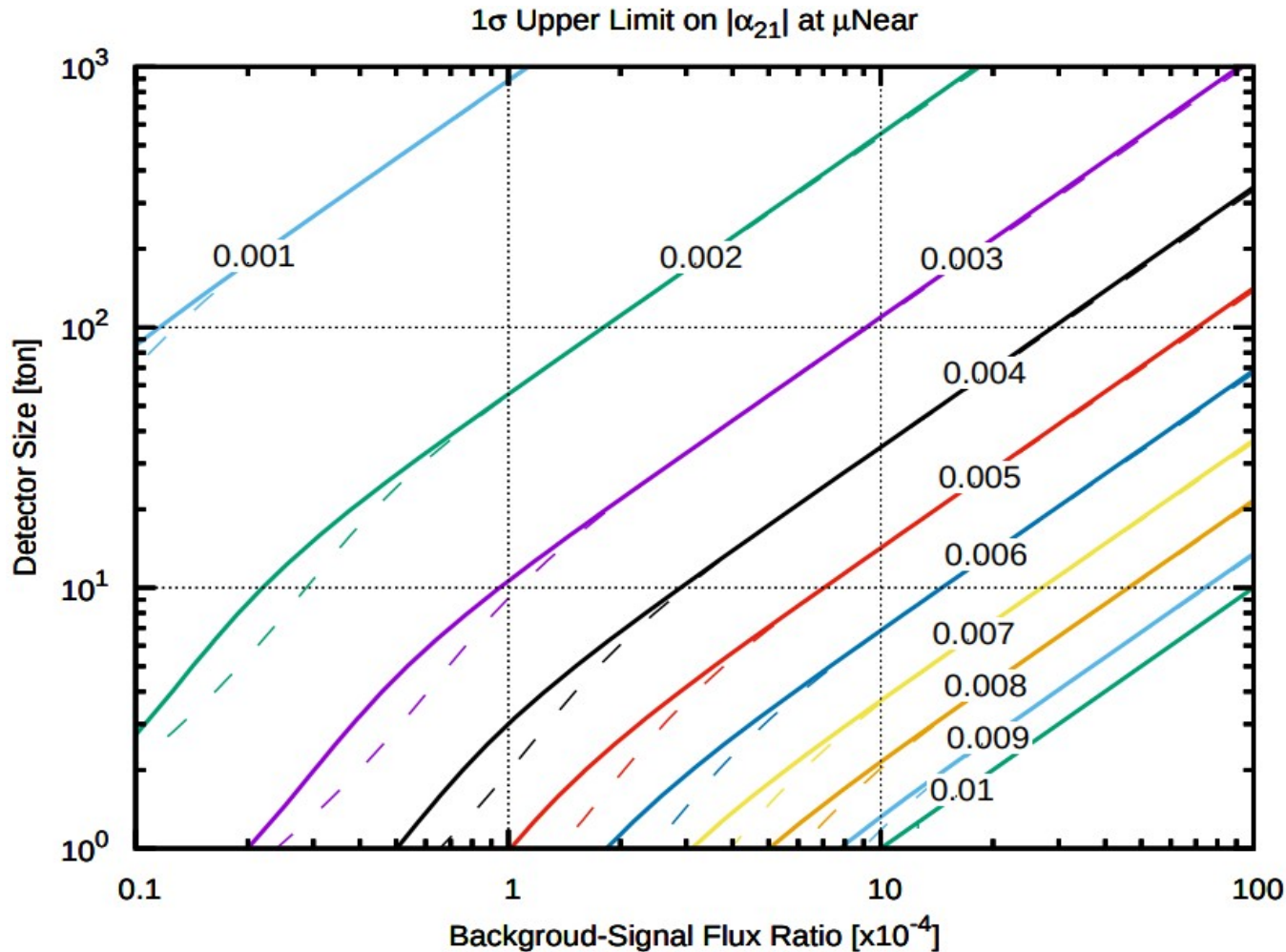
$$P_{\mu e}^{NP}(L \rightarrow 0) = \alpha_{11}^2 |\alpha_{21}|^2 P_{ee} \approx \alpha_{11}^2 |\alpha_{21}|^2 \approx |\alpha_{21}|^2$$



Transition @ $L \rightarrow 0$

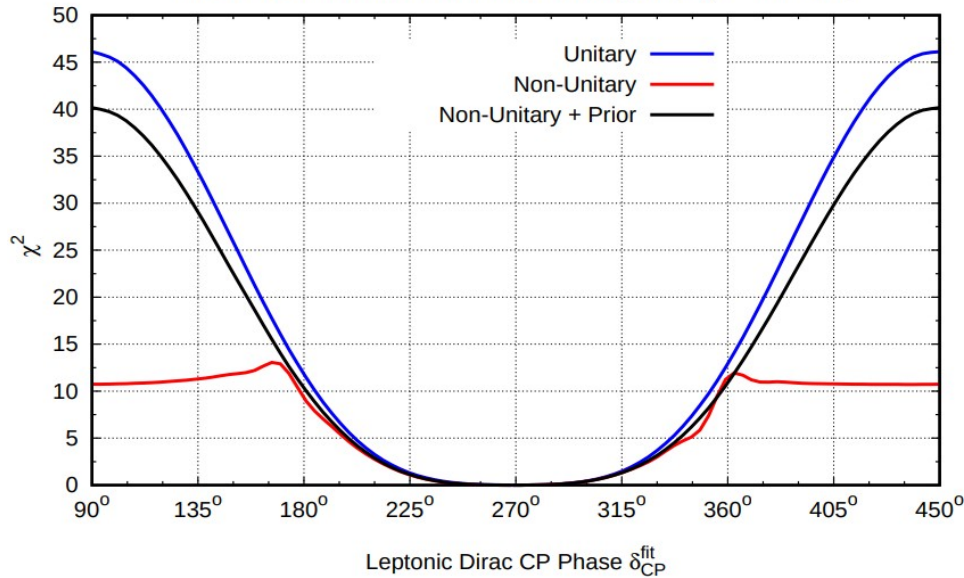
SFG, Pasquini, Tortola
& Valle [1605.01670]

$$P_{\mu e}^{NP}(L \rightarrow 0) = \alpha_{11}^2 |\alpha_{21}|^2 P_{ee} \approx \alpha_{11}^2 |\alpha_{21}|^2 \approx |\alpha_{21}|^2$$

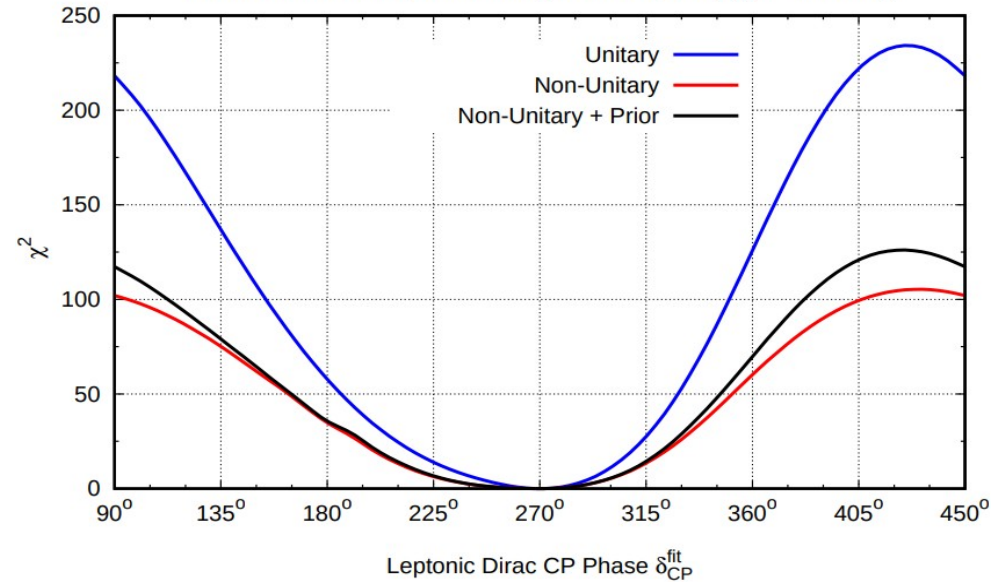


SFG, Pasquini, Tortola & Valle [1605.01670]

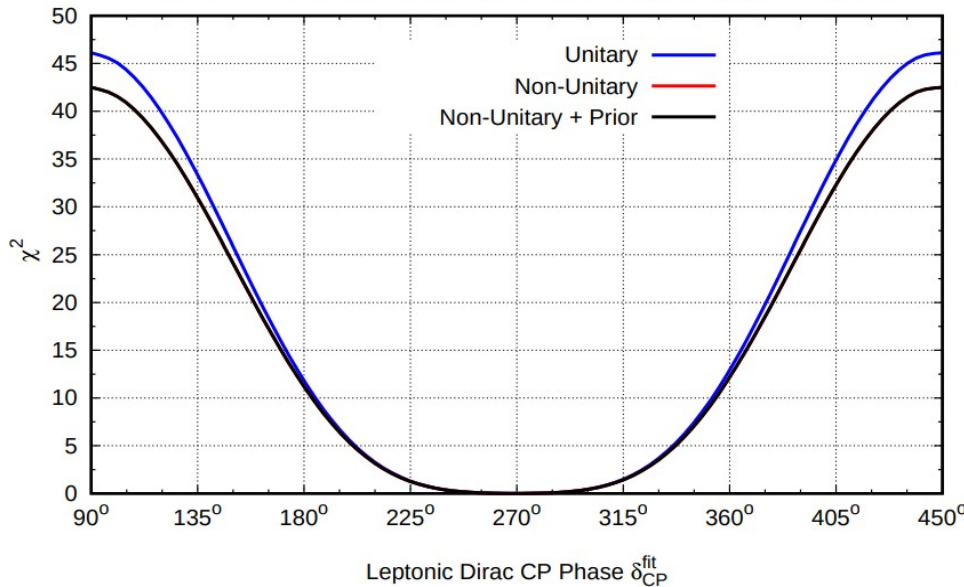
The effect of including non-unitarity at T2K+ μ SK [$\delta_{CP}^{true} = -90^\circ$, NH]



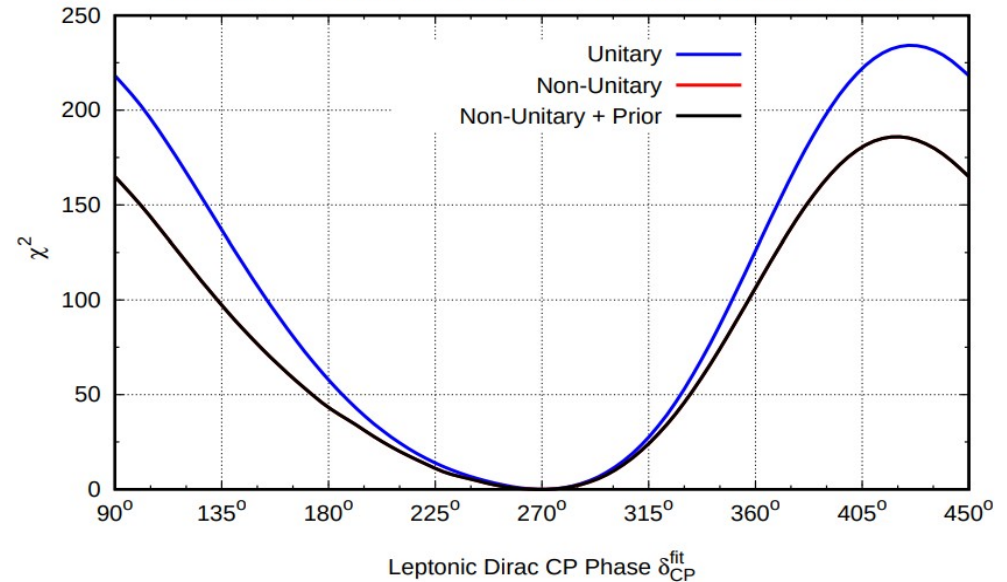
The effect of including non-unitarity at T2HK+ μ HK [$\delta_{CP}^{true} = -90^\circ$, NH]



The effect of including non-unitarity at T2K+ μ SK+ μ Near [$\delta_{CP}^{true} = -90^\circ$, NH]

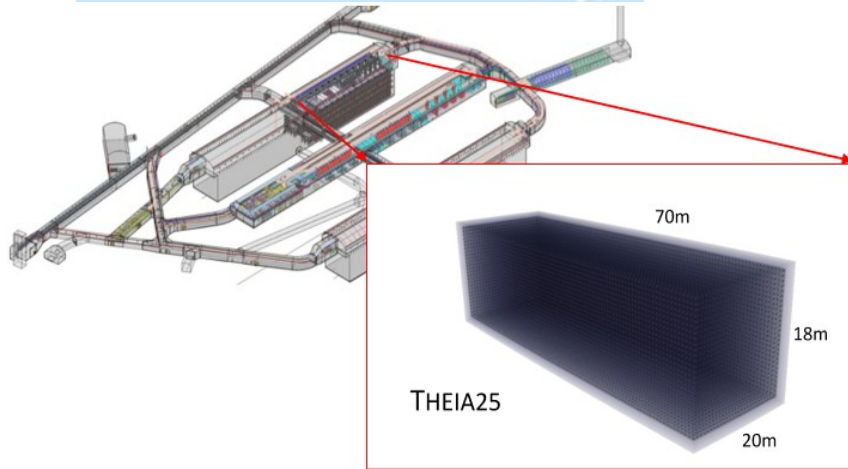


The effect of including non-unitarity at T2HK+ μ HK+ μ Near [$\delta_{CP}^{true} = -90^\circ$, NH]



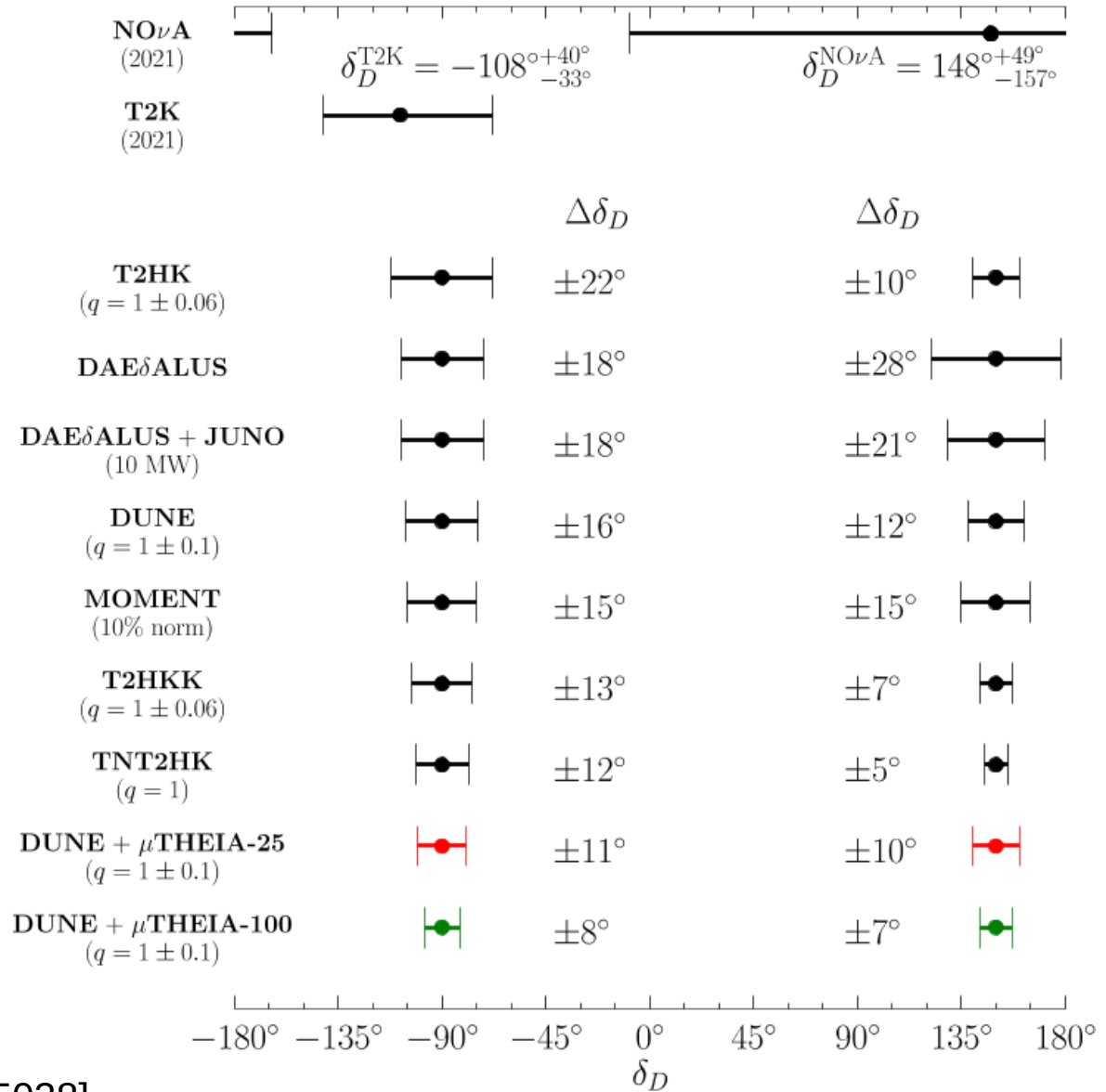
SFG, Pasquini, Tortola & Valle [1605.01670]

μ THEIA+DUNE



THEIA [1911.03501, 2202.12839]

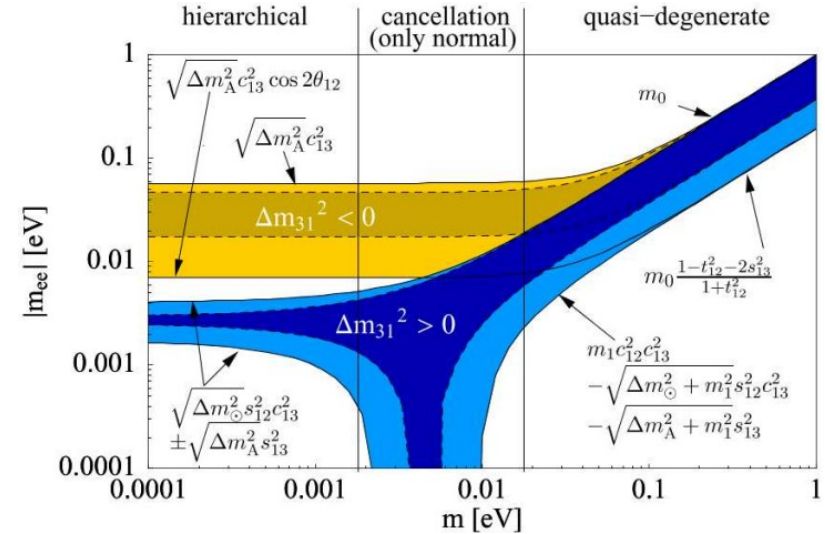
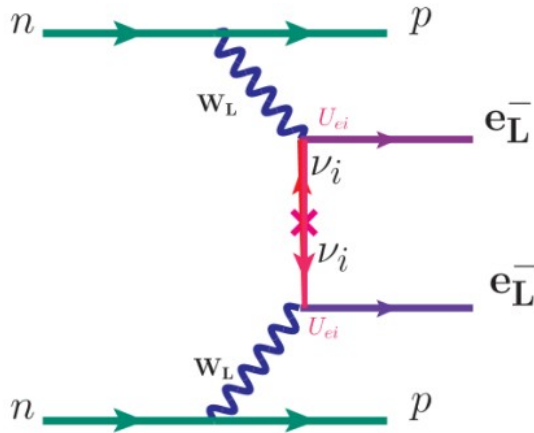
SFG, Kong, Pasquini [2202.05038]



- CP Violation in the Neutrino Sector
- Dirac CP Phase & New Physics
- Accelerator + μ DAR for better CP measurement
- **Majorana CP Phase**
- Neutrino Electromagnetic Properties

Neutrinoless Double Beta Decay

- Mediated by **Majorana Neutrino** + **Lepton # Violation**



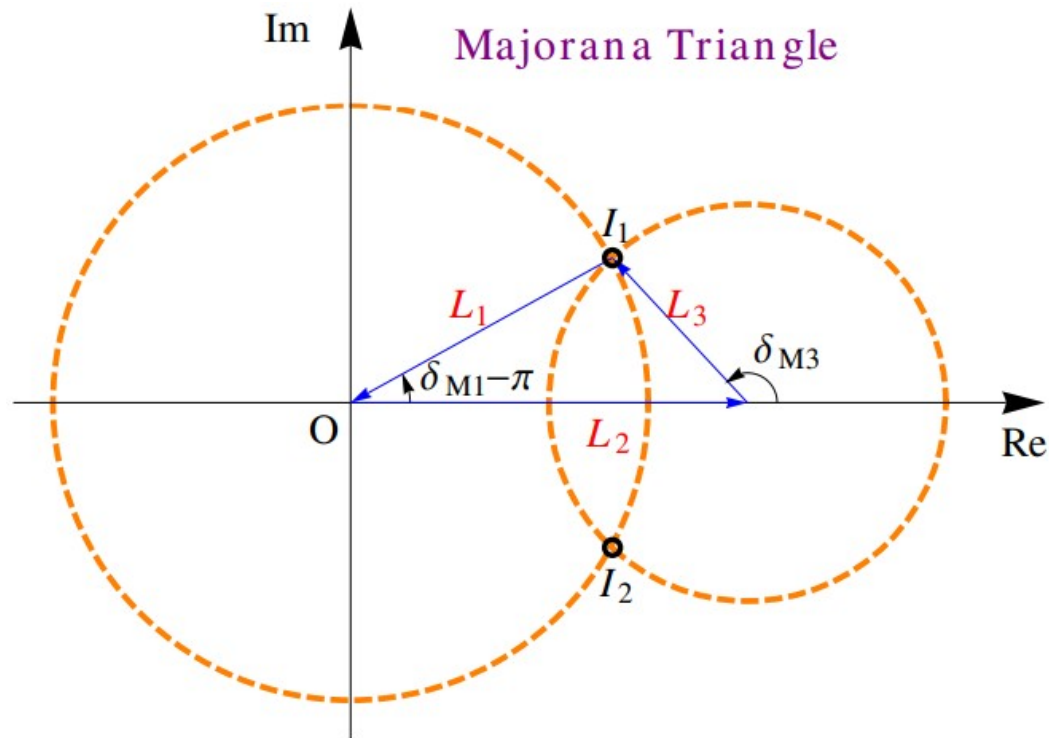
- Mass Suppression**

$$\mathcal{M} \propto \sum_i U_{ei} \frac{i}{p \not{=} - m_i} U_{ei} \approx \sum_i U_{ei} \frac{m_i}{p^2} U_{ei}$$

- Effective Mass**

$$\langle m \rangle_{ee} \equiv \left| \sum_i m_i U_{ei}^2 \right| = \left| c_s^2 c_r^2 m_1 e^{i\delta_{M1}} + s_s^2 c_r^2 m_2 + s_r^2 m_3 e^{i\delta_{M3}} \right|$$

Majorana Triangle



$$\langle m \rangle_{ee} \equiv \vec{L}_1 + \vec{L}_2 + \vec{L}_3$$

$$\vec{L}_1 \equiv m_1 U_{e1}^2 = m_1 c_r^2 c_s^2 e^{i\delta_{M1}},$$

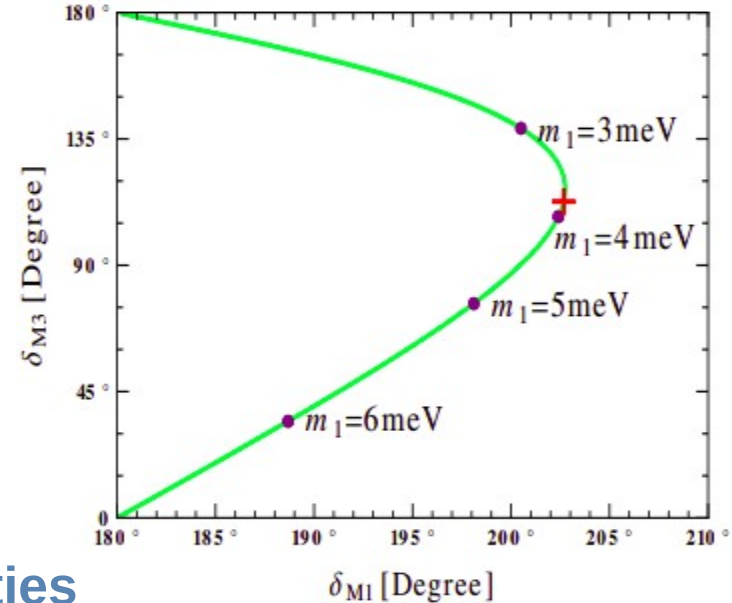
$$\vec{L}_2 \equiv m_2 U_{e2}^2 = \sqrt{m_1^2 + \Delta m_s^2} c_r^2 s_s^2,$$

$$\vec{L}_3 \equiv m_3 U_{e3}^2 = \sqrt{m_1^2 + \Delta m_a^2} s_r^2 e^{i\delta_{M3}}.$$

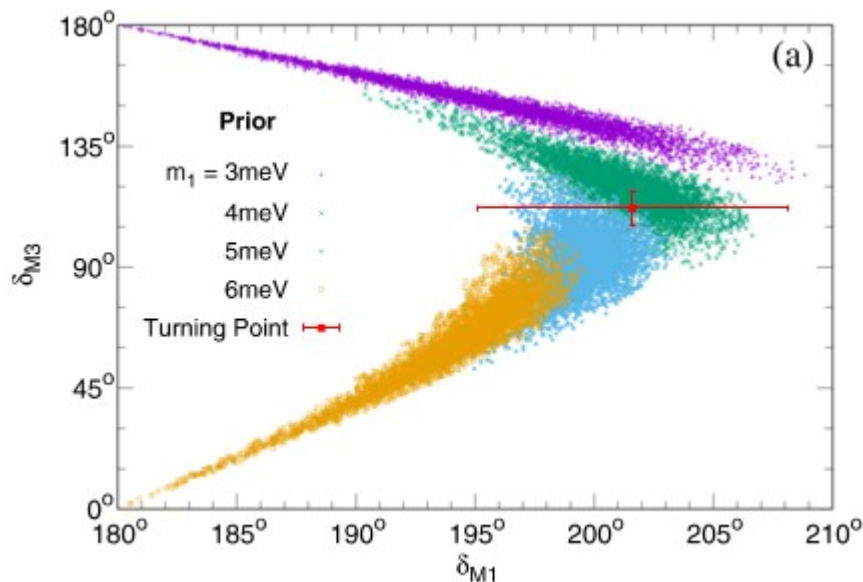
Catching 2 Majorana CP Phases

$$\cos \delta_{M1} = - \frac{m_1^2 c_r^4 c_s^4 + m_2^2 c_r^4 s_s^4 - m_3^2 s_r^4}{2m_1 m_2 c_r^4 c_s^2 s_s^2}$$

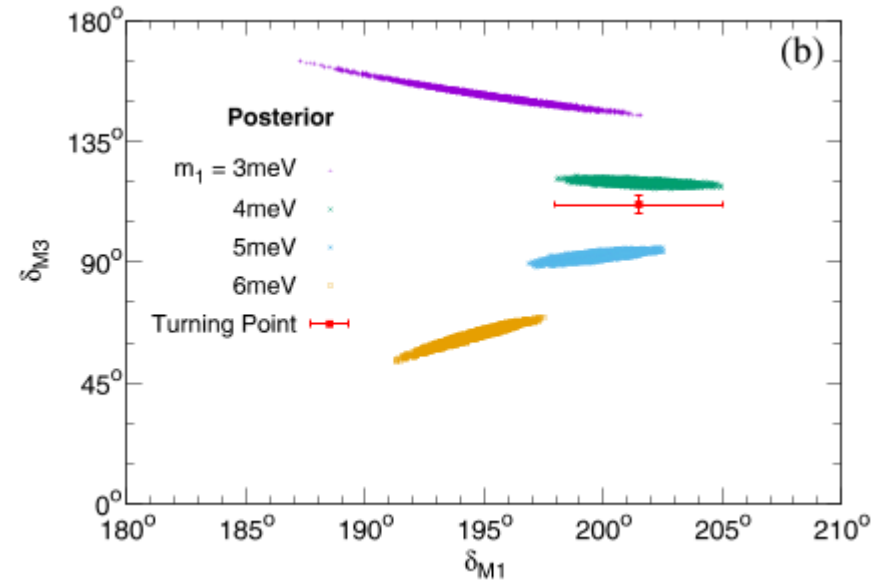
$$\cos \delta_{M3} = + \frac{m_1^2 c_r^4 c_s^4 - m_2^2 c_r^4 s_s^4 - m_3^2 s_r^4}{2m_2 m_3 c_r^2 s_r^2 s_s^2}$$



- JUNO can significantly reduce uncertainties

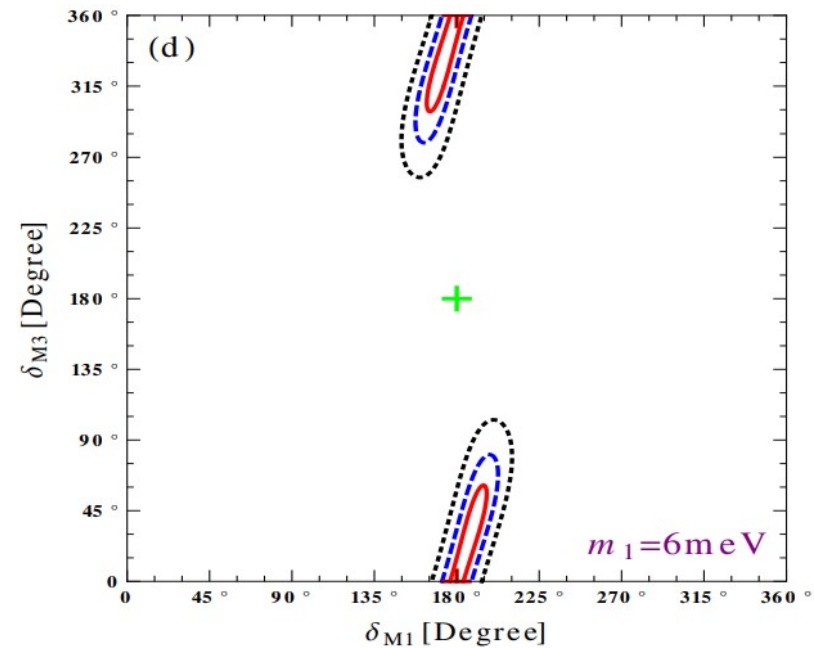
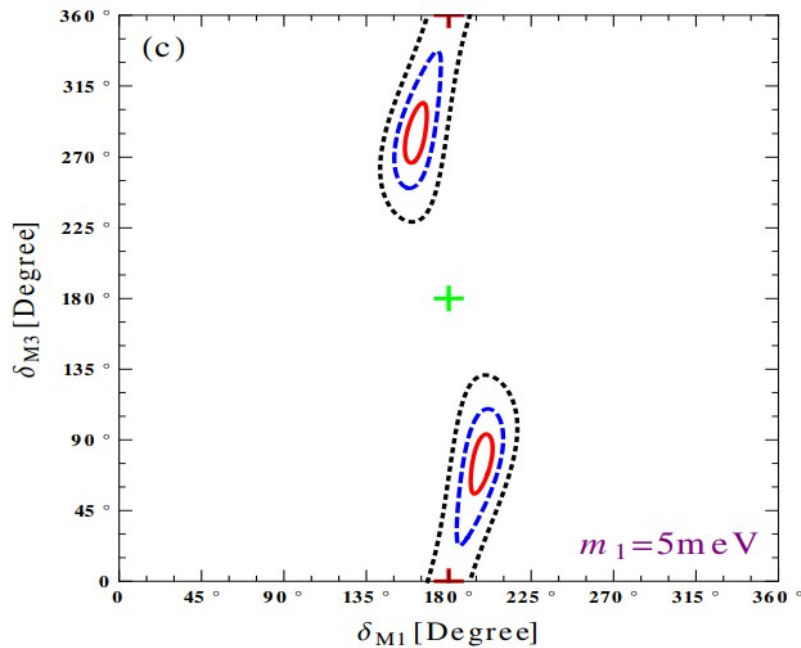
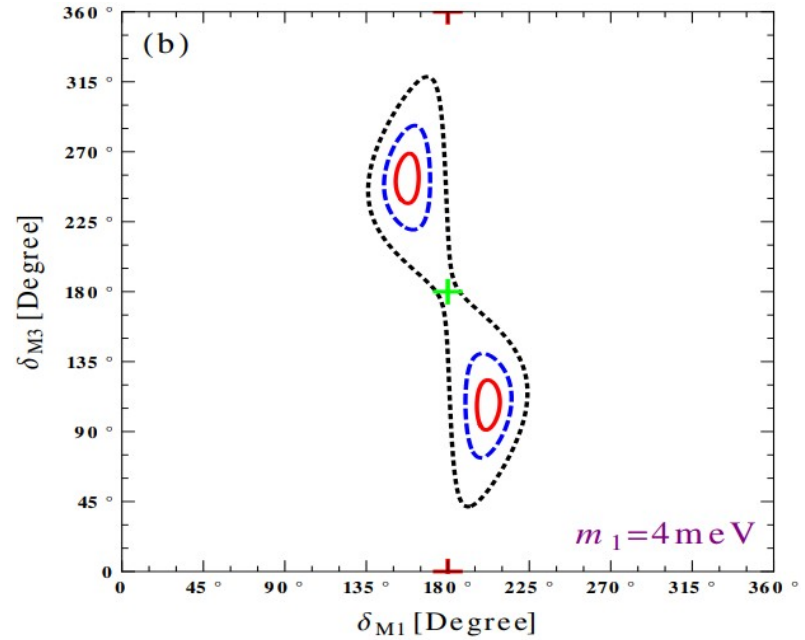
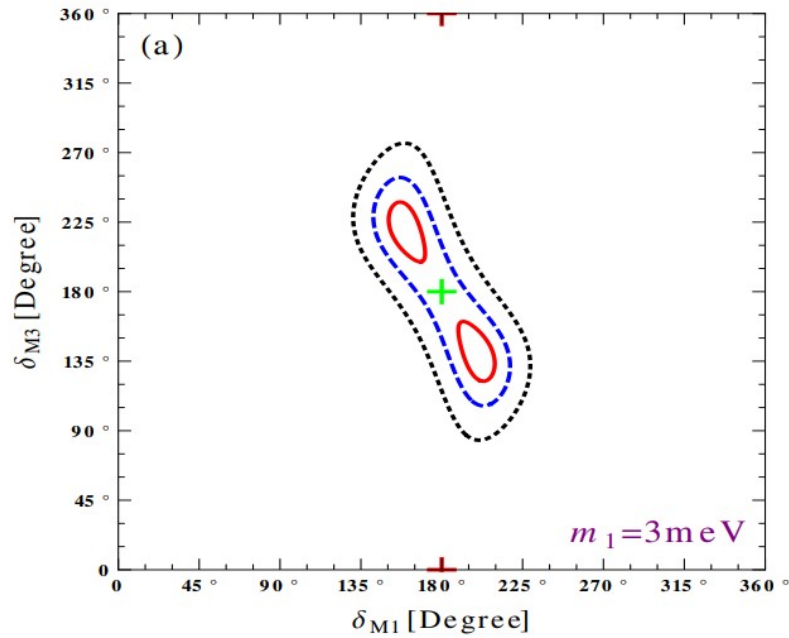


SFG & Manfred Lindner, [1608.01618]

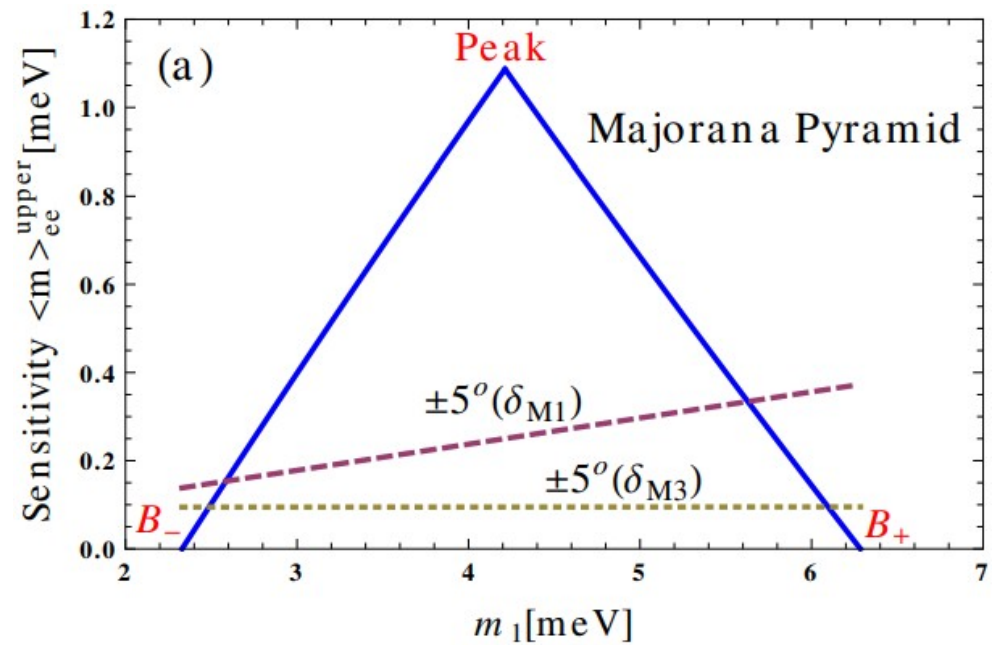
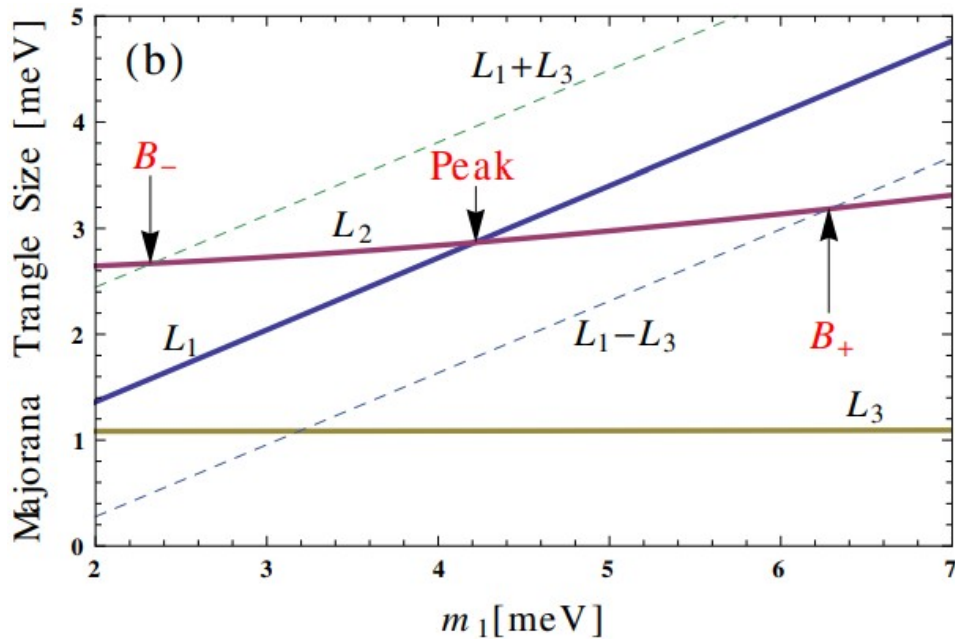
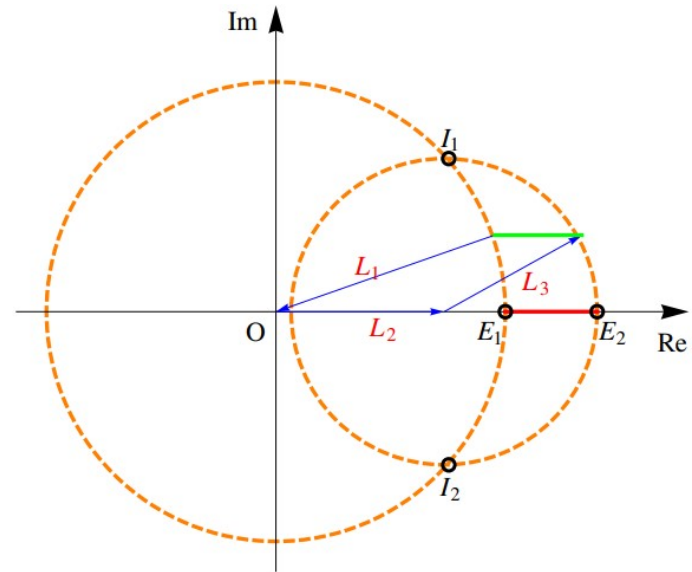
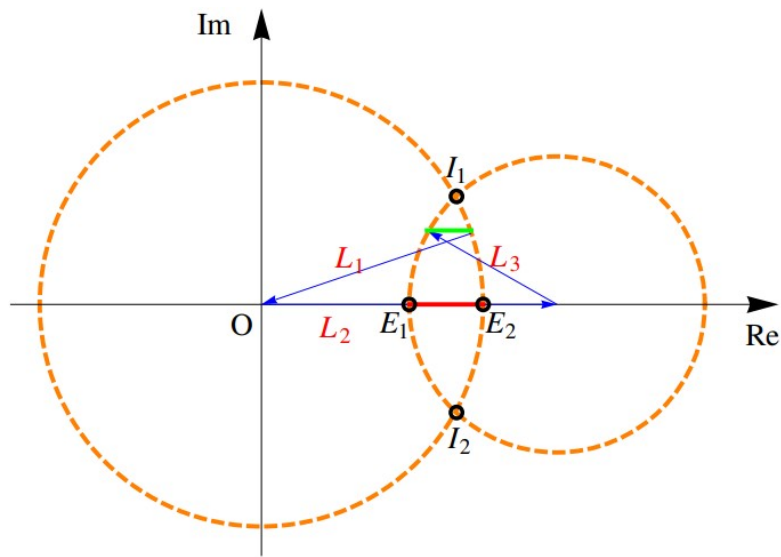


see also SFG & Rodejohann [1507.05514]

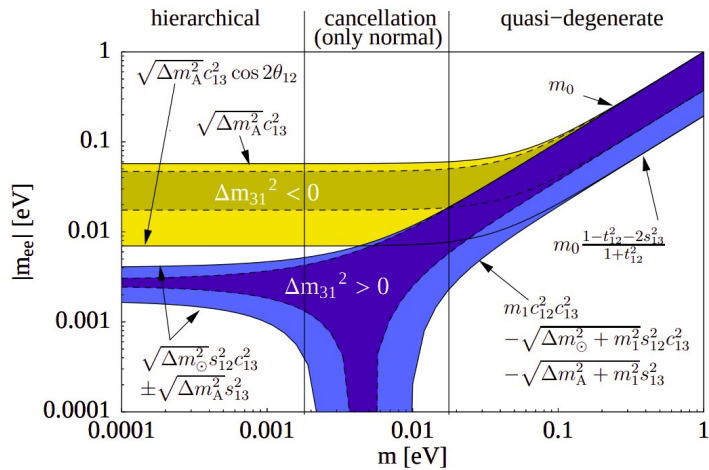
Uncertainty from m_{ee}



Majorana Pyramid



1meV frontier of $0\nu 2\beta$



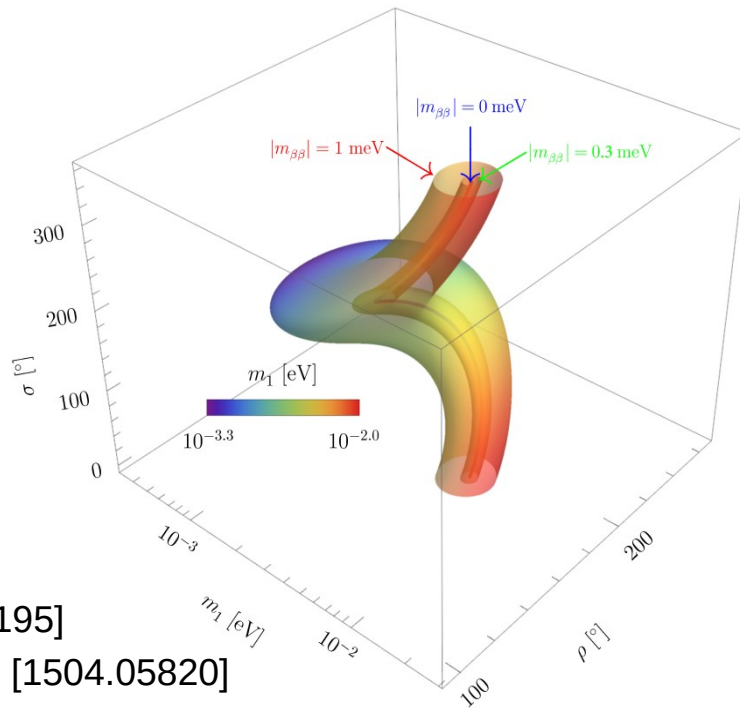
- Nonzero → only 1 Majorana CP

$$|m_{ee}| = f$$

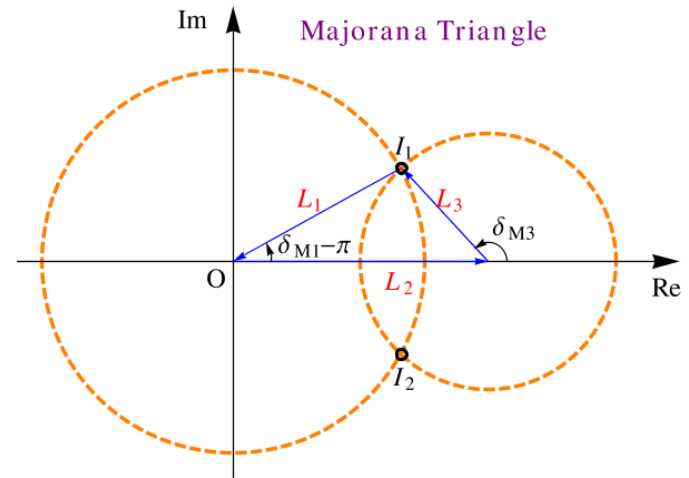
- Zero → 2 Majorana CP phases

$$|m_{ee}| = 0 \Rightarrow \mathbb{R}(m_{ee}) = \mathbb{I}(m_{ee}) = 0$$

$$|m_{ee}| < f \Rightarrow \mathbb{R}(m_{ee}) < f \quad \mathbb{I}(m_{ee}) < f$$



$$m_{ee} = c_r^2 c_s^2 m_1 e^{i\tilde{\delta}_{M1}} + c_r^2 s_s^2 m_2 + s_r^2 m_3 e^{i\tilde{\delta}_{M3}}$$



SFG & Manfred Lindner, [1608.01618]

Xing [hep-ph/0305195]

Xing, Zhao & Zhou [1504.05820]

Xing & Zhao [1612.08538]

Cao, Huang, Li, Wang, Wen, Xing, Zhao & Zhou [1908.08355]

Prey of Leptonic CP Phases



- CP Violation in the Neutrino Sector
- Dirac CP Phase & New Physics
- Accelerator + μ DAR for better CP measurement
- Majorana CP Phase
- **Neutrino Electromagnetic Properties**

- Being neutral, neutrino does not have charge
- But can have electric & magnetic moments

$$H_M = \bar{\nu} \left[-f_M(q^2) i \sigma_{\mu\nu} q^\nu + f_E(q^2) \sigma_{\mu\nu} q^\nu \gamma_5 \right] \nu A^\mu(q)$$

Magnetic Moment Electric Moment

1) Majorana neutrinos

$$(\mu_\nu)_{ij} = -(\mu_\nu)_{ji} \text{ and } (\epsilon_\nu)_{ij} = -(\epsilon_\nu)_{ji}$$

2) Dirac neutrinos

nonzero diagonal $(\epsilon_\nu)_{ii}$ also indicates CP violation

- Possible tests

$$(\mu_{\alpha\beta}^{\text{eff}})^2 \equiv \sum_j \left| \sum_k U_{\alpha k}^* [(\mu_\nu)_{jk} - i(\epsilon_\nu)_{jk}] \right|^2$$

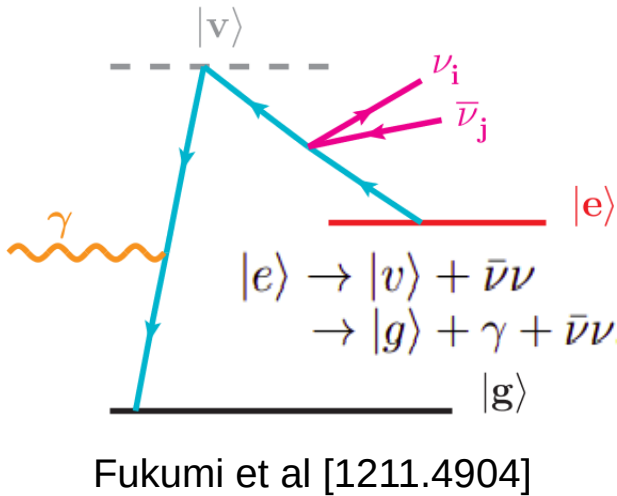
Scattering & recoil

$$(\mu_\nu^\odot)^2 \equiv \sum_{ij} |(\mu_\nu)_{ij}|^2 + |(\epsilon_\nu)_{ij}|^2$$

Stellar processes

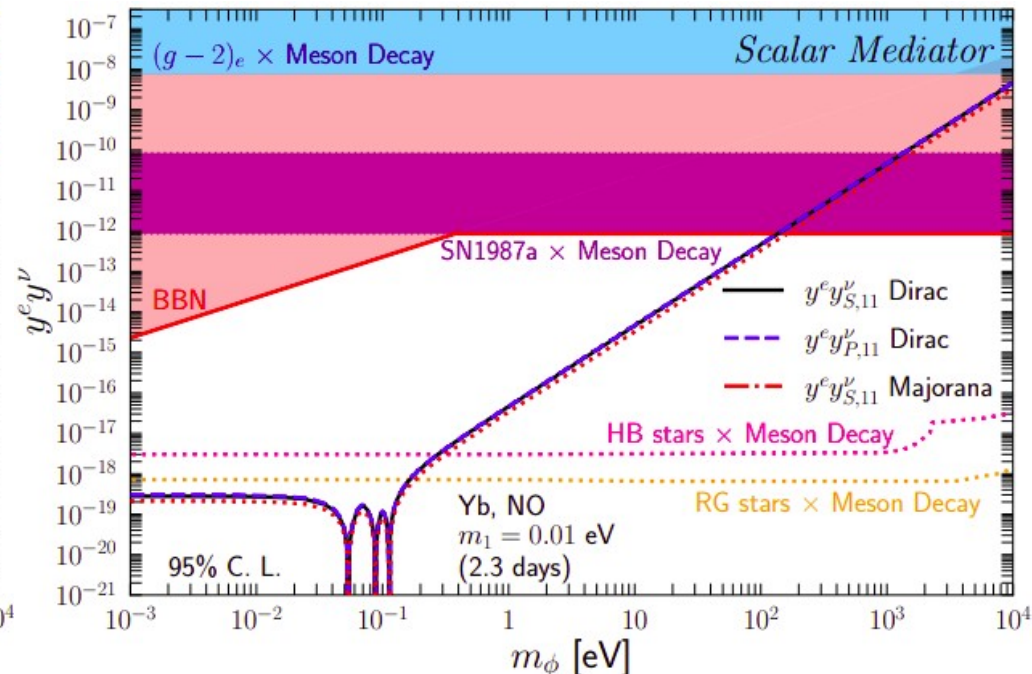
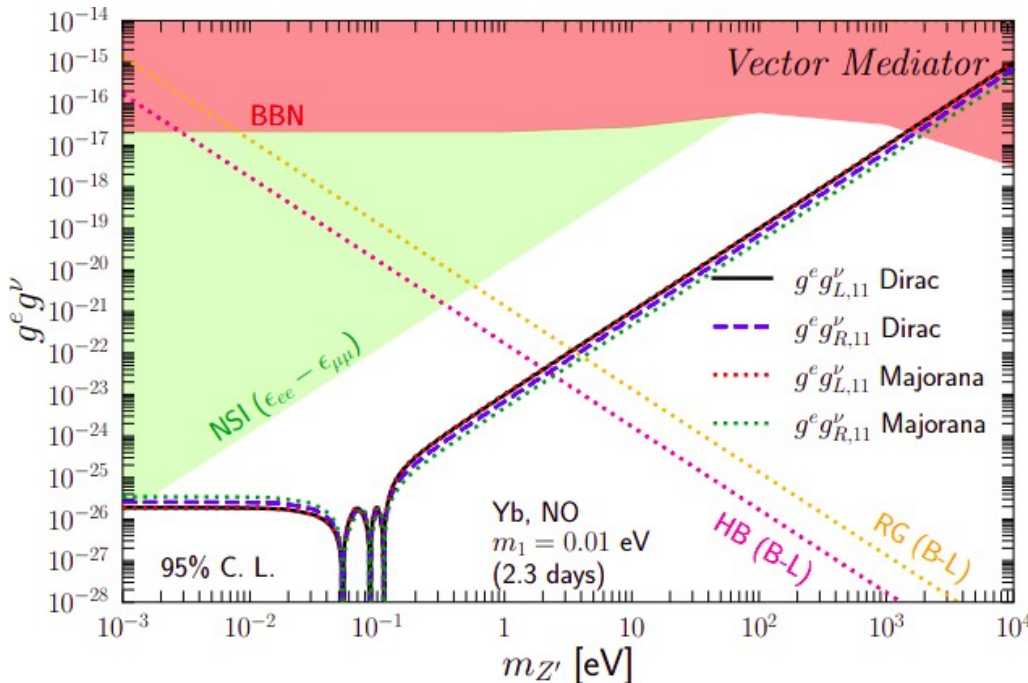
Giunt & Studenikin [1403.6344]

Light mediator @ ν Pair Emission



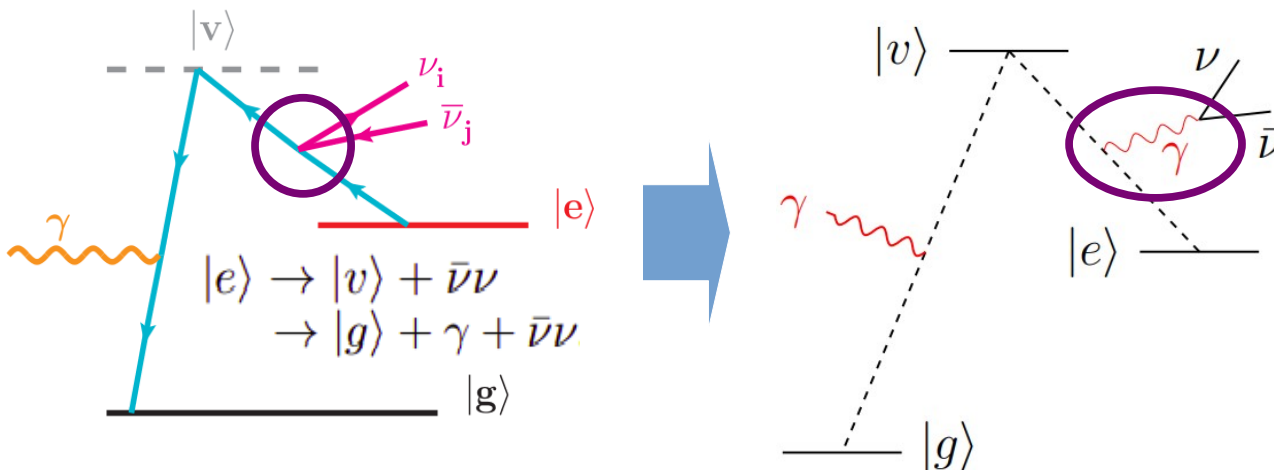
- 1) **Mass**
- 2) **CP phase**
- 3) **Dirac vs Majorana**
- 4) **Unitarity**
- 5) **Sterile neutrino**
- 6) **Axion**
- 7) **Dark photon**

- Song et al [1510.00421]; Zhang & Zhou [1604.08008]
 Yoshimura & Sasao [1506.08003]
 Dinh, Petcov, Sasao, Tanaka & Yoshimura [1209.4808]
 Huang, Sasao, Xing & Yoshimura [1904.10366]
 Dinh & Petcov [1411.7459]
 Huang & Zhou [1905.00367]
 Bhoonah, Bramante & Song [1909.07387]



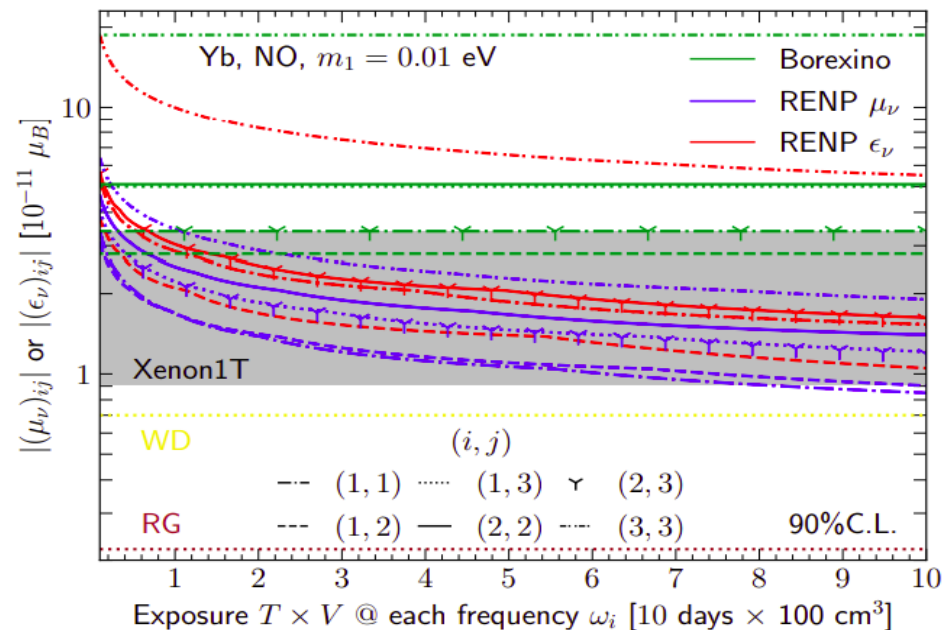
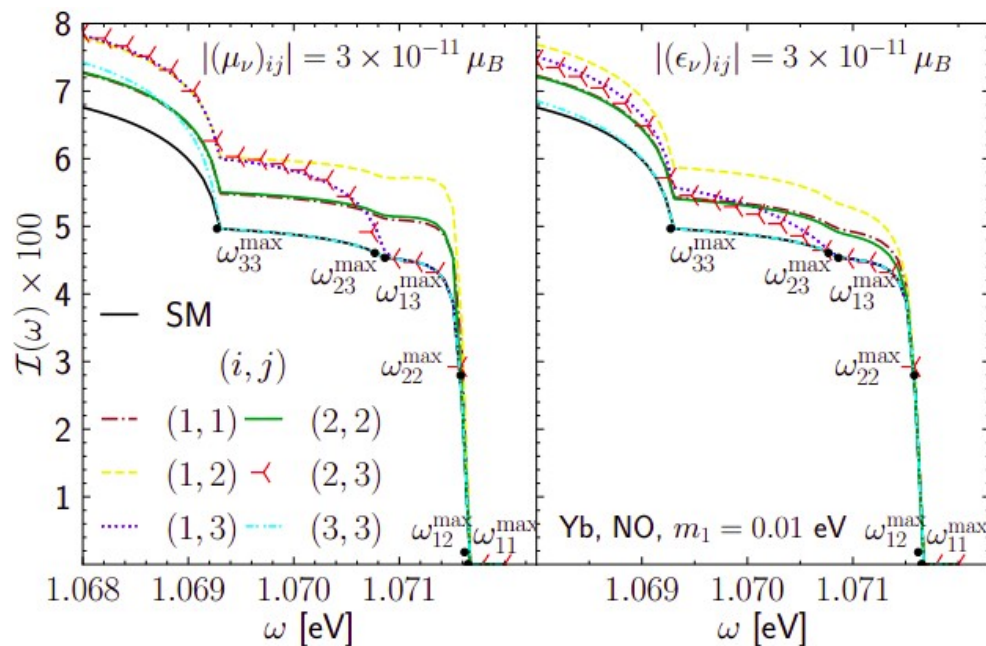
SFG & Pasquini [2110.03510]

ν Moments @ ν Pair Emission



$$\omega_{ij} = \frac{E_{eg}}{2} - \frac{(m_i + m_j)^2}{2E_{eg}}$$

$$\mathcal{I}_E^M \equiv \sum_{ij} \left(\frac{\mu_B}{G_F} \right)^2 \frac{\omega^2 \Delta_{ij} \Theta(\omega - \omega_{ij}^{\max})}{9(E_{vg} - \omega)^2} \times \left[1 + \frac{(m_i \pm m_j)^2 \pm 4m_i m_j}{q^2} - 2 \left(\frac{\Delta m_{ji}^2}{q^2} \right)^2 \right]$$

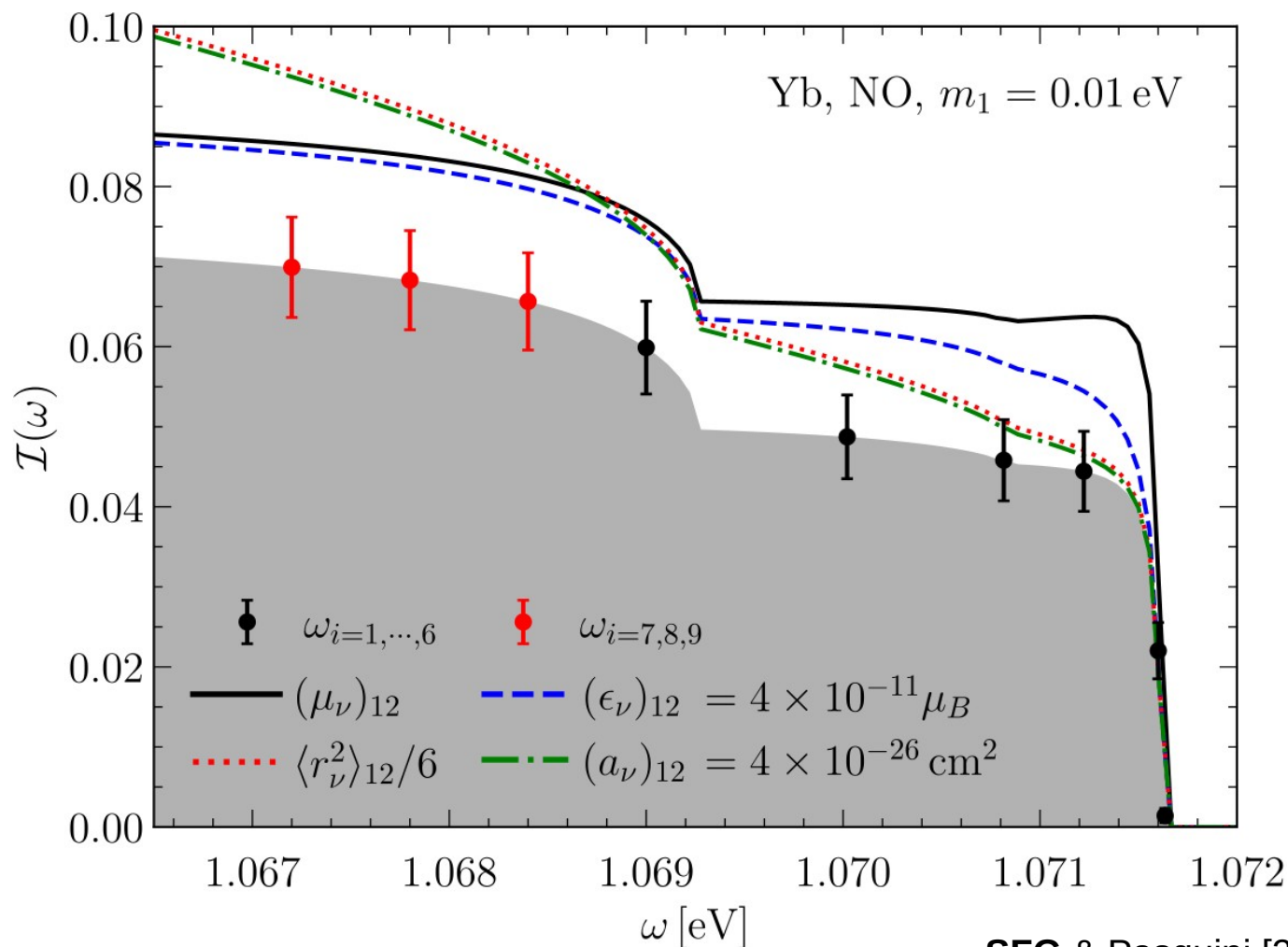


Can disentangle EDM & MDM elements!

SFG & Pasquini [2206.11717]

ν EM Properties @ ν Pair Emission

$$H_{EM} = \bar{\nu} \left[-i\sigma_{\mu\nu}q^\nu (\mu_\nu + i\epsilon_\nu\gamma_5) + \left(\gamma_\mu - \frac{q_\mu\not{q}}{q^2} \right) q^2 \left(\frac{\langle r_\nu^2 \rangle}{6} + a_\nu\gamma_5 \right) \right] \nu A^\mu(q),$$



SFG & Pasquini [2306.12953]

- **CP Violation in the Neutrino Sector**

Important consequence for explaining matter-antimatter asymmetry.

- **Dirac CP Phase & New Physics**

Needs to exclude multiple theoretical uncertainties.

- **Accelerator + μ DAR for better CP measurement**

Combination of high- and low-energy neutrino sources.

- **Majorana CP Phases**

No-loss for $0\nu 2\beta$ measurements.

Thank You

- **Neutrino Electromagnetic Properties**

Atomic RENP (radiative emission of neutrino pair) is a unique probe.



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Tsung-Dao Lee Institute

Backup Slides

- Oscillation probabilities @ Accelerator Neutrino Exps

$$P_{\nu_\mu \rightarrow \nu_e} \approx 4s_a^2 c_r^2 s_r^2 \sin^2 \phi_{31} \\ - 8c_a s_a c_r^2 s_r c_s s_s \sin \phi_{21} \sin \phi_{31} [\cos \delta_D \cos \phi_{31} \pm \sin \delta_D \sin \phi_{31}]$$

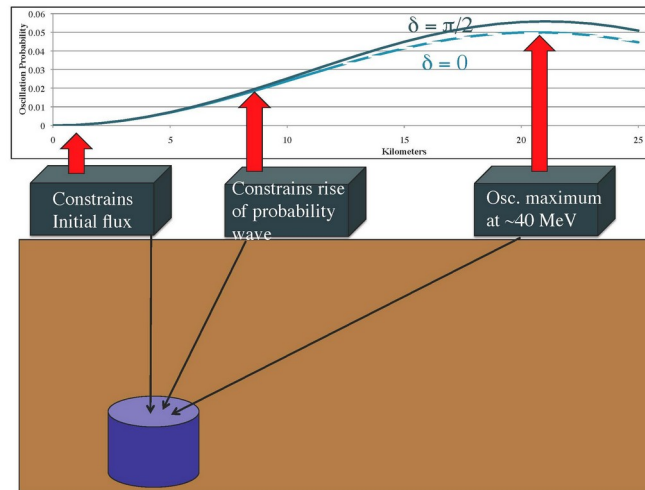
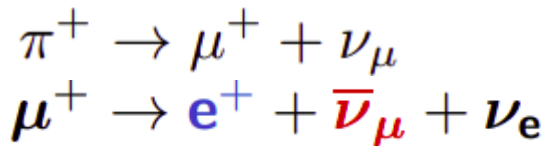
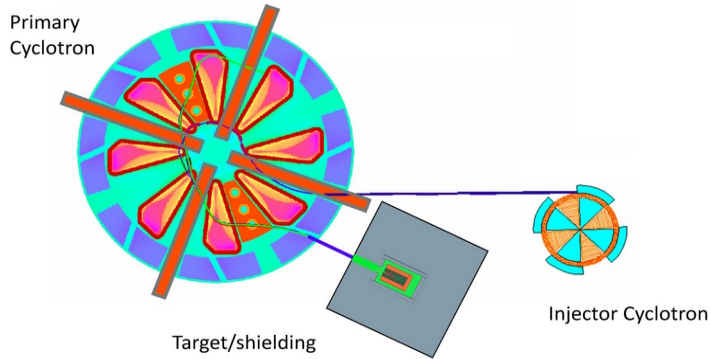
for ν & $\bar{\nu}$, respectively. [$\phi_{ij} \equiv \frac{\delta m_{ij}^2 L}{4E_\nu}$]

- Run both ν & $\bar{\nu}$ modes @ first peak [$\phi_{31} = \frac{\pi}{2}, \phi_{21} = \alpha \frac{\pi}{2}$], $\alpha = \frac{\delta M_{21}^2}{|\delta M_{31}^2|} \sim 3\%$

$$P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e} + P_{\nu_\mu \rightarrow \nu_e} = 2s_a^2 c_r^2 s_r^2,$$

$$P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e} - P_{\nu_\mu \rightarrow \nu_e} = \alpha \pi \sin(2\theta_s) \sin(2\theta_r) \sin(2\theta_a) \cos \theta_r \sin \delta_D.$$

中微子、反中微子的振荡概率之间的差别 \rightarrow 测量 CP 相角 δ_D 。



戴达罗斯 (希腊神话中的建筑师、工匠)

Disadvantages:

- The scattering lepton from IBD @ low energy is **isotropic**.
- **Cannot** distinguish $\bar{\nu}_e$ from different sources $\bar{\nu}_e + p \rightarrow e^+ n$
- Baseline **cannot be measured**.
- Cyclotrons **cannot** run simultaneously (20~25% duty factor).
- **Large** statistical uncertainty.
- **Higher intensity** is necessary.
- **Expensive** & Technically **challenging**.

Conrad & Shaevitz [0912.4079]

Agarwalla, Huber, Link & Mohapatra [1005.4055]

DAEdALUS [1006.0260, 1307.2949]

1 μ DAR source + **2** detectors

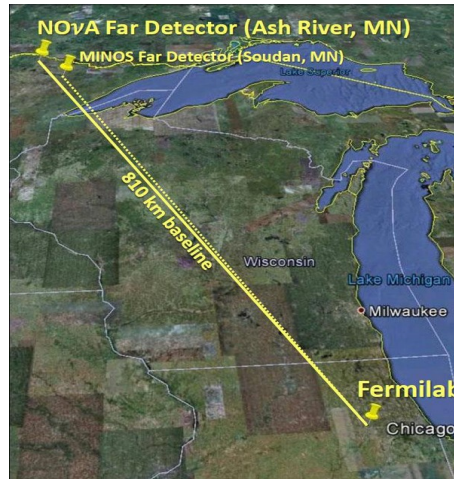
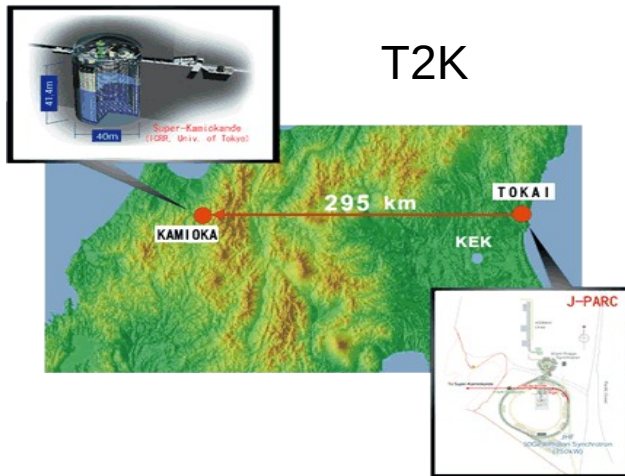
Advantages

- Full (**100%**) duty factor!
- **Lower** intensity: \sim **9mA** [\sim **4** \times lower than DAE δ ALUS]
- Not far beyond the current state-of-art technology of cyclotron [**2.2mA** @ Paul Scherrer Institute]
- MUCH **cheaper** & technically **easier**.
 - Only one cyclotron.
 - Lower intensity.

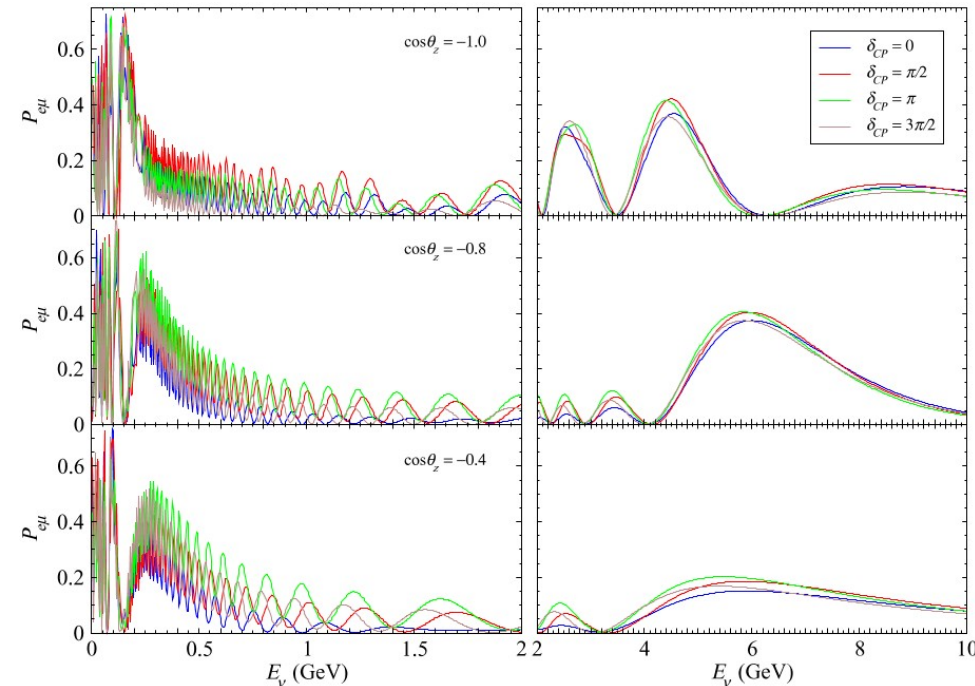
Disadvantage: A second detector!

- μ DAR with **Two Scintillators** (**μ DARTS**) [Ciuffoli, Evslin & Zhang, 1401.3977] also Smirnov, Hu, Li & Ling [1802.03677, 1808.03795]
- **Tokai 'N Toyama to(2) Kamioka** (**TNT2K**) [Evslin, Ge & Hagiwara, 1506.05023]

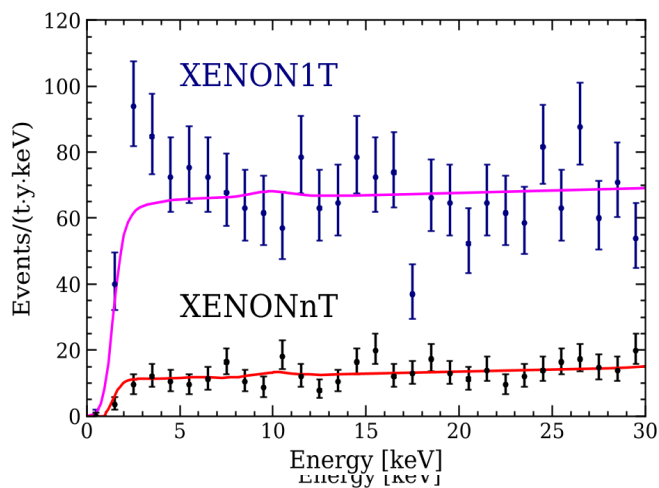
Future Accelerator Exps



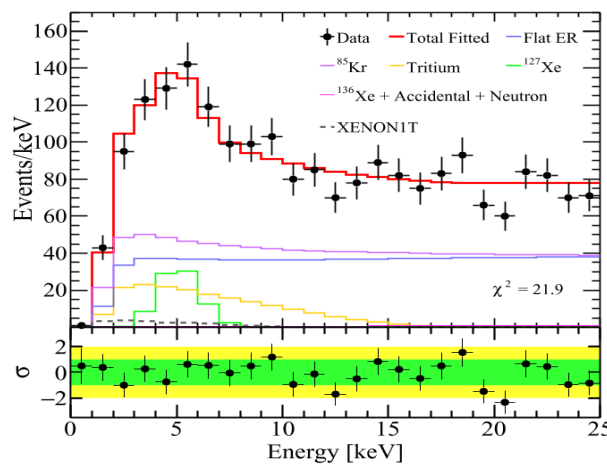
- Flux upgrade: **T2K-II**, **NOvA-II?**
- Detector upgrade: **T2HK**
- Baseline upgrade: **T2KK**, **T2KO**
- **MOMENT** Cao et al [1401.8125]
Tang et al [1909.01548]
- In china? Tang, Vihonen & Xu [2202.13595]
- Atmospheric
Super-PINGU: Razzaque & Smirnov [1406.1407]
JUNO: An et al [1507.05613]
DUNE: Kelly, Machado, Martinez-Soler, Parke & Perez-Gonzalez [1904.02751]
Super-ORCA: Hofestadt, Bruchner & Eberl [1907.12983]



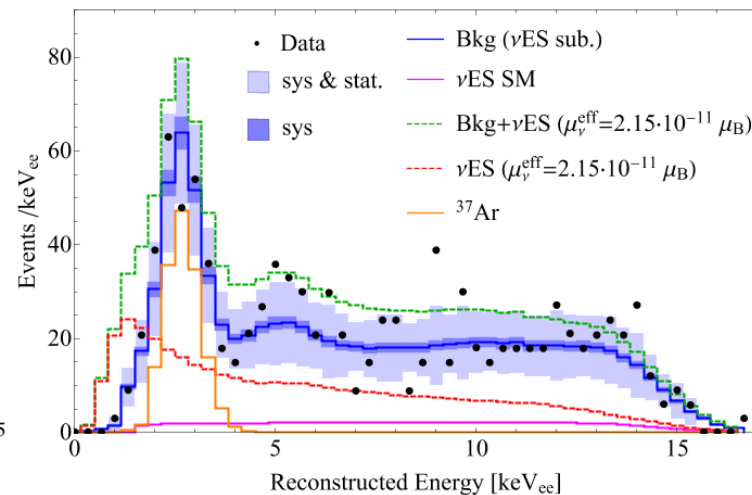
ν Magnetic Moments @ DM Exps



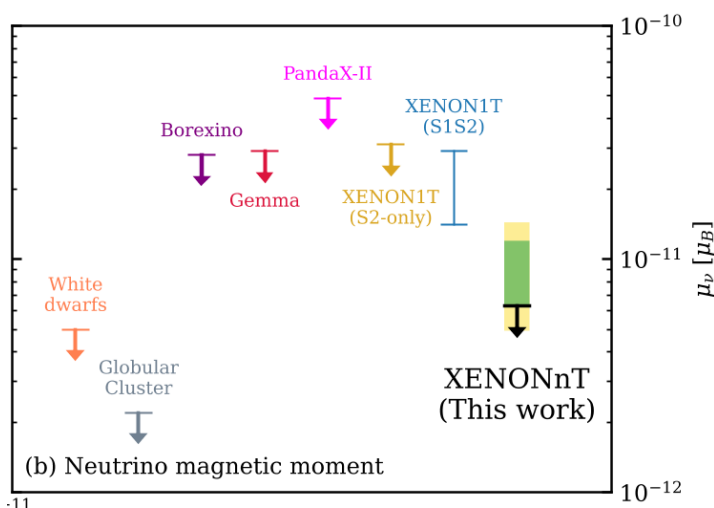
Xenon1T [2006.09721]



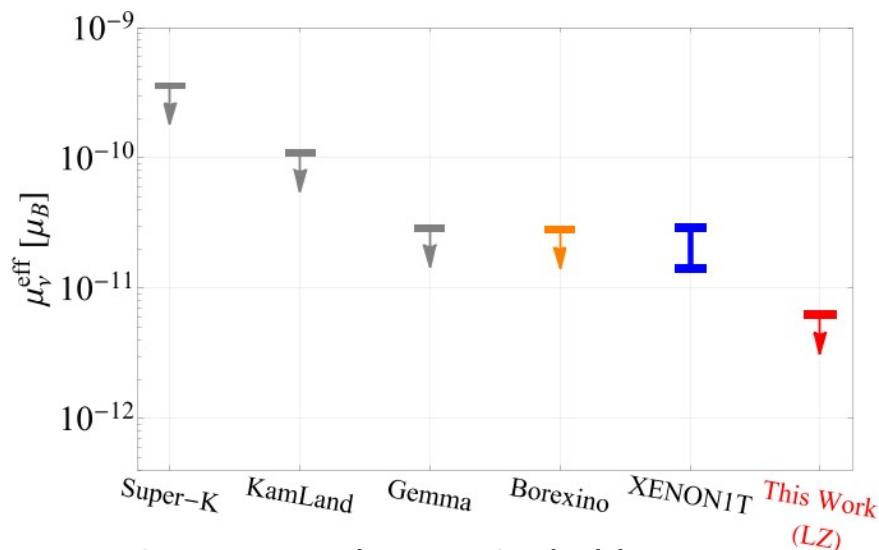
PandaX-II [2008.06485]



LZ [2207.03764]



XENONnT [2207.11330]



Corona, Bonivento, Cadeddu, Cargioli & Dordei [2207.05036]