



超子的自旋结构： 电子电子对撞/质子质子对撞自旋极化理论进展

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第一届Lambda超子自旋极化跨系统研讨会

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Outline

- **The Motivation**
- **Distribution and Fragmentation Functions**
- **Case Study:** $e^+ + e^- \rightarrow \vec{\Lambda} + X$
- **Case Study:** $\vec{l} + N \rightarrow \vec{\Lambda} + X$
- **Case Study: Λ production in proton+proton process**
- **Case Study: $\bar{\Lambda}$ Production in I+A Process**
- **Conclusions**

Our View of the Proton

with history

- Point-Like 1919
- Finite Size with Radius 1930s-1950s
- Quark Model 1960s
- QCD and Gluons 1970s
- Puzzles and Anomalies 1980s-present

- Quark Sea of the Nucleon
- Baryon-Meson Fluctuations
- Statistical Features
-

Surprises & Unknown about the Quark Structure of Nucleon: Sea

- **Spin Structure:** $\Sigma = \Delta u + \Delta d + \Delta s \approx 0.3$

“puzzle”: where is the proton’s missing spin

- **Strange Content** $\Delta s \neq 0$ $s(x) \neq \bar{s}(x)$

Brodsky & Ma, PLB381(96)317

- **Flavor Asymmetry** $\bar{u} \neq \bar{d}$

- **Isospin Symmetry Breaking** $\bar{u}_p \neq \bar{d}_n$ $\bar{d}_p \neq \bar{u}_n$

Ma, PLB 274 (92) 111

Boros, Londergan, Thomas, PRL81(98)4075

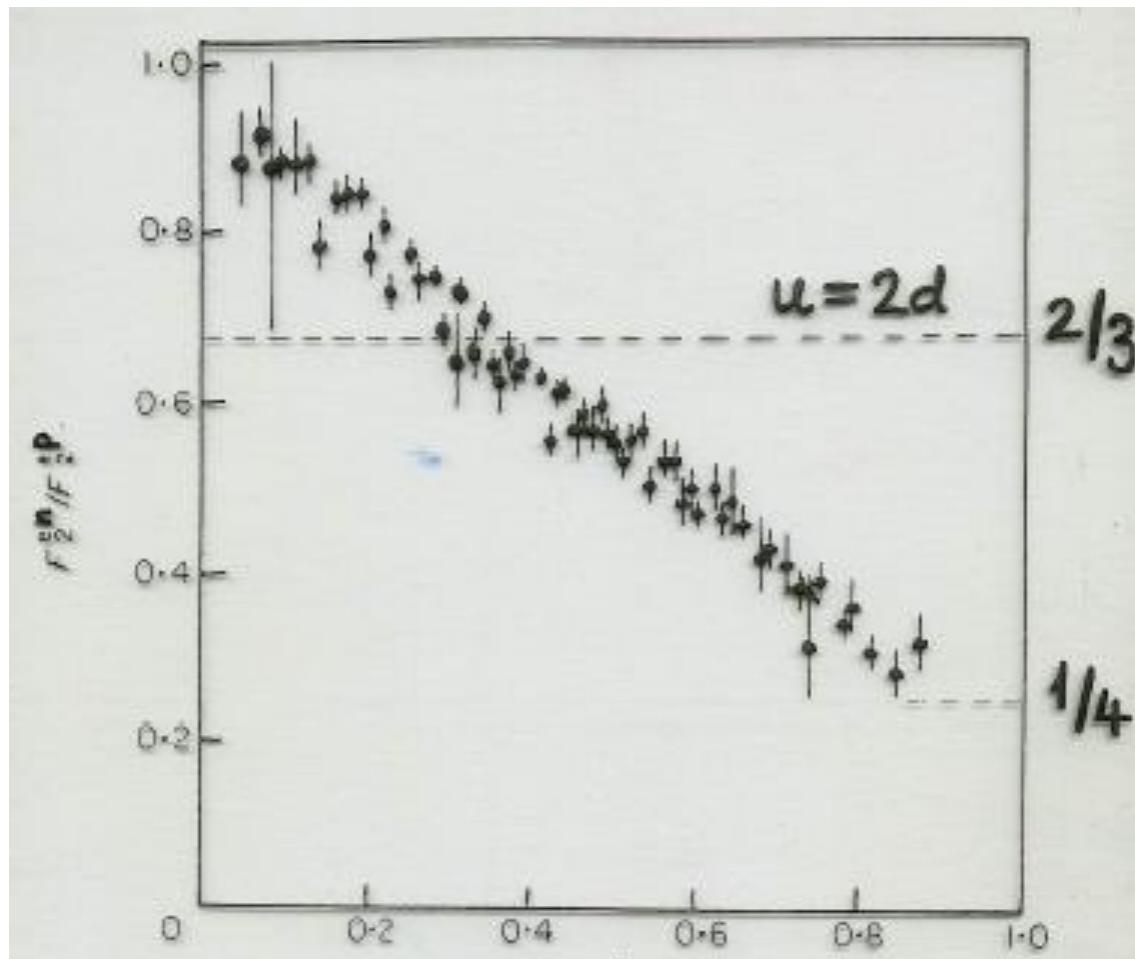
Unknown about the nucleon: valence

$x \rightarrow 1$ behaviors of flavor and spin

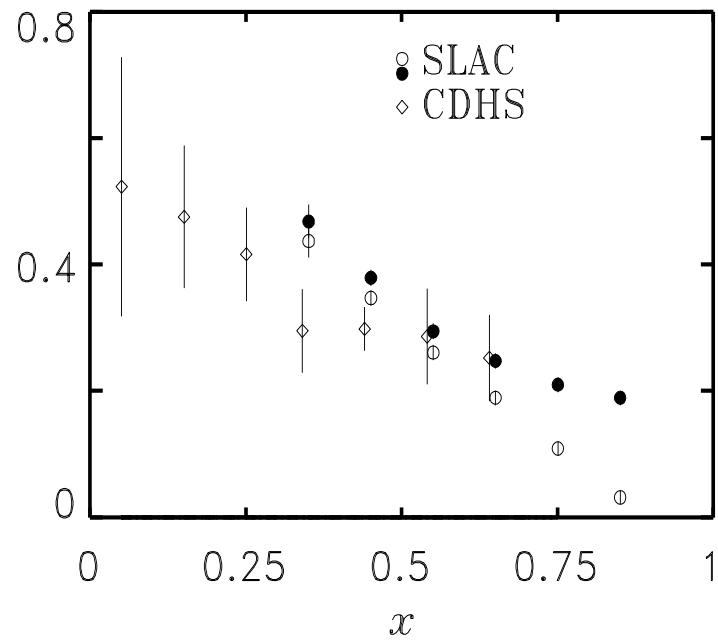
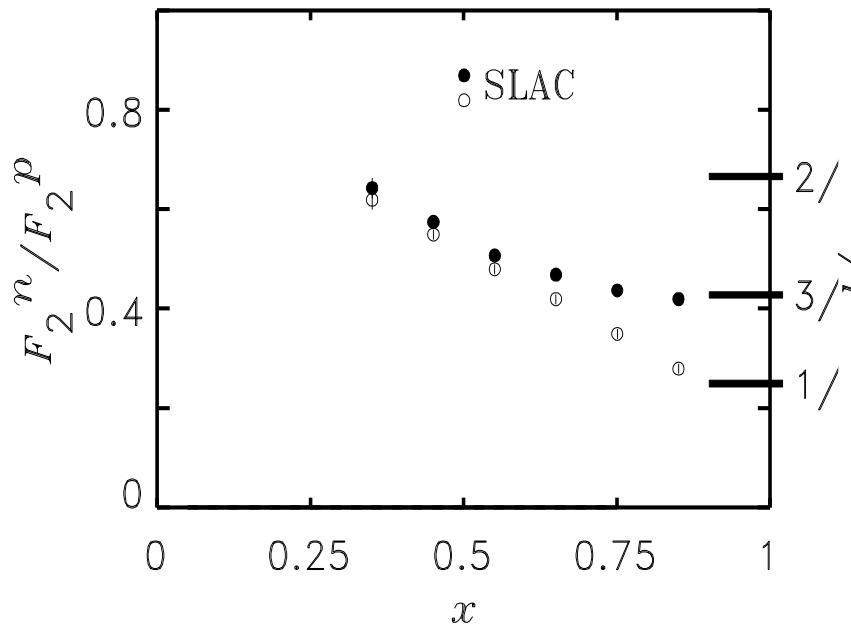
- Flavor

$\frac{d(x)}{u(x)}$	\rightarrow	0	Diquark Model
	\rightarrow	$\frac{1}{5}$	pQCD
$\frac{F_2^n(x)}{F_2^p(x)}$	\rightarrow	$\frac{1}{4}$	Diquark Model
	\rightarrow	$\frac{3}{4}$	pQCD
$\frac{d(x)}{u(x)}$	\rightarrow	$-\frac{1}{3}$	Diquark Model
	\rightarrow	1	pQCD

Ratio of Neutron/Proton Structure Functions

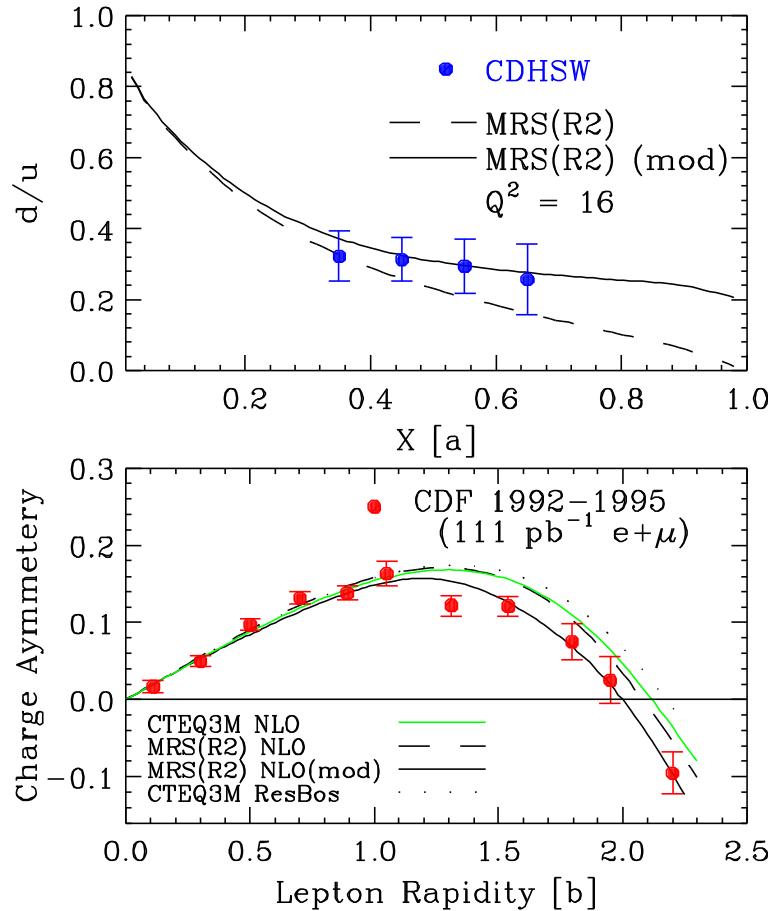


Flavor Content of the Proton with nuclear binding correction



W.Melnitchouk & A.W. Thomas
PLB 377(1996) 11

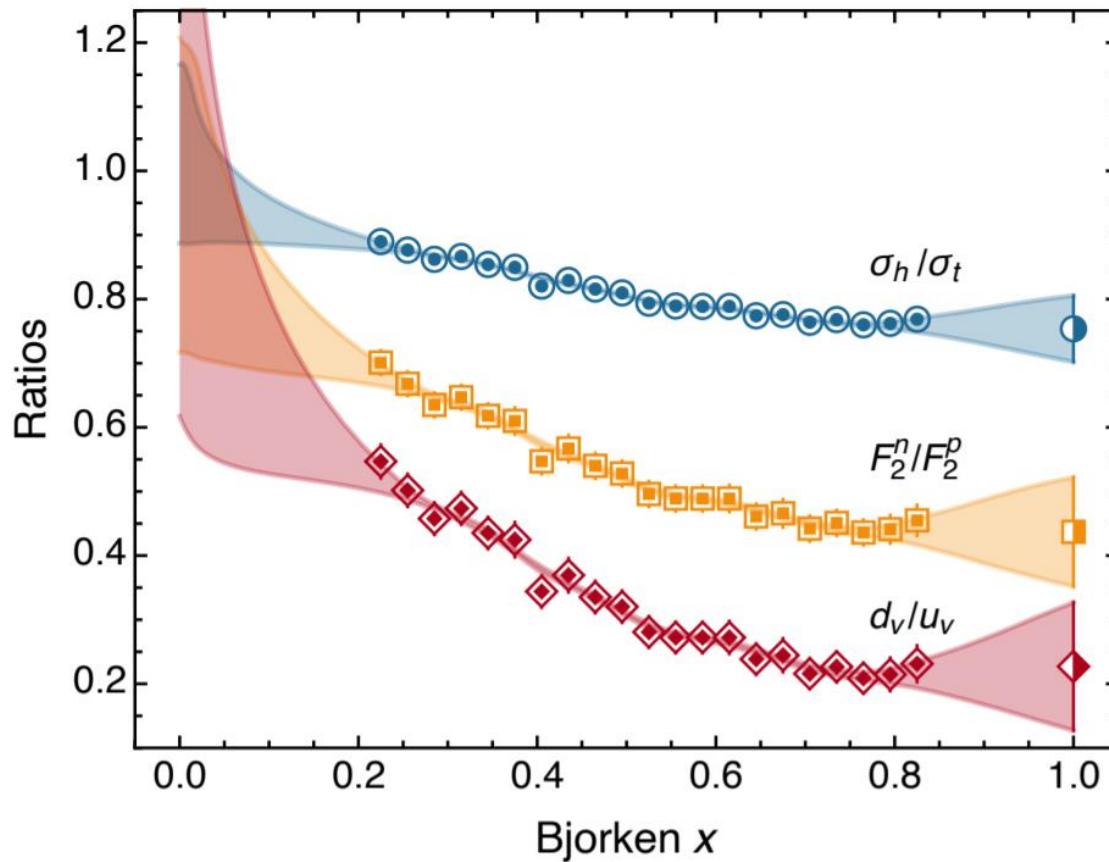
Flavor Content of the Proton from DIS neutrino data analysis



U.K. Yang & A. Bodek
PRL 82 (1999) 2467.

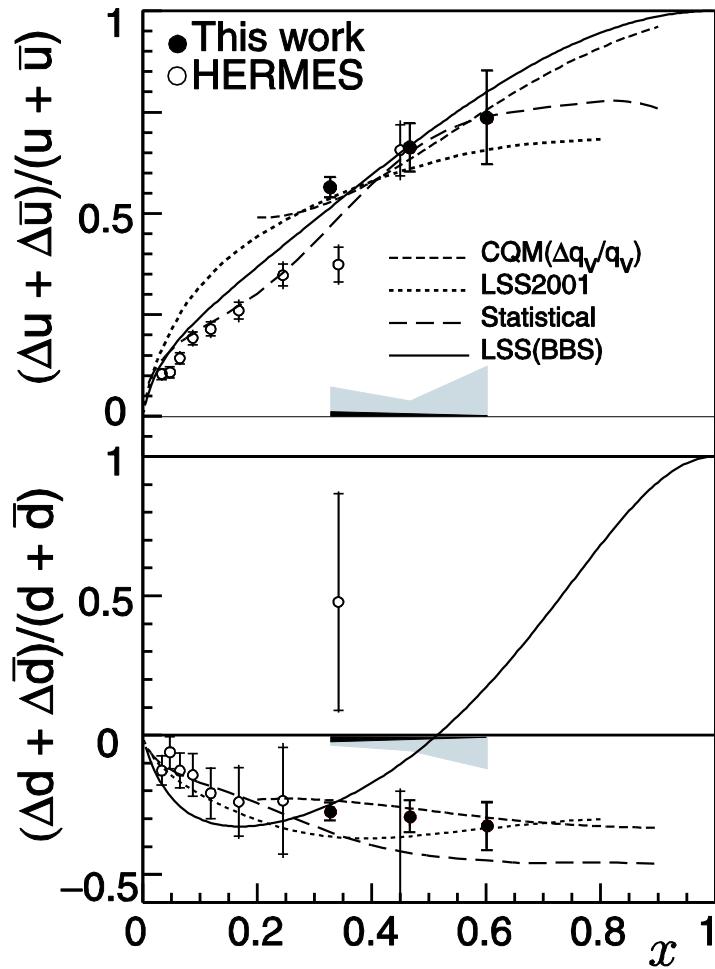
Flavor Content of the Proton

from DIS data of ${}^3\text{He}$ and ${}^3\text{H}$



Quark Helicity Distributions of Proton

Measurements at JLAB and HERMES



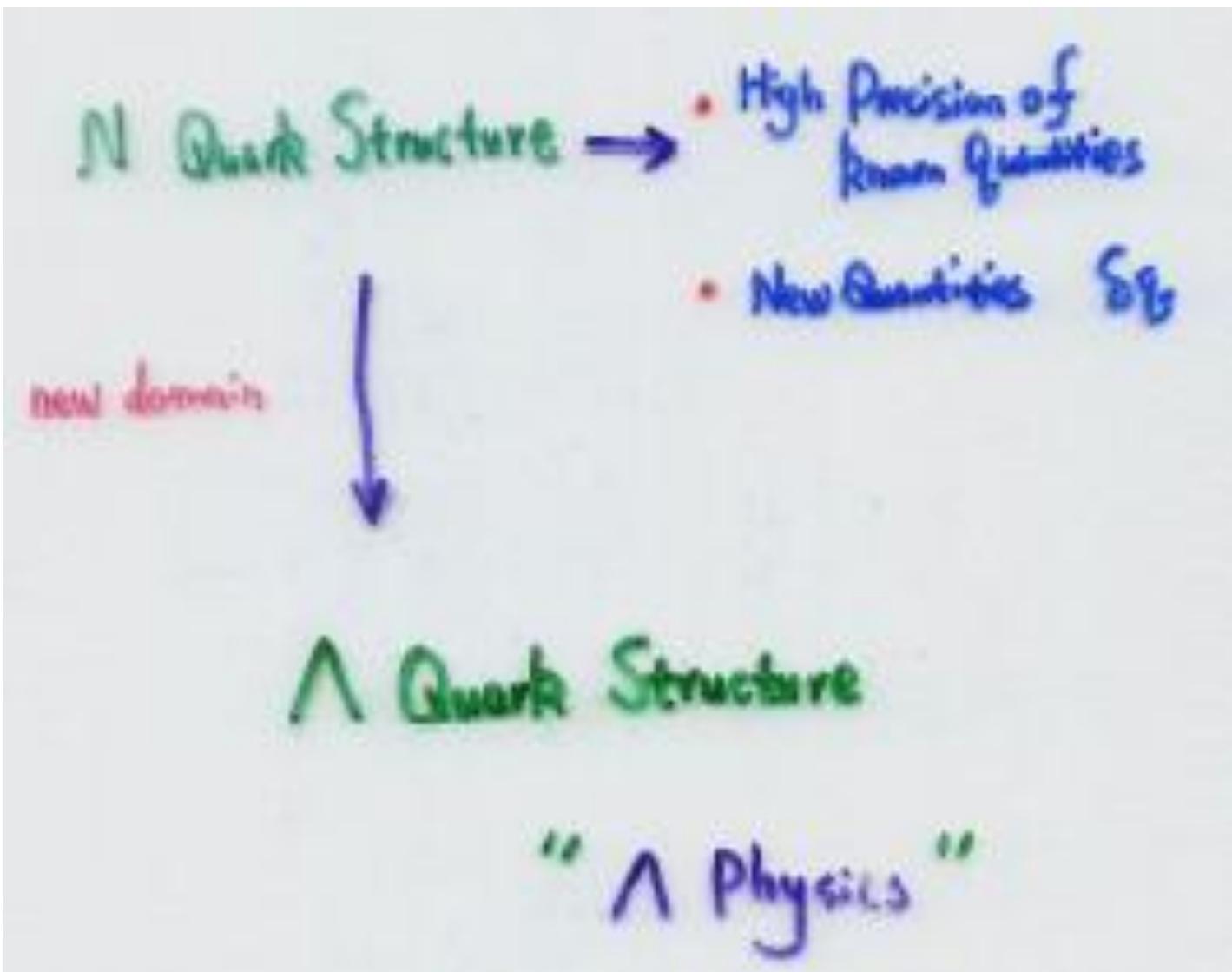
X. Zheng et al, JLab Hall A Collaboration
nucl-ex/0308011
PRL92 (2004) 012004.

Status of the Flavor and Spin Contents of the Proton

- **Flavor favors pQCD**
- **Spin favors diquark model**

Unclear & In Contradiction!

How to Test Various Theories?



SU(3) Symmetry together with Proton Spin Problem

PRL 70(1993) 2557

Burkhardt-Schiff SU(3) Argument:

$$\int_0^1 dx g_1^{u\Lambda}(x) = \frac{1}{18} (2\Sigma - D)$$

$$= \int_0^1 dx g_1^{d\Lambda}(x) - \frac{1}{18} (2D + 3F)$$

$$= -0.042 \pm 0.019$$

$$\Delta u^\Lambda = \Delta d^\Lambda = \frac{1}{3} (\Sigma - D) = -0.23 \pm 0.06$$

$$\Delta S^\Lambda = \frac{1}{3} (\Sigma + 2D) = 0.58 \pm 0.03$$

whereas the Quark Model predicts

$$\Delta u^\Lambda = \Delta d^\Lambda = 0$$

$$\Delta S^\Lambda = 1$$

The u,d sea of Lambda versus s sea of Nucleon

Ma-Soffer

PRL 81(98) 2250

u,d polarizations in Λ

is related to s polarization in N

$$\bullet \quad P(uudss) = \Lambda(uds) k^+(us)$$

$$\bullet \quad \left\{ \begin{array}{l} \Lambda(udsuu) = P(uud) k^-(su) \\ \Lambda(udsd\bar{d}) = P(udd) k(s\bar{d}) \end{array} \right.$$

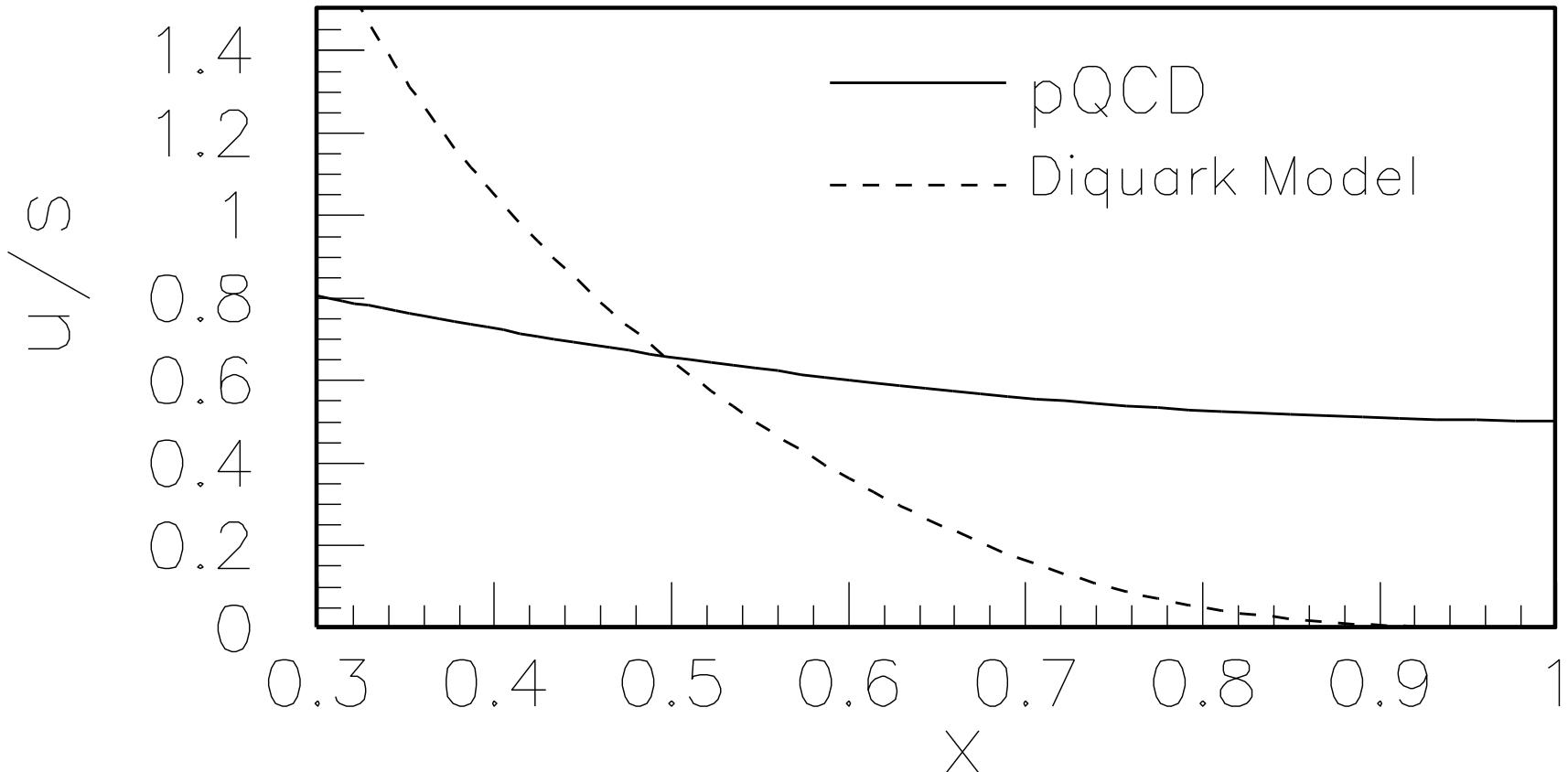
Different flavor & spin structure in different models

*Phys. Lett. B 477 (2000) 107
Ma-Schmidt-Yang $x \rightarrow 1$ behaviors*

- Flavor
 - $\frac{u(x)}{s(x)}$ $\rightarrow 0$ Diquark Model
 - $\frac{1}{2}$ pQCD
- $\frac{\Delta S(x)}{S(x)} \rightarrow 1$
- $\frac{\Delta u(x)}{u(x)} \rightarrow 1$

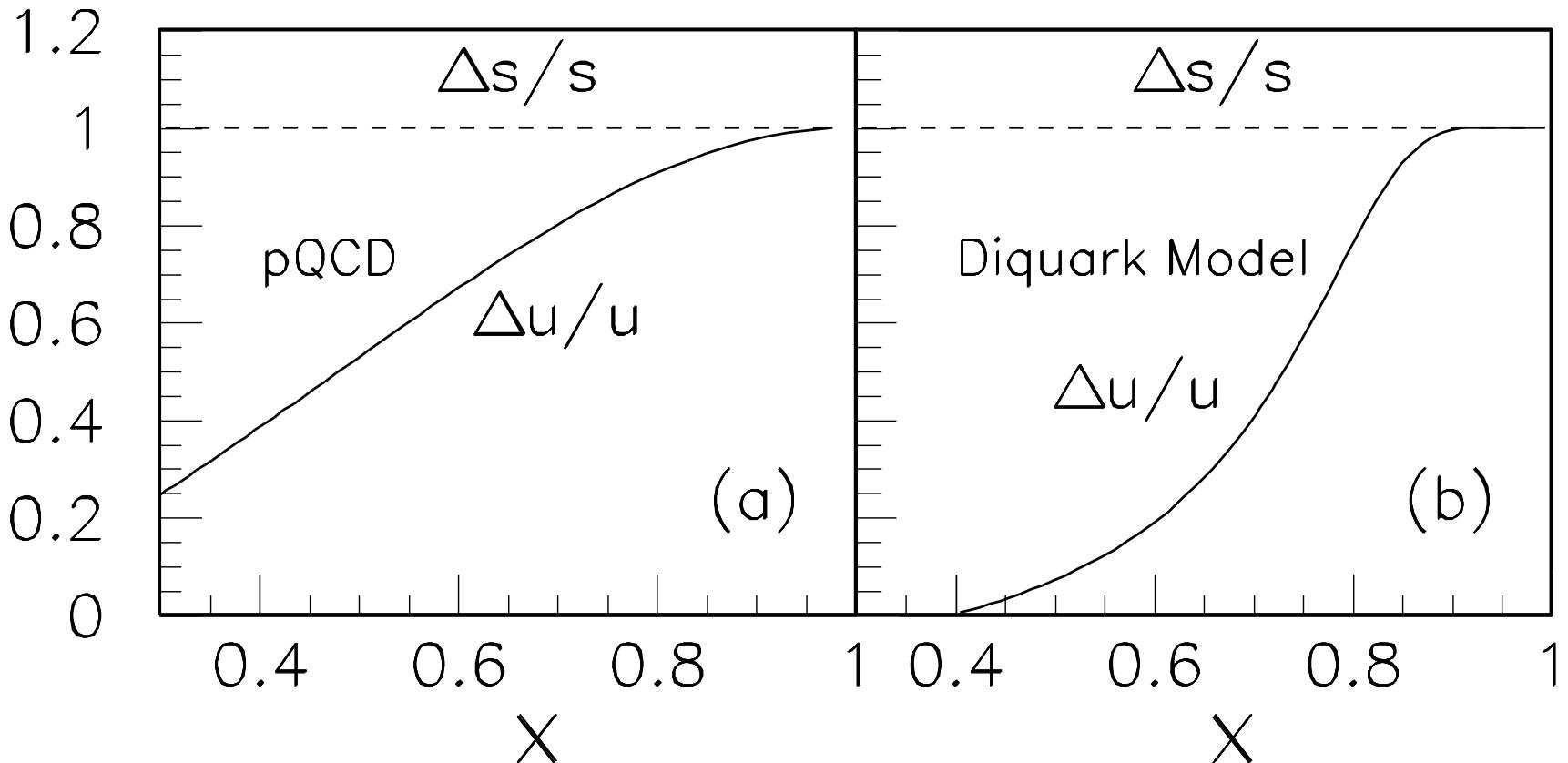
$\Delta u(x) = \Delta d(x) > 0 \quad \text{at large } x$
 $\int_0^1 dx \Delta u(x) = \int_0^1 dx \Delta d(x) \leq 0$

Flavor structure in two different models



B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Lett. B 477 (2000) 107

Spin structure in two different models



B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Lett. B 477 (2000) 107

An intuitive argument

Quark-Diquark Model

$S(u\bar{u}d\bar{d}) S$	$U^+ S(d\bar{s}) \quad M_D = M_S$
$D(\bar{d}s) u$	$U^+ V(\bar{d}s) \quad M_D = M_U > M_S$

$$\Psi_{D^*} \approx e^{-\omega \left(\frac{\vec{k}_1^2 + M_S^2}{r^2} + \frac{\vec{k}_2^2 + M_D^2}{r'^2} \right)}$$

at $x \rightarrow 1 \quad \Psi_{D^*}(x) \ll \Psi_S(x)$

pQCD Analysis

$\Psi_{D^*}(x) \sim (1-x)^p \quad p = m-d+2/\alpha_S + 1$

$$\Delta S_d = S_q - S_p$$
$$S_q = S_p \quad |\delta S_d| = 0$$
$$S_q \neq S_p \quad |\delta S_d| = 1 \quad \text{Suppressed}$$

Significantly different predictions of Λ Structure

- naive quark model predicts:

$$\Delta u = \Delta d = 0, \quad \Delta S = 1$$

- Jaffe-Burkhardt predict:

$$\Delta u = \Delta d = -0.2 \quad \Delta S = 0.6$$

- We predict:

$$\frac{\Delta u}{u} = \frac{\Delta d}{d} \rightarrow 1 \quad \text{at } x \rightarrow 1$$

in Both quark-antiquark model
and pQCD analysis

Connections between structure functions and fragmentation functions

How to Measure $q^A(x)$, $D_q^A(z)$?

$$q^A(x) \propto D_q^A(z)$$

space-like

$$z = \frac{\theta^2}{\theta^2}$$

time-like

$$z = \frac{\theta^2}{\theta^2}$$

- space-like
- time-like
- The Gubis-Lipatov reciprocity relation
- parton distribution by parton fragmentation duality

• S.J. Brodsky, B.-Q. Ma,
PLB 392 (1997) 452.

• V. Barone, A. Drago, B.-Q. Ma,
PRC 62 (2000) 062201 (R).

• B.-Q. Ma, I. Schmidt, J. Soffer,
J.-J. Yang,
PLB 547 (2002) 245.

Advantage of Λ Physics:

Self-Analyzing Property of Λ

- **Polarization of Λ can be measured through the self-analyzing process of Λ decay:**



Various Processes of Polarized Fragmentation

Various Processes to Measure $D_f^\wedge(\theta)$, $\Delta D_f^\wedge(\theta)$

- $e^+e^- \rightarrow \vec{\Lambda} + x$ M. Burkhardt, R.L. Jaffe
PRD 74 (2006) 034013

- $\vec{t} N \rightarrow \vec{\Lambda} + x$ R.L. Jaffe, PRD 54 (2001) 074003

- $p\bar{p} \rightarrow \vec{\Lambda} + x$ D. de Flaviis, M. Stevenson, A. Wetzler, W. Vogelsang, PRD 61 (2000) 074020

- $\nu N \rightarrow \vec{\Lambda} + x$ $\frac{\Delta D_u^\wedge(\theta)}{D_u^\wedge(\theta)}$ Kuhnlein-Pohl, -de Hass, EPJC 51 (2007) 121

Flavor separation of fragmentation functions

B.-Q. Ma, J. Soffer, PRL 82 (1999) 2250

Complete flavor separation of

$$D_g^{\wedge}(z), \Delta D_g^{\wedge}(z), D_{\bar{q}}^{\wedge}(z), \Delta D_{\bar{q}}^{\wedge}(z)$$

- $\nu N \rightarrow \mu^- \vec{\Lambda} X$
- $\bar{\nu} N \rightarrow \mu^+ \vec{\Lambda} X$
- $\nu N \rightarrow \mu^- \vec{\Lambda} X$
- $\bar{\nu} N \rightarrow \mu^+ \vec{\Lambda} X$

Extension to Sigma Hyperon

Ma-Schmidt-Yang hep-ph/9907556

Nucl. Phys. B 574 (2000) 331

Σ^+ Valence: $x \rightarrow 1$

$\frac{s}{u} \rightarrow 0$ Diquark Model

$\frac{1}{s} \rightarrow \frac{1}{5}$ pQCD

$\frac{as}{s} \rightarrow -\frac{1}{3}$ Diquark Model

$\rightarrow 1$ pQCD

$\frac{\delta u}{u} \rightarrow 1$

- Bigger difference at middle X

The advantage of using Sigma hyperons

The Advantage of Σ^\pm

charged, us beam

$$\Sigma^\pm + N \rightarrow \mu^\pm \mu^- + X$$

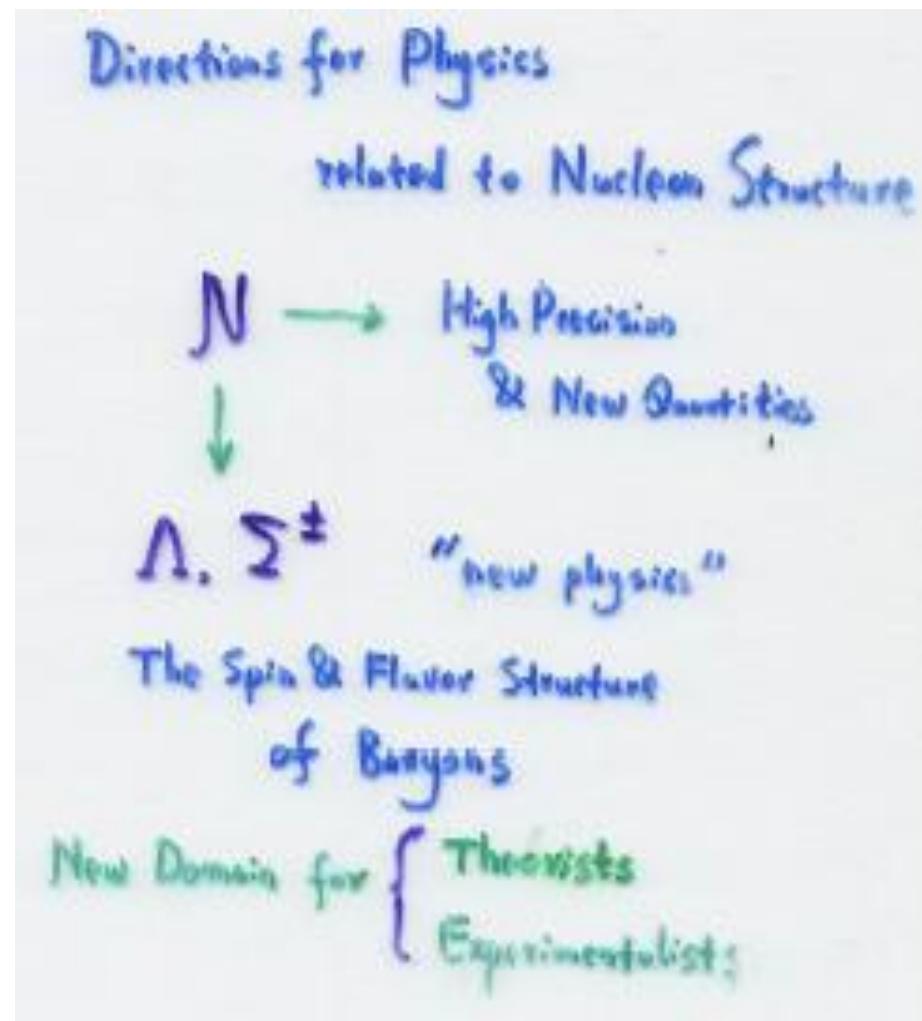
A Feynman diagram showing the annihilation of a Sigma baryon (Σ^\pm) and a nucleon (N). The Σ^\pm baryon decays via an s-channel exchange of a gluon (s) into a muon (μ^\pm) and an intermediate state. This intermediate state then decays via a t-channel exchange of a gluon (t) into a muon (μ^-) and another intermediate state. The final state consists of a muon (μ^\pm), a muon (μ^-), and an additional particle (X). The quark distributions can be measured.

Alberg et al. (Nucl. Phys. B 373)

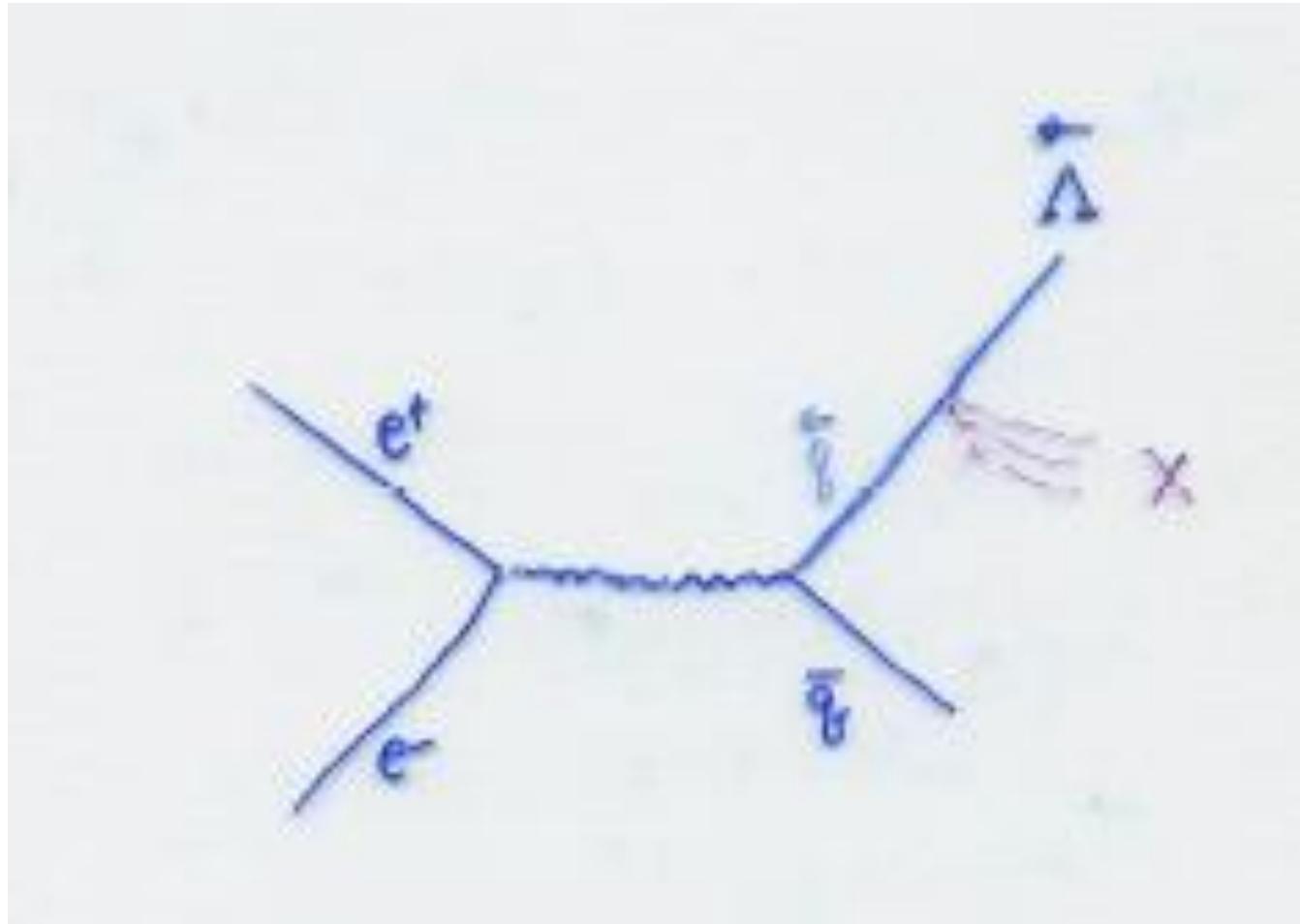
Mo, Schmidt, Yang, hep-ph/9501017

The quark distributions can be measured

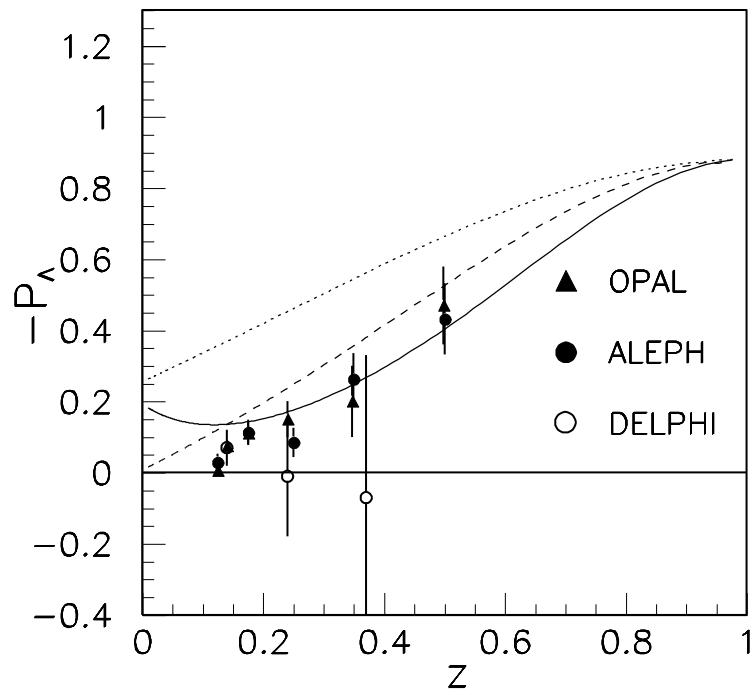
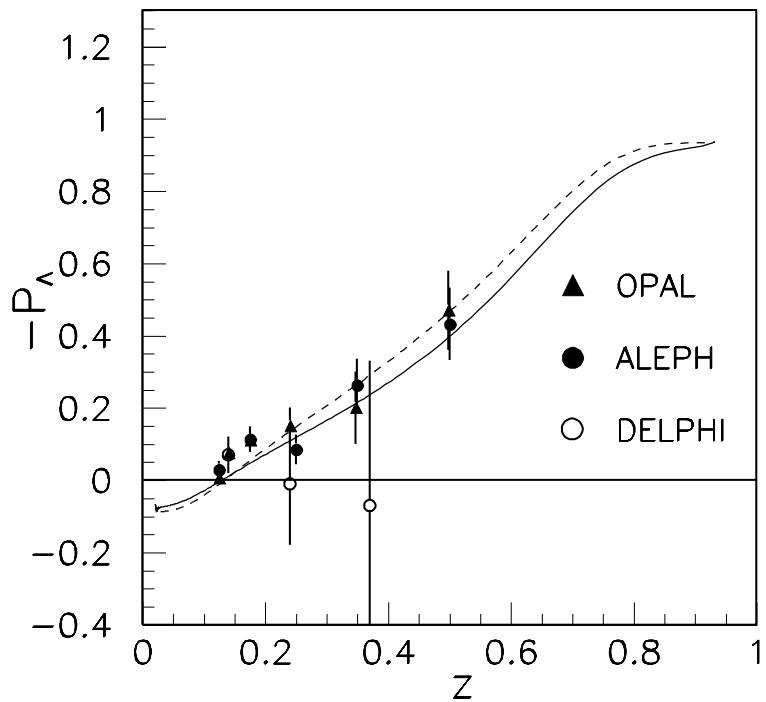
New Domain for Theorists and Experimentlists



Spin structure of Lambda from Lambda polarization in Z^0 decay

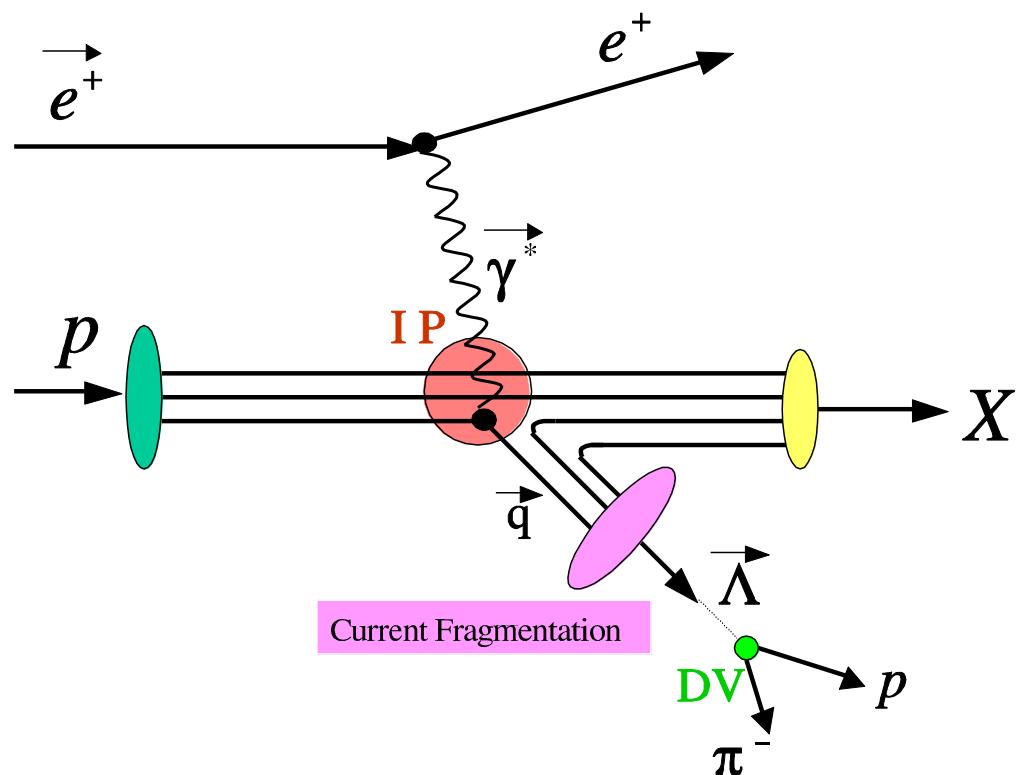


Diquark model and pQCD results

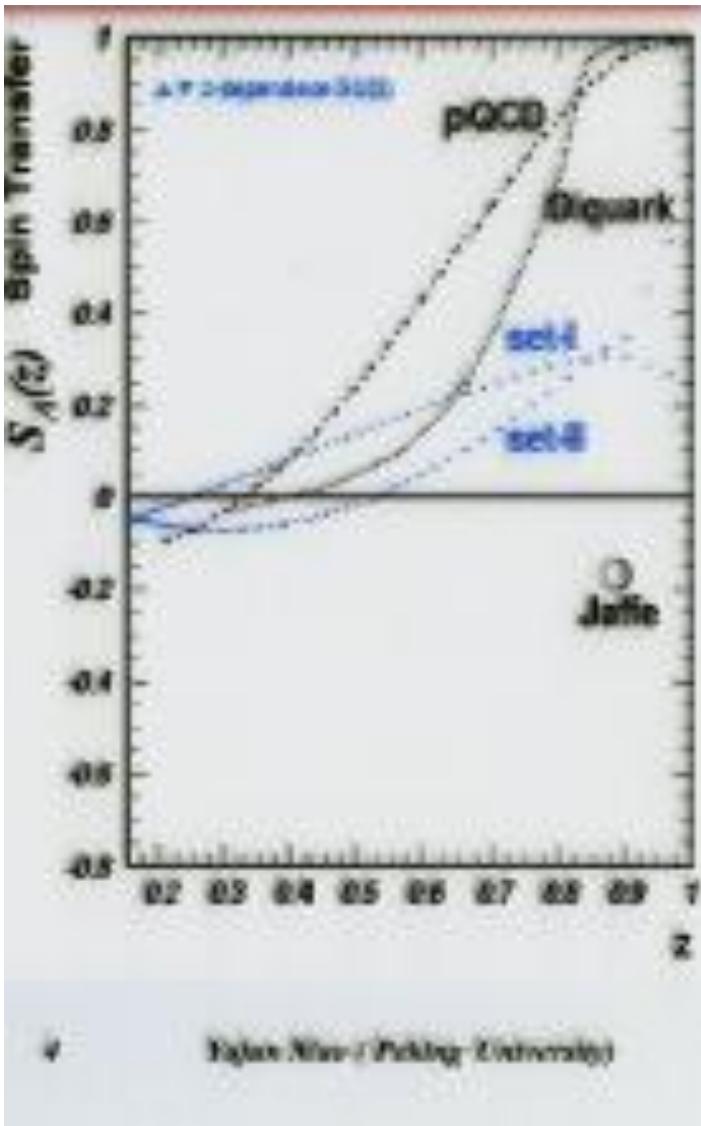


B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Rev. D 61 (2000) 034017

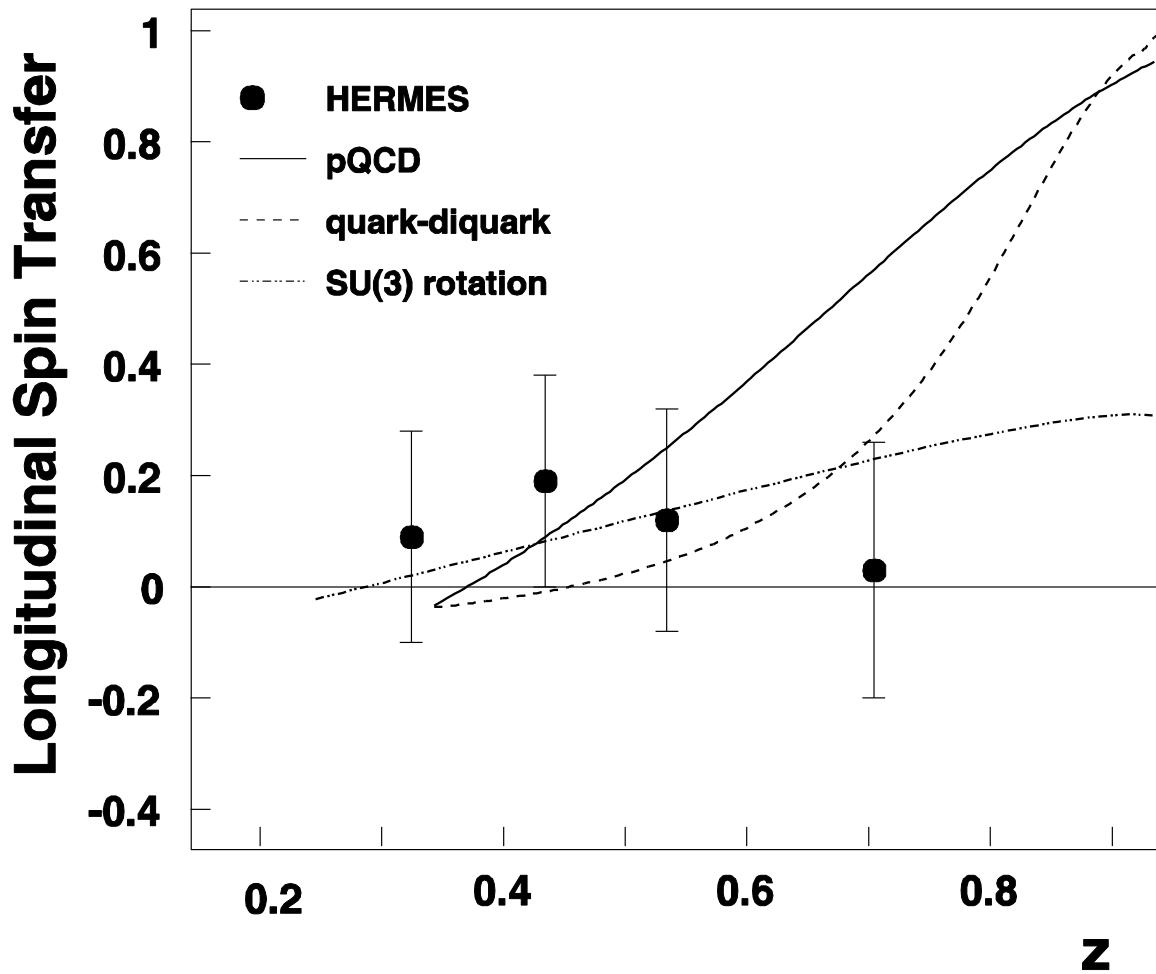
Spin Transfer to Λ in Semi-Inclusive DIS



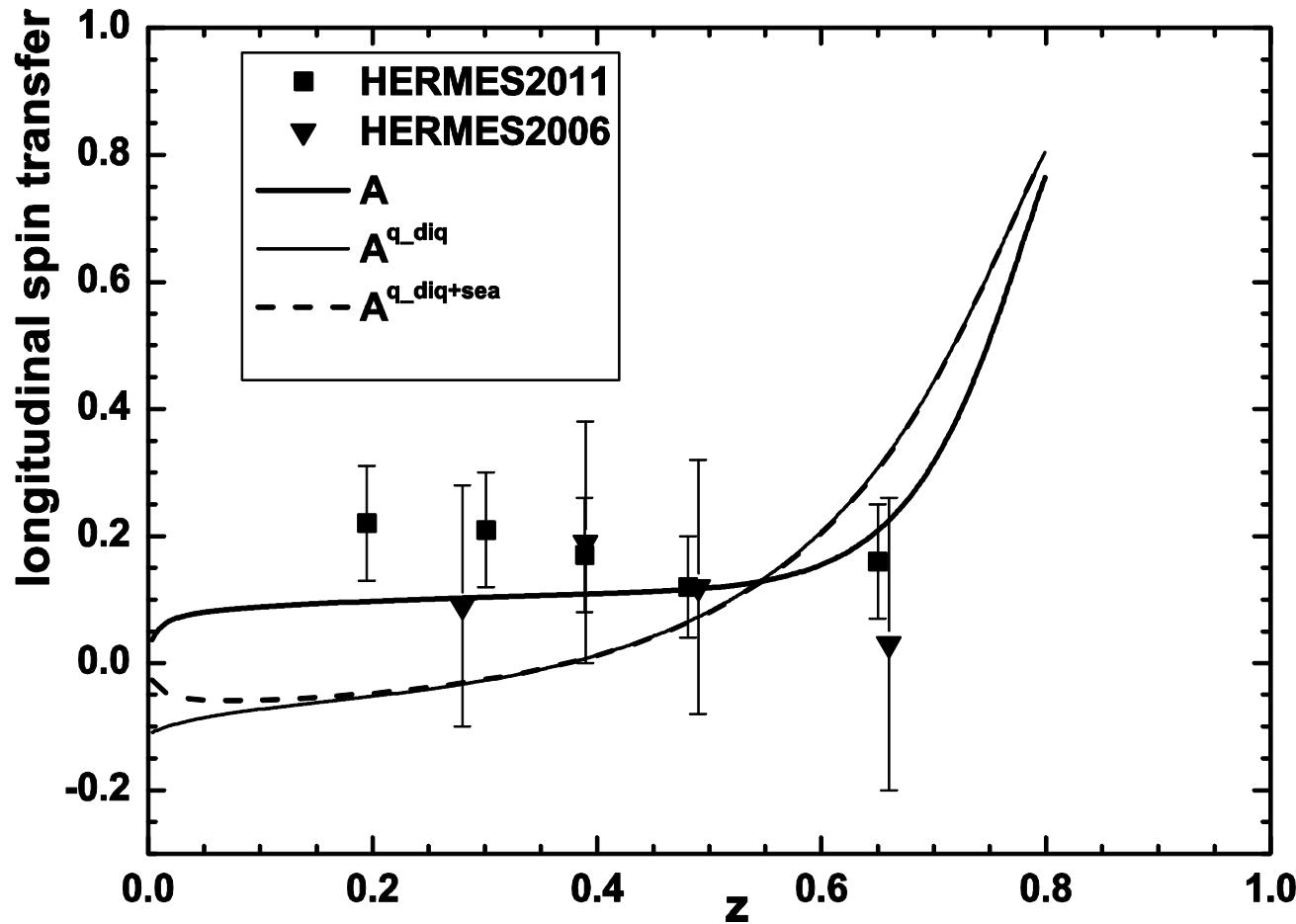
Different predictions



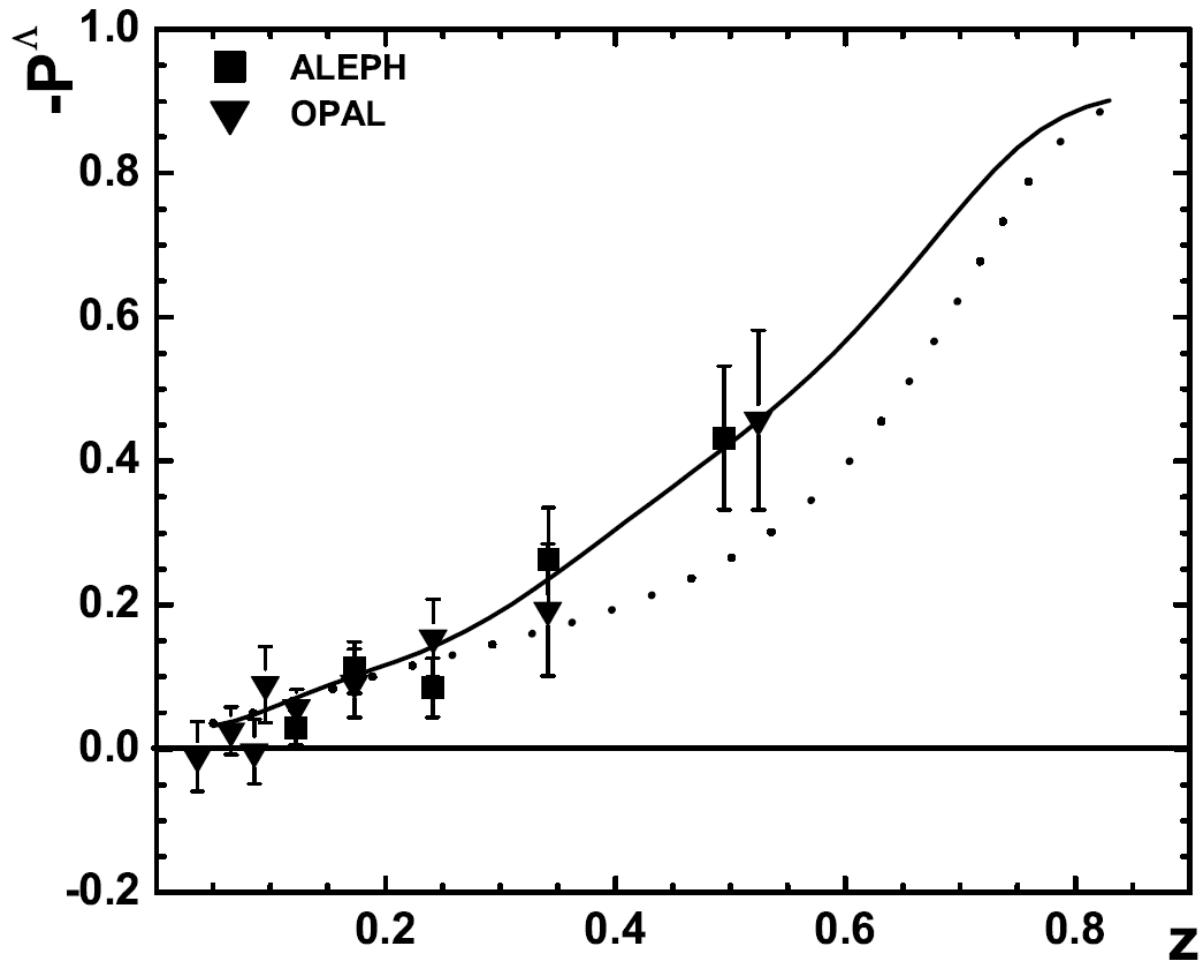
Comparison with data



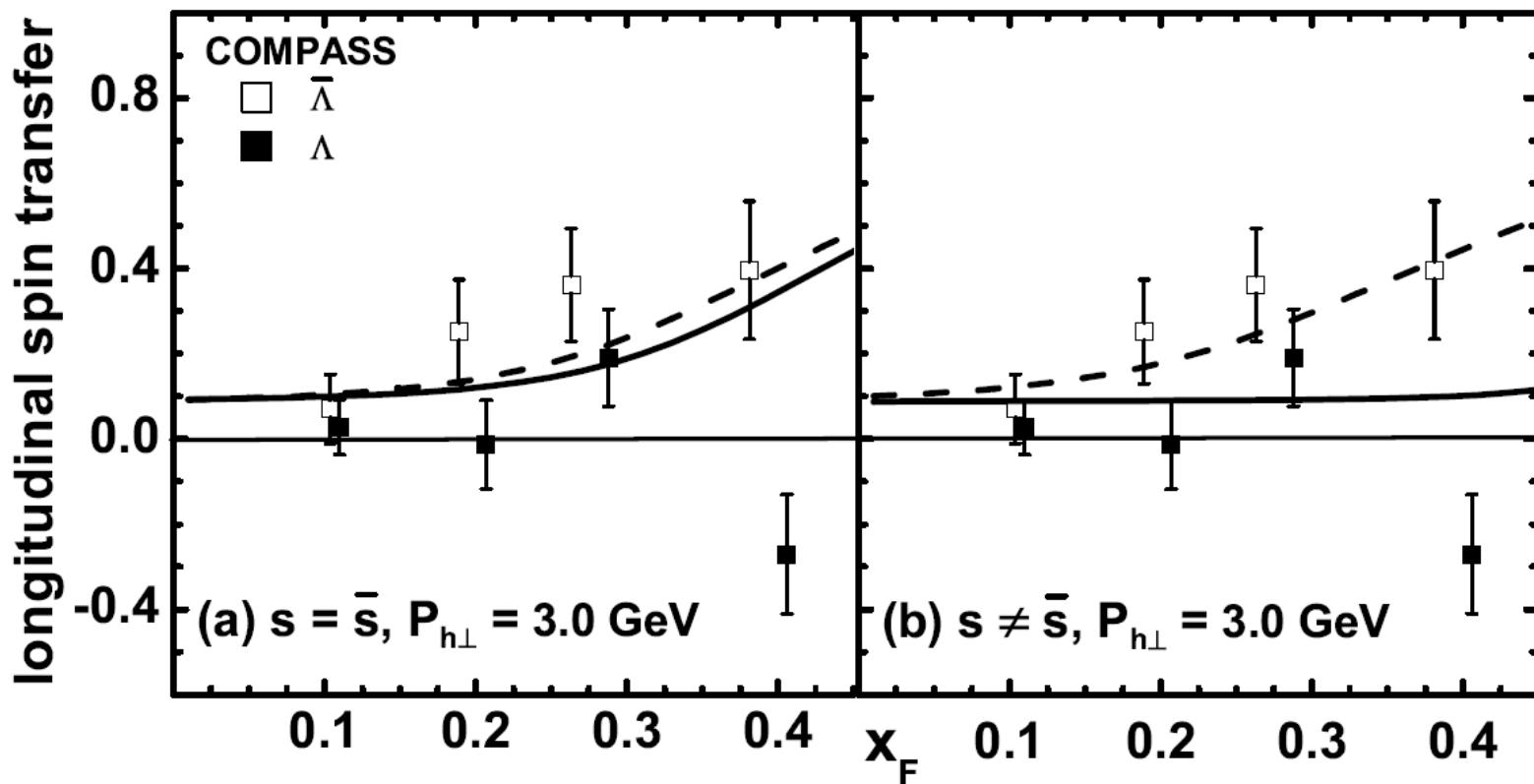
New results including both unfavored and indirect decays: SIDIS



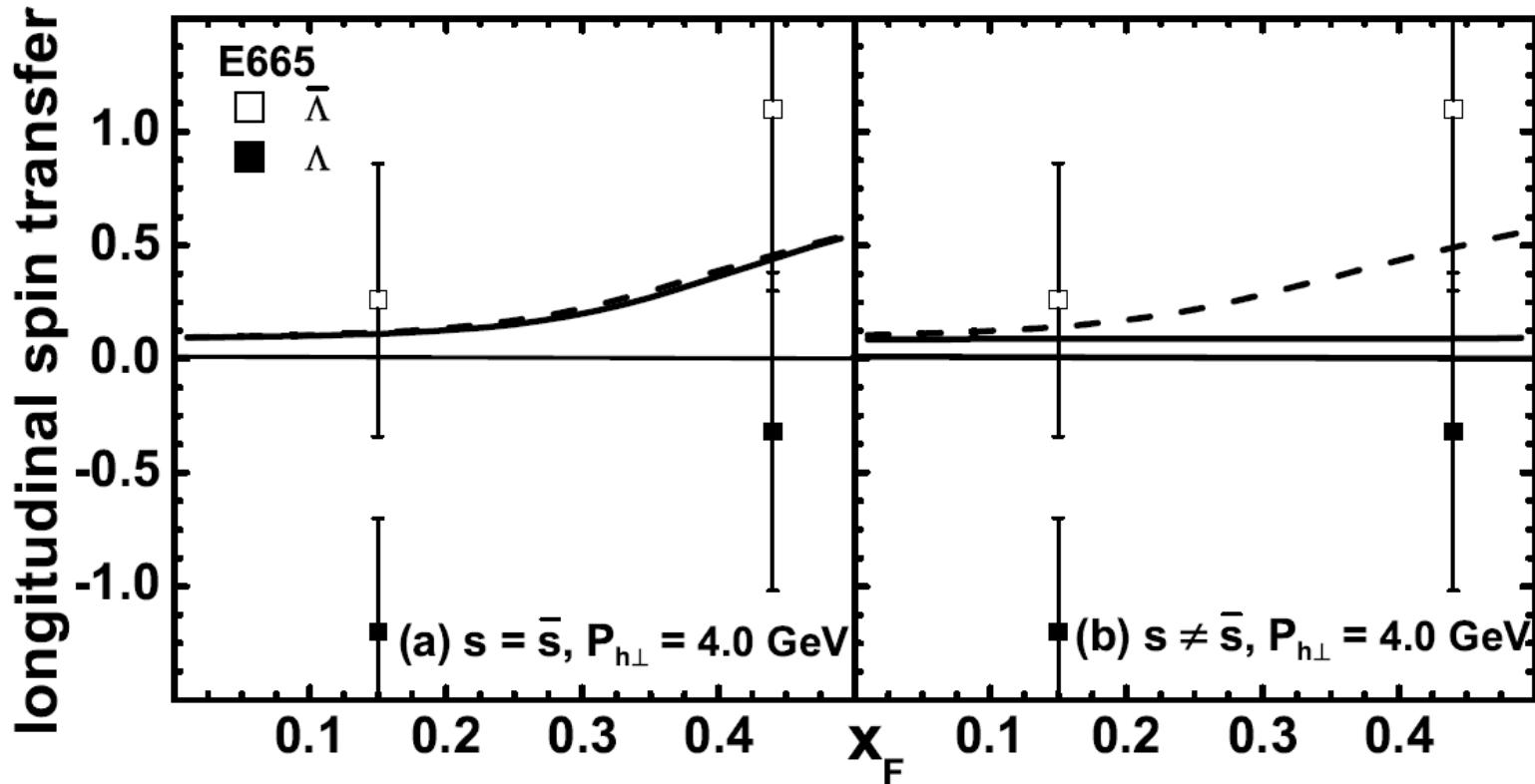
Results with new parametrization: Z-pole



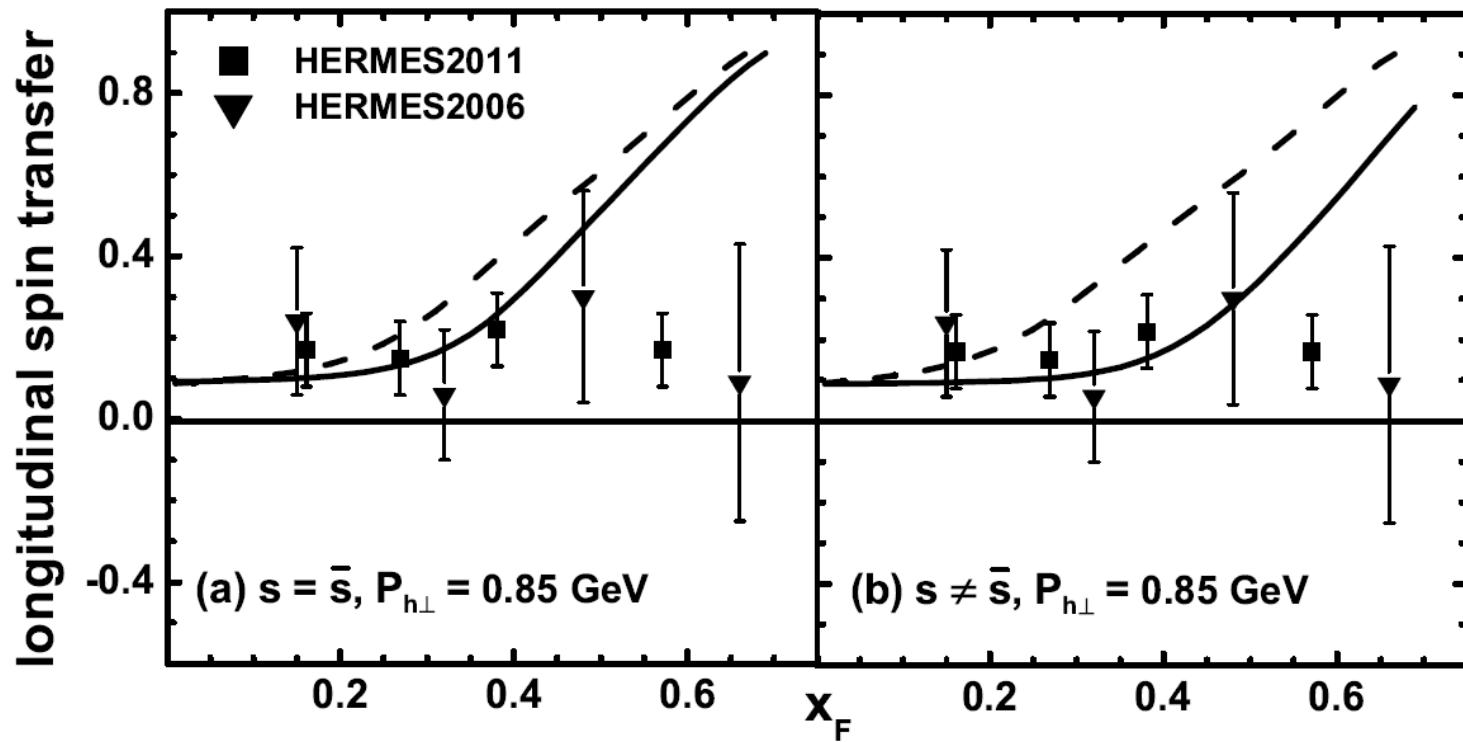
Difference between Lambda and anti-Lambda spin transfers with the COMPASS data



*Difference between Lambda and anti-Lambda spin transfers
with s-sbar asymmetry for E665*



Difference between Lambda and anti-Lambda spin transfers with s-sbar asymmetry for HERMES

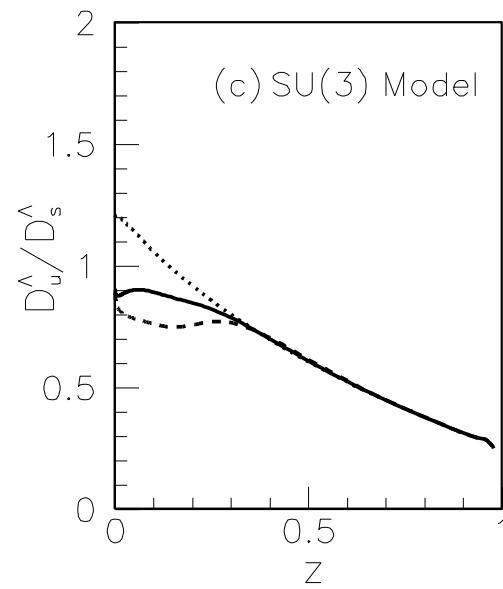
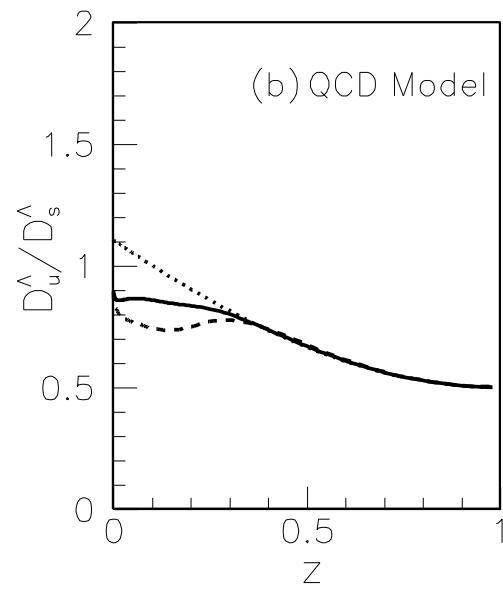
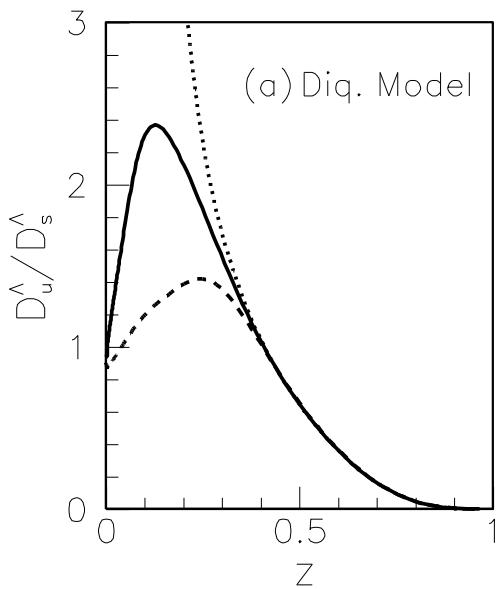


$\bar{\Lambda}/\Lambda$ Ratio in DIS Production

- A sensitive quantity that can provides information about the flavor structure of Λ hyperon.

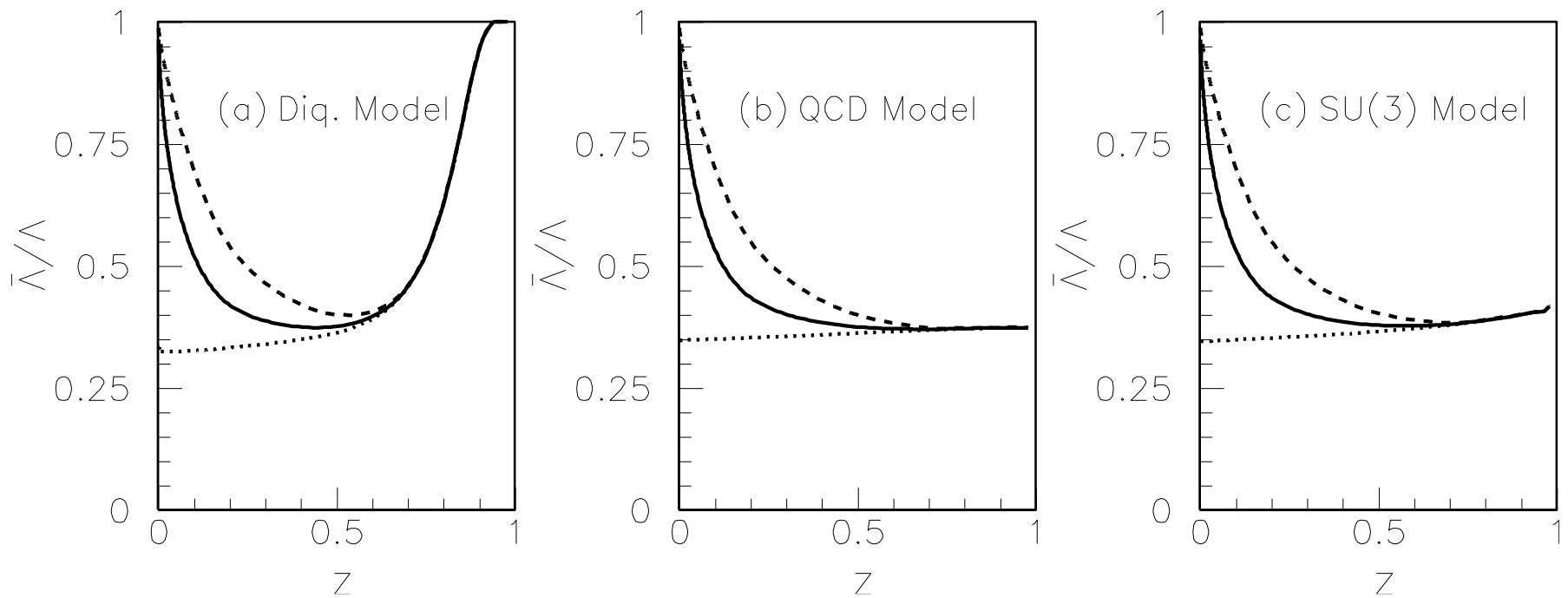
B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

The flavor structure of Lambda u/s ratio with x-dependence



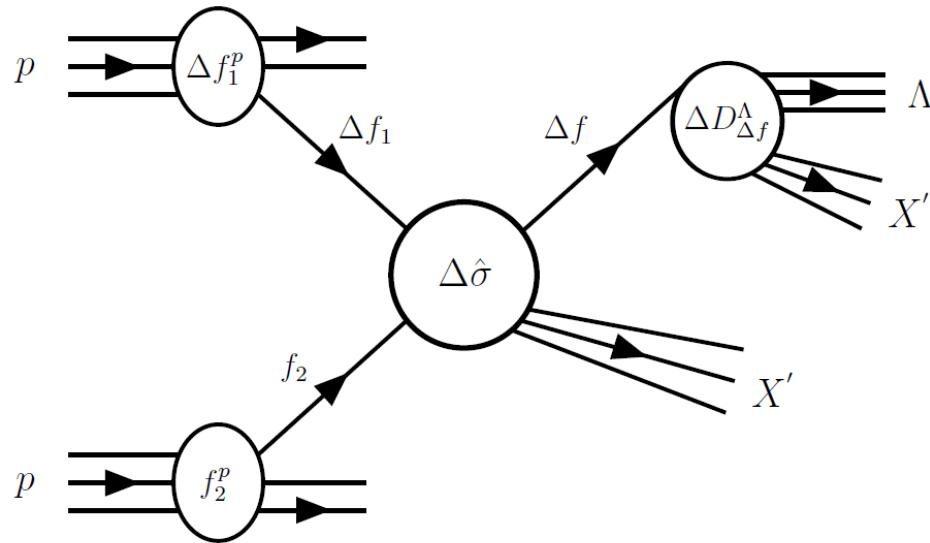
B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

Different predictions



B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

$$\vec{p}p \rightarrow \vec{\Lambda}X$$

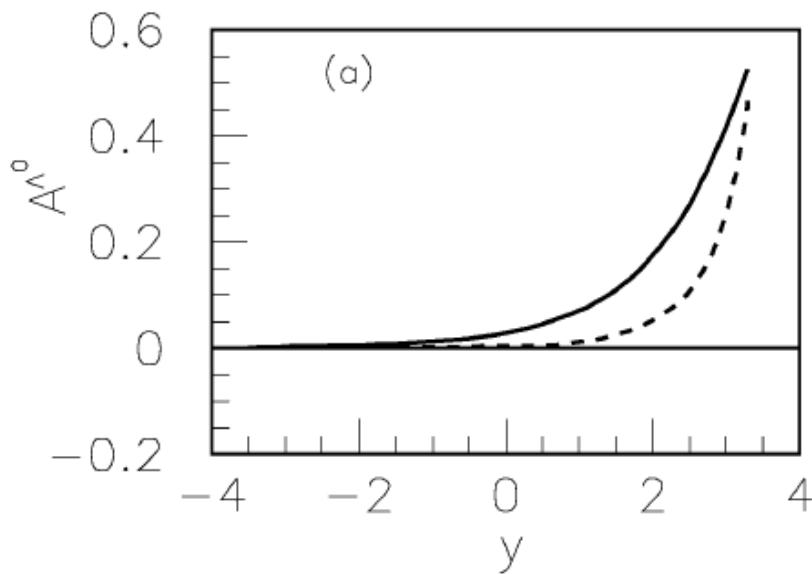


Providing information about

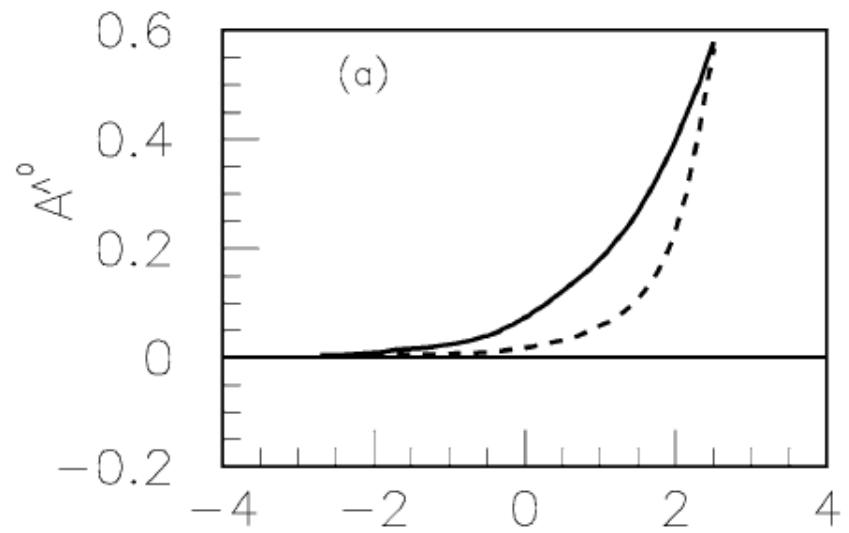
- the inclusive production of hadrons
- the strange and antistrange quark polarizations of the proton.

Spin transfer for

$\vec{p} p \rightarrow \vec{\Lambda} X$ at RHIC-BNL



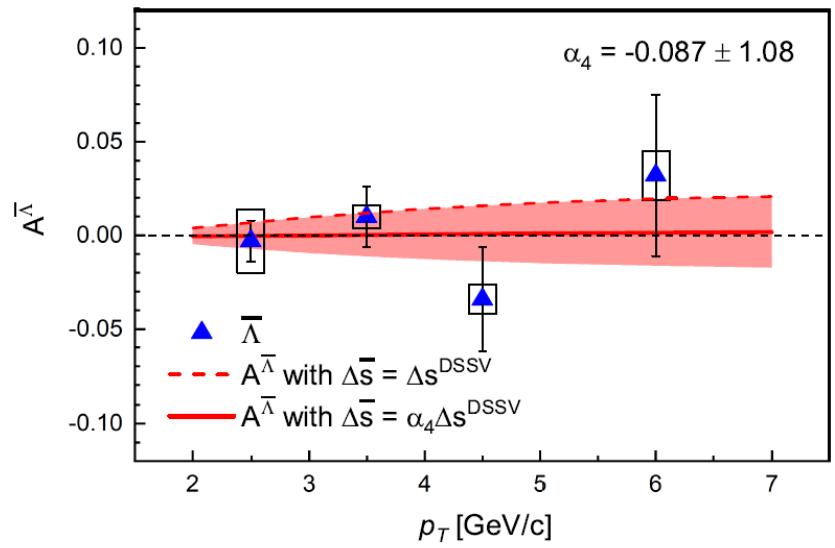
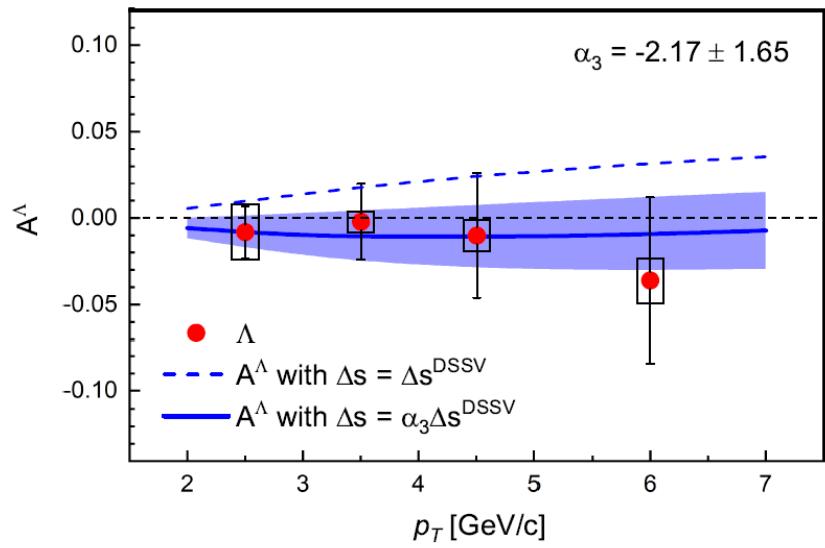
$\sqrt{s} = 500$ GeV



$\sqrt{s} = 200$ GeV

B.-Q. Ma, I. Schmidt, J.-J. Yang, J.Soffer, Nucl. Phys. A 703 (2002) 346

Fitting to STAR DATA



Results from fitting STAR data

Table: Fitting results of α_i and calculated results of Δs and $\Delta \bar{s}$.

	value	Δs	$\Delta \bar{s}$	χ^2_{\min}
α_1	-1.20 ± 1.31	-0.014 ± 0.015		0.37
α_2	-0.24 ± 0.49		-0.003 ± 0.005	2.48
α_3	-2.17 ± 1.65	-0.025 ± 0.019		0.42
α_4	-0.087 ± 1.08		-0.001 ± 0.012	2.24

Two options: with/without gluon polarization

Comparison with Predictions & Results

The central values of the fitting results are basically compatible with

- the light-cone meson-baryon fluctuation model²⁴ prediction $\Delta s(x) \approx -0.05$ to -0.01 and $\Delta \bar{s}(x) \approx 0$.
- the recent lattice QCD determination²⁵, $\Delta s^+ = -0.02(1)$ at $Q^2 \approx 7\text{GeV}^2$.
- the results from Jefferson Lab Angular Momentum (JAM) Collaboration²⁶ $\Delta s^+(Q_0^2) = -0.03(10)$.

²⁴S. J. Brodsky and B.-Q. Ma, Phys. Lett. B 381, 317 (1996).

²⁵G. S. Bali et al. [QCDSF Collaboration], Phys. Rev. Lett. 108, 222001 (2012)

²⁶J.J.Ethier, N.Sato and W.Melnitchouk, Phys. Rev. Lett. 119, 132001 (2017) 

Relating Λ -production with nucleon strangeness

- The spin transfer process of $\vec{p}p \rightarrow \vec{\Lambda}X$ is feasible to study strange-antistrange polarizations of the nucleon.
- The fitting to STAR data suggests:

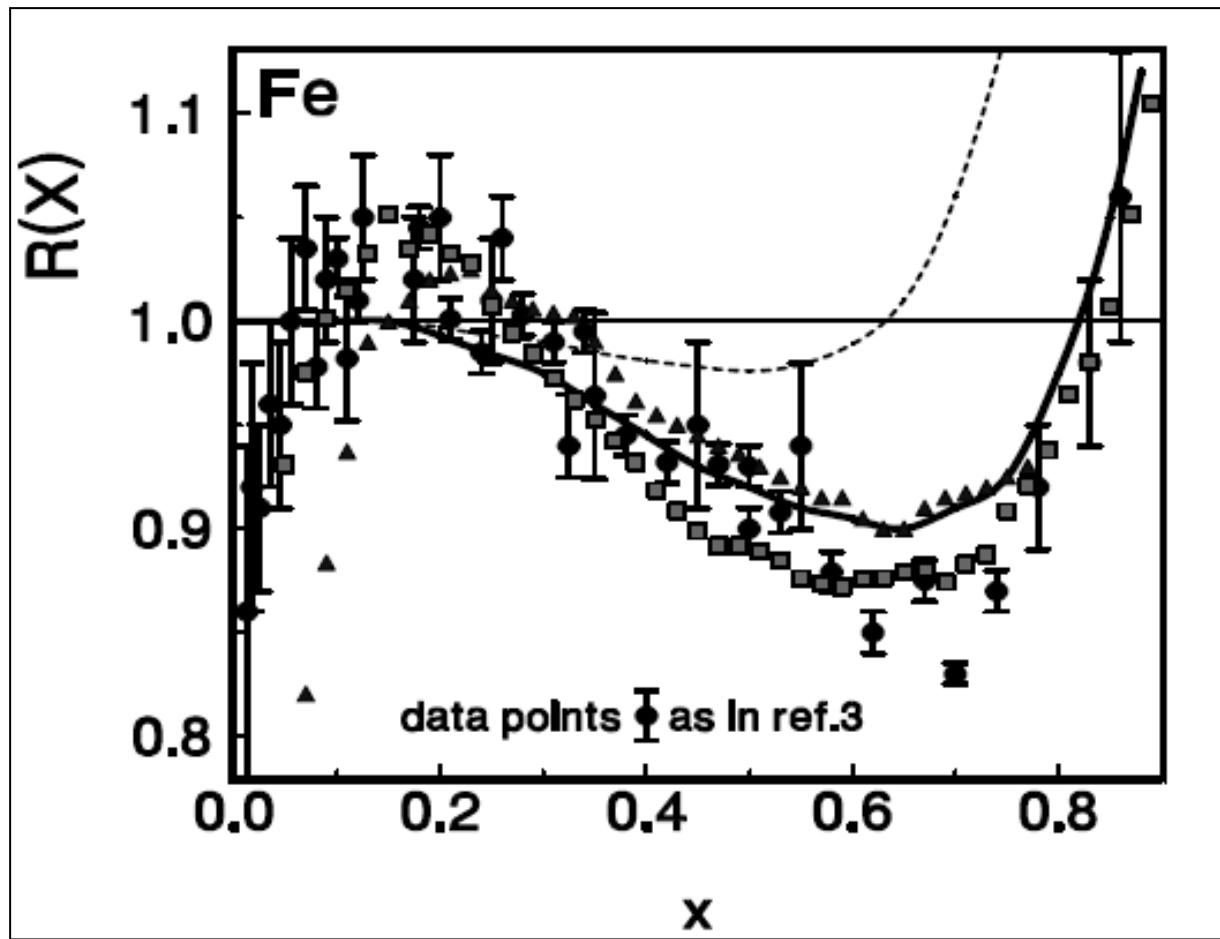
$$\Delta s \neq \Delta \bar{s}$$

$$\Delta s \approx -0.025 \pm 0.019$$

$$\Delta \bar{s} \approx -0.001 \pm 0.012$$

- The results are compatible with the light-cone baryon-meson fluctuation model prediction.

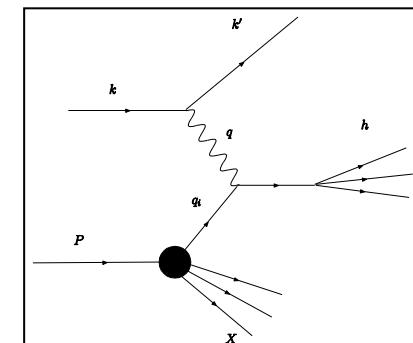
Nuclear EMC Effect



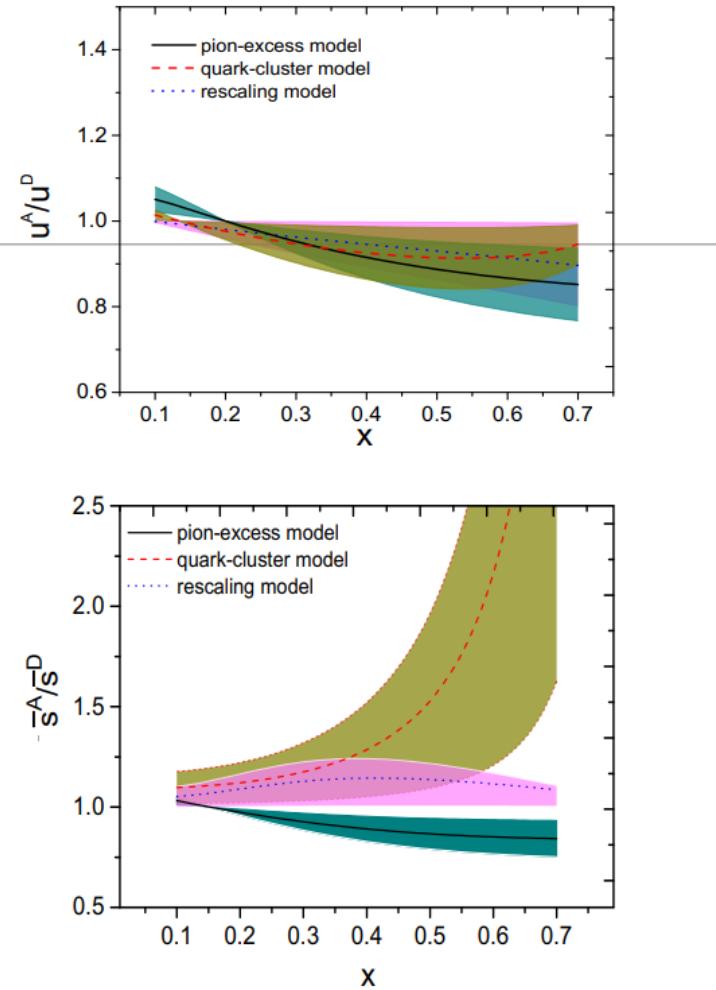
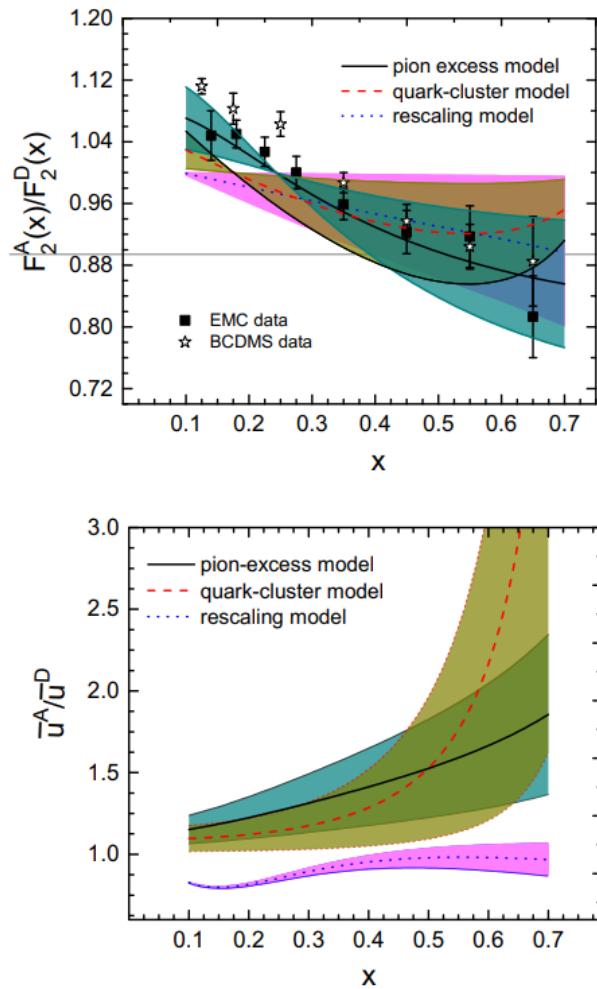
Anti-Lambda production as a probe of the nuclear sea structure?

- 有三类原子核模型可以解释EMC效应: 团簇模型, Pi盈余模型, 重新标度模型.
- 通过考察原子核与核子荷电轻子半单举过程中末态强子anti-Lambda的产率比对x的依赖性, 我们发现 anti-Lambda能够区分定性描述原子核EMC效应的三类不同原子核结构模型.

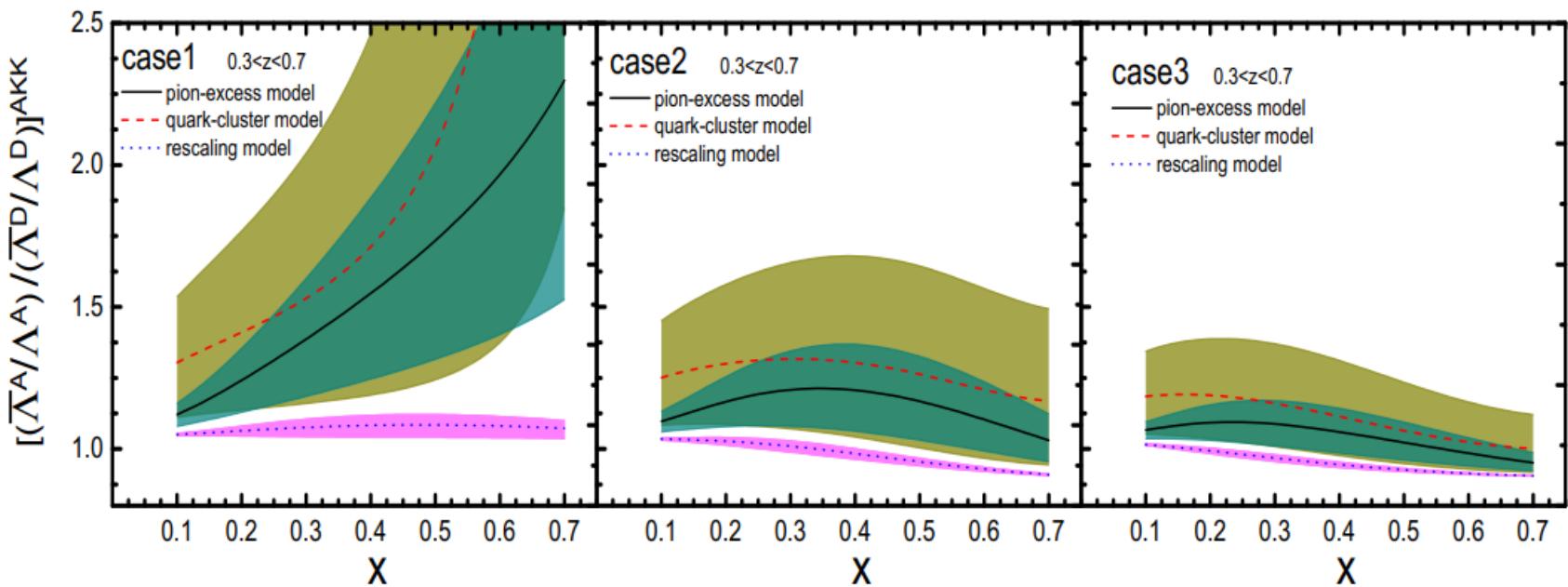
B.Lu, B.-Q. Ma, Phys.Rev.C74 (2006) 055202
C.Gong, B.-Q. Ma, Phys.Rev.C97 (2018) 065207



Different sea behaviors in nuclei



Different predictions of anti-lambda production



Conclusion 1: Lambda Physics

Λ Quark Structure

- High Precision of known Quantities

- New Quantities S_{fit}

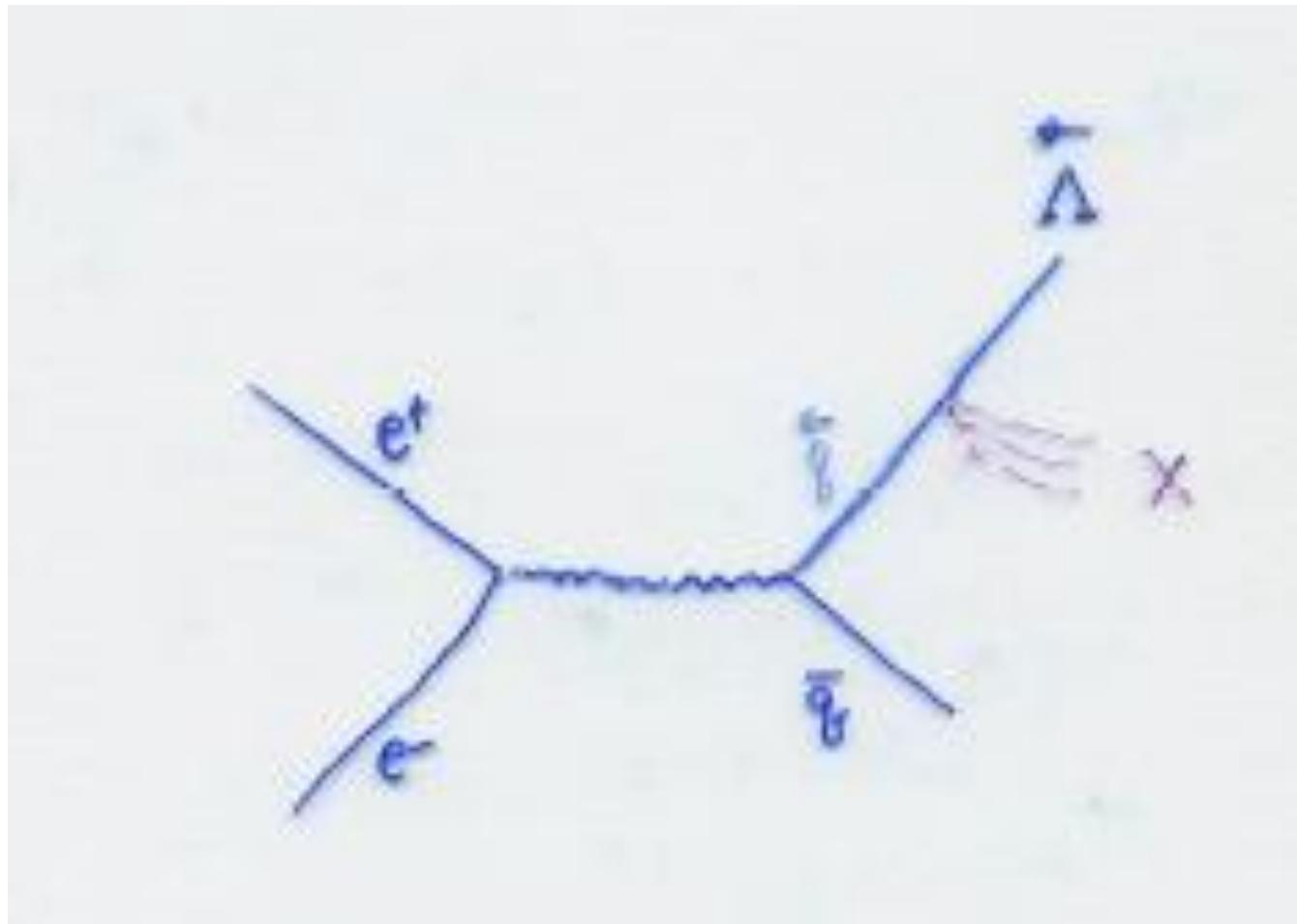
new domain



Λ Quark Structure

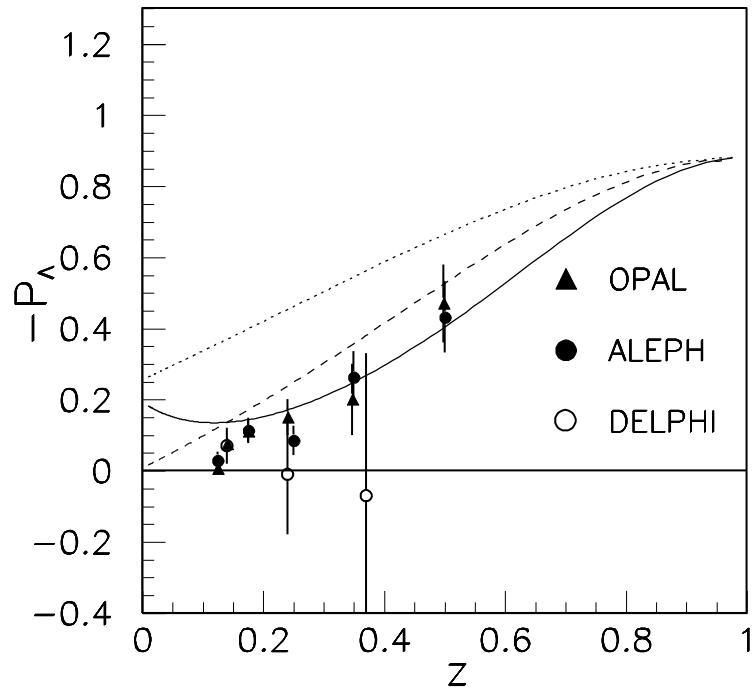
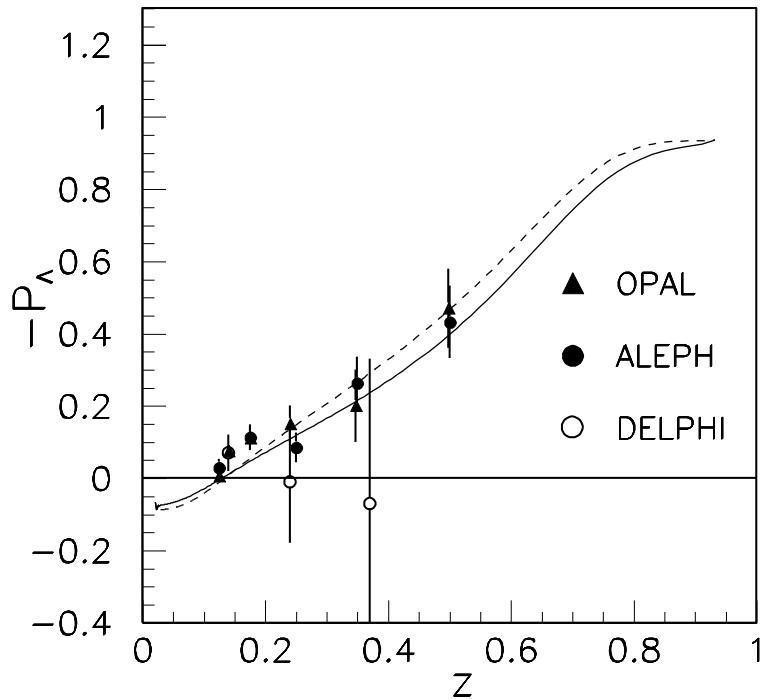
" Λ Physics"

Spin structure of Λ from Λ polarization in Z^0 decay



Conclusion 2: Spin Structure of Λ

Diquark model and pQCD results



**B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Rev. D 61 (2000) 034017**

Conclusion 3: Relating Λ -production with nucleon strangeness

- The spin transfer process of $\vec{p}p \rightarrow \vec{\Lambda}X$ is feasible to study strange-antistrange polarizations of the nucleon.
- The fitting to STAR data suggests:

$$\Delta s \neq \Delta \bar{s}$$

$$\Delta s \approx -0.025 \pm 0.019$$

$$\Delta \bar{s} \approx -0.001 \pm 0.012$$

- The results are compatible with the light-cone baryon-meson fluctuation model prediction.

Conclusion 4: Anti-Lambda Production for Nuclear Physics

Anti-Lambda production charged lepton semi-inclusive deep inelastic scattering off nuclear target is ideal to figure out the nuclear sea content, which is differently predicted by different models accounting for the nuclear EMC effect.

B.Lu, B.-Q. Ma, Phys.Rev.C74 (2006) 055202

C.Gong, B.-Q. Ma, Phys.Rev.C97 (2018) 065207