

Production of Light Nuclei in Au+Au Collisions with the STAR BES-II Program

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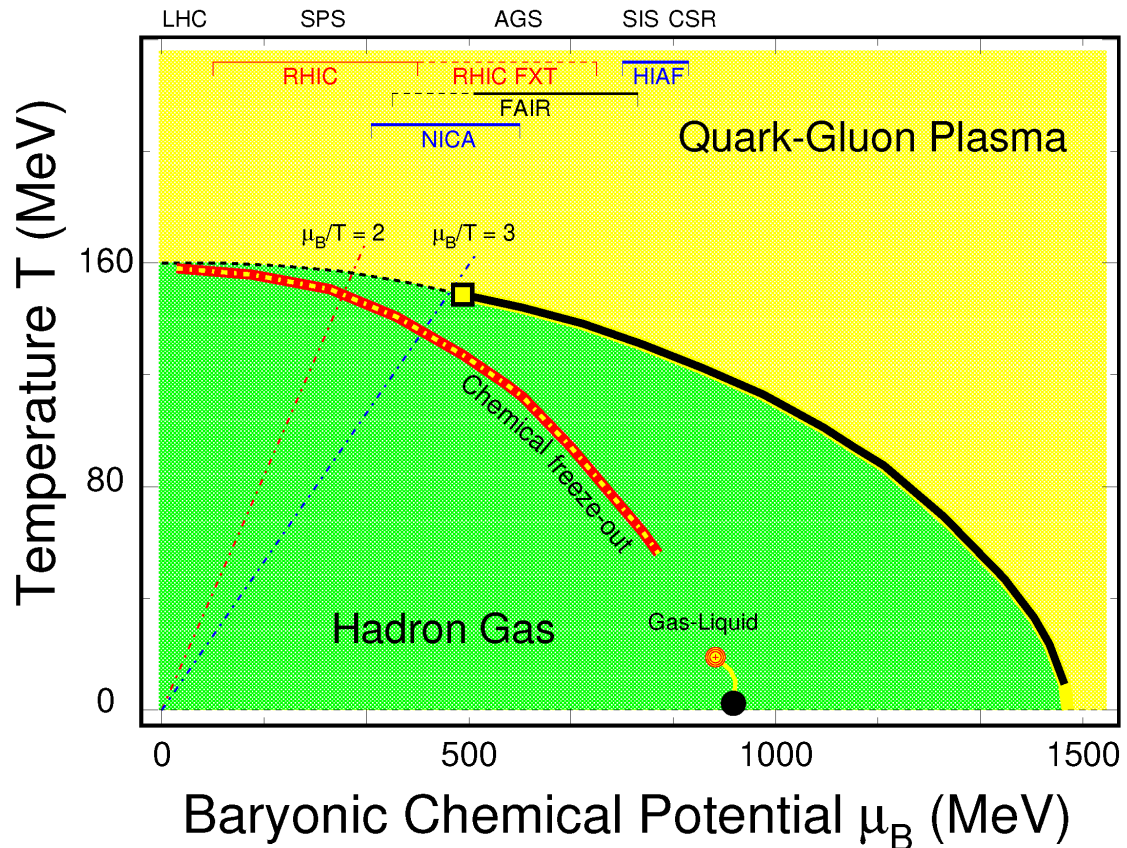
Central China Normal University

October 12, 2024

Outline

- Motivation
- The STAR Experiment
 - Dataset and Particle Identification
- Results in BES-I, FXT and Isobar
- Results in BES-II
 - Transverse Momentum Spectra
 - Particle Yields and Ratios
 - Coalescence Parameters
 - Nuclear Modification Factors
- Summary and Outlook

Motivation – QCD Phase Diagram and HIC



Beam Energy Scan Program at RHIC:

- ❖ Control beam energy and centrality to vary initial T and μ_B .
- ❖ Create QGP and explore its properties.
- ❖ Map out the crossover and/or 1st order QCD phase boundary.
- ❖ Search for the signatures of possible QCD critical point.

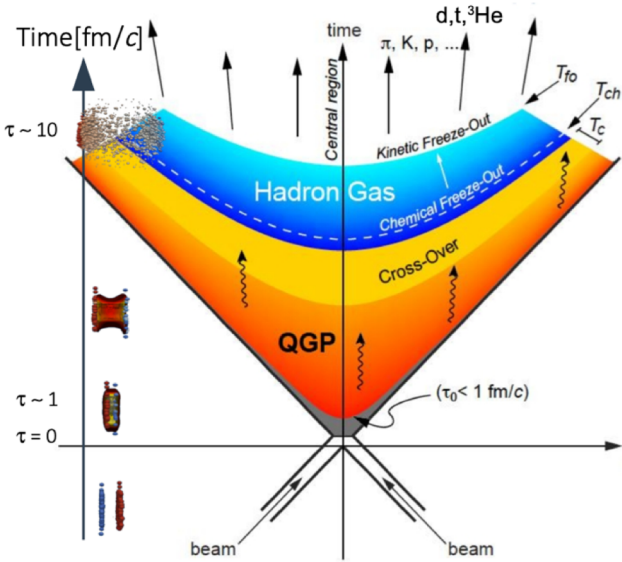
A. Bzdak, S. Esumi, V. Koch, J. F. Liao, M. Stephanov, and N. Xu. *Physics Reports*, 853 (2020) 1–87.

X. Luo, N. Xu, *Nucl. Sci. Tech.* 28, (2017) 112

X. Luo, Q. Wang, N. Xu, P. F. Zhuang. *Properties of QCD Matter at High Baryon Density*. Springer, 2022.

<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0598>

Motivation – Light Nuclei



1. Why Light Nuclei?

- ❖ May carry information about local baryon density fluctuations.
- ❖ Provide an effective probe to study 1st order phase boundary and the QCD critical point.

2. Observable : Yield ratio of light nuclei ($N_t \times N_p / N_d^2$)

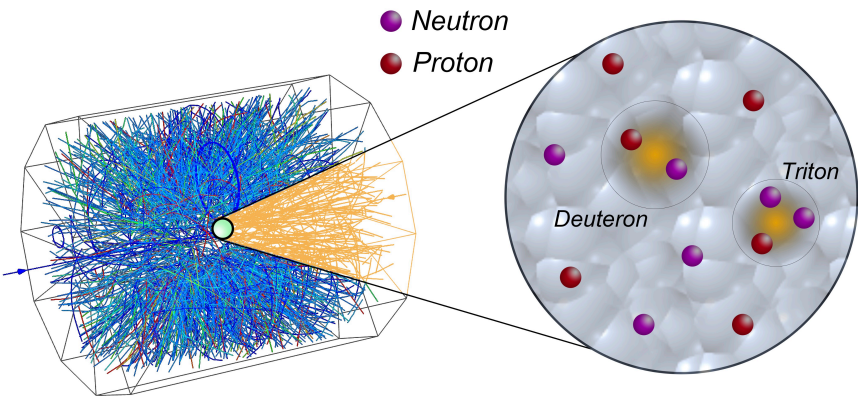
- ❖ Based on coalescence model:

$$N_A = g_c \int d\Gamma \rho_s(\{x_i, p_i\}) \times W_A(\{x_i, p_i\})$$

- ❖ The yield ratio is related to neutron density fluctuations:

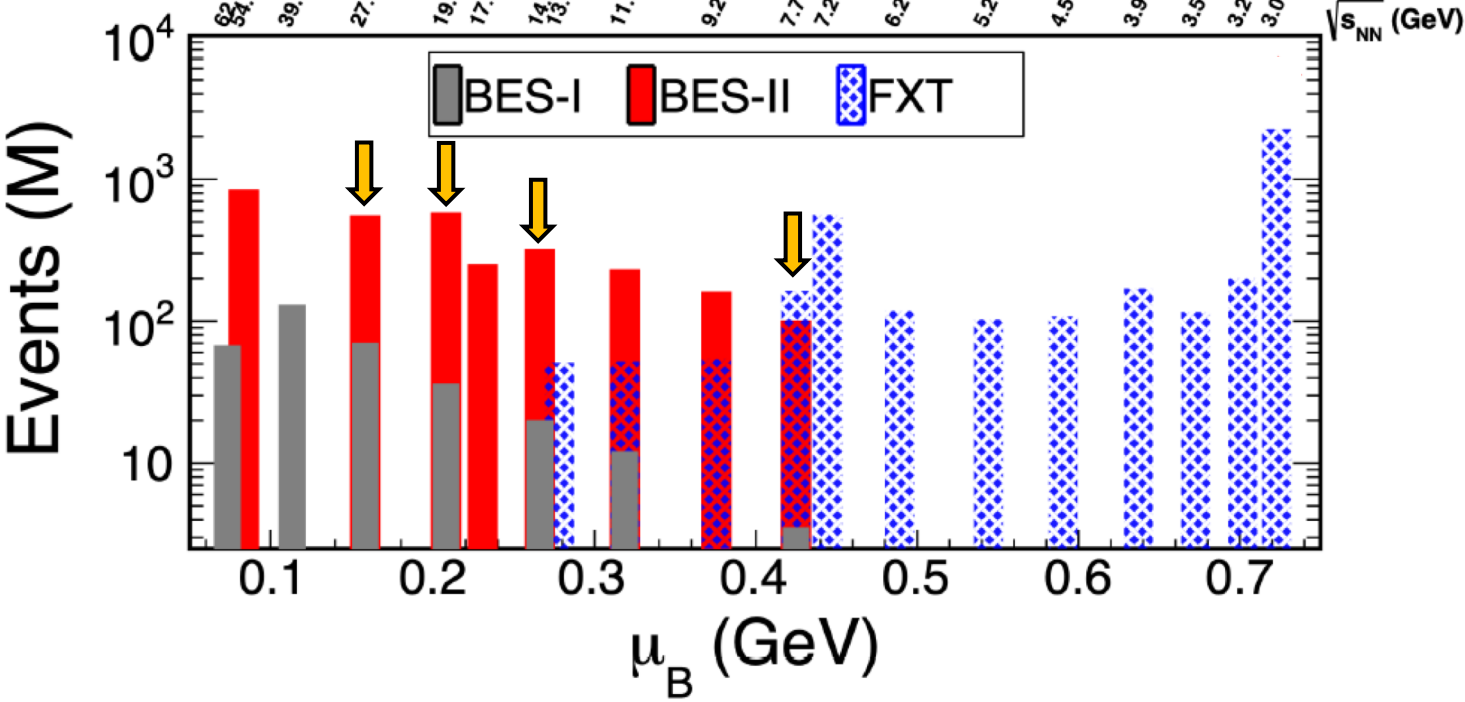
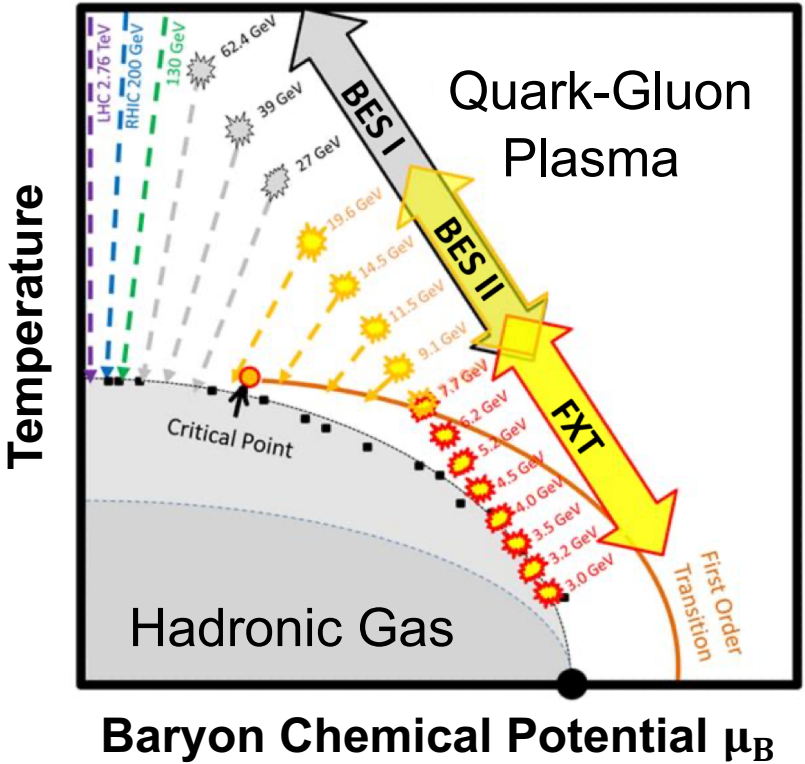
$$N_t \times N_p / N_d^2 \approx g(1 + \Delta n)$$

factor $g = \frac{1}{2\sqrt{3}}$ comes from the thermal equilibrium assumption of nucleon abundances.



K. Sun et al. Phys.Lett.B 774 (2017) 103-107
E. Shuryak et al. Phys.Rev.C 101 (2020) 3, 034914

RHIC Beam Energy Scan Program

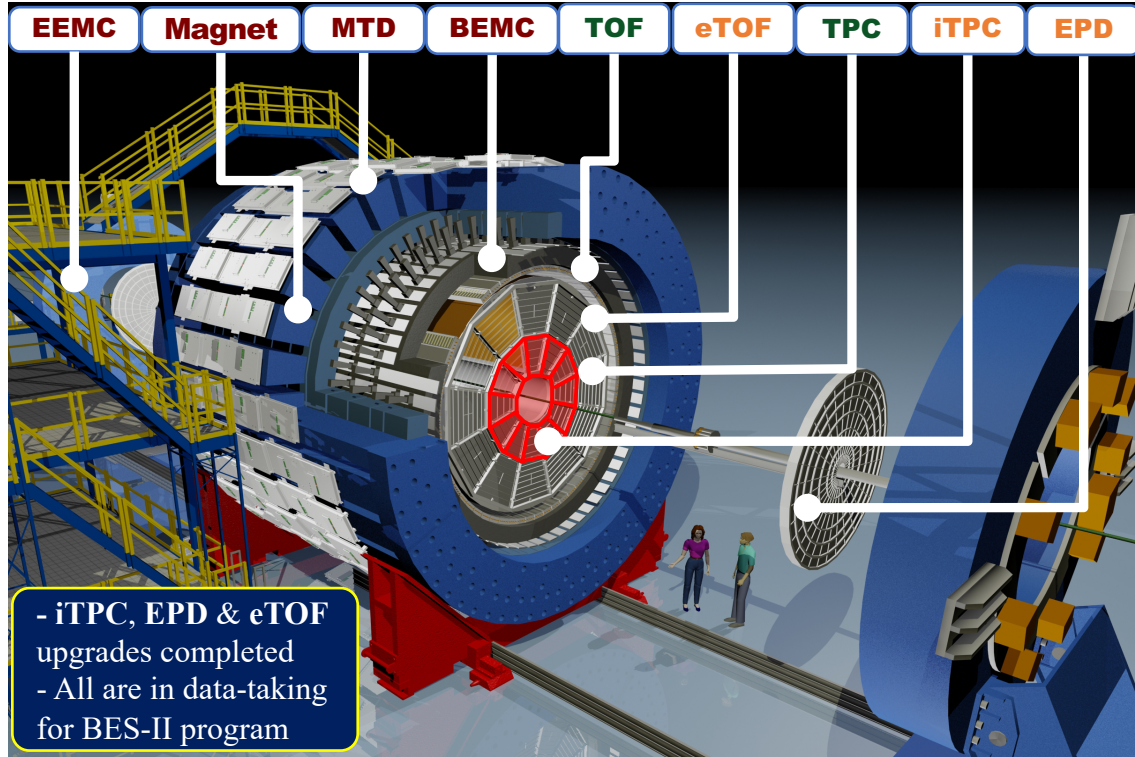


- ❖ STAR has completed BES-II data-taking with factors of 10 – 20 more statistics compared to BES-I.
- ❖ BES-II: 8 collider energies ($\sqrt{s_{NN}} = 7.7 - 54 \text{ GeV}$) / 12 FXT energies ($\sqrt{s_{NN}} = 3.0 - 13.7 \text{ GeV}$)
- ❖ μ_B coverage : $25 < \mu_B < 750 \text{ MeV}$.

STAR Collaboration, arXiv:1007.2613
<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0493>
<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0598>

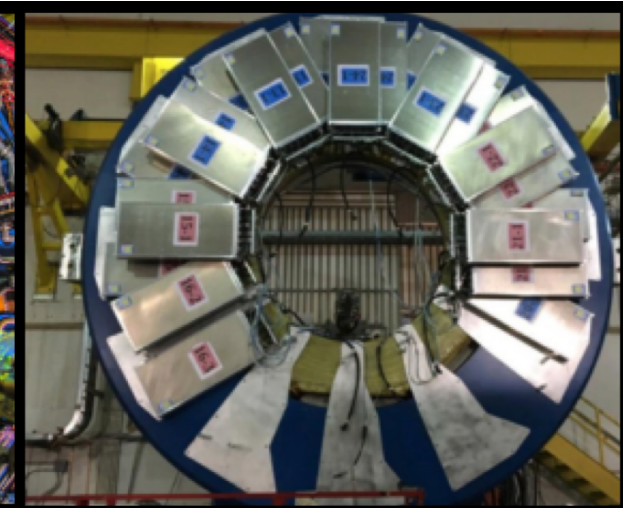
The Solenoidal Tracker At RHIC (STAR)

Detector upgrades for STAR BES-II



iTPC:

- Improves dE/dx
- Extends η coverage from 1.0 to 1.5
- Lower p_T cut-in from 125 to 60 MeV/c
- Ready in 2019



eTOF:

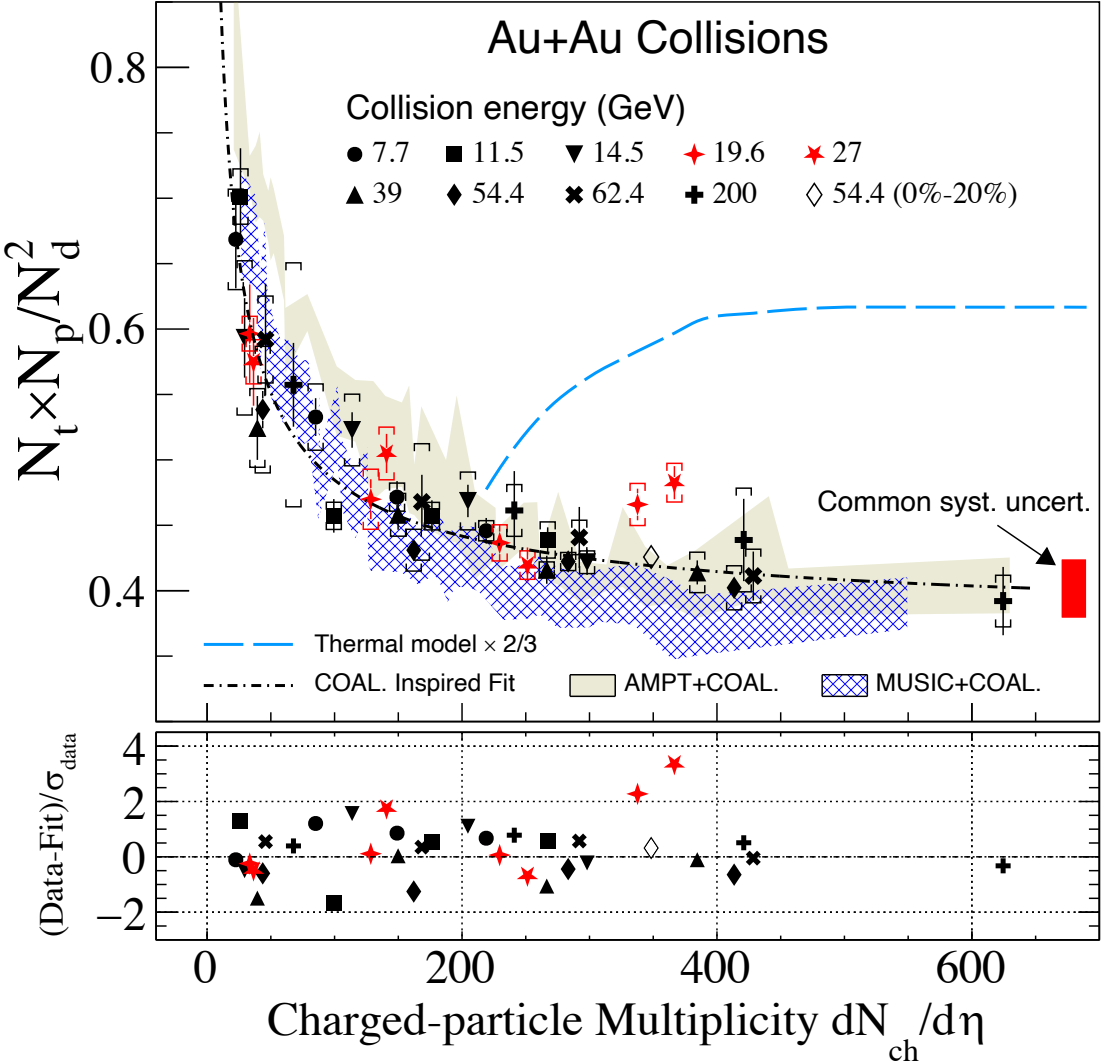
- Forward rapidity coverage
- η coverage from 1.0 to 1.5
- Borrowed from CBM-FAIR
- Ready in 2019

iTPC: <https://drupal.star.bnl.gov/STAR/starnotes/public/sn0619>

eTOF: [STAR and CBM eTOF group, arXiv: 1609.05102](https://arxiv.org/abs/1609.05102)

- ❖ Enlarge the rapidity acceptance
- ❖ Improve particle identification
- ❖ Lower p_T cut-in to reduce uncertainty in spectra extrapolation

Results from RHIC BES-I



STAR Collaboration, Phys.Rev.Lett. 130 (2023) 202301

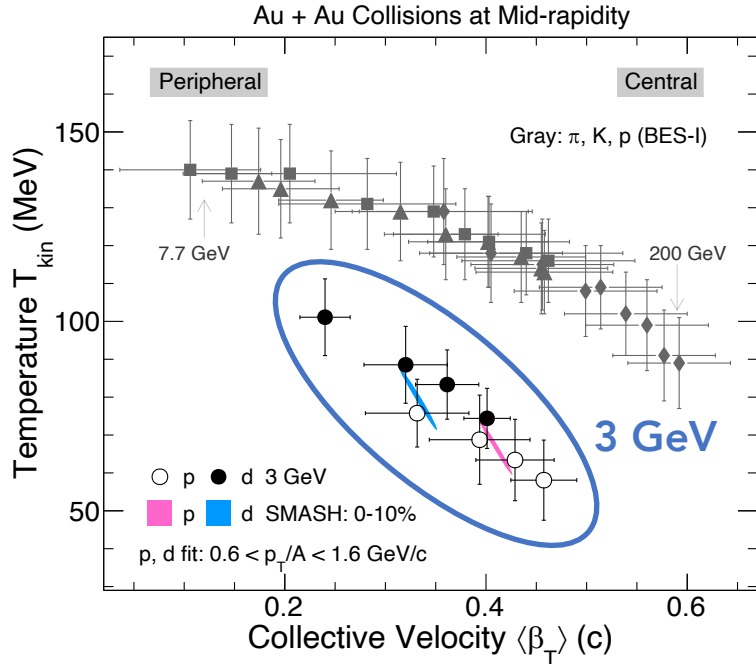
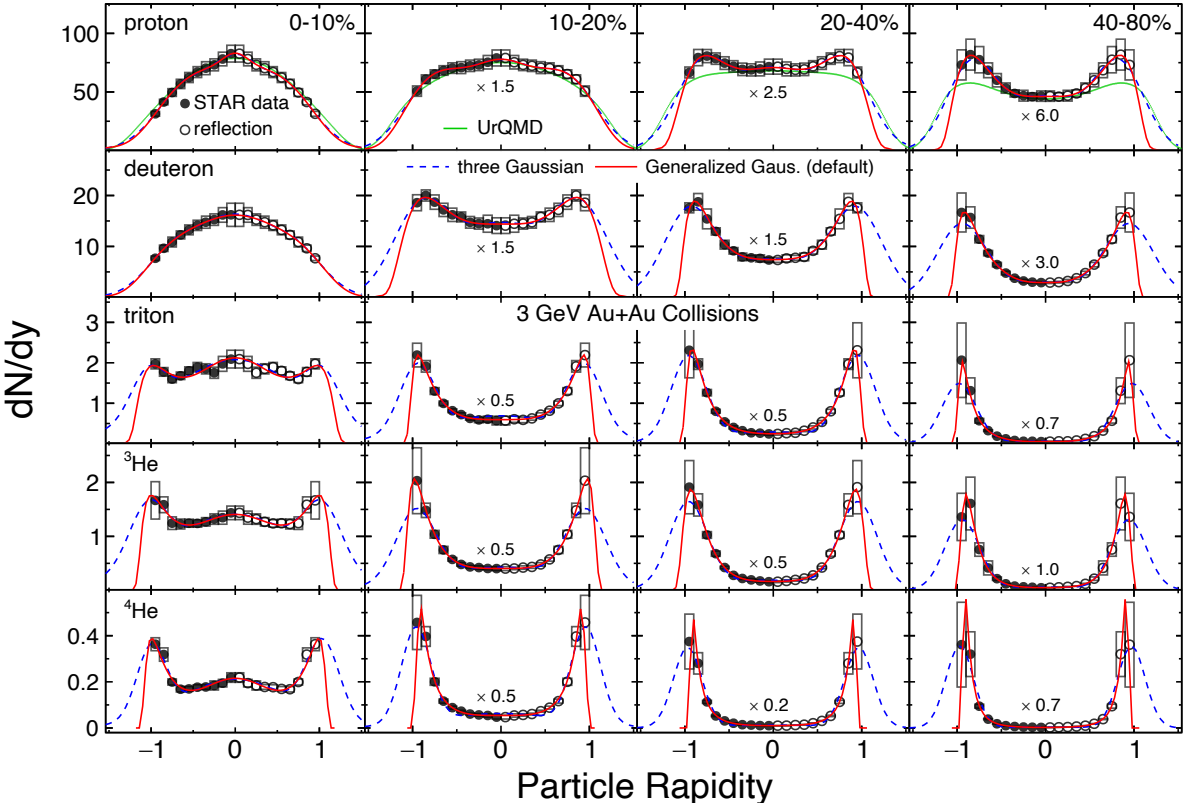
- The yield ratio exhibits a scaling behavior: A trend driven by the interplay between the size of light nuclei and the size of the fireball created in HIC.

$$\frac{N_t \times N_p}{N_d^2} = p_0 \times \left(\frac{R^2 + \frac{2}{3}r_d^2}{R^2 + \frac{1}{2}r_t^2} \right)^3, \text{ where } R \propto (dN_{ch}/d\eta)^{1/3}.$$

W. Zhao, K. J. Sun, C. M. Ko and X. Luo, Phys. Lett. B 820 (2021) 136571

- The ratios at $\sqrt{s_{NN}} = 19.6$ and 27 GeV in 0-10% centrality show enhancements with respect to the coalescence baseline with a combined significance of 4.1σ .
- The thermal model overestimates the experimental data and shows a clear difference compared to the coalescence model.

Results from RHIC FXT at $\sqrt{s_{NN}} = 3 \text{ GeV}$



$$\text{Blast-Wave: } \frac{1}{2\pi p_T} \frac{d^2N}{dp_T dy} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}} \right)$$

STAR Collaboration, arXiv: 2311.11020
 K.J. Sun, W.H. Zhou, L.W. Chen, C.M. Ko, F. Li, R. Wang, J. Xu, arXiv: 2205.11010

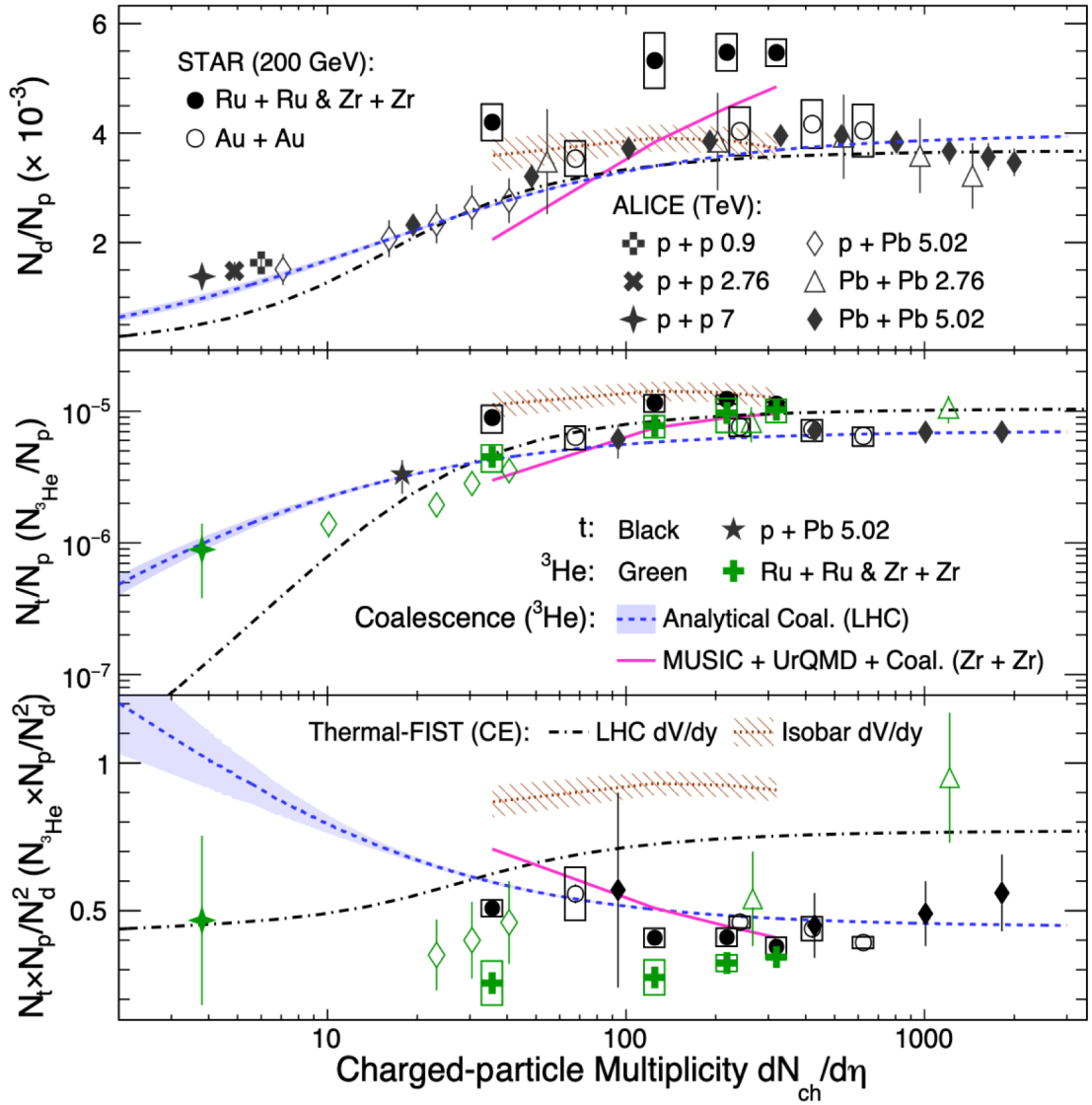
- The centrality dependence of rapidity density is attributed to the interplay between baryon stopping and the spectators' contribution.
- At 3 GeV Au+Au collisions, the kinetic freeze-out parameters (T_{kin} , $\langle \beta_T \rangle$) show different trend compared to that of higher energy collisions.
- The freeze-out parameter (T_{kin}) of deuteron is systematically higher than that of proton at 3 GeV, which is different from higher energies.

➔ Indicate a different equation of state (EoS)

Similar trend seen in SMASH Model

$$T_{kin}^d > T_{kin}^p$$

Results from RHIC Isobar

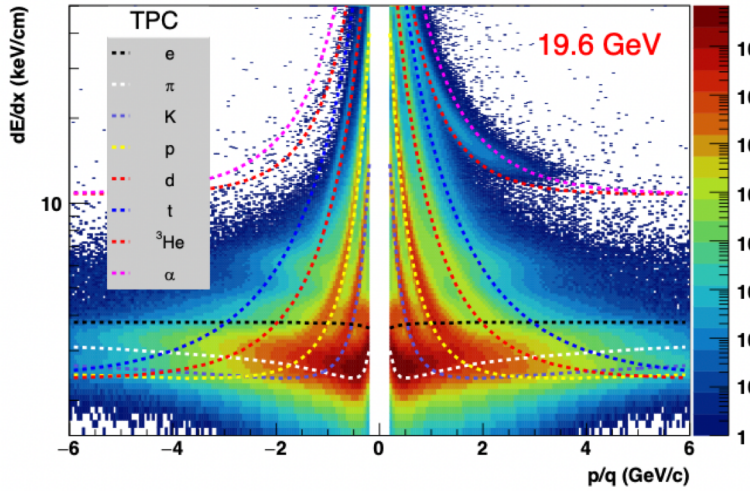


- ❖ Deuteron production are enhanced in Ru\Zr + Ru\Zr collisions compared with other systems.
- ❖ Triton yield deviate from ³He yield in peripheral collision which may effected by neutron skin.
- ❖ $N_t \times N_p / N_d^2$ are overestimated by the thermal model.

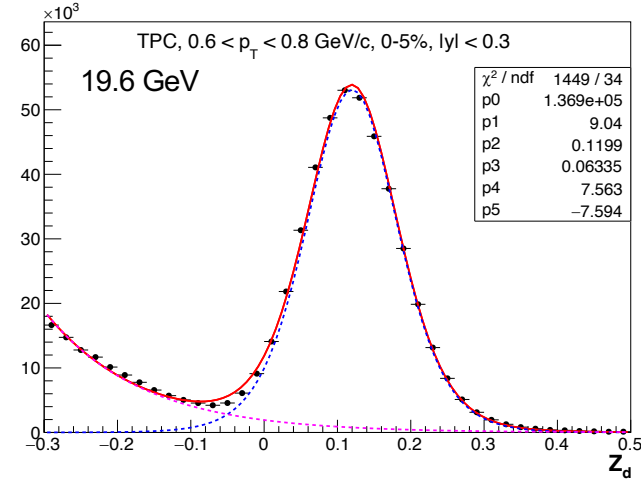
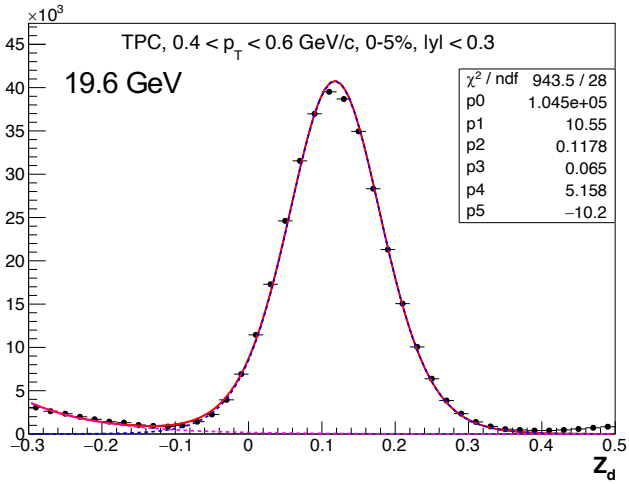
V. Vovchenko, et al. *Phys.Lett.B* 785 (2018) 171-174
 K-J. Sun, et al. *Phys.Lett.B* 792 (2019) 132-137; arXiv:2404.02701
 [STAR Collaboration] *Phys.Rev.C* 99 (2019) 6, 064905; *Phys.Rev.Lett.* 130 (2023) 202301
 [ALICE Collaboration] *Phys.Rev.C* 88 (2013) 044910; *Phys.Rev.C* 93 (2016) 2, 024917;
Phys.Rev.C 107 (2023) 6, 064904

Particle Identification and Signal Extraction

$$\text{TPC: } z = \log \left(\frac{\langle dE/dx \rangle_{\text{measure}}}{\langle dE/dx \rangle_{\text{Bichsel}}} \right)$$



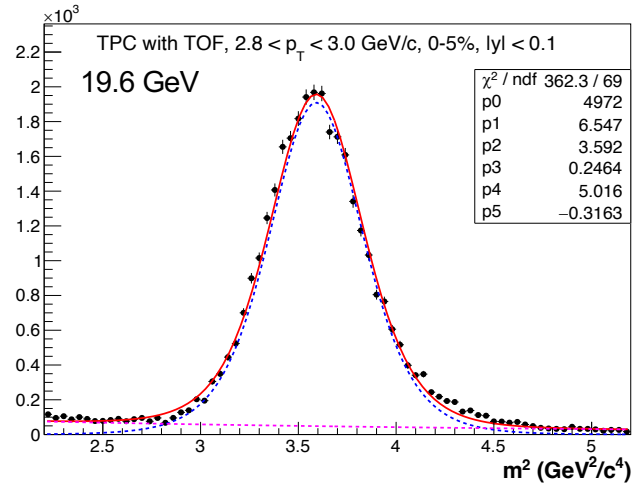
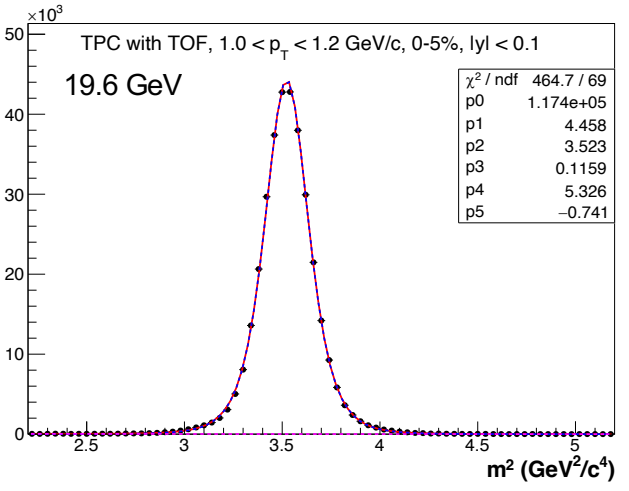
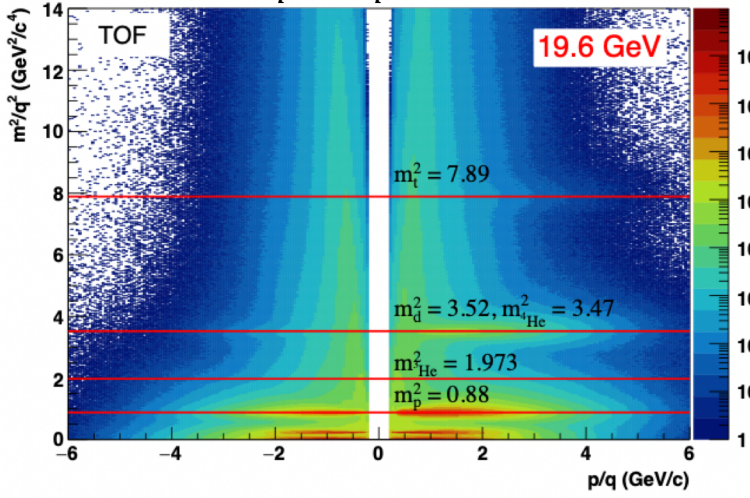
Signal Extraction Examples (19.6 GeV, deuteron, 0-5%)



Low p_T : TPC

Signal: Gaussian (Blue)
BG: Gaussian (Magenta)
Total: (Red)

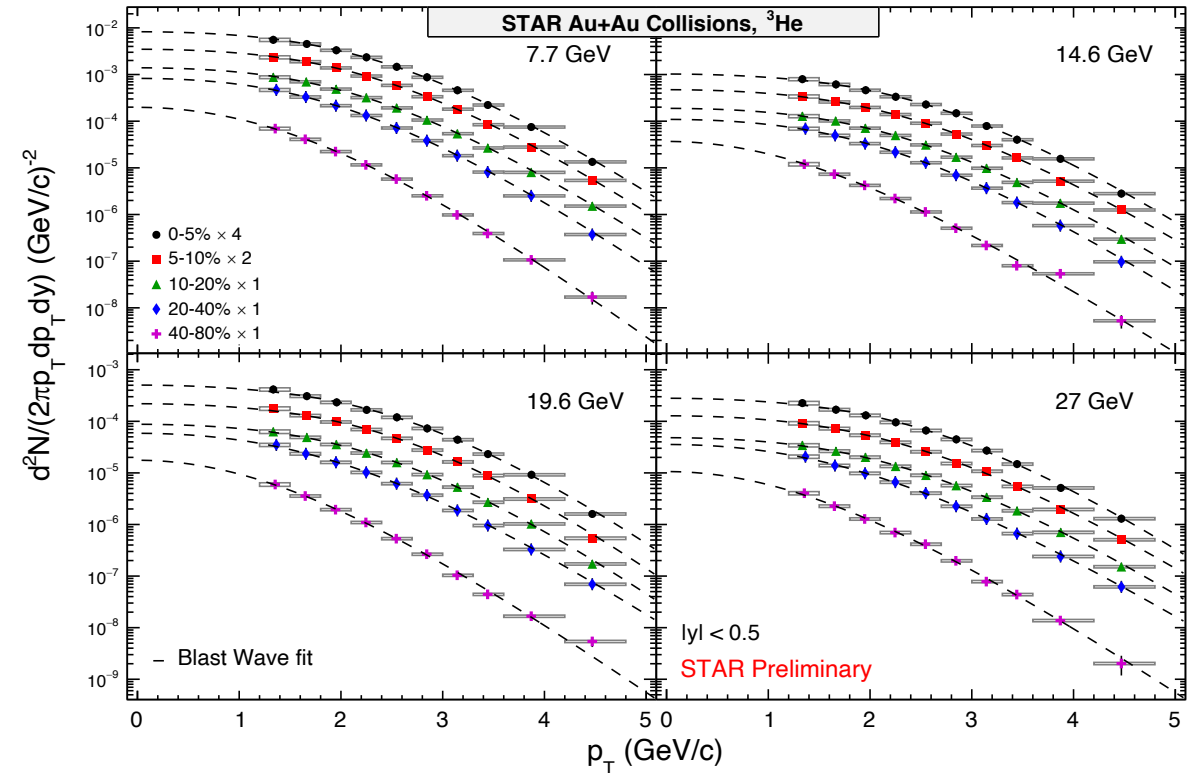
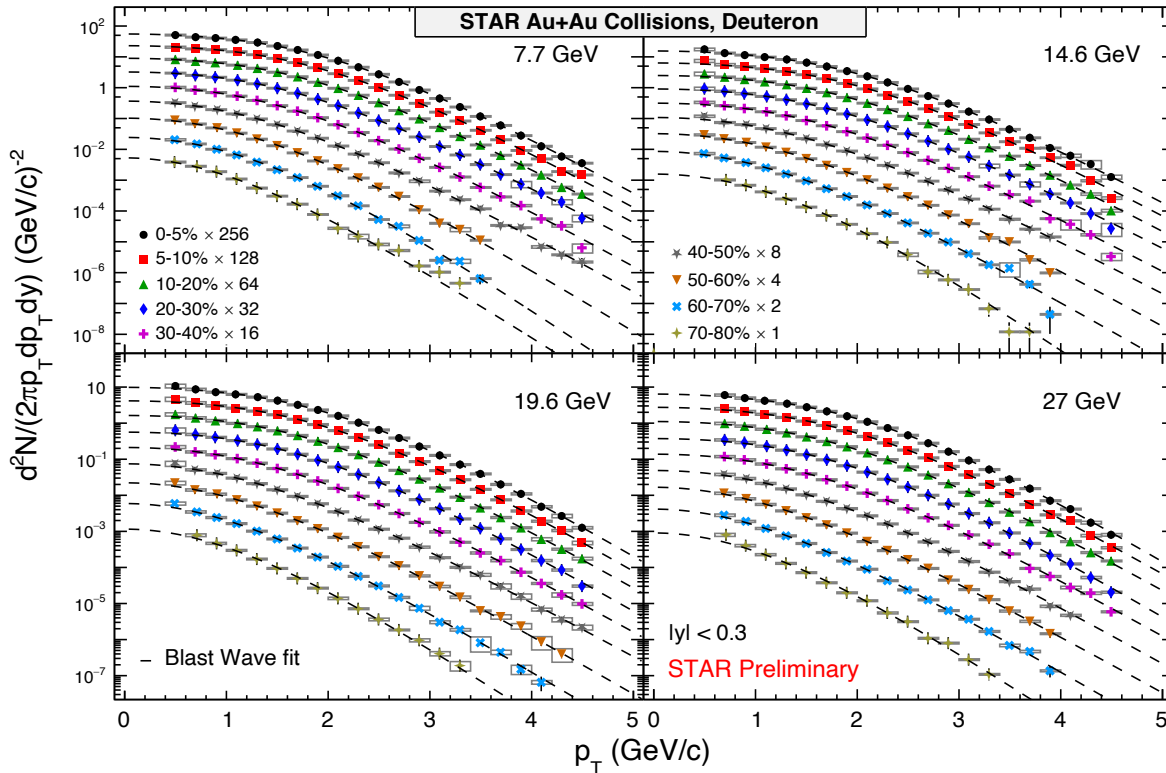
$$\text{TOF: } \frac{m^2}{q^2} = \frac{p^2}{q^2} \left(\frac{c^2 t^2}{L^2} - 1 \right)$$



High p_T : TPC with TOF

Signal: Student-t (Blue)
BG: Gaussian (Magenta)
Total: (Red)

Transverse Momentum Spectra



❖ Blast-Wave Function

$$\frac{1}{2\pi p_T} \frac{d^2N}{dp_T dy} \propto \int_0^R r dr m_T I_0 \left(\frac{p_T \sinh \rho}{T_{kin}} \right) K_1 \left(\frac{m_T \cosh \rho}{T_{kin}} \right)$$

$$\rho = \tanh^{-1} \beta_r, \beta_r(r) = \beta_T \left(\frac{r}{R} \right)^n, n \text{ fix to } 1 \text{ in this analysis}$$

Freeze-out parameters:

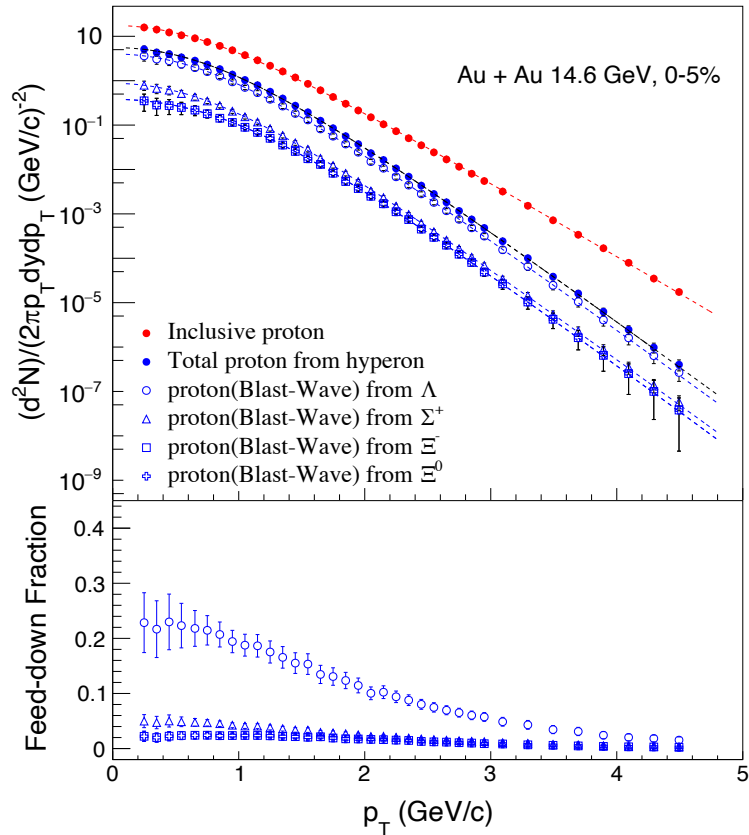
T_{kin} : kinetic freeze-out temperature

$\langle \beta_T \rangle$: average radial flow velocity

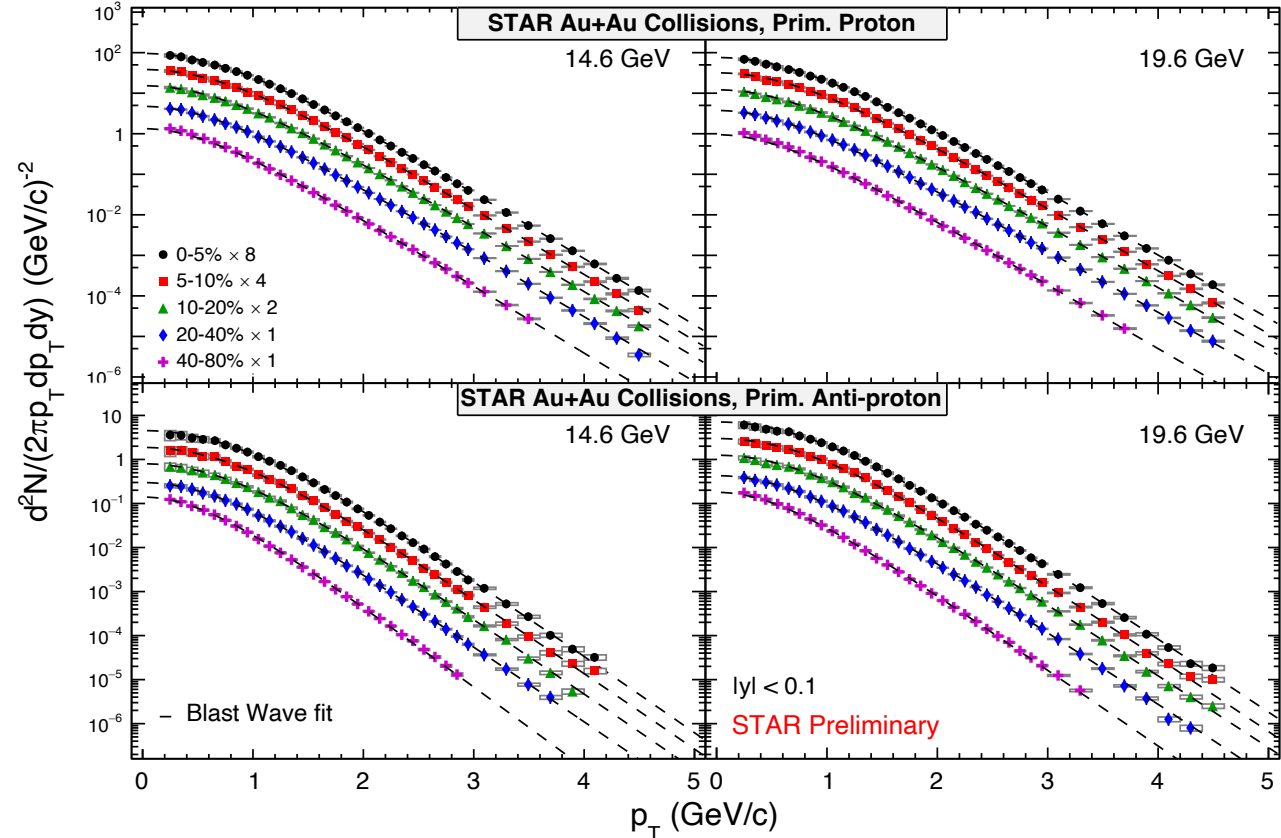
I_0 and K_1 : from Bjorken Hydrodynamic assumption

- All efficiencies and corrections are included.
- The statistical and systematic uncertainties are shown as vertical lines and boxes, respectively.
- The p_T ranges are extended in BES-II, which lead to smaller systematic uncertainties in particle yields.

(anti-)Proton Weak Decay Feed-down Correction

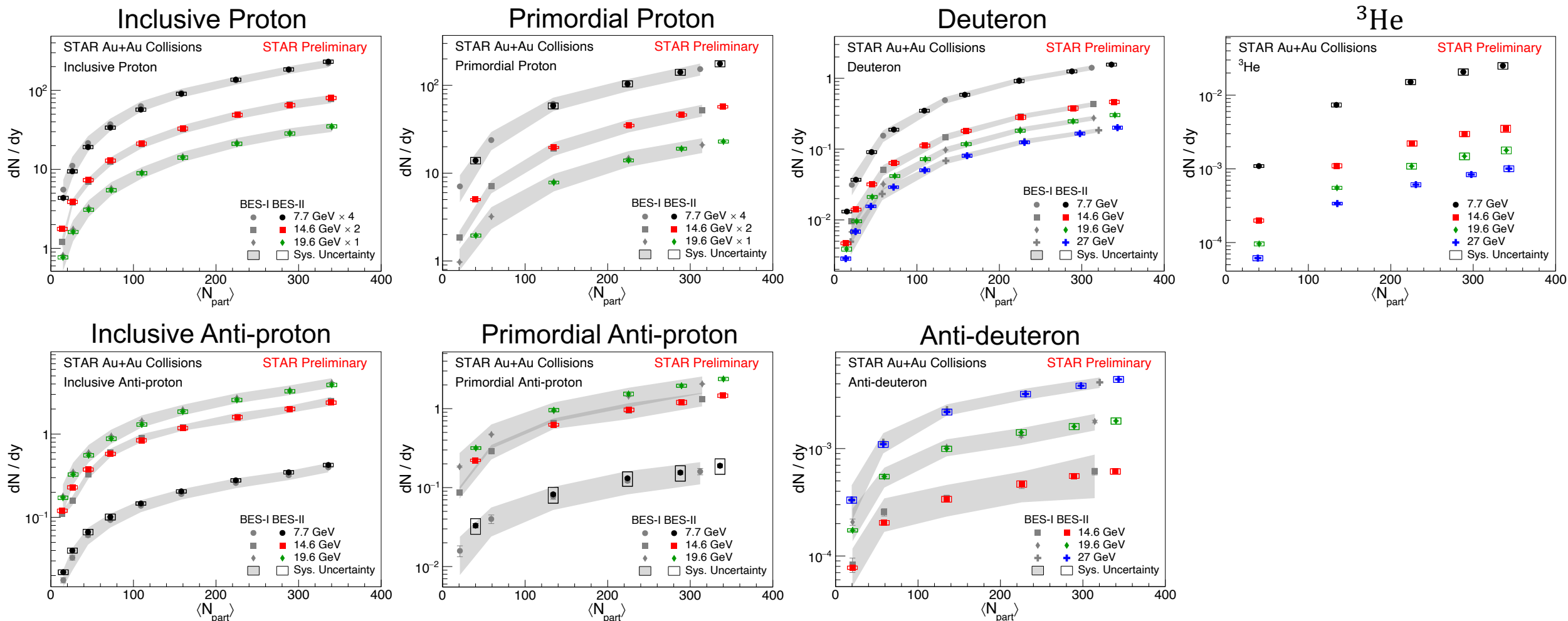


- $\Lambda \rightarrow p + \pi^-$, branching ratio = 63.9 %
- $\Sigma^+ \rightarrow p + \pi^0$, branching ratio = 51.57 %
- $\Xi^- \rightarrow \Lambda + \pi^-$, branching ratio = 99.887 %
- $\Xi^0 \rightarrow \Lambda + \pi^0$, branching ratio = 99.524 %



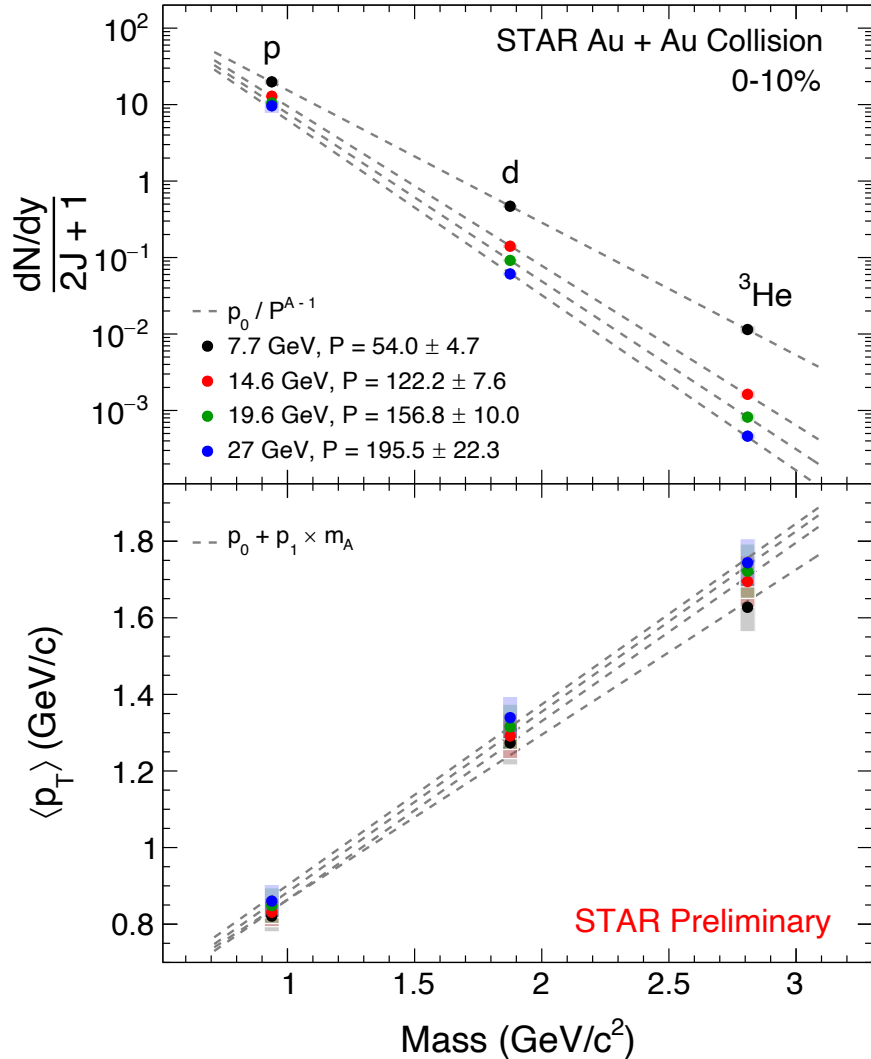
- The primordial spectra were obtained by subtracting the (anti-)proton weak decayed from strange hadrons.
- Data driven method: Use STAR published strange particle ($\Lambda, \Sigma^+, \Xi^-, \Xi^0$) yields and embedding simulation samples.
- The spectra of Σ^+ : Obtained by multiplying the Λ spectra by a factor of 0.27.
- The spectra of Ξ^0 : Assumed to be the same as those of Ξ^- .

Particle Yields and Ratios



- The systematic uncertainties are reduced in BES-II because the p_T ranges are extended.
- Yields for light nuclei increase from peripheral to central collisions.
- dN/dy for positive particles decrease with increasing energy, while the behavior is opposite for antiparticles.

Particle Yields and Ratios

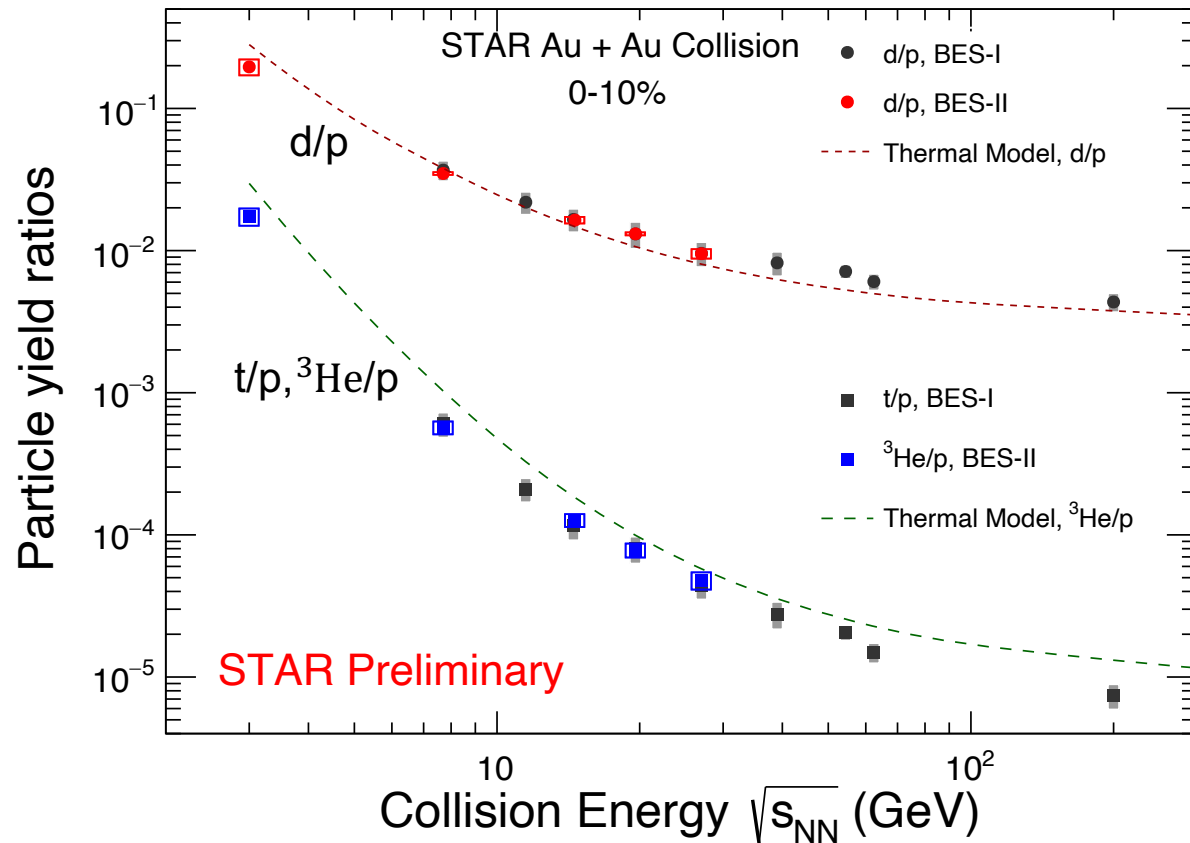


- Mass dependence of $\frac{dN/dy}{2J+1}$:
 - Fitted by an exponential function form p_0/P^{A-1} , where P is the penalty factor and can be determined by the Boltzmann factor $e^{(m_N - \mu_B)/T}$ in thermal model.
 - The production of light nuclei are proportional to the spin degeneracy.
 - The penalty factor is larger at higher beam energy, which indicates that it is harder to form high-mass objects.

E864 Collaboration, Phys.Rev.Lett. 83 (1999) 5431-5434
STAR Collaboration, Phys.Rev.Lett. 130 (2023) 202301

- Mass dependence of $\langle p_T \rangle$:
 - Fitted by an linear function form $p_0 + p_1 \times m_A$.
 - The $\langle p_T \rangle$ increases linearly with increasing mass of the particles.

Particle Yields and Ratios



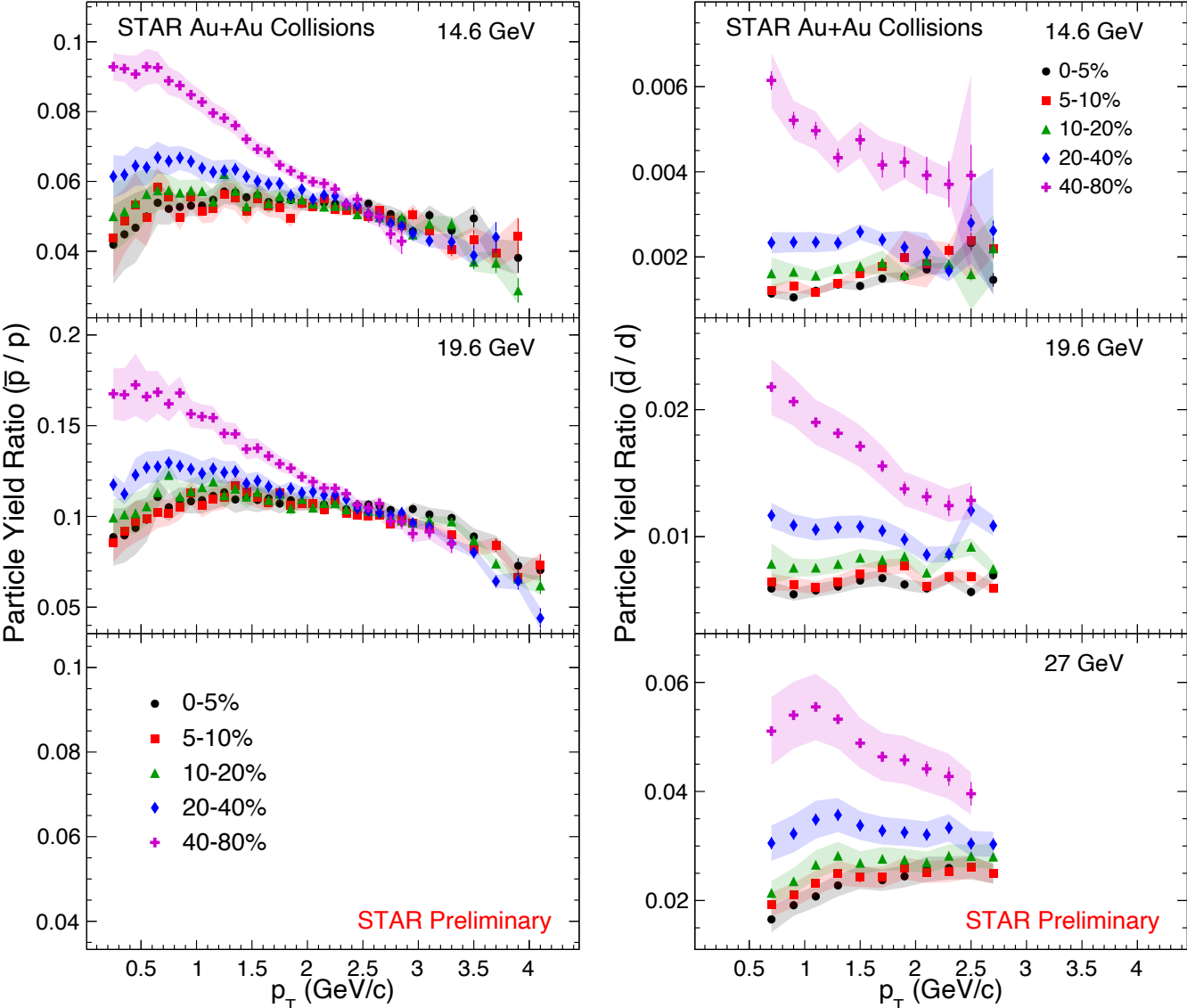
STAR, *Phys. Rev. Lett.* 130, 202301 (2023)

- The particle ratios are consistent with BES-I.
- The particle ratios show a monotonic decrease with collision energy.
- d/p ratio can be described well by the thermal model.
- The thermal model overestimates t/p and $^3\text{He}/p$ by a factor of approximately two.
 - The hadronic re-scatterings may play a crucial role during the hadronic expansion phase which lead to this discrepancy.

V. Vovchenko, B. Dönigus, B. Kardan, M. Lorenz, and H. Stoecker, *Phys. Lett. B*, (2020) 135746;

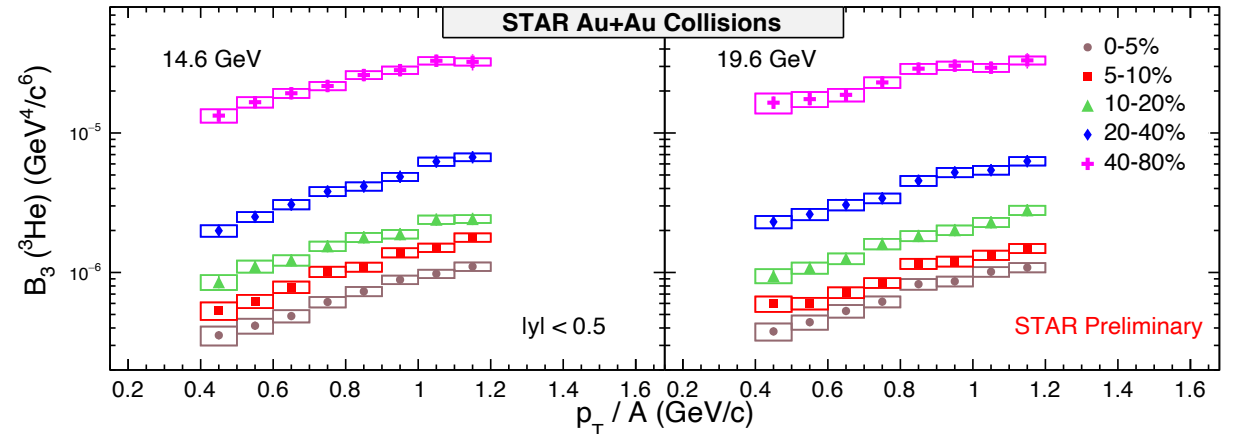
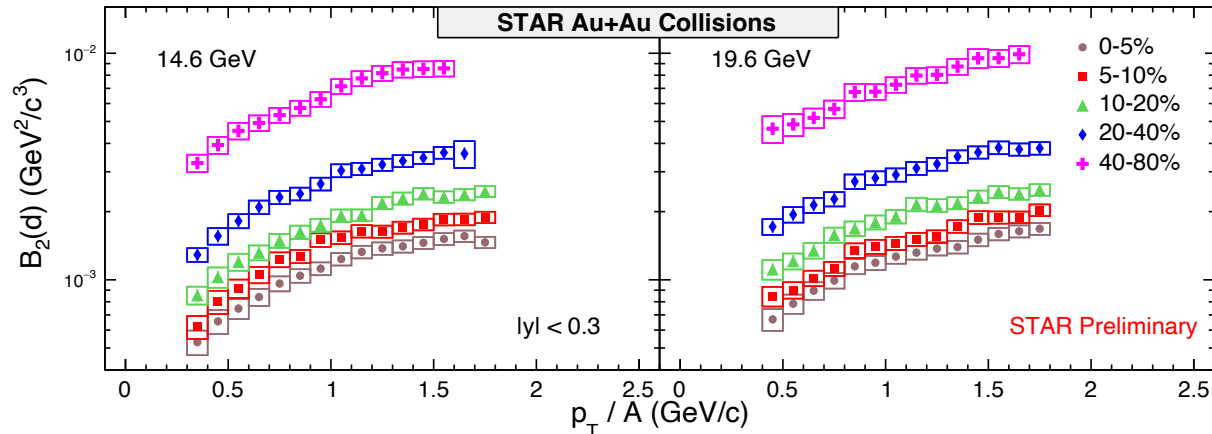
K.J. Sun, R. Wang, C. M. Ko, Y.G. Ma, and C. Shen, *Nature Commun.*, 15 (2024) 1, 1074

Particle Yields and Ratios



- The statistical and systematic uncertainties are shown as vertical lines and color bands, respectively.
- \bar{p}/p and \bar{d}/d ratios show strong centrality dependence.
 - This could be due to the annihilation between the particles and antiparticles.

Coalescence Parameter



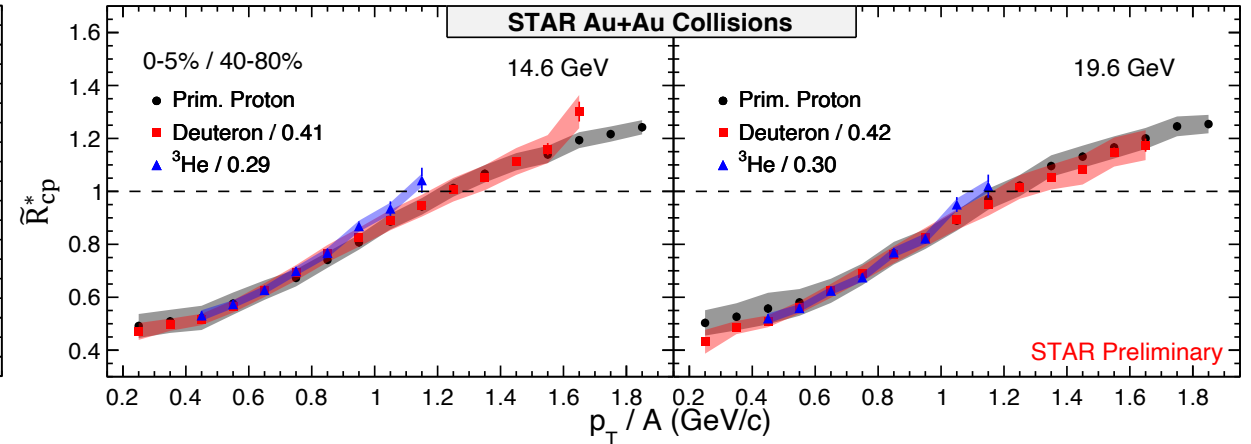
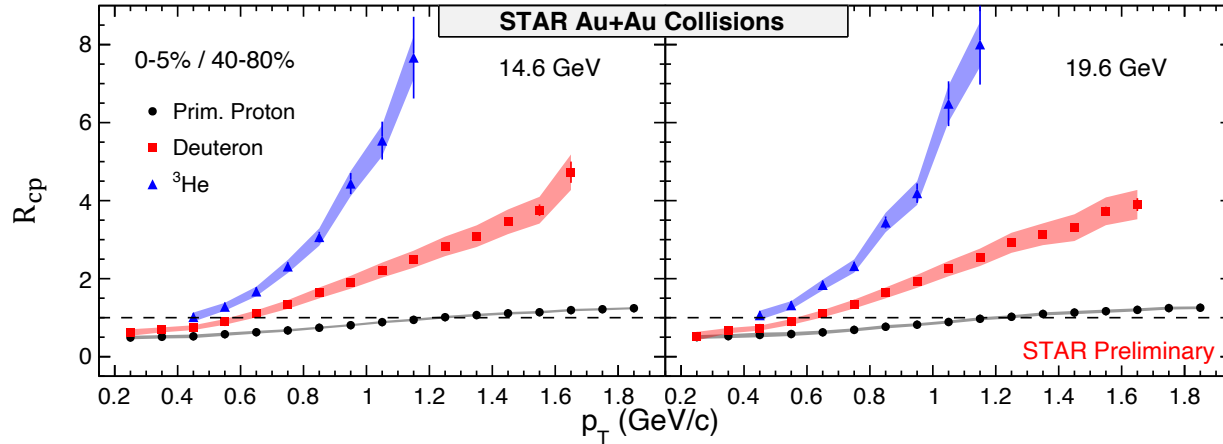
- ❖ In the coalescence picture, the invariant yield of light nuclei is proportional to the invariant yield of nucleons. The coalescence parameter B_A reflects the probability of nucleon coalescence, which is related to the local nucleon density.

$$E_A \frac{d^3 N_A}{d^3 p_A} = B_A \left(E_p \frac{d^3 N_p}{d^3 p_p} \right)^Z \left(E_n \frac{d^3 N_n}{d^3 p_n} \right)^{A-Z} \approx B_A \left(E_p \frac{d^3 N_p}{d^3 p_p} \right)^A.$$

*R. Scheibl and U. Heinz Phys.Rev.C 59 (1999) 1585-1602
STAR Collaboration, Phys.Rev.C 99 (2019) 6, 064905*

- B_A increase with increasing p_T which might indicate an expanding collision system.
- B_A increase from central to peripheral collisions, which can be explained by a decreasing source volume.

Nuclear Modification Factors (R_{cp})



$$\diamond R_{cp}(p_T) = \frac{\langle N_{coll} \rangle_p}{\langle N_{coll} \rangle_c} \times \frac{d^2 N_{AA}^c / dp_T dy}{d^2 N_{AA}^p / dp_T dy}, \quad \langle N_{coll} \rangle \text{ is the average number of binary nucleon-nucleon collisions per event.}$$

❖ Number of constituent nucleon (NCN) scaling for $R_{cp}(p_T)$:

$$R_{cp}^*(p_T) = \left(\frac{B_{A, \text{Central}}}{B_{A, \text{Peripheral}}} \right)^{-1/A} \left(R_{cp}(Ap_T) \right)^{1/A} \left(\frac{\langle N_{coll} \rangle_c}{\langle N_{coll} \rangle_p} \right)^{1/A-1} \equiv \left(\frac{B_{A, \text{Central}}}{B_{A, \text{Peripheral}}} \right)^{-1/A} \tilde{R}_{cp}^*(p_T)$$

C. S. Zhou, Y. G. Ma, and S. Zhang. Eur.Phys.J.A 52 (2016) 12, 354

➤ R_{cp} for different particles exhibit distinct trends at all energies.

➤ In coalescence picture, the \tilde{R}_{cp}^* of light nuclei will follow a common trend when scaled by a constant factor, which is determined by $\tilde{R}_{cp}^* \text{ Nuclei} / \tilde{R}_{cp}^* \text{ Proton}$ at $p_T/A = 0.65 \text{ GeV}/c$.

Summary and Outlook

Summary:

- We report the light nuclei productions (p , d , ${}^3\text{He}$, \bar{p} and \bar{d}) in Au+Au collisions at $\sqrt{s_{\text{NN}}} = 7.7 - 27$ GeV from RHIC STAR BES-II.
- The particle ratios N_d/N_p and $N_{{}^3\text{He}}/N_p$ have been measured.
 - The particle ratios show a monotonic decrease with collision energy.
 - The thermal model over-predicts t/p and ${}^3\text{He}/p$ by a factor of about 2.
- The coalescence parameter B_A have been measured.
 - Collective expansion leads to an increase in B_A from low to high p_T
 - Decreasing source volume results in a rise of B_A from central to peripheral collisions.
- The nuclear modification factor R_{cp} of light nuclei shows a scaling behavior, which is consistent with a nucleon coalescence mechanism of light nuclei production.

Outlook:

- ❖ Working on the compound ratios ($N_p \times N_t/N_d^2$ and $N_p \times N_{{}^3\text{He}}/N_d^2$) in BES-II.
- ❖ Continue the analysis of other BES-II (collider + FXT) energies.
- ❖ Working on the production of ${}^4\text{He}$ in BES-II.

Thank you for your attention!

Backup

Au+Au collisions at RHIC (Collider mode)				Au+Au collisions at RHIC (Fixed-Target)				
$\sqrt{s_{NN}}$ (GeV)	nEvents (M)	μ_B (MeV)	Time	$\sqrt{s_{NN}}$ (GeV)	E_{beam} (GeV)	nEvents (M)	μ_B (MeV)	Time
200	380	25	Run-10, 19	13.7	100	50	280	Run-21
62.4	46	75	Run-10	11.5	70	50	320	Run-21
54.4	1200	85	Run-17	9.2	44.5	50	370	Run-21
39	86	112	Run-10	7.7	31.2	260	420	Run-18, 19, 20
27	585	156	Run-11, 18	7.2	26.5	470	440	Run-18, 20
19.6	595	206	Run-11, 19	6.2	19.5	120	490	Run-20
17.3	256	230	Run-21	5.2	13.5	100	540	Run-20
14.6	340	262	Run-14, 19	4.5	9.8	110	590	Run-20
11.5	235	316	Run-10, 20	3.9	7.3	120	633	Run-20
9.2	160	372	Run-20	3.5	5.75	120	670	Run-20
7.7	104	420	Run-10, 21	3.2	4.59	200	699	Run-19
				3.0	3.85	2300	750	Run-18, 21

- ❖ STAR has completed BES-II data-taking with factors of 10-20 more statistics compared to BES-I.
- ❖ BES-II: 8 collider energies (7.7 – 54 GeV) / 12 FXT energies (3.0 - 13.7 GeV)
- ❖ μ_B coverage : $25 < \mu_B < 750$ MeV.

STAR, arXiv:1007.2613

<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0493>

<https://drupal.star.bnl.gov/STAR/starnotes/public/sn0598>