

# QCD临界点附近的动力学研究

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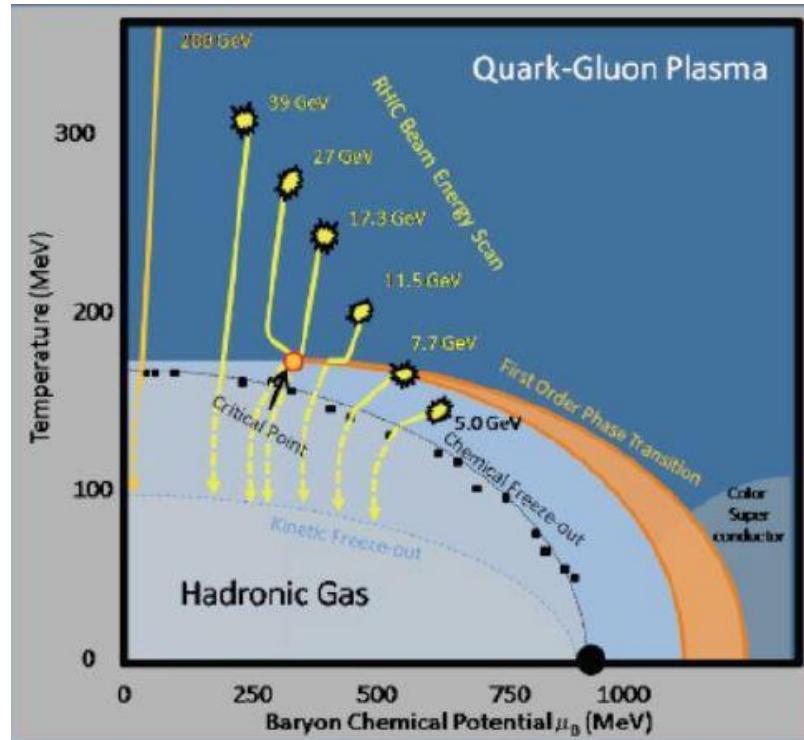
## STAR区域研讨会

◆会议时间◆

2024年10月10日-10月15日

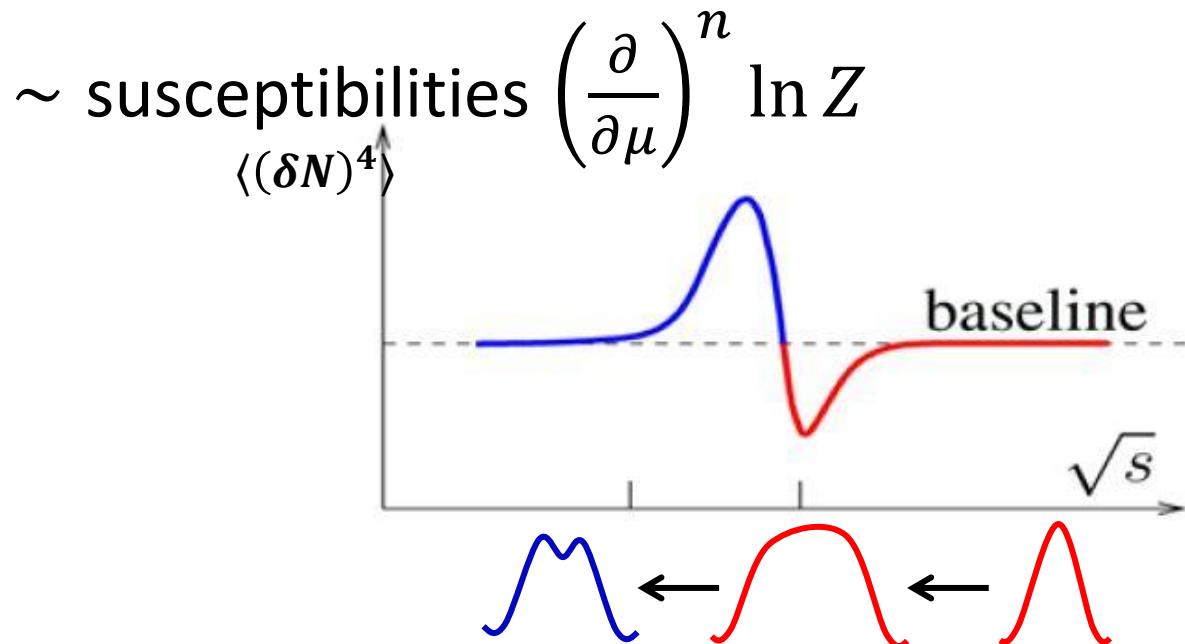
# QCD phase diagram

- **Lattice QCD** (small  $\mu_B$  finite  $T$ ):
    - Crossover
  - **Effective models** (large  $\mu_B$ )
    - 1<sup>st</sup> order phase transition
- **Critical point**
- Lattice QCD: sign problem at large  $\mu_B$
  - Effective models: parameters dependent
- **Heavy-ion collisions :**
- tuning  $\sqrt{s_{NN}}$ , mapping  $T - \mu$  phase diagram:  
RHIC(BES),NICA,FAIR,J\_PARC,HIAF....

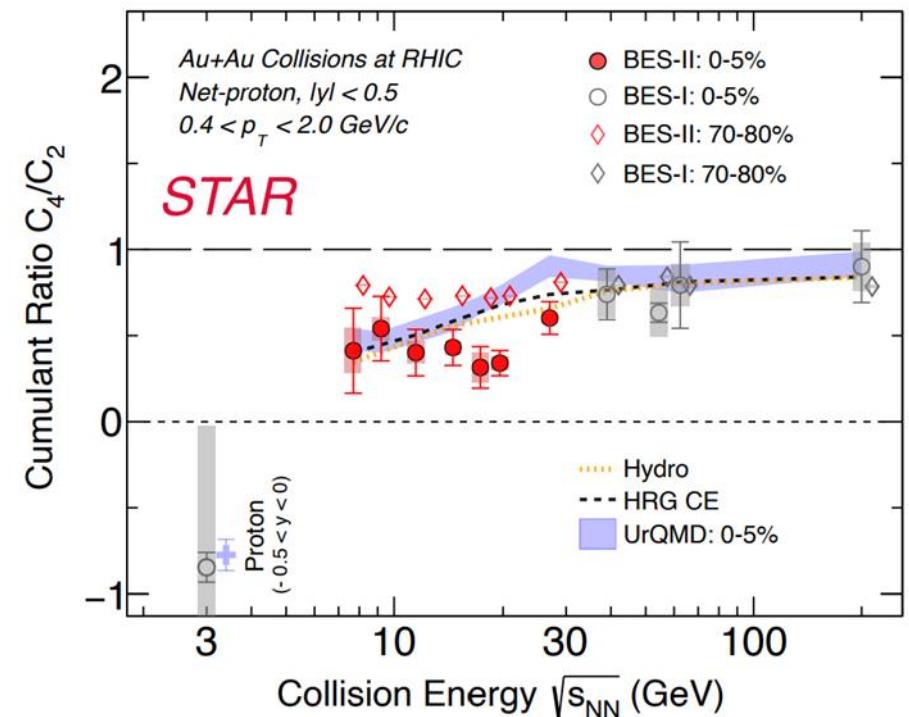


# Net-proton fluctuations near critical point

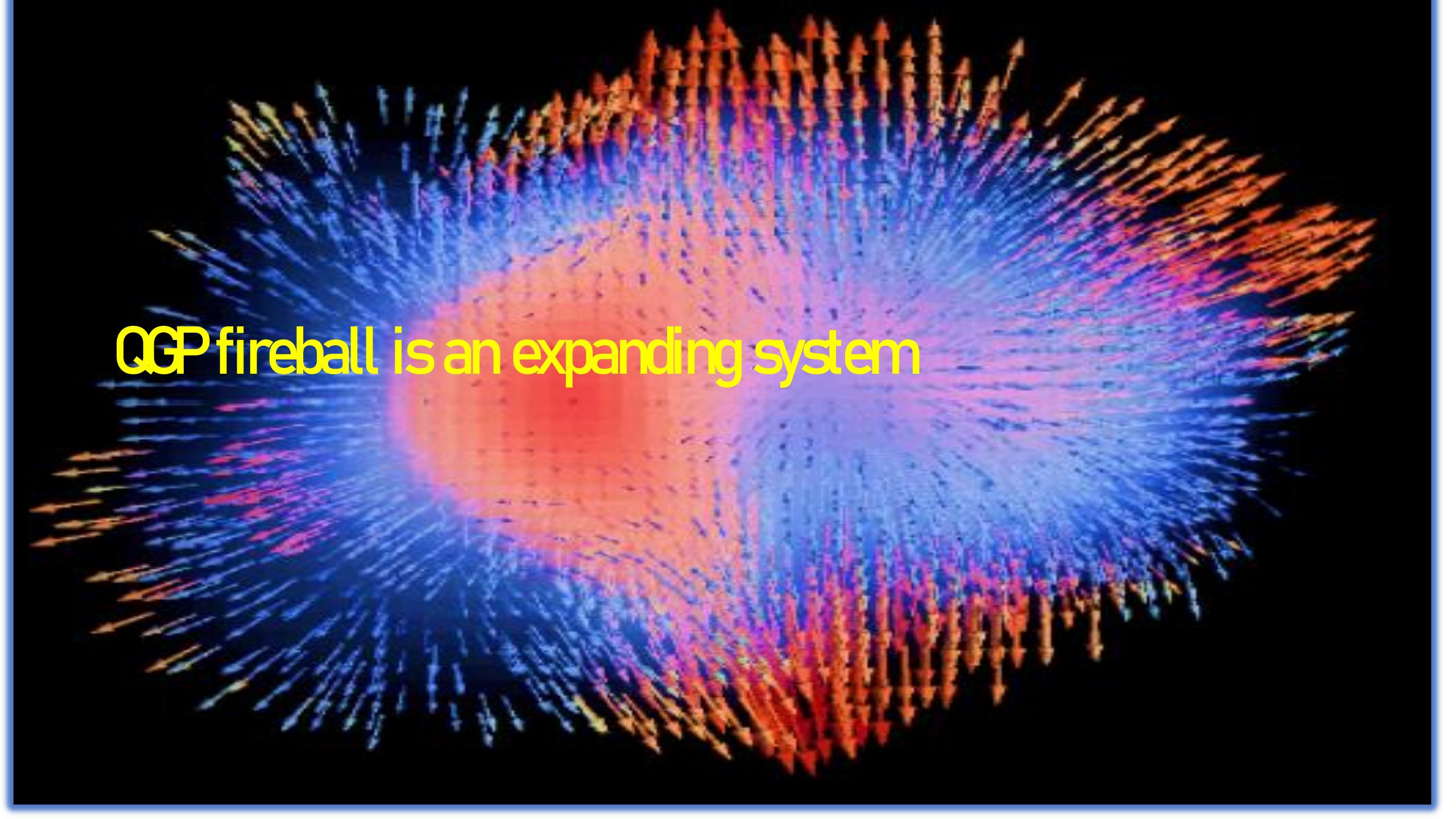
- Characteristic feature of critical point:
  - long range correlation
  - large fluctuations
- Non-monotonicity of Net-Proton Cumulant



M.Stephanov, PRL 107,052301



STAR, CPOD 2024

A 3D vector field plot visualizing the expansion of a Quark-Gluon Plasma (QGP) fireball. The plot is set against a black background and shows a central, luminous orange-red region that tapers off into a complex, multi-colored pattern of blue, purple, and yellow. Numerous small, colored arrows point radially outward from the center, indicating the direction and speed of expansion. The overall shape is roughly spherical but flattened at the top and bottom.

**QGP fireball is an expanding system**

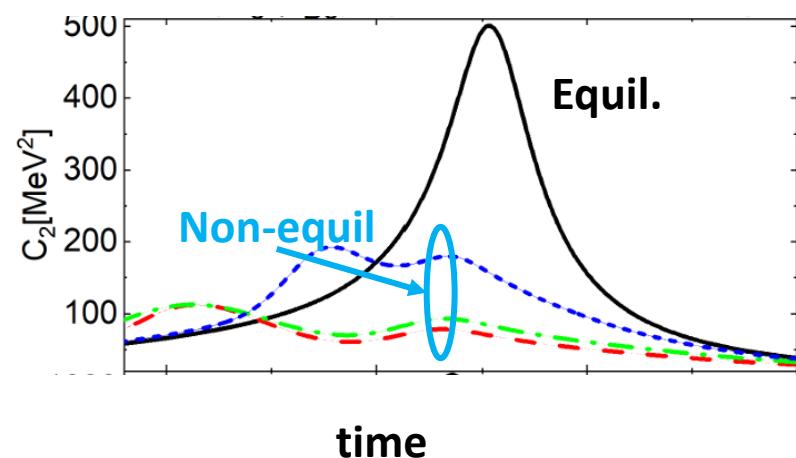
# Dynamical effects modifies the critical fluctuations

- Expanding QGP fireball forces the fluid cell swipes the critical regions with finite time  $\tau_{expand}$ ;
- Long range correlation in the system near the critical point requires long time to correlate with each other  $\tau_{relax} \sim \xi^z$ ;
- In heavy-ion collision, near the QCD critical point,

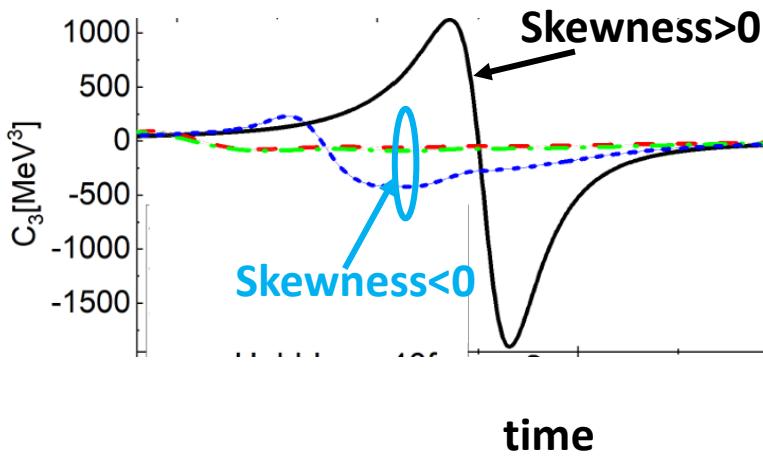
Shian Tang, Shanjin Wu, Huichao Song, PRC(2021)

$$\tau_{relax} \gg \tau_{expand}$$

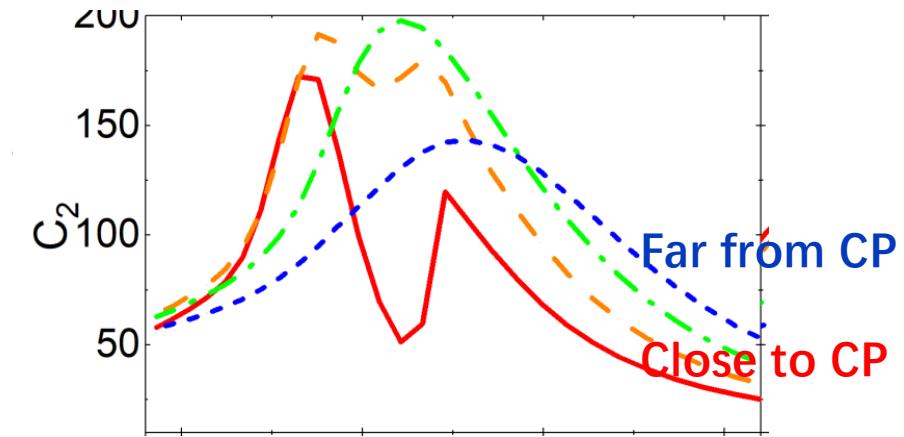
- Expansion drives the system has no enough time to relax to equilibrium



Suppression of the fluctuations



Reverse the sign of skewness



Largest  $C_2 \neq$  closet to CP

# Dynamical models near QCD critical point

- **QCD critical point belongs to model H** (Symmetry and dynamical analysis): D.T.Son,Stephanov(04)
- **Model H** (conserved baryon  $\mathbf{n} \approx \mathbf{n} + \boldsymbol{\sigma}$  + transverse momentum density):

$$\frac{\partial \mathbf{n}}{\partial t} = \Gamma[\mathbf{n}, \boldsymbol{\pi}^T] + \text{noise}, \quad \frac{\partial \boldsymbol{\pi}^T}{\partial t} = \boldsymbol{\eta}[\mathbf{n}, \boldsymbol{\pi}^T] + \text{noise}$$

Numerical simulation is time consuming.

- **Simplified models:**

- **Model A** (order parameter field) S.Mukherjee et al15' 16', L.Jiang et al17', S.Wu et al 19', S.Tang et al 23',

$$\frac{\partial \boldsymbol{\sigma}}{\partial t} = \Gamma[\boldsymbol{\sigma}] + \text{noise}$$

- **Model B** (conserved field) M.Sakaida et al 17', S.Wu et al 19', M.Nahrgang et al 19', G.Pihan et al 22' ...

$$\frac{\partial \mathbf{n}}{\partial t} = \nabla^2 \Gamma[\mathbf{n}] + \text{noise}$$

# Dynamical models near QCD critical point

More efficient modeling in heavy-ion collisions: **Hydrodynamics + Critical fluctuations**

- **Non-equilibrium chiral hydrodynamics** M. Nahrgang et al 11'12'14'16'19'

$$\partial_\mu T_{fluid}^{\mu\nu} = \partial_\mu T_\sigma^{\mu\nu}, \quad \frac{\partial \sigma}{\partial t} = \Gamma[\sigma] + \text{noise}$$

- **Fluctuating hydrodynamics** J.Kapusta et al 12',12', K.Murase et al 13', X.An et al 19',21'...

$$\partial_\mu T_{fluid}^{\mu\nu} = \text{noise} , \partial_\mu N^\mu = \text{noise}$$

- **Hydro+, hydro++...** (hydro + slow modes) M. Stephanov et al 18'19'20', N. Abbasi et al 22', L. Du et al 20',.....

$$\partial_\mu T_*^{\mu\nu} = 0, \partial_\mu N^\mu = 0, \quad \frac{\partial \phi}{\partial t} = \Gamma[\phi]$$

- **Hydro-kinetics** D.Teaney et al 17'18'19'22'...

$$\partial_\mu T^{\mu\nu} = 0, \partial_\mu N^\mu = 0, \quad \partial_t(\text{noise correlator}) = \#$$

See reviews: e.g. Shanjia Wu, et al.,2104.13250; Lipei Du et al. 2402.10183; Xin An et al., 2108.13867; Marcus Bluhm et al., 2001.08831; Adam Bzdak et al.,1906.00936;M.Asakawa et al.,1512.05308

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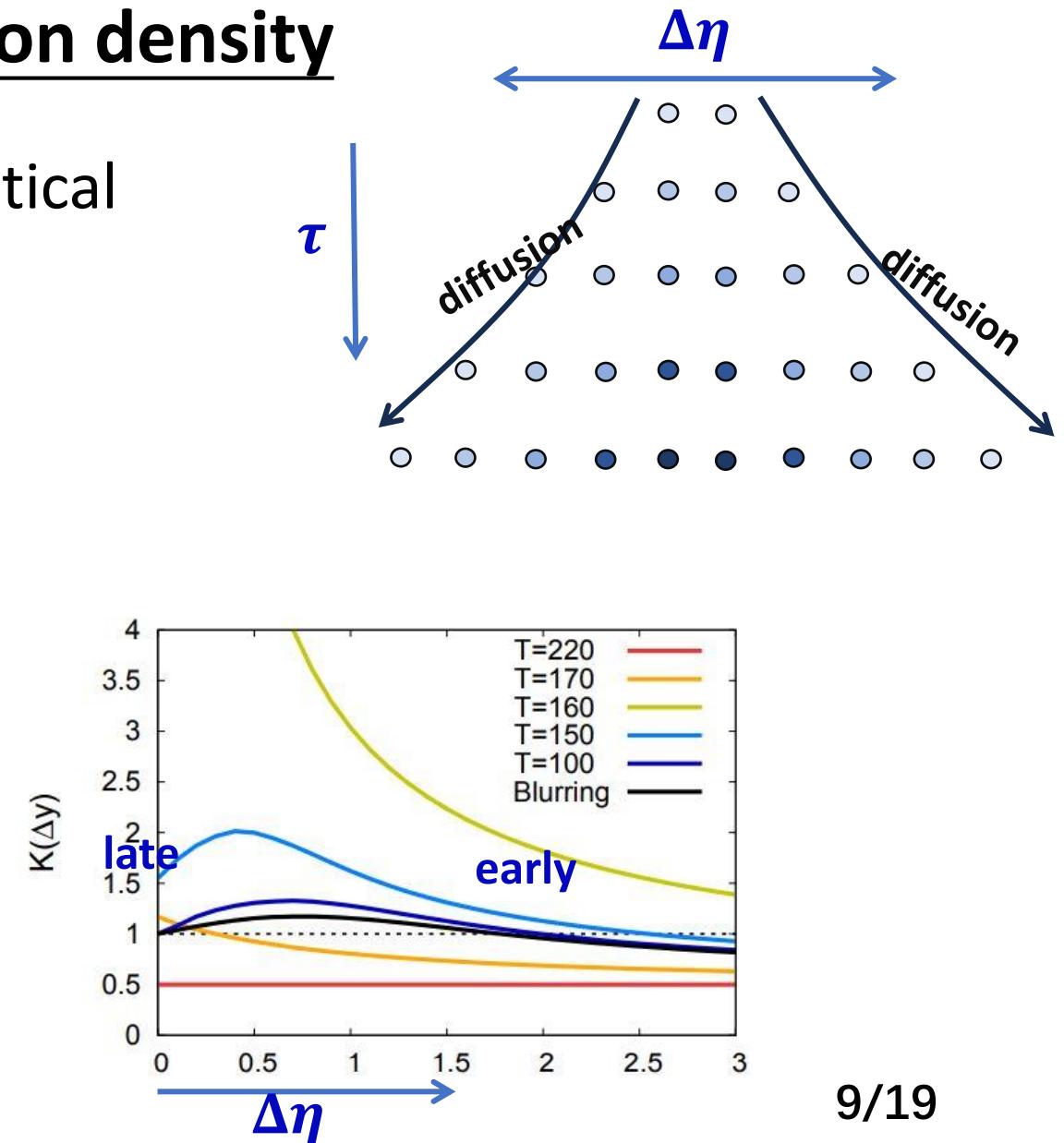
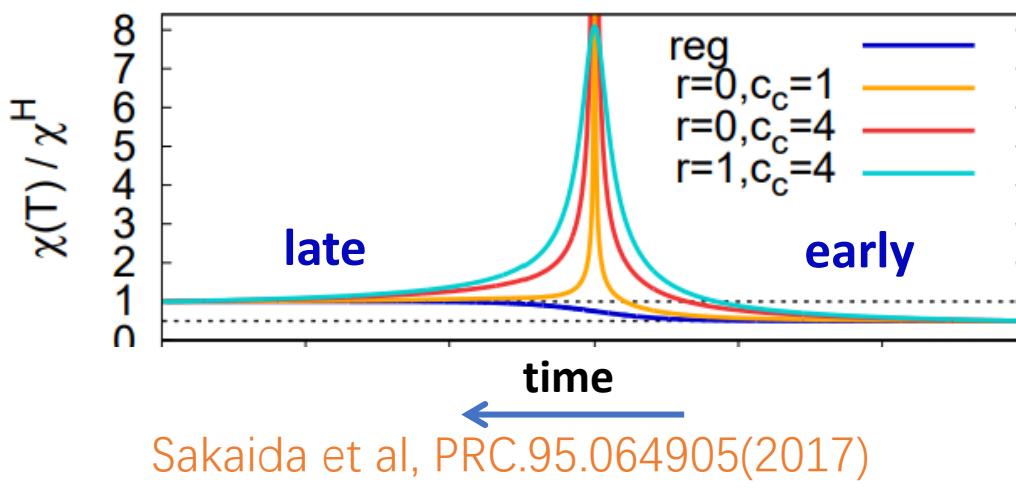
$$\frac{\partial \mathbf{n}}{\partial t} = \nabla^2 \Gamma[\mathbf{n}] + \text{noise}$$

# Dynamics of conserved net-baryon density

- Diffusion of conserved baryon near critical point:

$$\partial_\tau n = \nabla^2 n + \text{noise}$$

- The process of diffusion consumes time.
- Larger  $\Delta\eta \sim$  the early evolution.



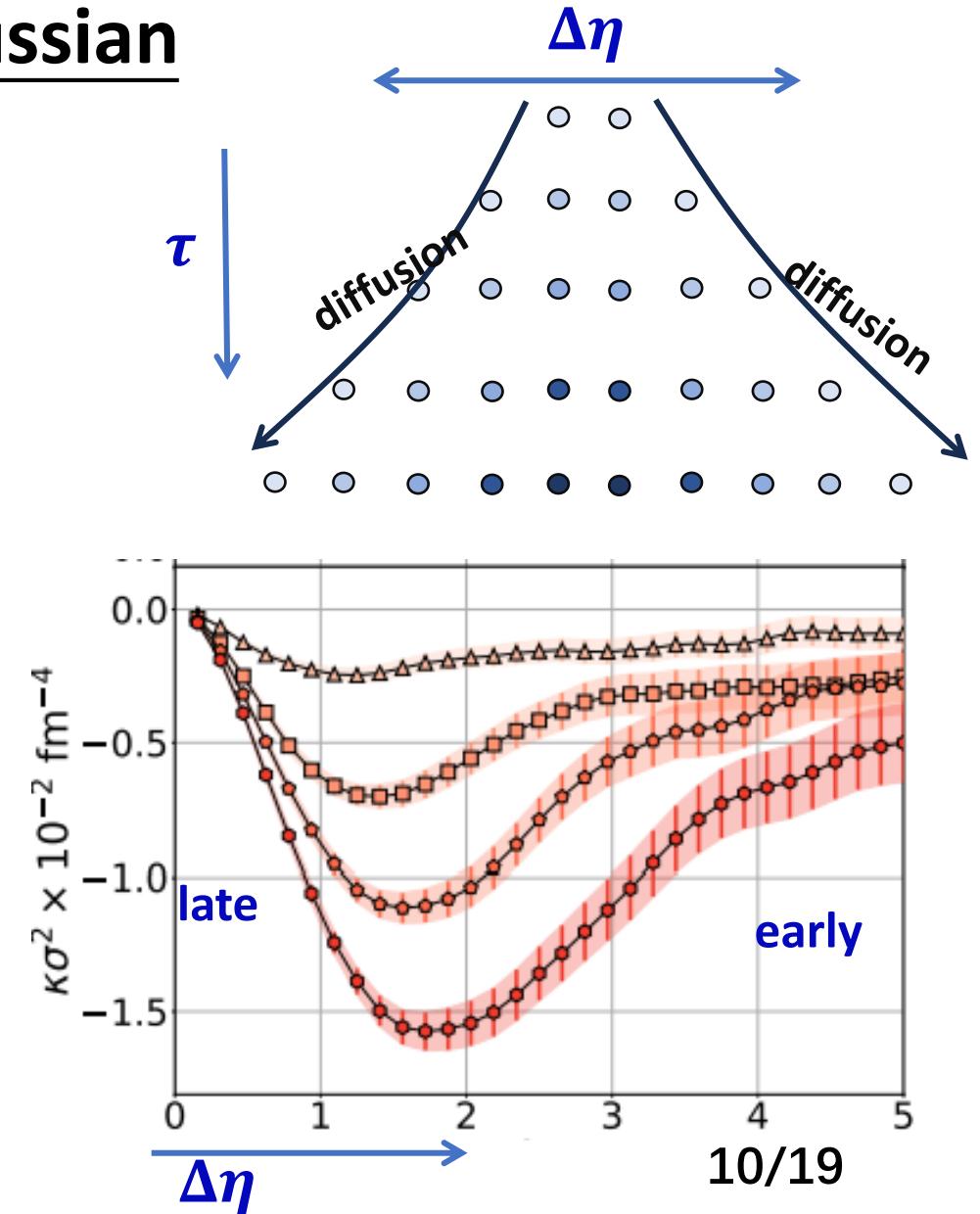
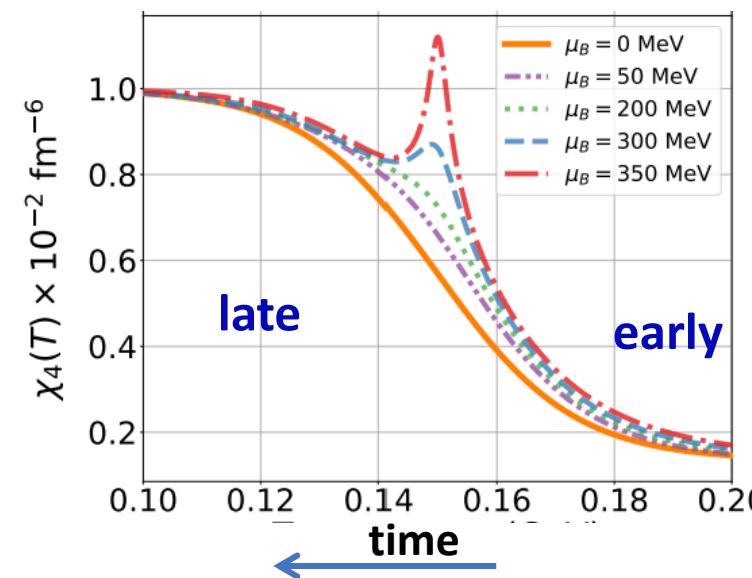
# Conserved net-baryon with non-Gaussian

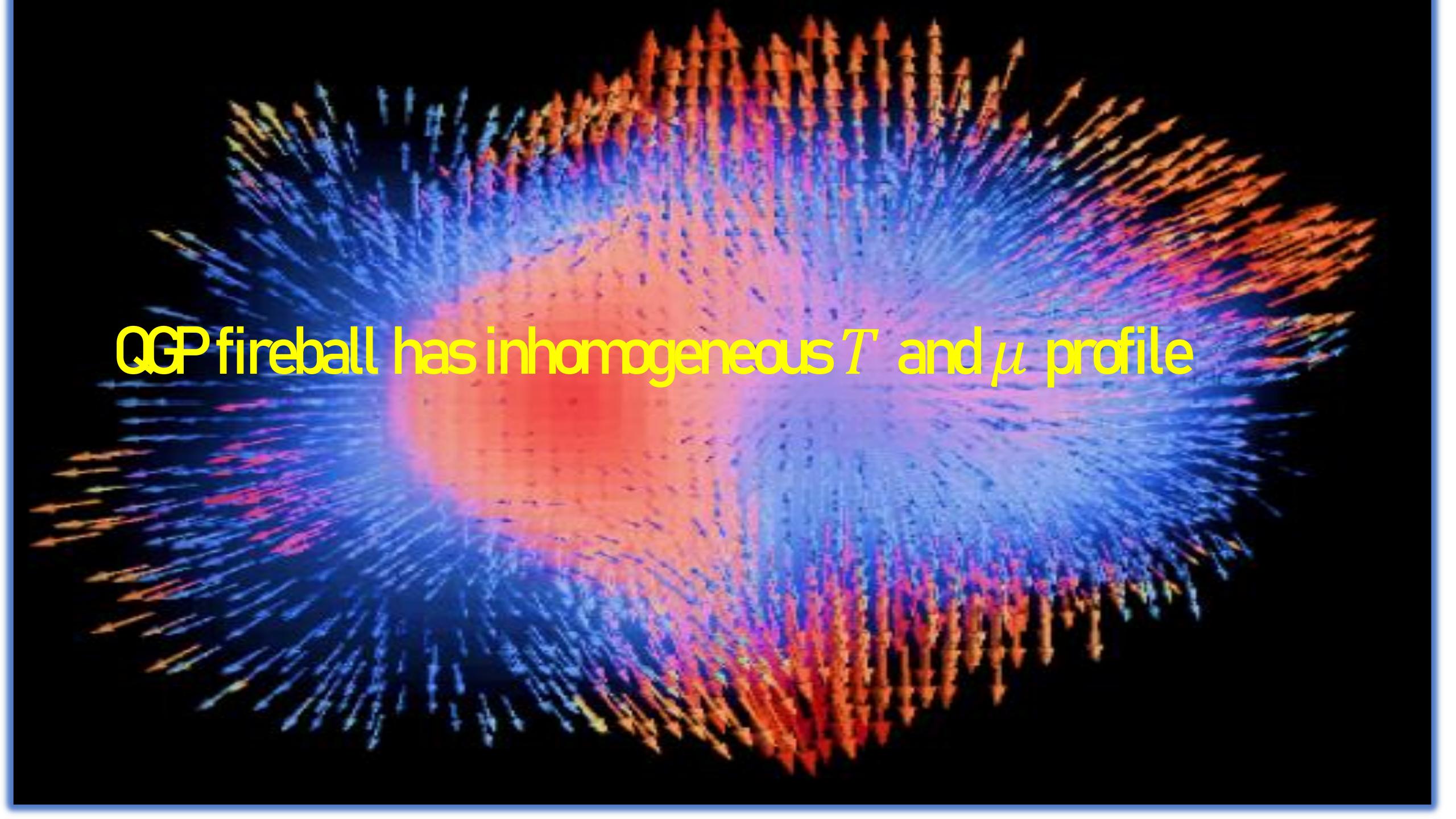
- Diffusion of conserved baryon near critical point:

$$\partial_\tau n = \nabla^2(n + n^2 + n^3) + \text{noise}$$

- The process of diffusion consumes time.
- Larger  $\Delta\eta \sim$  the early evolution.

G.Pihan et al.,  
PRC.107.014908(2022)

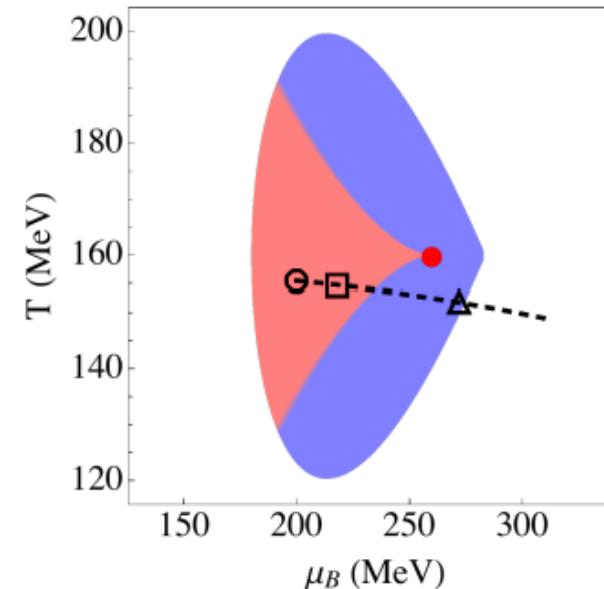
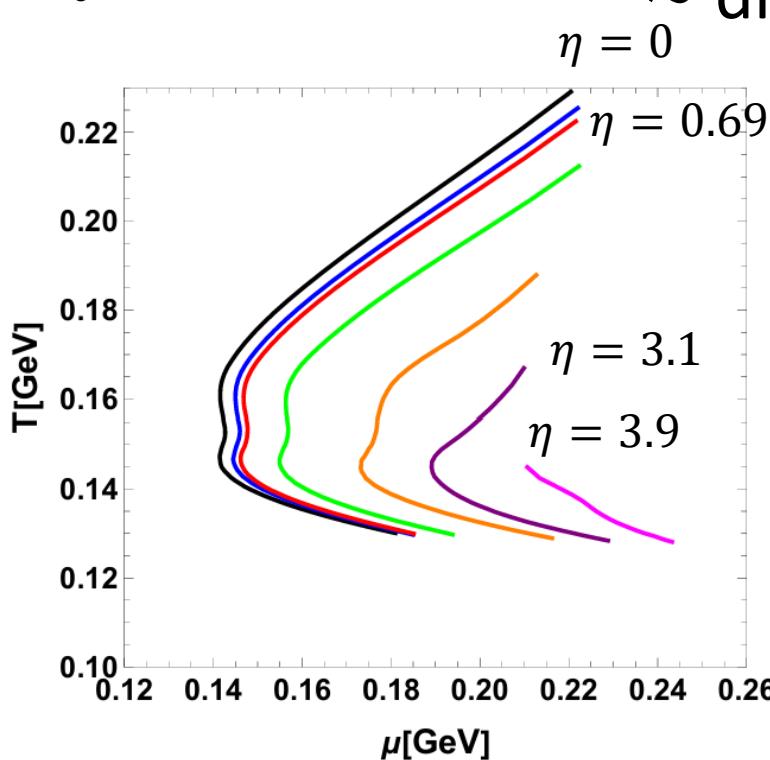




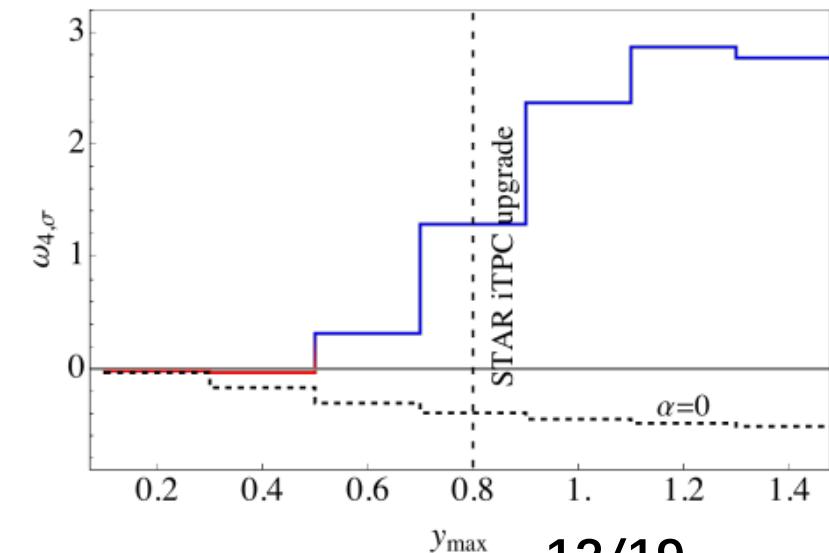
**QGP fireball has inhomogeneous  $T$  and  $\mu$  profile**

## Inhomogeneous QGP profile

- Different rapidity  $\sim$  different trajectories
- $\sim$  detect different region of critical region
- $\sim$  different critical behavior



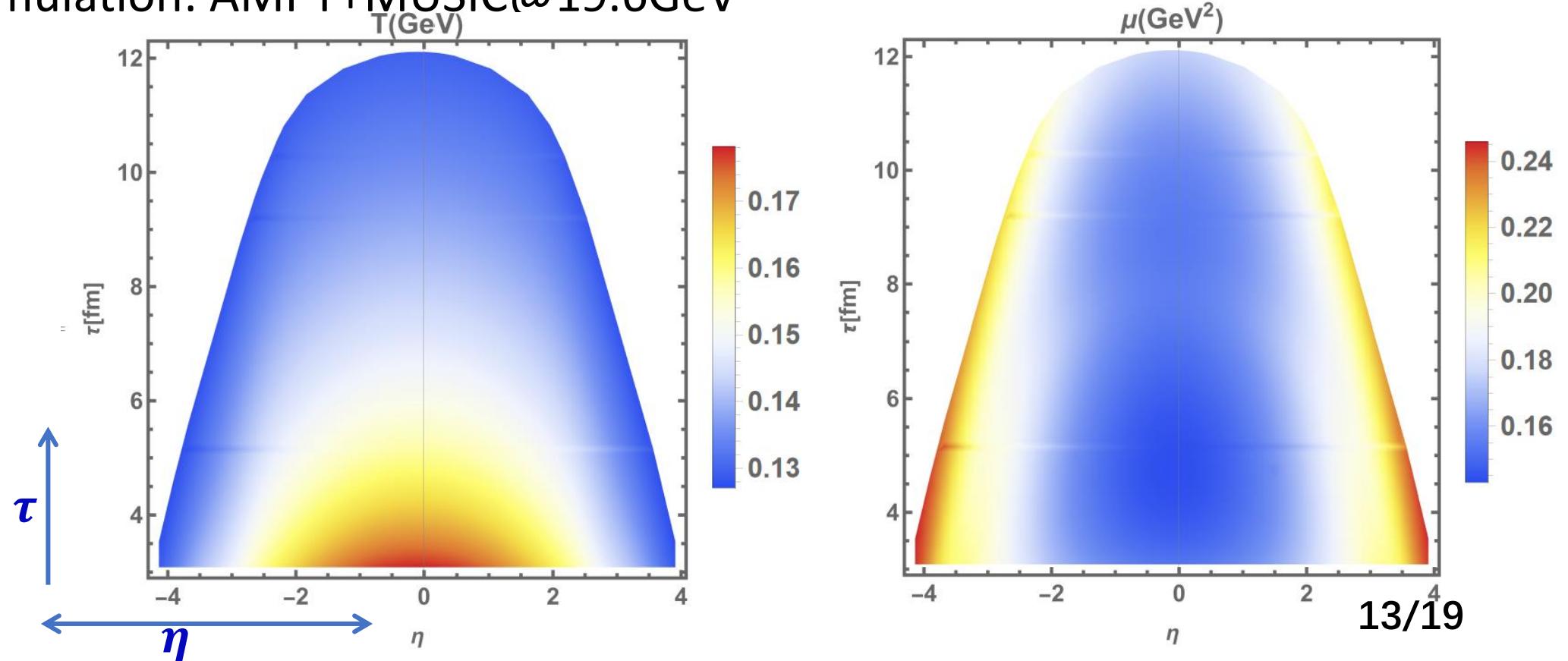
J.Brewer et al., PhysRevC.98.061901



# Inhomogeneous T and $\mu$ profile from hydro simulation

- This talk aims to study the inhomogeneous QGP profile effects on the diffusion of net-baryon
- Hydro simulation: AMPT+MUSIC@19.6GeV

Shanjin Wu 2406.12325

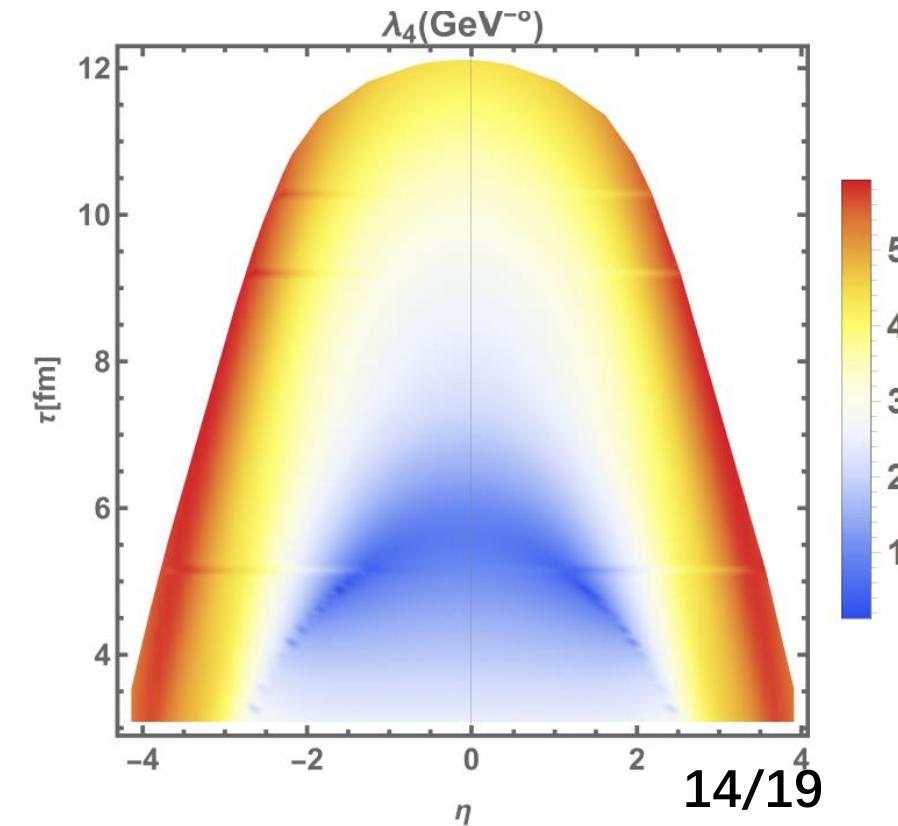
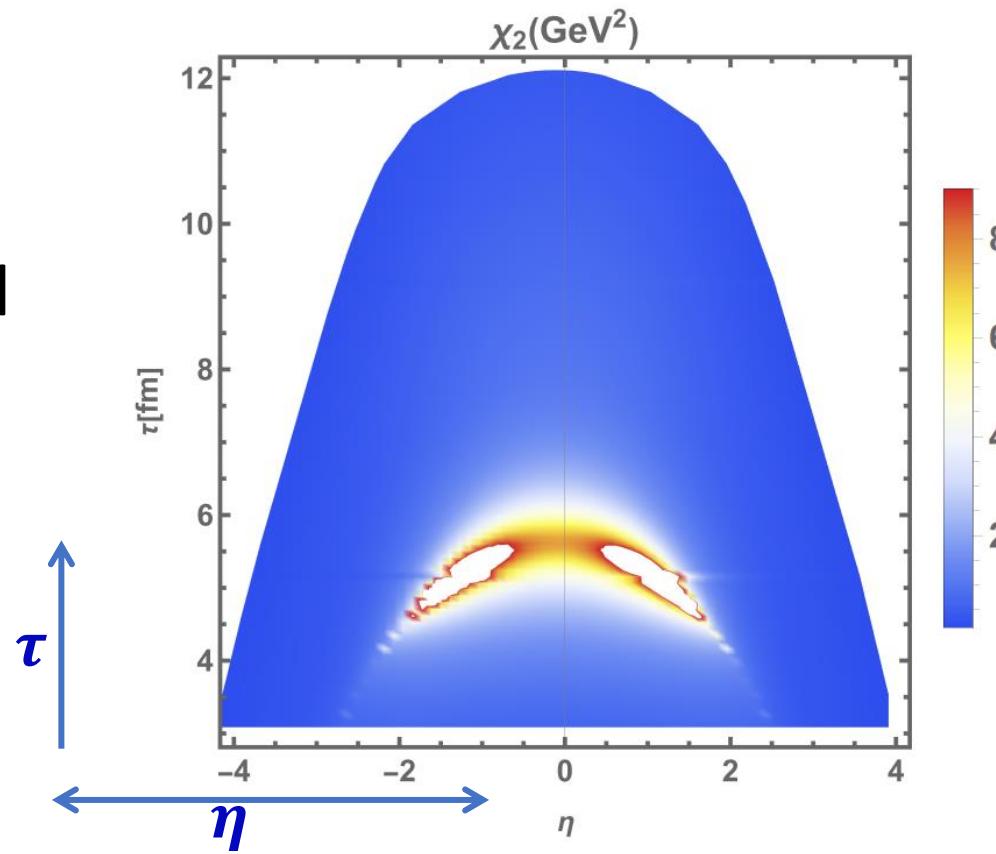


# Conserved net-baryon with inhomogeneous T and $\mu$ profile

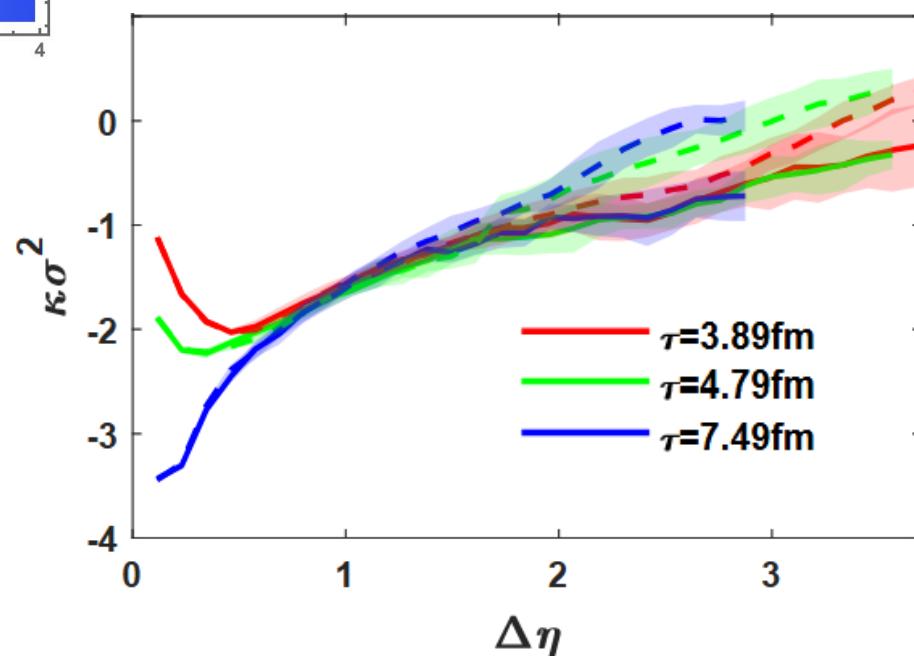
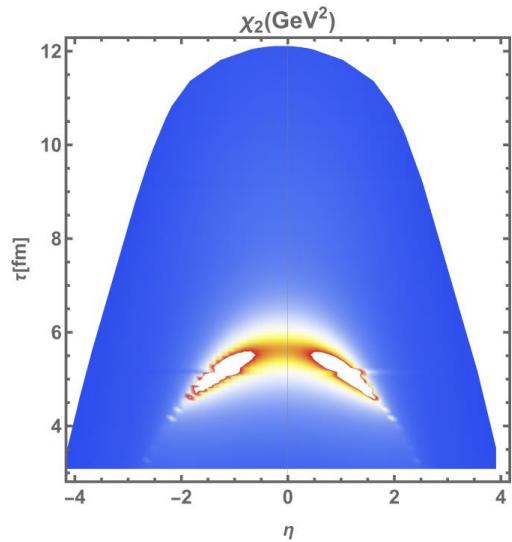
- Diffusion of conserved baryon near critical point: Shanjin Wu 2406.12325

$$\partial_\tau n = \nabla^2 \left( \frac{n}{\chi_2} + \lambda_3 n^2 + \lambda_4 n^3 \right) + \text{noise}$$

$\chi_2, \lambda_4$  from Ising model



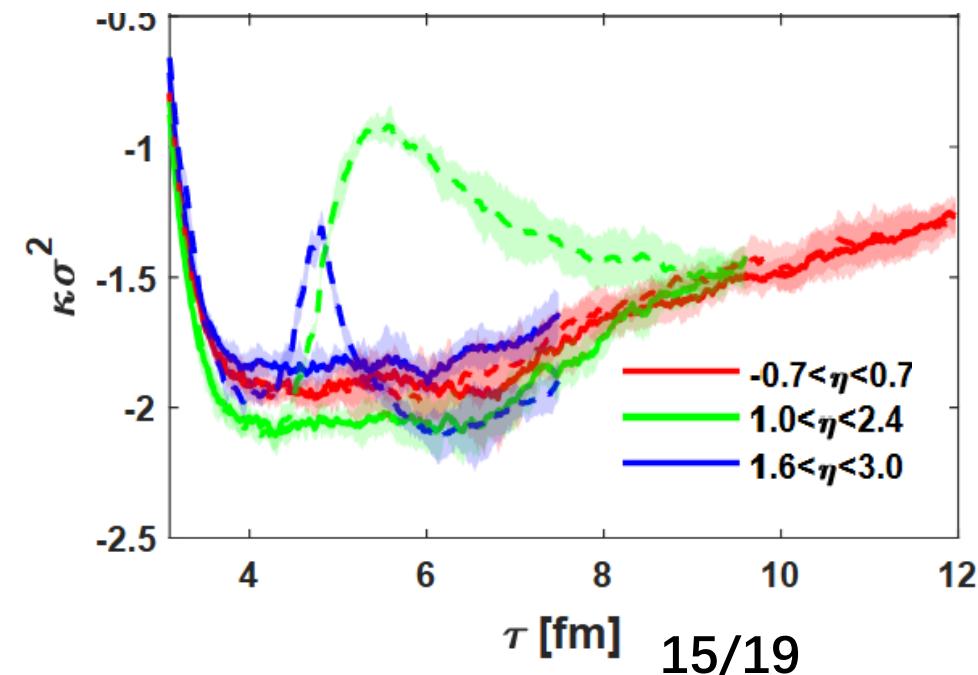
# Conserved net-baryon fluctuations



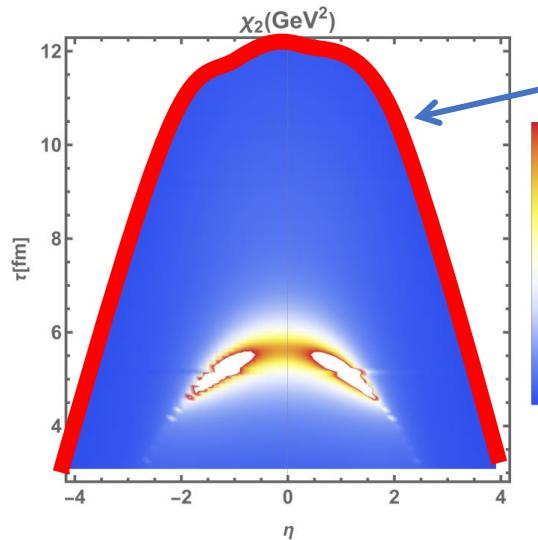
Shanjin Wu 2406.12325

Inhomogeneous profile effects is significant at large rapidity

Solid: uniform profile; Dashed: inhomogeneous profile



# Conserved net-baryon fluctuations at freeze-out surface



Freeze-out surface

Shanjin Wu 2406.12325

Net-baryon number at freeze-out surface

$$N_B = g \int \frac{d^3 p}{(2\pi)^3} \frac{1}{p^0} \int d\sigma_\mu p^\mu f(\mathbf{x}, \mathbf{p})$$

In Bjorken limit

$$= \frac{2gA}{(2\pi)^2} \int_{-\Delta\eta}^{\Delta\eta} d\eta \tau_f \exp\left(\frac{\mu}{T}\right) T^3 \left(\frac{m}{T}\right)^2 K_2\left(\frac{m}{T}\right)$$

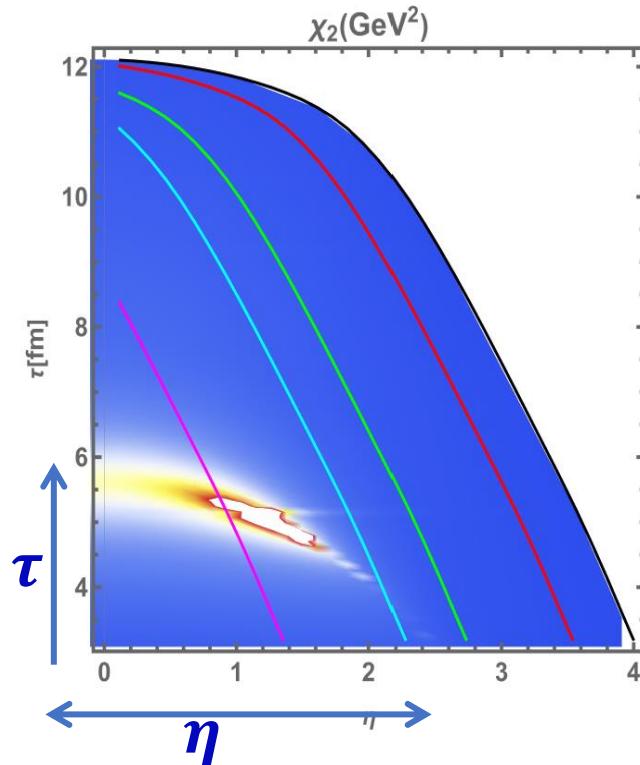
Net-baryon fluctuations at freeze-out surface

$$\delta N_B = \frac{2gAm^2}{(2\pi)^2} \int_{-\Delta\eta}^{\Delta\eta} d\eta \exp\left(\frac{\mu}{T}\right) \frac{\delta n_B(\tau_f, \eta)}{\chi} K_2\left(\frac{m}{T}\right)$$

# Conserved net-baryon fluctuations at freeze-out surface

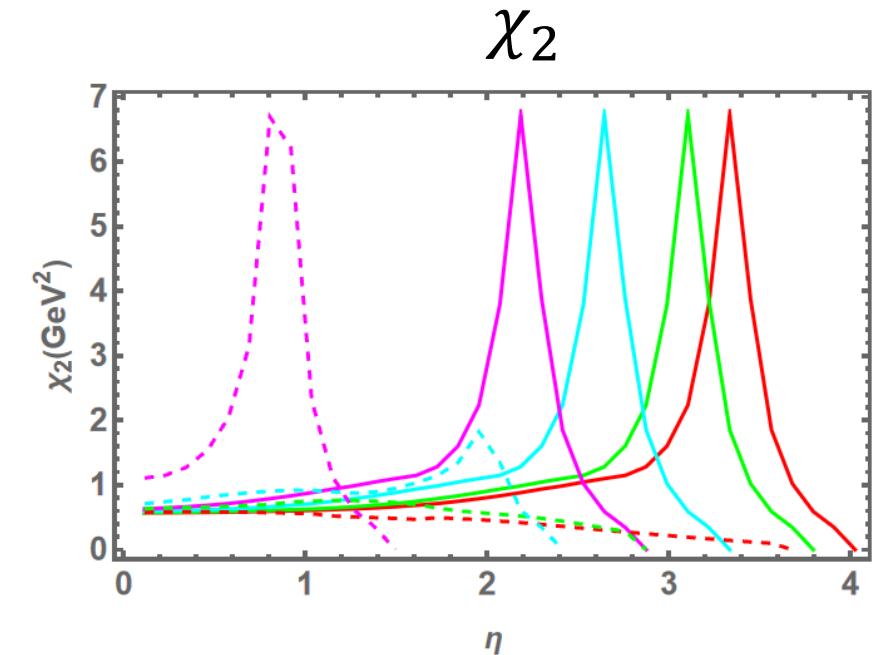
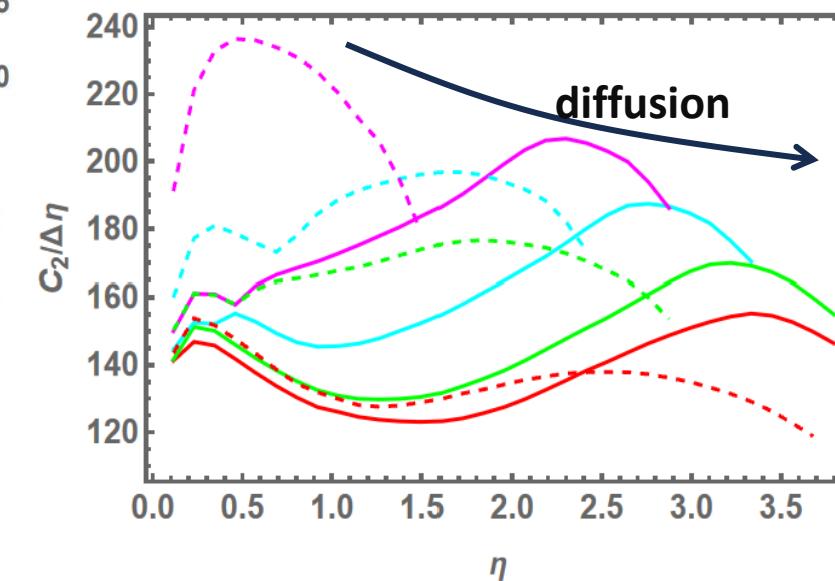
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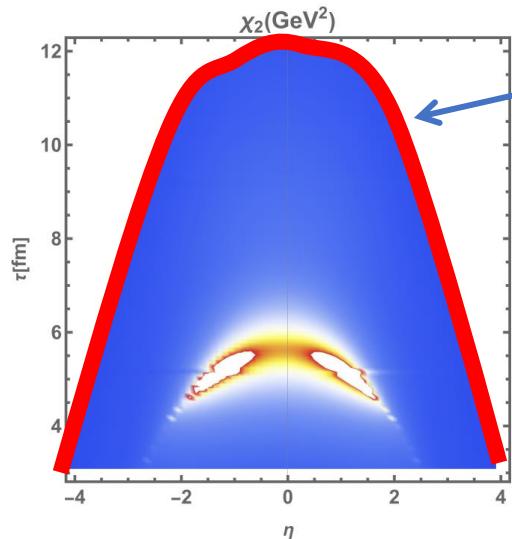
Critical slowing down effects of the fluctuations

$$C_2 = \langle (\delta n_B)^2 \rangle$$



Solid: uniform profile; Dashed: inhomogeneous profile

# Conserved net-baryon fluctuations at freeze-out surface

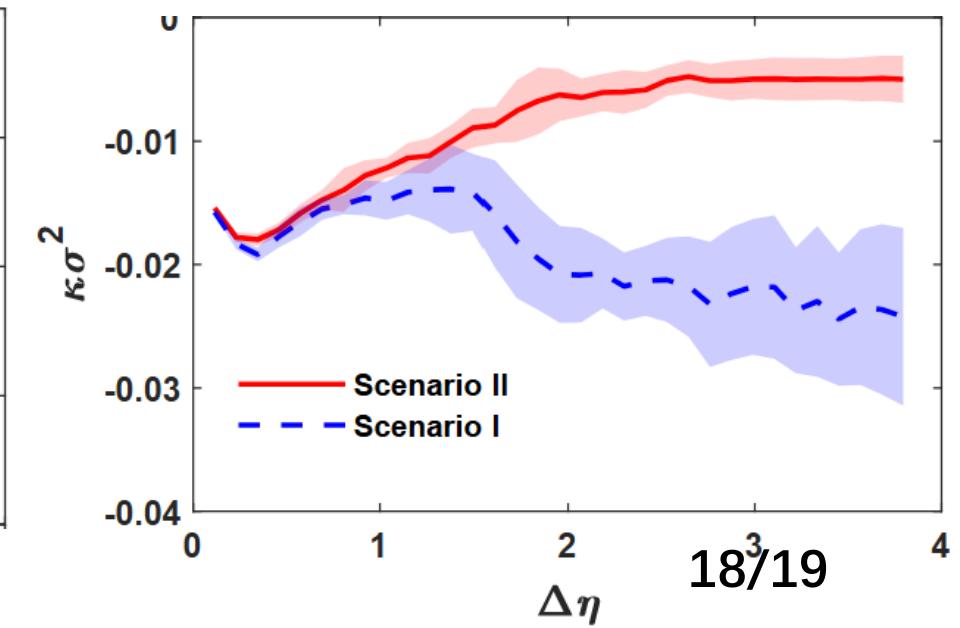
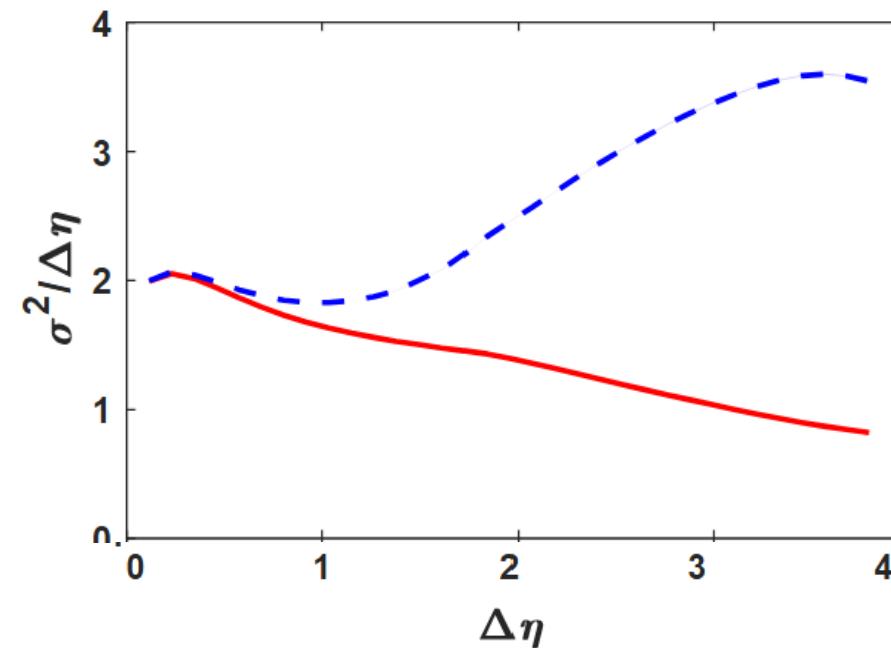


Freeze-out surface

Shanjin Wu 2406.12325

Critical slowing down => Significant inhomogeneous profile effects at large rapidity

Solid: uniform profile; Dashed: inhomogeneous profile



## Summary

- Dynamical modeling the QGP evolution near the QCD critical point is essential for the study of fluctuations in heavy-ion experiments;
- The diffusion of conserved net-baryon density preserves the early evolution history and behaves non-monotonically with increasing rapidity;
- Considering the inhomogeneous  $T$  and  $\mu$  profile has significant effects at large rapidity.
- Multiplicative noise

Thank you!