

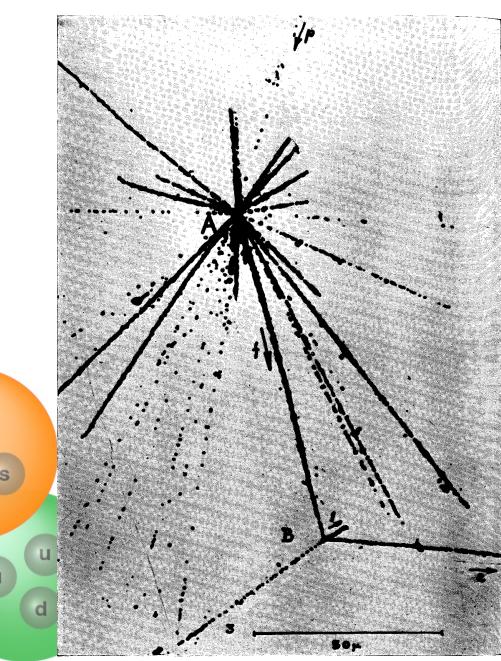
Collision Energy Dependence of Hypertriton Production in Au+Au Collisions at RHIC

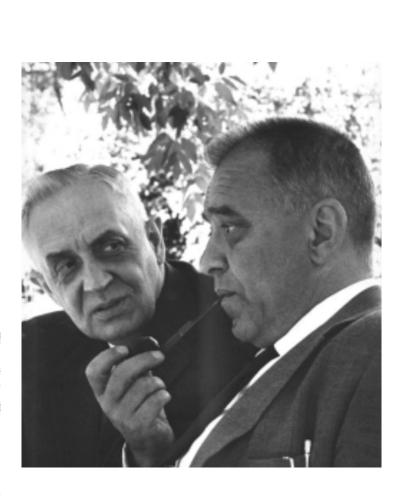
Xiujun Li (李 秀君) University of Science and Technology of China Oct 12, 2024

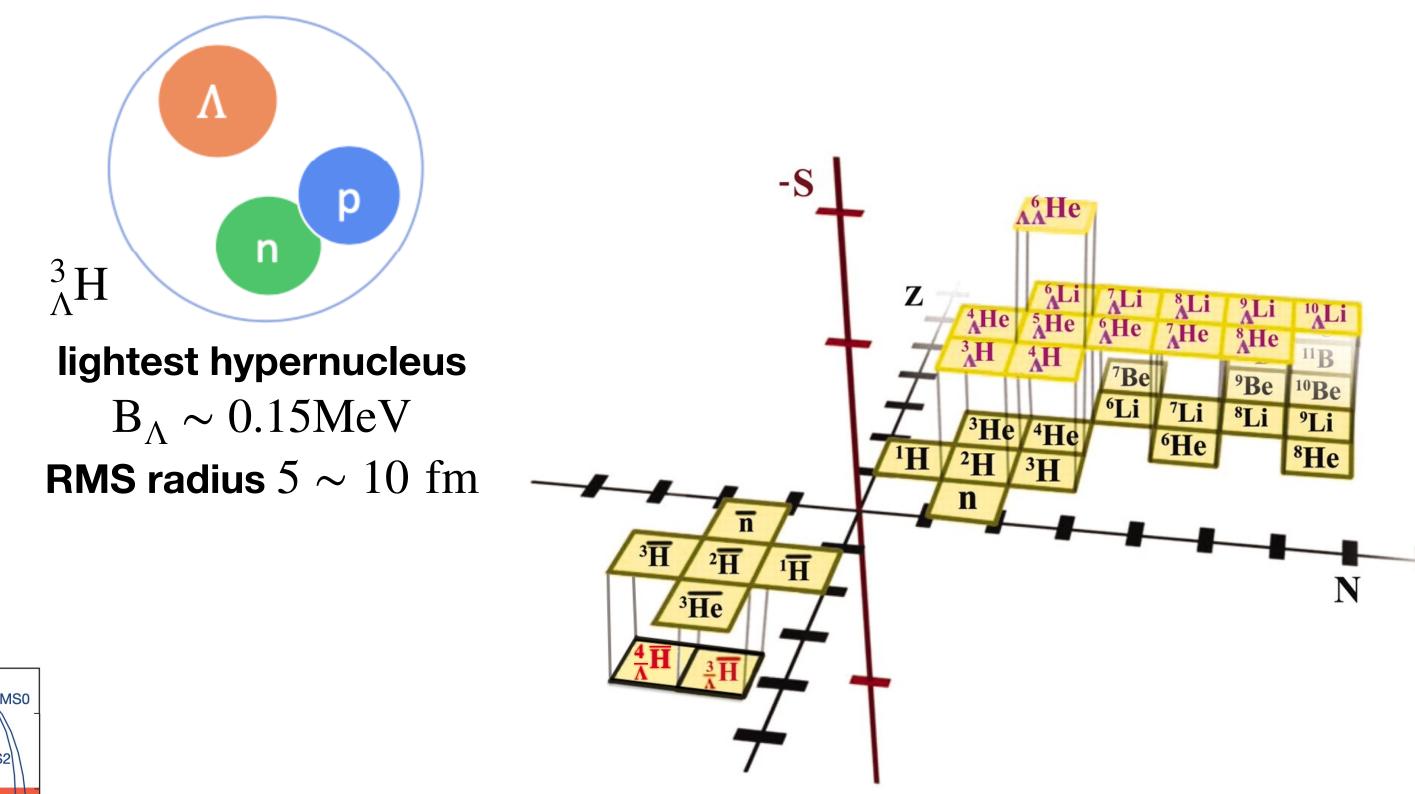


Introduction: hypernuclei

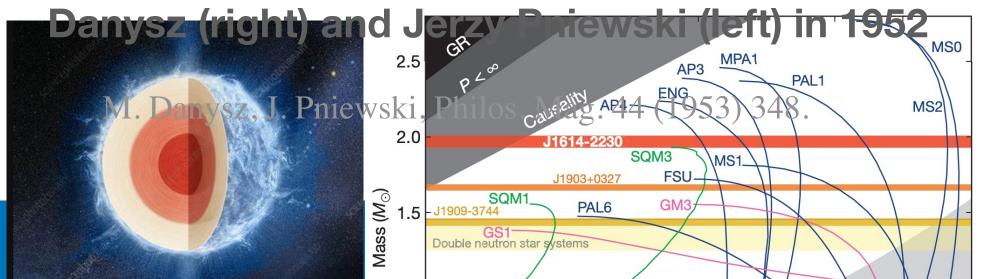
-Natural hyperon-baryon correlation system







The first discovery of hypernucleus by Marian



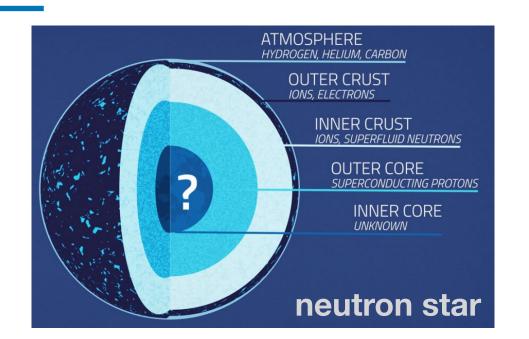
R Regional Meeting, Chongqing University, Oct.12, 2024

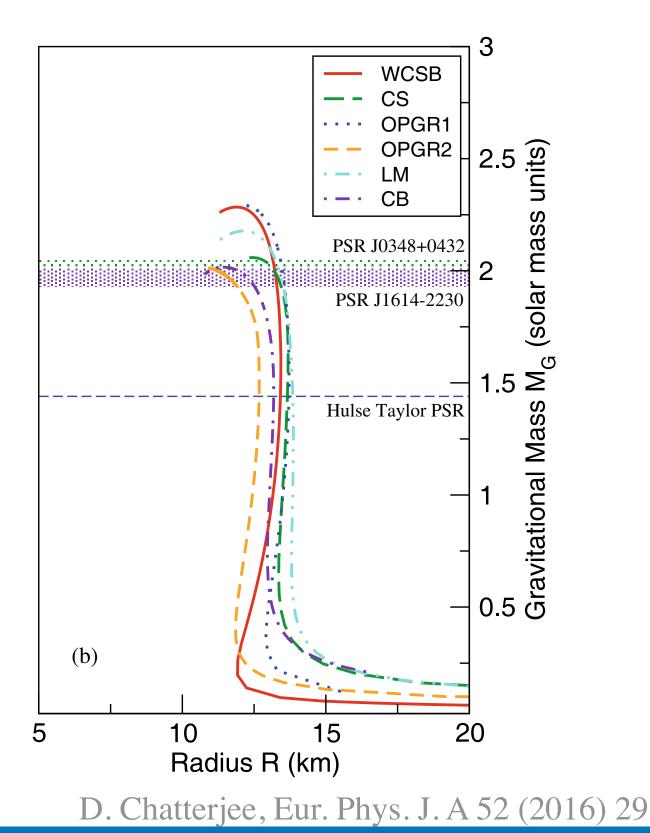
Hypernuclei: bound nuclear systems of non-strange and strange baryons



Introduction: YN interaction in dense matter

- Hypernuclei serve as a laboratory to study the hyperon-nucleon (YN) interaction
 - YN interaction is essential in probing dense inner core of neutron star
 - Hyperon puzzle: do hyperons exist in the dense inner core of neutron stars?
 - No direct method to probe inner core
 - Rely on theoretical models compared with observations
 - Lack of experimental data of YN, YNN, YY interactions to constrain theoretical models of the dense matter equation of state (EoS)



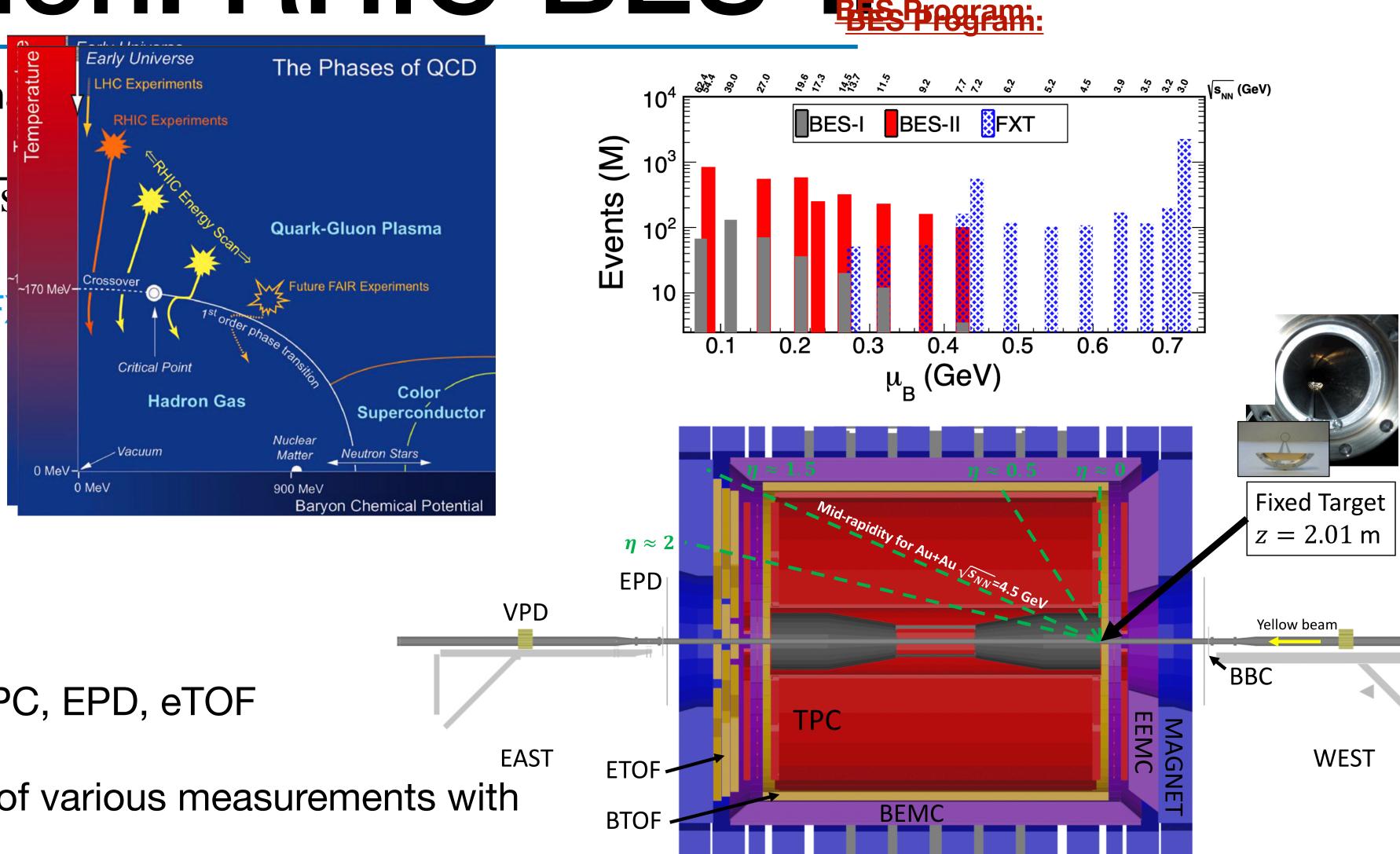






Introduction: RHIC BES-I

- RHIC beam energy scan Ph
 - Specific focus on low $\sqrt{2}$
 - Include fixed target (F) energies, increase $\mu_{\rm B}$ ~700 MeV
 - High statistics data
 - Improve systematics ullet
 - Detector upgrade: iTPC, EPD, eTOF
 - Enhances the capability of various measurements with excellent precision

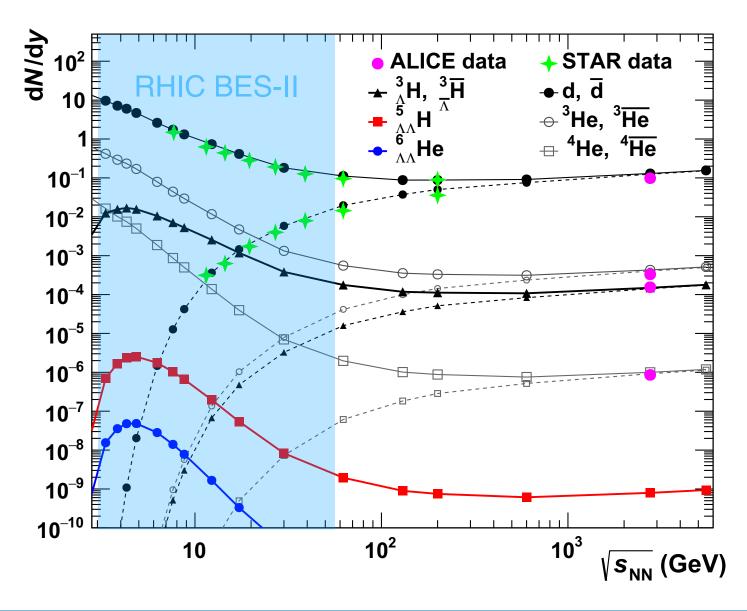




Introduction: hypernuclei in HI collisions

Production mechanism of hypernuclei is still not well understood.

Hypernuclei formation process in relativistic heavy-ion (HI) collisions



B. Dönigus, Eur. Phys. J. A (2020) 56:280 A. Andronic et al. PLB (2011) 697:203–207

can be studied through measurements related to spectra and collective flow.

Hypernuclei measurements are scarce in HI collision experiments

At low beam energies, hypernuclei production is expected to be enhanced due to high baryon density

RHIC BES-II offers great opportunity for hypernuclei measurements.

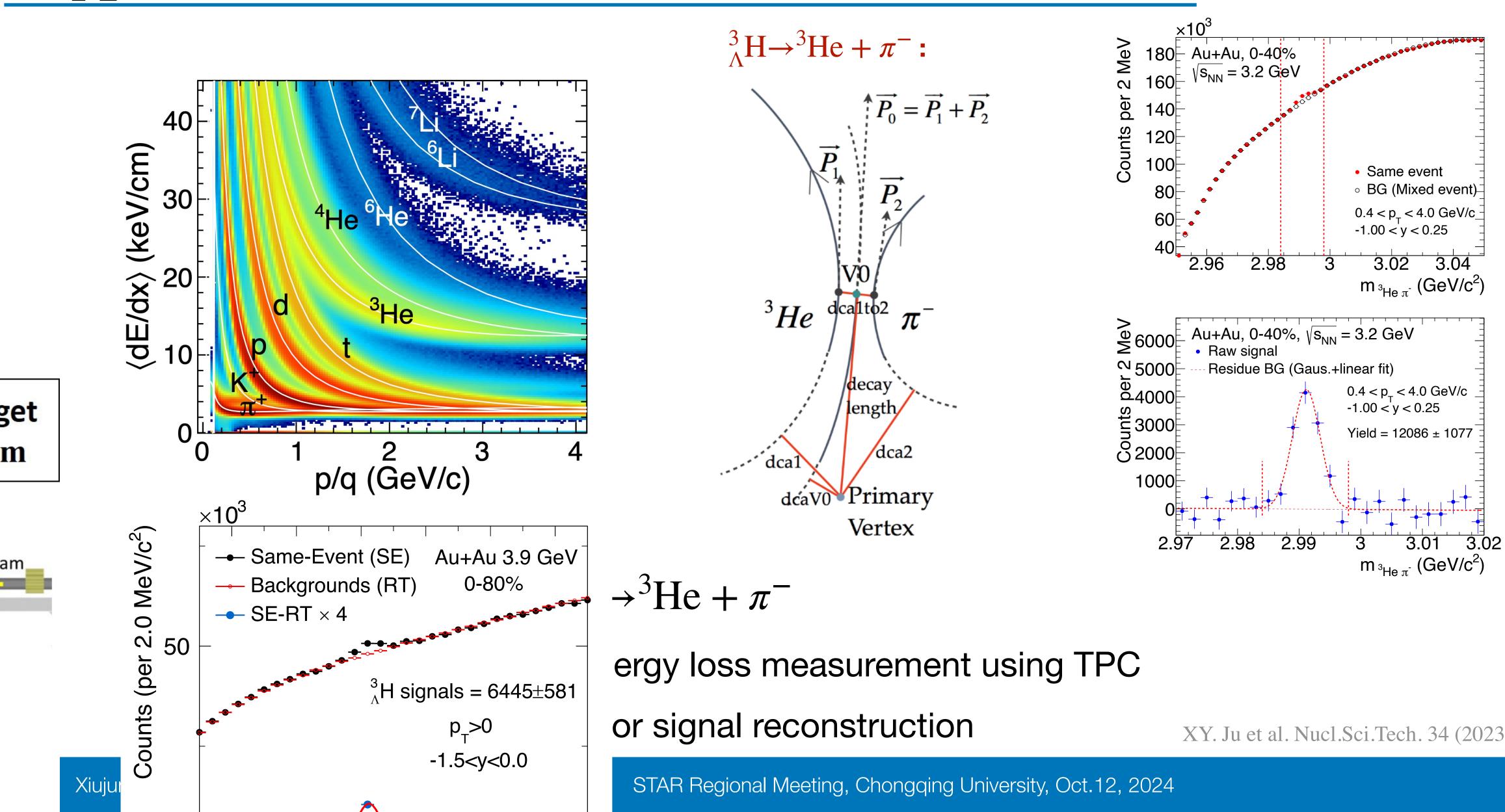
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³H reconstruction

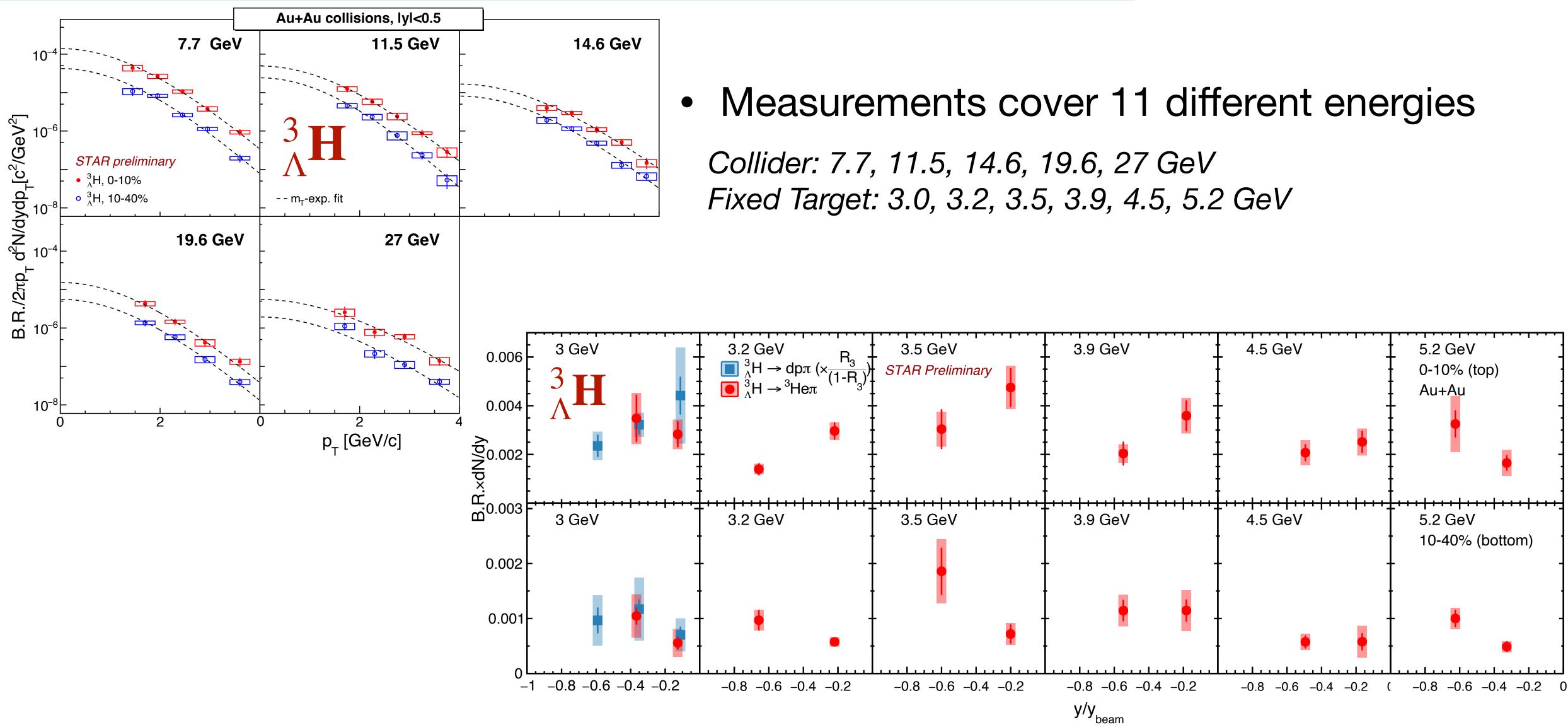


XY. Ju et al. Nucl.Sci.Tech. 34 (2023) 10, 158





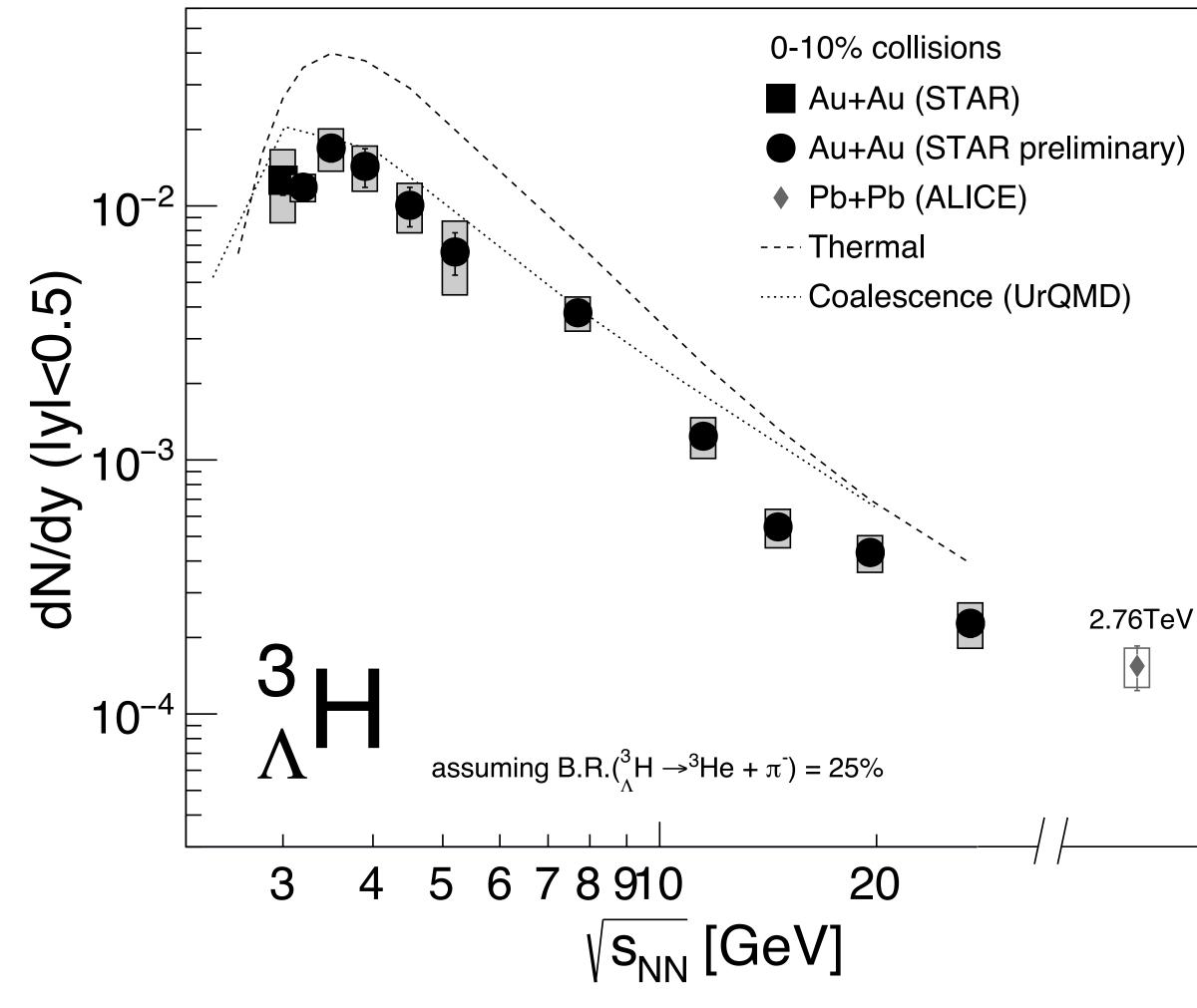
3 H rapidity and p_{T} spectra



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Energy dependence of ${}^3_{\Lambda}H$ production



STAR, PRL 128 (2022) 202301 ALICE, PLB 754 (2016) 360 T. Reichert, et al, PRC 107 (2023) 014912

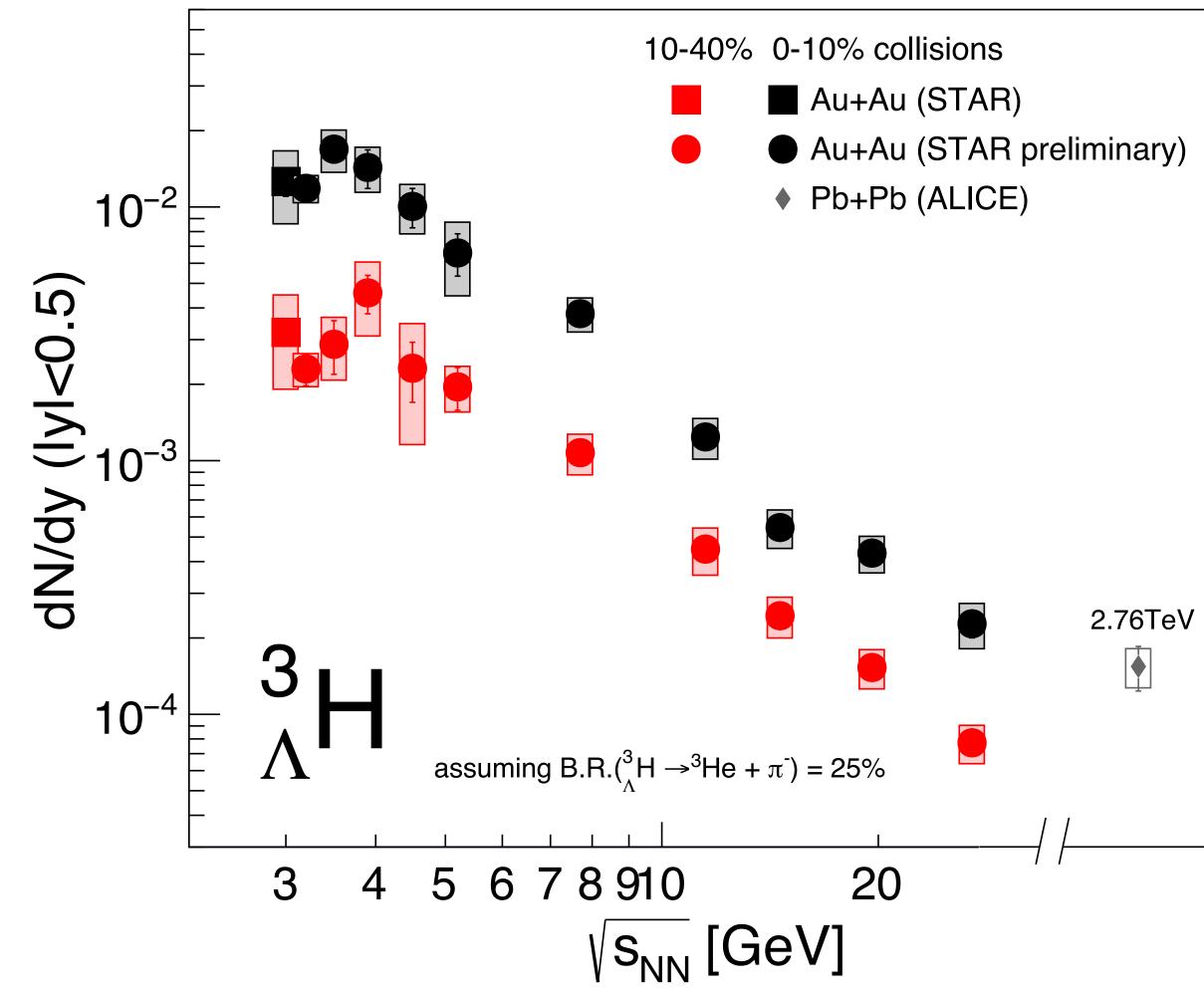
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- Yields increase strongly from $\sqrt{s_{NN}} =$ 27 GeV to ~4 GeV
- Peak at 3-4 GeV
- Hadronic transport + coalescence models qualitatively describe the data
- Thermal model overestimates the data





Centrality dependence of ${}^3_{\Lambda}H$ production



STAR, PRL 128 (2022) 202301 ALICE, PLB 754 (2016) 360

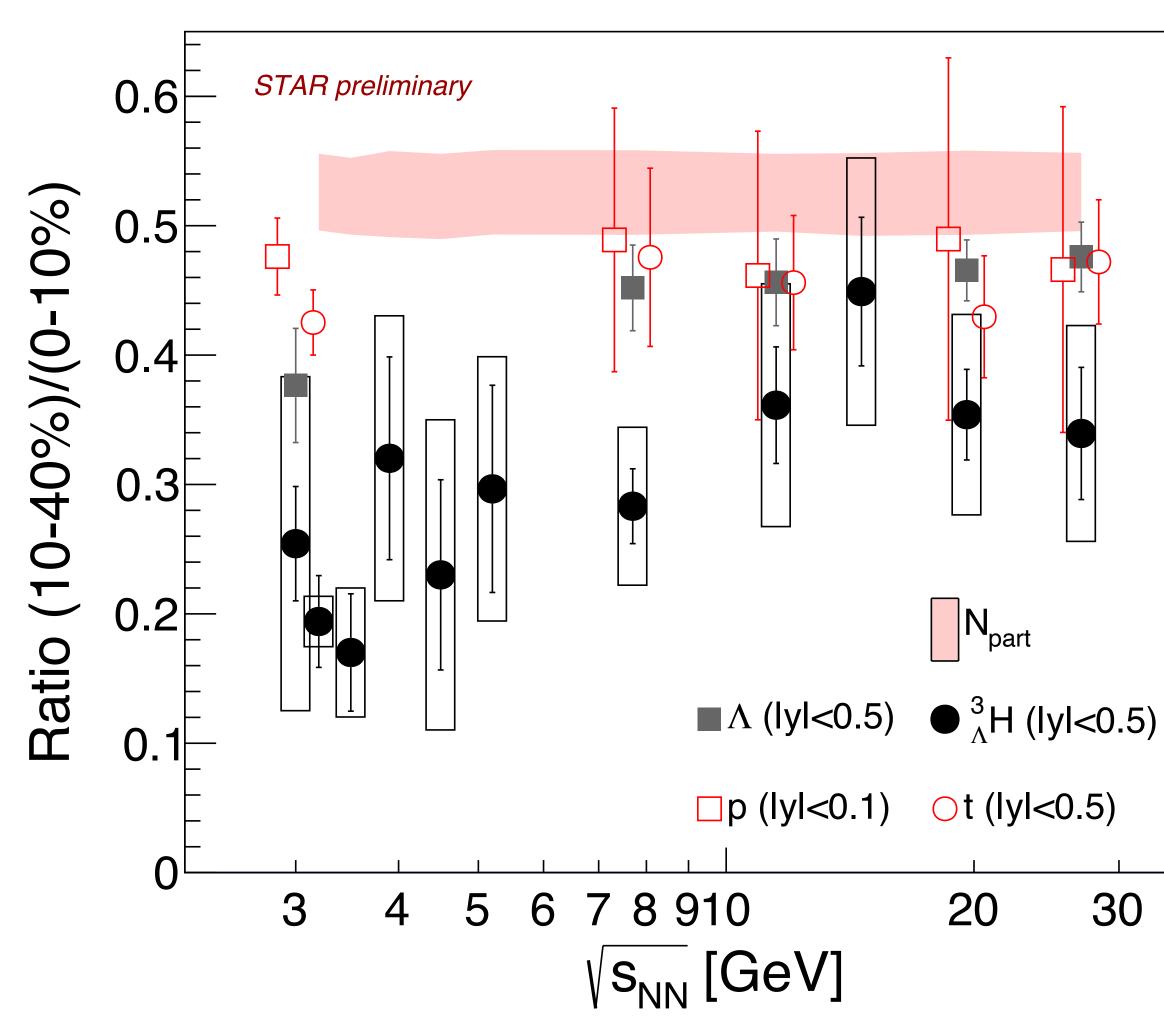
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 Similar trend in central (0-10%) and mid-central (10-40%) collisions



Centrality dependence of ${}^{3}_{\Lambda}H$ production

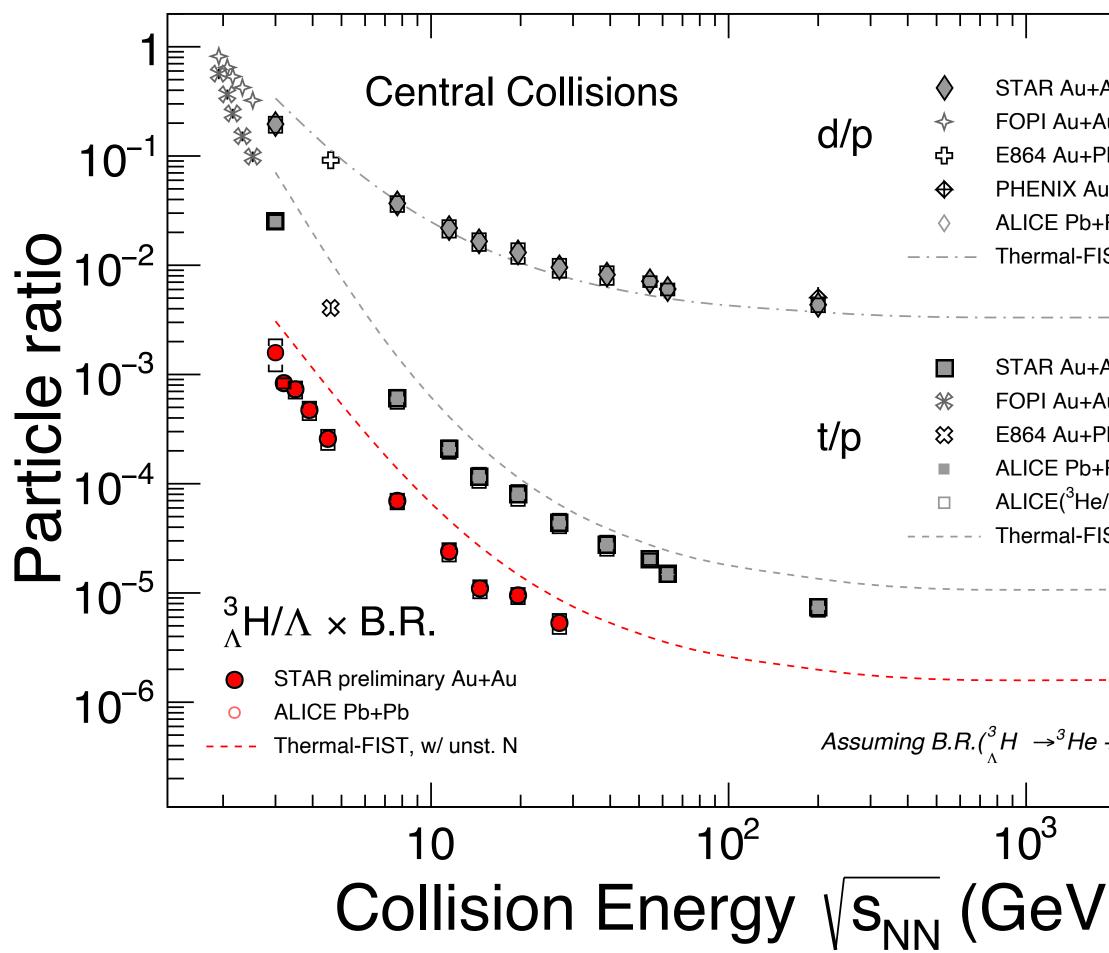


- Suppression of mid-central/central $^{3}_{\Lambda}H$ yield ratio w.r.t N_{part} , seems more apparent below $\sqrt{s_{NN}} = 7.7 \text{ GeV}$
- ${}^{5}_{\Lambda}H$ yield ratio tends to increase more steeply than proton, Λ , triton at low energies

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Nuclei-to-hadron ratios



STAR, PRL 130 (2023) 202301STAR, arXiv: 2311.11020T. Reichert, et al, PRC 107 (2023) 014912

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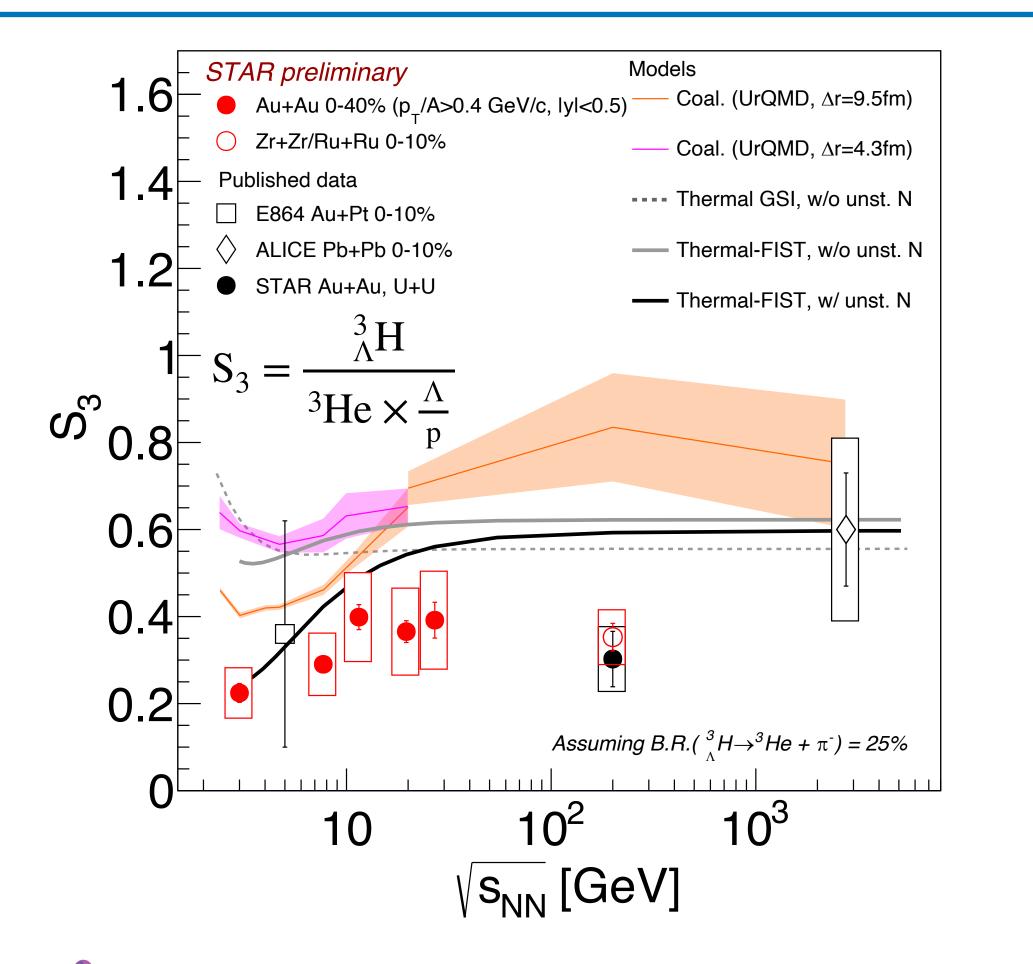
- Thermal model assumes that chemical freeze-out of light/hypernuclei happens at same time with hadrons
 - Particle yield ratio is independent of volume. ${}^3_{\Lambda}H/\Lambda$ yield ratio is dependent of strangeness correlation length
- d/p consistent with thermal model
- ${}^3_{\Lambda}H/\Lambda$, as well as t/p, overestimated by thermal model by a factor of ~2

Suggest ${}^3_\Lambda H$ and t yields are not in equilibrium and fixed at chemical freeze-out simultaneously with other hadrons



Energy dependence of strangeness population factor S_3

ullet



Science 328 (2010) 58 STAR, arXiv: 2310.12674 ALCE, PLB 754 (2016) 360 E864, PRC 70 (2004) 024902

A. Andronic et al, PLB 697 (2011) 203 (Thermal (GSI)) S. Zhang, PLB 684 (2010) 224 (Coal.+AMPT) T. Reichert, et al, PRC 107 (2023) 014912 (UrQMD, Thermal-FIST)

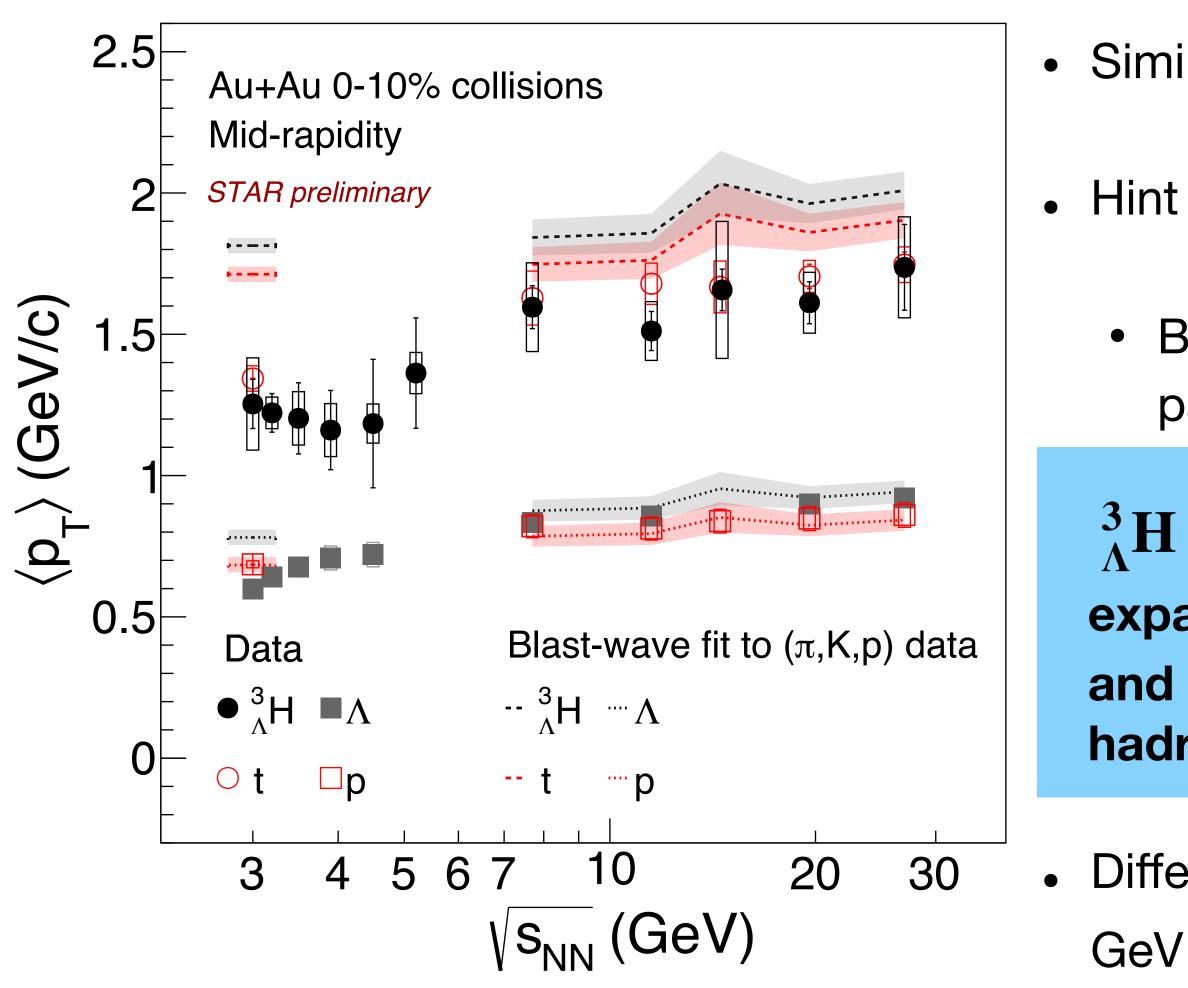
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- A prominent enhancement of S_3 was proposed as a probe S. Zhang et al. PLB 684 (2010) 224–227 for deconfinement
- Data shows a mild increasing trend from $\sqrt{s_{NN}} = 3.0$ GeV to 2.76 TeV
 - For coalescence(UrQMD) models, the energy dependence is sensitive to the **source radius** (Δr). Data favor larger radius.
 - Due to the difficulty in forming ${}^3_{\Lambda}H$ of large radius in small systems
- Thermal-FIST, which includes **feed-down** from unstable nuclei to stable p, 3 He, describes the S₃ data better
 - Possible feed-down should be accounted





Energy dependence of ${}^{3}_{\Lambda}H \langle p_{T} \rangle$



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- Similar $\langle p_T \rangle$ for ${}^3_{\Lambda}$ H and t
- Hint of ${}^3_{\Lambda}$ H and t $\langle p_T \rangle < \langle p_T \rangle^{BW}$ at $\sqrt{s_{NN}} > 7.7$ GeV
 - Blast-wave fit using measured kinetic freeze-out parameters from light hadrons (π , K, p)
 - $^{3}_{\Lambda}$ H and t might do not follow same collective expansion as light hadrons. Can be interpreted as $^3_{\Lambda}{
 m H}$ and t decoupling at different times compared to light hadrons
- Different trend for $\sqrt{s_{NN}}$ = 3-4.5 GeV and $\sqrt{s_{NN}}$ = 7.7-27
 - Suggest different expansion dynamics?



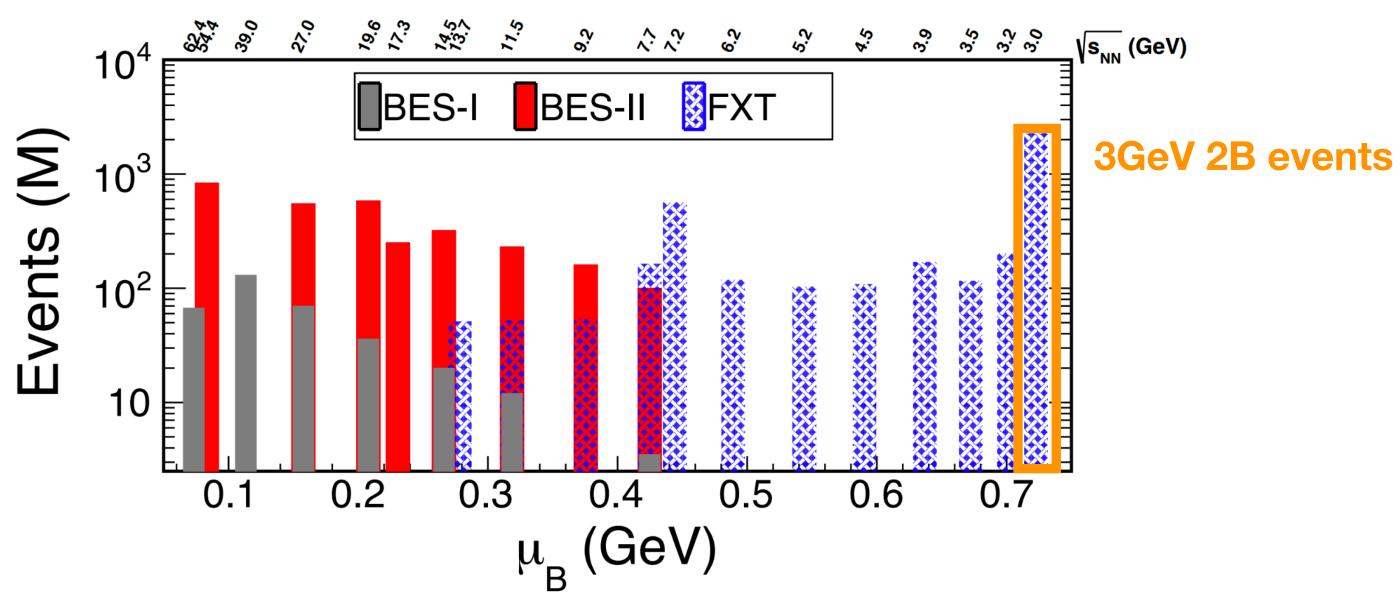
Summary

- ${}^{3}_{\Lambda}H$ yields and ${}^{3}_{\Lambda}H/\Lambda$ ratio in 0-10% collisions overestimated by thermal model, assuming chemical freeze-out of light/hypernuclei happens at same time with hadrons, by a factor of ~ 2
- ${}^{3}_{\Lambda}H\langle p_{T}\rangle$ overestimated by Blast-wave fit parameterization from light hadrons
 - time compared to the light hadrons
- Suppression of ${}^{3}_{\Lambda}H$ in 10-40% collisions at low collisions energies observed
- Energy dependence of S_3 suggests feed-down from unstable nuclei

• ${}_{\Lambda}^{3}$ H are likely formed at or decouples from the system at a different



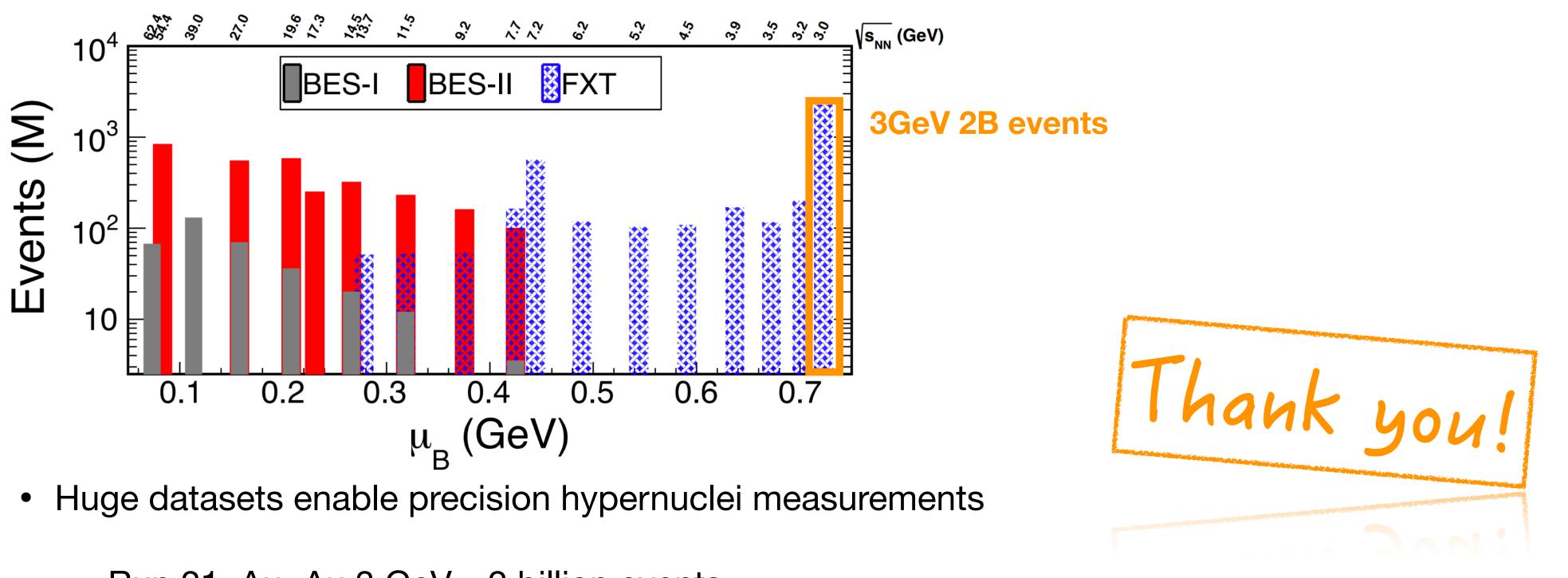
Outlook



- Huge datasets enable precision hypernuclei measurements
 - Run 21, Au+Au 3 GeV, ~2 billion events
 - Run 18, Isobar 200 GeV, ~6 billion events ullet
 - Run 23-25, Au+Au 200 GeV, ~18 billion events
- Opportunities for heavier hypernuclei: ${}^{4}_{\Lambda}H$, ${}^{4}_{\Lambda}He$, ${}^{5}_{\Lambda}He$, ${}^{6}_{\Lambda}H$, ${}^{A}_{\Lambda\Lambda}H$, ${}^{A}_{\Lambda\Lambda}He$



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