HIGGS POTENTIAL AND BSM OPPORTUNITIES



NNLO corrections to *W***-pair and** H^+H^- **production**

Zhe Li (黎哲) ShanDong University

Based on [JHEP 12 (2024) 038, JHEP 12 (2024) 136, 2409.08879]

In collaboration with Ren-You Zhang, Wen-Jie He, Zhi-Xing Zhang, Shu-Xiang Li and Xiao-Feng Wang

NNLO corrections to *W*-pair and H^+H^- production

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2024.12.21

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Outline





NNLO corrections to *W*-pair and H^+H^- production

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Higgs Discvover











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NNLO QCD-EW corrections to W-pair production at electron-positron colliders

Based on [JHEP 12 (2024) 038]

With Shu-Xiang Li and Xiao-Feng Wang, Ren-You Zhang, et. al

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$$e^+e^- \rightarrow W^+W^-$$
@ QCD-EW



Calculation setup

Consider the following process

$$e^+(p_1,\lambda_1) + e^-(p_2,\lambda_2) \to W^+(p_3,\lambda_3) + W^-(p_4,\lambda_4),$$

Unpolarized differential cross sections

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = \frac{\beta}{64\pi^2 s} \frac{1}{4} \sum_{\lambda_{1,\dots,4}} \left| \mathcal{M}(s,t,\lambda_{1,\dots,4}) \right|^2$$

Perturbative expansion



 $e^+e^- \rightarrow W^+W^-$ @ QCD-EW



Workflow of calculations

I. Generating Feynman amplitudes

- FeynArts
- QGRAF

II. Scattering amplitudes calculation

- 1. Dirac algebra, Lorentz algebra, color algebra evaluation
- 2. Tensor reduction
- 3. IBP reduction (Laporta algorithm)
- 4. MI calculation (CDEs)

III. UV divergence

- On-shell renormalization scheme
- MS renormalization scheme

IV. IR divergence

Phase-space slicing method

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Subtraction method

V. Phase-space integrals

- Monte-Carlo
- Reverse unitarity

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NNLO QCD \otimes EW corrections

• $\mathcal{O}(\alpha \alpha_s)$ corrections to amplitudes

$$\delta \mathcal{M}^{\mathcal{O}(\alpha\alpha_{s})} = \delta \mathcal{M}^{\mathcal{O}(\alpha\alpha_{s})}_{e\nu_{e}W} + \delta \mathcal{M}^{\mathcal{O}(\alpha\alpha_{s})}_{eeV} + \delta \mathcal{M}^{\mathcal{O}(\alpha\alpha_{s})}_{\gamma/Z-S.E.} + \delta \mathcal{M}^{\mathcal{O}(\alpha\alpha_{s})}_{VWW}$$
Feynman diagrams
$$\times \mathcal{O}(\alpha_{s}) \qquad \bigotimes \mathcal{O}(\alpha\alpha_{s})$$

$$\xrightarrow{e\nu_{e}W \text{ vertex}} 2\text{-loop massive triangle FIs}$$

$$\gamma/Z \text{ self-energy}$$

eeV vertex

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VWW vertex

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Integral family

Two-loop triangle integrals

$$F(\alpha_1,\ldots,\alpha_7) = \int \mathcal{D}^d l_1 \mathcal{D}^d l_2 \frac{1}{D_1^{\alpha_1} \ldots D_7^{\alpha_7}}$$

Propagators

Full mass dependence!

$$D_{1} = l_{1}^{2} - m_{t}^{2}, \quad D_{2} = l_{2}^{2} - m_{t}^{2}, \quad D_{3} = (l_{1} - p_{3})^{2} - m_{b}^{2}, \quad D_{4} = (l_{2} - p_{3})^{2} - m_{b}^{2}$$
$$D_{5} = (l_{1} - p_{3} - p_{4})^{2} - m_{t}^{2}, \quad D_{6} = (l_{2} - p_{3} - p_{4})^{2} - m_{t}^{2}, \quad D_{7} = (l_{1} - l_{2})^{2}.$$

Topologies On-shell W: this talk
Off-shell W: X.F Wang's talk





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General solutions

See Prof. YZ's talk

Path-ordered iterated integral:

$$\mathbf{I}(\vec{x},\epsilon) = \mathcal{P}\exp\left\{\epsilon\int_{\gamma}d\mathbb{A}\right\}\mathbf{I}\left(\vec{x}_{0},\epsilon\right)\left\{\mathbf{Boundary}\right\}$$
Boundary

Order by order

$${f I}^{(n)}(ec x) = egin{cases} & {f I}^{(0)}(ec x_0) & n = 0, \ & {f I}^{(n)}(ec x_0) + \int_\gamma d{\mathbb A}(ec x){f I}^{(n-1)}(ec x) & n > 0. \end{cases}$$

Iterated integrals

K(t) = dA/dt

$$\int_{\gamma} \underbrace{\mathrm{d}A \cdots \mathrm{d}A}_{n \text{ times}} = \int_{0}^{1} \mathrm{d}t_{1} \underbrace{K(t_{1})}_{0} \int_{0}^{t_{1}} \mathrm{d}t_{2} \underbrace{K(t_{2})}_{0} \cdots \int_{0}^{t_{n-1}} \mathrm{d}t_{n} \underbrace{K(t_{n})}_{\text{[K. Chen, `77]}}$$

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Rationalization

Three square roots

$$\lambda_1 = \sqrt{s(s - 4m_t^2)}, \qquad \lambda_2 = \sqrt{s(s - 4m_W^2)},$$

$$\lambda_3 = \sqrt{(m_W^2 - m_t^2 - m_b^2)^2 - 4 m_t^2 m_b^2}.$$

Change of variables

$$\frac{s}{m_t^2} = -\frac{(1-x)^2}{x}, \quad \frac{m_W^2}{m_t^2} = -\frac{(1-x)^2 z}{x(1+z)^2}, \quad \frac{m_b^2}{m_t^2} = \left(1 + \frac{y}{z+1}\right) \left(1 - \frac{(1-x)^2 z}{xy(z+1)}\right)$$

Rationalized simultaneously

$$\lambda_1' = \frac{(1-x)(1+x)}{x}, \quad \lambda_2' = \frac{(1-x)^2(1-z)}{x(1+z)}, \quad \lambda_3' = \frac{(1-x)^2z + xy^2}{xy(1+z)}$$

$$e^+e^- \rightarrow W^+W^-$$
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Multiple polylogarithms

See Prof. YZ's talk

$$\mathbf{I}^{(n)}(\vec{x}) = \mathbf{I}^{(n)}(\vec{x}_0) + \int_{\gamma} d\mathbb{A}(\vec{x}) \mathbf{I}^{(n-1)}(\vec{x}) \quad d\mathbb{A}(\vec{x}) = \sum_{i} \mathbb{M}_i d\log \eta_i$$

* MPLs

$$\begin{aligned} \mathbf{G}(a_1, ..., a_n; z) = \int_{0}^{z} \frac{1}{t - a_1} G(a_2, ..., a_n; t) dt, \\ \text{with} \\ G(a_1; z) = \int_{0}^{z} \frac{1}{t - a_1} dt \ a_1 \neq 0, \quad G(\underbrace{0, ..., 0}_{n \text{ times}}; z) = \frac{\log^n(z)}{n!} \end{aligned}$$

* eMPLs

[Goncharov, `98,`01]

Unrationalizable square root See X.F Wang's talk

$$e^+e^- \rightarrow W^+W^-$$
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Boundary conditions

• Vanishing inputs ($s = m_W^2 = 0, m_b^2 = m_t^2$)

$$\mathbf{I}_{i}(\vec{x}_{0},\epsilon) = \begin{cases} 1, & i = 1, 2, 3, \\ 0, & i \neq 1, 2, 3. \end{cases}$$

♦ Vacuums $\left\{ \overbrace{I_{1,2,3}}^{2} \qquad \overbrace{I_{4,5,6,7,11,12}}^{2} \qquad \overbrace{I_{4,5,6,7,11,12}}^{2} \qquad \overbrace{I_{8,9,10,13,14}}^{2} \\
\overbrace{I_{5,17,20,23,26,30}}^{2} \qquad \overbrace{I_{16,18,19}}^{2} \qquad \overbrace{I_{21,22,24,25,27,28,31}}^{2} \qquad \overbrace{I_{29,32}}^{2} \\$ NNLO corrections to W-pair and H^+H^- production Zhe Li (SDU) 2024.12.21

$$e^+e^- \to W^+W^-@ \ QCD-EW$$

$$interpretation = 0$$

$$for the error of th$$

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Scale dependence



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Symmetry relation

 $\frac{\mathrm{d}\sigma}{\mathrm{d}\cos\theta_{W^-}} = \frac{\mathrm{d}\sigma}{\mathrm{d}\cos\theta_{W^+}}\Big|_{\theta\to\pi-\theta}\,,$

Strongly peak in the forward direction!!!

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Two-loop planar master integrals for NNLO QCD corrections to W-pair production at hadron colliders

Based on [JHEP 12 (2024) 136]

With Wen-Jie He, Shu-Xiang Li, Xiao-Feng Wang, et. al

Integral family

$$q(p_1) + \bar{q}(p_2) \to W^+(p_3) + W^-(p_4)$$

Plannar box-type topologies

Two-loop master integrals

 $S_3 \equiv (1, 1, 1, 1, 0, 0, 1, 1, 0)_F$

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$$F(n_1, \dots, n_9) = \int \mathcal{D}^d l_1 \mathcal{D}^d l_2 \frac{1}{D_1^{n_1} \dots D_9^{n_9}}$$

Propagators

$$D_{1} = l_{1}^{2}, \qquad D_{2} = (l_{1} + p_{1})^{2}, \qquad D_{3} = (l_{1} - p_{2})^{2},$$

$$D_{4} = l_{2}^{2} - m_{t}^{2}, \qquad D_{5} = (l_{2} + p_{1})^{2} - m_{t}^{2}, \qquad D_{6} = (l_{2} - p_{2})^{2} - m_{t}^{2},$$

$$D_{7} = (l_{1} - l_{2})^{2} - m_{t}^{2}, \qquad D_{8} = (l_{1} - p_{2} - p_{3})^{2}, \qquad D_{9} = (l_{2} - p_{2} - p_{3})^{2}.$$

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Linear basis

$$g_{1} = \epsilon^{2} f_{1}$$

 $g_{2} = \epsilon^{2} f_{2} x$
 $g_{3} = \epsilon^{2} f_{3} y$
 $g_{4} = \epsilon^{2} f_{4} z$
 $g_{5} = \epsilon^{3} f_{5} r'_{3}$
 $g_{6} = \epsilon^{3} f_{6} x y$
 $g_{7} = \epsilon^{2} f_{7} y$
 $g_{8} = \epsilon^{2} f_{8} r'_{2} + 1/2 f_{7} r'_{2}$

Linear form DEs

$$g_{9} = \epsilon^{3} f_{9} x$$

$$g_{10} = \epsilon^{2} f_{10} r'_{1} - \epsilon^{3} f_{9} r'_{1}$$

$$g_{11} = \epsilon^{3} f_{11} (y - z)$$

$$g_{12} = \left[\epsilon^{2} f_{12} z + \epsilon^{3} f_{11} (y - z) - \epsilon^{2} f_{8} y \right] r'_{4} / (y - z)$$

$$g_{13} = \epsilon^{3} \frac{\pi r'_{3}}{\Psi_{1}} f_{13}$$

$$g_{14} = \epsilon^{3} f_{14} r'_{3}$$

$$g_{15} = \frac{1}{\epsilon} \frac{\Psi_{1}^{2}}{2\pi i W_{y}} \frac{\partial g_{13}}{\partial y}$$

$$e. Square roots$$

$$r'_{1}^{2} = x(x - 4), r'_{2}^{2} = y(y + 4), r'_{3}^{2} = x(x - 4z), r'_{4}^{2} = (z - y)(y - z + 4).$$

$$f. Lower triangular$$

$$I. Cover triangular$$

$$A^{(0)}(\vec{x}) + \epsilon \, dA^{(1)}(\vec{x}) g(\vec{x}, \epsilon)$$

$$I. Adams, `18]$$

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 $d\mathbf{g}(\vec{x},\epsilon) =$

Validation

Good agreements!!!

Branch	\mathbf{G}	Results (Analytic / AMFlow)	
\mathcal{T}_{3F} -	g_{13}	$3.2548275606982903057\epsilon^3$	(Analytic)
		$3.2548275606982903051\epsilon^3$	(AMFlow)
	g_{14}	$-2.33756055601827687\epsilon^3$	(Analytic)
		$-2.33756055601827691\epsilon^3$	(AMFlow)
	g_{15}	$-0.094132461909778403737363\epsilon^2$	(Analytic)
		$-0.0715655314006887\epsilon^3$	
		$- 0.094132461909778403737361 \epsilon^2$	(AMFlow)
		$-0.0715655314006885\epsilon^3$	

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Mixed QCD-EW corrections to THDM charged Higgs-pair production at lepton colliders

Based on [2312.17207]

With Zhi-Xing Zhang, Shu-Xiang Li, Ren-You Zhang, et. al

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 $e^+e^- \rightarrow H^+H^- @ QCD-EW$

THDM background

- Simplest extensions of SM Higgs sector
- Scalar sector of many other models (SUSY)

 $\rightarrow H^+H^-$ @ QCD-EW

 $e^+e^- \rightarrow H^+H^-$ @ QCD-EW

 $e^+e^- \rightarrow H^+H^- @ QCD-EW$

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 m_{H^+} scan plots

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$$e^+e^- \rightarrow H^+H^- @ \text{QCD-EW}$$

$$(aa_s) \text{ corrections density plots}$$

$$(aa_s) \text{ corrections den$$

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NNLO QCD-EW corrections to W-pair production at lepton colliders (~1%)

Planar MIs for NNLO QCD corrections to W-pair production at hadron colliders (elliptic integrals)

NNLO QCD-EW corrections to charged Higgspair production at lepton colliders (~1%)

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Thanks

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Backup

NNLO corrections to *W*-pair and H^+H^- production

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$$e^+e^- \rightarrow W^+W^-$$
@ QCD-EW

NLO EW corrections

EW corrections

 $\Delta \sigma_{\rm EW} = \sigma_{\rm virtual} + \sigma_{\rm real} + \sigma_{\rm h.o.ISR}$

Virtual corrections

$$\frac{\mathrm{d}\sigma_{\mathrm{virtual}}}{\mathrm{d}\Omega} = \frac{\beta}{64\pi^2 s} \frac{1}{4} \sum_{\lambda_{1,\dots,4}} 2\operatorname{Re}\mathcal{M}_0^* \delta\mathcal{M}^{\mathcal{O}(\alpha)} \xrightarrow{\text{On-shell}}_{\text{[A.Denner `97]}}$$

Real photon radiation

×

 $\sigma_{\text{real}} = \sigma_{\text{S}}(\delta_s) + \sigma_{\text{HC}}(\delta_s, \delta_c) + \sigma_{\text{H}\bar{\text{C}}}(\delta_s, \delta_c) \frac{\text{Slicing method}}{\text{[B. W. Harris, `02]}}$

$$\sigma_{\rm h.o.ISR} = \sigma_{\rm ISR-LL} - \sigma_{\rm LL,sub}$$
 Structure functions

[W. Beenakker, `96]

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$$e^+e^- \rightarrow W^+W^-$$
 (MIs)

Analytic continuation

Euclidean region

$$s < 0 \land m_W^2 < 0$$

$$\frac{s + i0}{m_W^2 - i0}$$

$$s + i0^+, m_W^2 + i0^+$$

 $n_t^2 - i0^+, m_b^2 - i0^+$

Physical region
$$s > 4 m_W^2 > 0$$

Inverse transfrom

$$\begin{aligned} x &= \frac{\sqrt{4 \, m_t^2 - s} - \sqrt{-s}}{\sqrt{4 \, m_t^2 - s} + \sqrt{-s}} \,, \qquad z &= \frac{\sqrt{-s} - \sqrt{4 \, m_W^2 - s}}{\sqrt{-s} + \sqrt{4 \, m_W^2 - s}} \,, \\ y &= (1+z) \, \frac{1}{2 \, m_t^2} \left[\, m_b^2 - m_t^2 - m_W^2 + \sqrt{\lambda(m_W^2, \, m_t^2, \, m_b^2)} \, \right] \,. \end{aligned}$$

Family	Kinematic configuration	x	z	y
${\cal F}$	$2m_W < \sqrt{s} < 2m_t$	$e^{iartheta}$	z	$y - i 0^+$
	$2 m_t < \sqrt{s}$	$x + i 0^+$	z	$y-i0^+$
\mathcal{F}^*	$2 m_W < \sqrt{s}$	$x + i 0^+$	z	$y-i0^+$

$$e^+e^- \rightarrow W^+W^-$$
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Kinematic distributions

Symmetry relation

$$\frac{\mathrm{d}\sigma}{\mathrm{d}p_{T,W^-}} = \frac{\mathrm{d}\sigma}{\mathrm{d}p_{T,W^+}}$$

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Boundary conditions

Expansion @ $\vec{x}_0 = (x, 1, 0)$ $A_x^{(0,1)}(\vec{x}) = A_{x,(0,0)}^{(0,1)}(x), \quad A_y^{(0,1)}(\vec{x}) = \frac{A_{y,(-1,0)}^{(0,1)}}{\vec{x}}, \quad A_z^{(0,1)}(\vec{x}) = \frac{A_{z,(0,-1)}^{(0,1)}}{\vec{x}}$ Lowest-order approximation $d\tilde{\mathbf{g}}(\vec{x},\epsilon) = \epsilon \left[A_{x,(0,0)}^{(1)}(x) dx + \frac{1}{v} A_{y,(-1,0)}^{(1)} \right] \tilde{\mathbf{g}}(\vec{x},\epsilon)$ $\tilde{g}(\vec{x}) = g(\vec{x}_0)$ Degeneration $r_1' = r_3' = \sqrt{x(x-4)}$ $\Psi_1 = \frac{\pi}{2}\sqrt{x-4}\,,$ $r_2' = 0$, $\partial_y \Psi_1 = \frac{\pi}{8} \left(2 - x \right) \sqrt{x - 4} \,.$ $r'_{\scriptscriptstyle A} = \sqrt{3}$, Boundary • Known integrals: g_1, \dots, g_8 [S. Di Vita, 17] • Symmetry relation: $g_{11} = -g_9$, $g_{12} = ig_{10}$ (a) (x, y, z) = (1, 0, 1)

• Regularity condition:

 $\bigstar x \to 0: g_9 \qquad \bigstar x \to 3: g_{13} \qquad \bigstar x \to 4: g_{14}$

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THDM types

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Constraints on THDM parameters

Experiment constraints:

- Direct searches for BSM Higgs bosons [M. Aiko, `10, G. Abbiendi, `13]
- Flavor constraints: $B \rightarrow X_s \gamma$ [M. Misiak, `20, V. Khachatryan, `14, R. Aaij, `17]
- Anomalous magnetic moment of the muon [J. Haller, `18]
- SM Higgs boson coupling measurements [G. Aad, `15]
- Oblique parameters [W. Grimus, `15, W. Grimus, `08, R. Workma, `22]

Theoretical constraints:

- Vacuum stability [S. Nie, `99, A. Barroso, `13, V. Branchina, `18,]
- Perturbative unitarity [S. Kanemura, `15]

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$\tan\beta$ scan plots

