

Dec 20th, 2024, Hefei

OutLine

- Introduction
- Some technical details
- Numerical results
- Summary



- Experimental evidence
- Previous calculations

Experimental evidence

3.4 *σ*

Evidence for the Higgs Boson Decay to a Z Boson and a Photon at the LHC

ATLAS and CMS Collaborations • Georges Aad (Marseille, CPPM) Show All(5264) Sep 7, 2023

32 pages Published in: Phys.Rev.Lett. 132 (2024) 2, 021803 Published: Jan 11, 2024 e-Print: 2309.03501 [hep-ex] DOI: 10.1103/PhysRevLett.132.021803 (publication) PDG: $H \rightarrow Z\gamma$ Show All(2) Report number: CERN-EP-2023-157 Experiments: CERN-LHC-ATLAS, CERN-LHC-CMS View in: CERN Document Server, HAL Science Ouverte, OSTI Information Bridge Server, ADS Abstract Service **Γ**⁴ cite 🗟 claim D pdf @ links datasets reference search \Rightarrow 57 citations -7

ATLAS and CMS, PRL 132 (2024) 021803 [2309.03501]

Evidence (2309.03501)

 obtained from a
 combination of ATLAS
 (2005.05382) and CMS
 (2204.12945).

Previous searches: 1806.05996, 1402.3051, 1307.5515, 0806.0611 ...

Experimental evidence



Branching Ratio: $(3.4 \pm 1.1) \times 10^{-3}$ Signal strength: $\mu = 2.2 \pm 0.7$ $(\mu = \sigma_{exp} / \sigma_{SM})$

Motivation: provide SM prediction as precise as possible.

ATLAS and CMS, PRL 132 (2024) 021803 [2309.03501]

Previous calculations

• Leading order.

J.R. Ellis, M. K. Gaillard, D. V. Nanopoulos, Nucl. Phys. B 106 (1976) 292 R.N. Cahn, M. S. Chanowitz, N. Fleishon, Phys. Lett. B 82 (1979) 113-116

• QCD corrections:

M. Spira, A. Djouadi, P. M. Zerwas, Phys. Lett. B 276 (1992) 350-353

T. Gehrmann, S. Guns, D. Kara, JHEP 09 (2015) 038 [1505.00561]

R. Bonciani, V. D. Duca, H. Frellesvig, et al, JHEP 08 (2015) 108 [1505.00567]

Previous calculations Why NLO EW?

The QCD corrections amount to 0.22% of the LO width. It is resonable to expect that the EW corrections will yield a larger contribution.

6.68 keV

The theoretical uncertainty has not been fully discussed

 $\Gamma = \frac{G_F^2 \alpha m_W^2}{4 m_W^3 (m_W^2 - m_Z^2)} |A|^2$ 7.09 keV

1505.00567: $\Gamma_{H \to Z\gamma} = \frac{G_F \alpha^2}{64\sqrt{2\pi^3}m_H} \frac{(m_H^2 - m_Z^2)^3}{m_H^2} |\mathcal{F}|^2$

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1505.00561:

Previous calculations

Signal background interference: F. Buccioni, F. Devoto, A. Djouadi, J. Ellis, et. al, Phys. Lett. B 851 (2024) 138596 [2312.12384]





Interference effect (Higgs resonance & nonresonance) is negative and small (about -3%)...



- Lorentz structure
- Feynman integral
- Electroweak coupling

Some technical details

Lorentz structure



$$\begin{split} H &\to Z(p_1) \gamma(p_2) \\ \mathcal{M} &= T^{\mu\nu} \varepsilon^*_\mu(p_1) \varepsilon^*_\nu(p_2) \end{split}$$

 $T^{\mu\nu} = T_1 p_1^{\mu} p_1^{\nu} + T_2 p_2^{\mu} p_2^{\nu} + T_3 p_1^{\mu} p_2^{\nu} + T_4 p_2^{\mu} p_1^{\nu} + T_5 g^{\mu\nu} + T_6 \epsilon^{\mu\nu\rho\sigma} p_{1\rho} p_{2\sigma}$ Constraint:

 $T_1 = 0, T_5 = -p_1 \cdot p_2 T_4$ (Gauge invariance).

 T_2, T_3, T_6 do not contribute to $|\mathcal{M}|^2$.

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Feynman integrals

About 5000 Feynman diagrams (Feynman gauge)



We work in the on-shell renormalization scheme.

• $\alpha(0) \approx 1/137$, *eey*-coupling in the zero momentum transfer limit (p' = p)



For non-zero momentum transfer (like $(p - p')^2 = M_Z^2$), $\ln\left(\frac{m_f}{M_Z}\right)$ terms survive.

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$$\alpha(M_Z^2) \approx 1/128$$
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• $\alpha_{G_{\mu}} = \sqrt{2}G_{\mu}M_{W}^{2}/\pi(1 - M_{W}^{2}/M_{Z}^{2}) \approx 1/132$, defined through muon

decay, $\alpha_{G_{\mu}} = \alpha(0) (1 + \Delta r^{(1)}).$

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no $\ln(m_f)$ terms

We use five different electroweak coupling schemes:

- 1. $\alpha(0)$ scheme: all to $\alpha(0)$;
- 2. $\alpha(M_Z^2)$ scheme: all to $\alpha(M_Z^2)$;
- 3. $\alpha_{G_{\mu}}$ scheme: all to $\alpha_{G_{\mu}}$;
- 4. mixed 1: the one connecting to the external photon to $\alpha(0)$, others to $\alpha(M_Z^2)$;
- 5. mixed 1: the one connecting to the external photon to $\alpha(0)$, others to $\alpha_{G_{u}}$.





- Numerical results
- Other potential uncertainty

sources

TABLE I. The LO and NLO decay widths of $H \rightarrow Z\gamma$ under different coupling schemes. The relative EW corrections are also given.

Scheme	Input parameters	Γ ^{LO} (keV)	$\Gamma_{\rm EW}^{\rm NLO}$ (keV)	δ_{EW} (%)
$\alpha(0)$	$\alpha(0), m_f$	5.920	6.234	5.3
$\alpha(m_Z^2)$	$\alpha(m_Z^2), m_f$	7.273	6.303	-13
G_{μ}	G_{μ}, m_f	6.599	6.343	-3.9
Mixed 1	$\alpha(0), \alpha(m_Z^2)$	6.791	6.316	-7.0
Mixed 2	$\alpha(0), G_{\mu}$	6.364	6.316	-0.75

 $\Gamma^{\text{LO}} = 6.364^{+0.909}_{-0.444} \text{ keV}, \Gamma^{\text{NLO}}_{\text{EW}} = 6.316^{+0.027}_{-0.082} \text{ keV}.$

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The EW corrections

- more significant than the QCD corrections
- greatly suppress the theoretical uncertainty

Combining EW, QCD, *b*-mass corrections: $\Gamma^{H \to Z\gamma} = 6.348^{+0.028}_{-0.085}$ keV



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Other potential uncertainty sources

Finite width effect / proper definition of $H \rightarrow Z\gamma$ —— subtle in theory



Other potential uncertainty sources

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phenomenologically:

- *H*-width effect: about -3%
 [PLB 851 (2024) 138596]
- Z-width effect: may reach 10%
 [PLB 727 (2013) 424, PRD 89
 (2014) 3, 033013]

larger than quick estimation $\mathcal{O}(\Gamma/M)$

Alternative way: study the $H \rightarrow \ell \overline{\ell} \gamma$ process, experimentally and theoretically.

Summary

- We obtain the most accurate prediction $Br(H \rightarrow Z\gamma) = 1.56^{+0.07}_{-0.08} \times 10^{-3}$ by including the NLO EW corrections.
- The EW corrections are more significant than the QCD corrections, and can greatly suppress the theoretical uncertainty.
- The state-of-the-art SM prediction is significantly lower than the measured value. This could probably be attributed to the underestimated experimental uncertainties or the new physics beyond the SM.

Cross check:

W. L. Sang, F. Feng, Y. Jia, Phys. Rev. D 110 (2024) 5, L051302 [2405.03464]

Thank You!