

半导体量子芯片

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中国科学技术大学

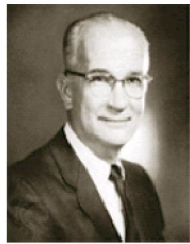
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- 半导体普适量子逻辑门
- 半导体多量子比特扩展
- 量子数据线和量子比特的长程耦合
- 类石墨烯单电子晶体管和量子点

半导体芯片起源

20世纪50—60年代物理学家针对电子管面临的物理极限问题，研发了晶体管和集成电路，实现了信息技术从电子管到晶体管芯片的第一次跨越式变革。

1956年度诺贝尔物理奖



肖克利 (William Shockley)

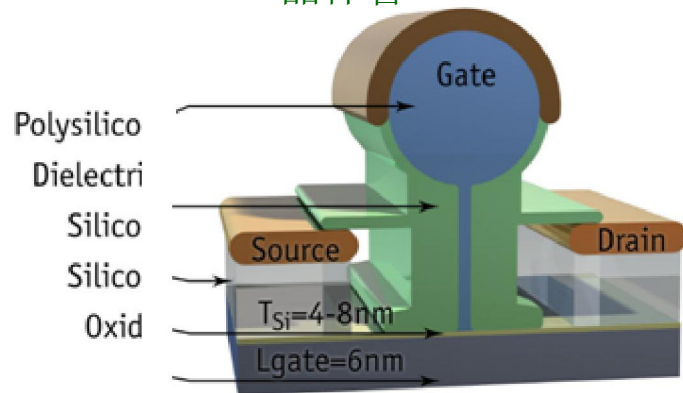


巴丁 (John Bardeen)



布拉坦 (Walter Brattain)

晶体管

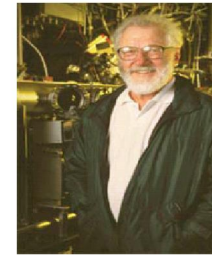


电流导通 1状态 比特(经典信息单元)
电流截止 0状态

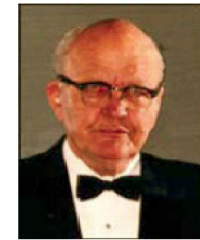
2000年度诺贝尔物理奖



若尔斯-阿尔费罗夫 (ZHORES I. ALFEROV)



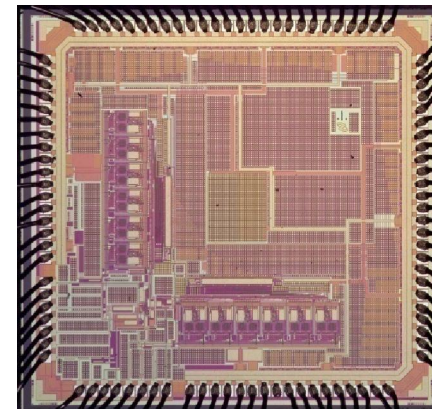
赫伯特-克勒默 (HERBERT KROEMER)



杰克-S-基尔比 (JACK S. KILBY)

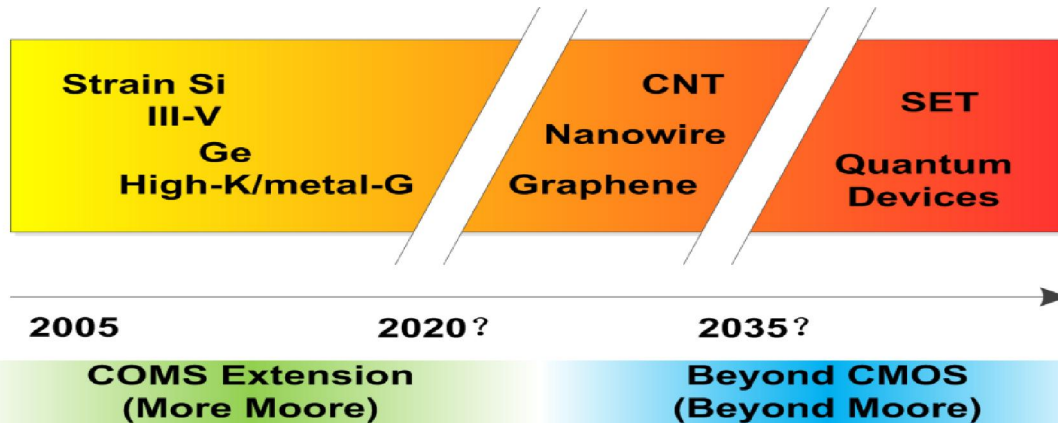
芯片(集成电路)

集成



挑战和机遇

随着摩尔定律的发展，尺寸等导致的量子效应越来越明显，发热和功耗也越来越严重。



半导体信息技术第二次
变革的浪潮即将来临!



International Technology Roadmap for Semiconductors

NATURE | NEWS FEATURE

عربي

No Moore 2016 ITRS

The chips are down for Moore's law

The semiconductor industry will soon abandon its pursuit of Moore's law. Now things could get a lot more interesting.

M. Mitchell Waldrop

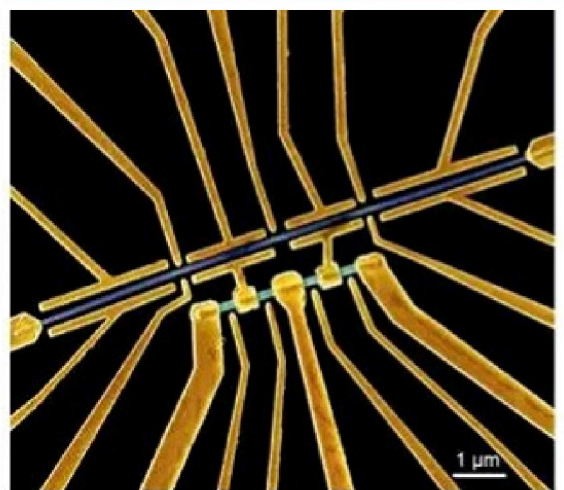
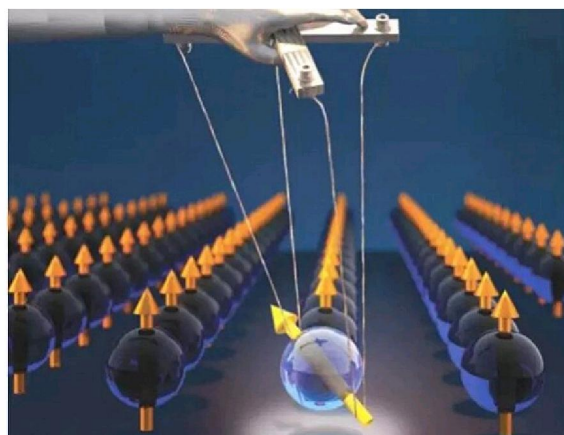
09 February 2016



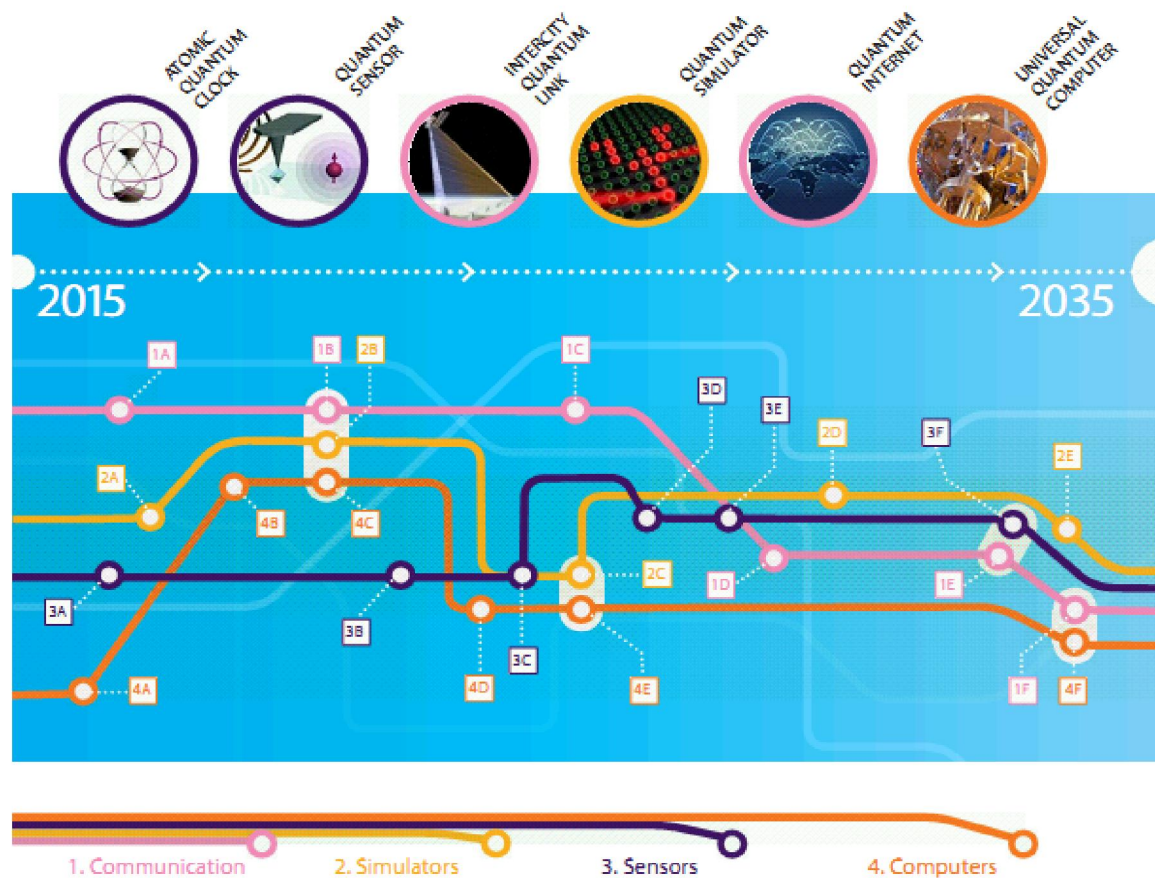
国内外进展与趋势



欧洲“量子宣言”
10亿欧元



Quantum Technologies Timeline



国内外进展与趋势



UNSW AUSTRALIA | Newsroom

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Prime Minister hails UNSW's quantum computing research as the world's best

22 APR 2016 | UNSW MEDIA

There is no bolder idea than quantum computing, said Prime Minister Malcolm Turnbull, hailing UNSW's research in the transformative technology as the "best work in the world".



Commonwealth Bank Chief Information Officer David Whiteing said: "Commonwealth Bank is proud to support the University of New South Wales' world-leading quantum computing research team and join the Australian Government in providing tangible support for their National Innovation and Science Agenda.

澳大利亚政府成立**硅基半导体量子芯片实验室**，抢占半导体量子芯片发展的制高点。

研究背景：国内外进展与趋势



“核、高、基”专项

首页 工作动态 通知公告 政策文件 专项介绍 专题研究 办事指南 工作简报 媒体聚焦 联系方式

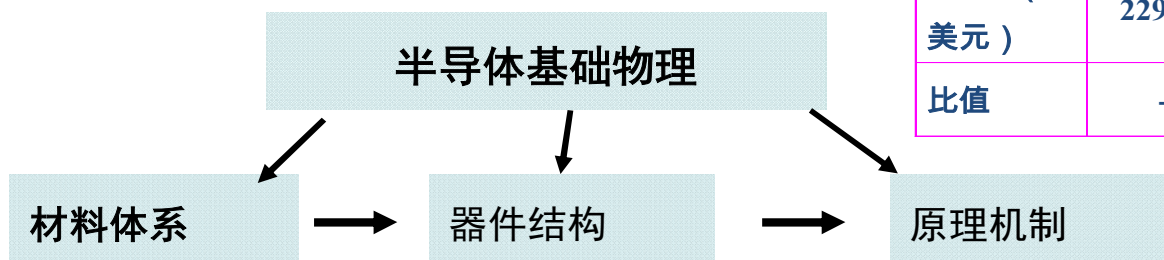
首页 > 专项介绍 > 核心电子器件、高端通用芯片及基础软件产品

错失半导体技术第一次变革缺乏核心专利！

2015年我国集成电路和部分产品的进口额比较

	集成电路	铁矿砂	石油（含原油、成品油）	农产品
进口额（亿美元）	2299.3	576.2	1344.5	1168.8
比值	-	3.99	1.71	1.97

核心电子器件、高端通用芯片及基础软件产品



技术、平台和队伍等基础薄弱

- 纳米结构制备技术和平台
- 极低温量子输运技术和平台

研究背景：国内外进展与趋势



中华人民共和国科学技术部

Ministry of Science and Technology of the People's Republic of China

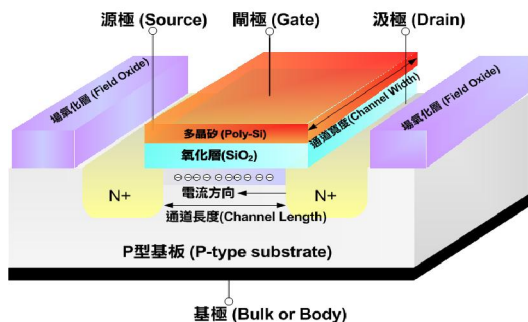
2011年“十二五”导向性重大项目（超级973）“固态量子芯片研究”
要求2015年前实现比特数3的量子芯片。

JF *** 半导体量子芯片***

2016年“十三五”重点研发计划“半导体量子芯片研究”，要求2020
年前获得品质因子1000、比特数6的量子芯片。

研究背景：传统芯片和量子芯片比较

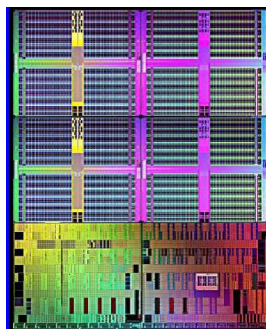
半导体晶体管



电流导通 1状态
电流截止 0状态

比特 (经典信息单元)

半导体芯片
(处理器-CPU)



避免克服量子效应

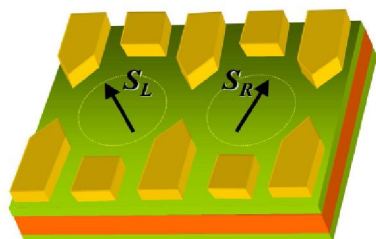
计算机



- 可逆运算解决功耗和发热问题
- 指数加速的计算能力

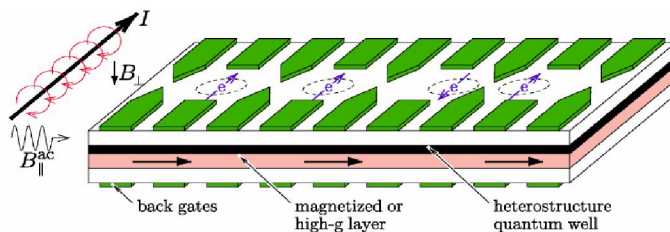
信息的量子化

单电子晶体管



电子自旋向上 1状态
电子自旋向下 0状态

半导体量子芯片
(量子处理器-Quantum CPU)



开发利用量子效应

量子计算机



传统芯片和量子芯片比较

半导体晶体管

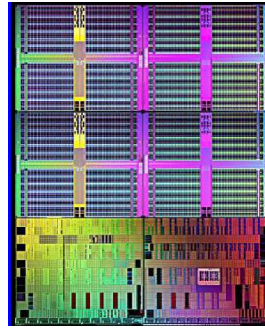
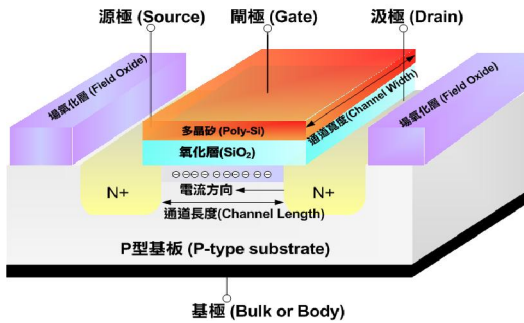


半导体芯片

(处理器-CPU)



计算机



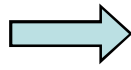
避免克服量子效应

对N个量子存储器实行一次操作，相当于对经典存储器进行2的N次方个操作。

电流导通 1状态
电流截止 0状态
比特(经典信息单元)



单电子晶体管

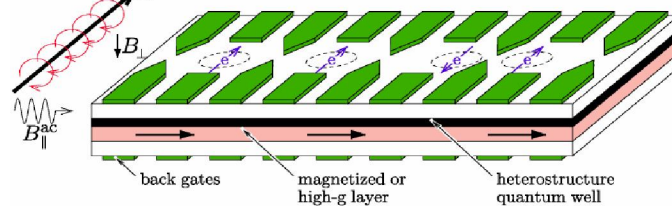
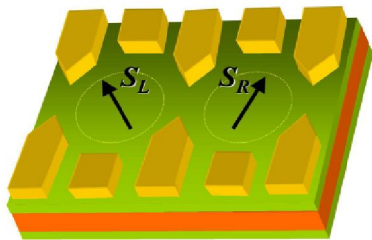


半导体量子芯片

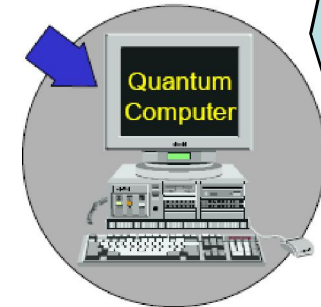
(量子处理器-Quantum CPU)



量子计算机



开发利用量子效应



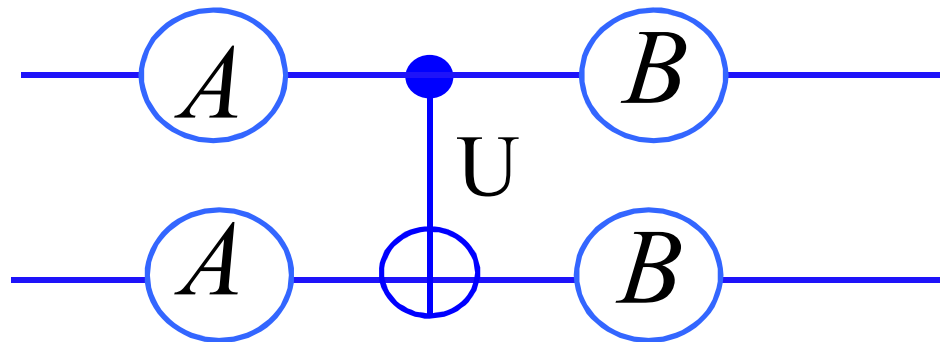
电子自旋向上 1状态
电子自旋向下 0状态
量子比特(量子信息单元)

量子芯片核心任务

普适量子计算：

单比特逻辑门+两比特逻辑门 (CNOT)

A、B 是两个不平行的单比特逻辑门



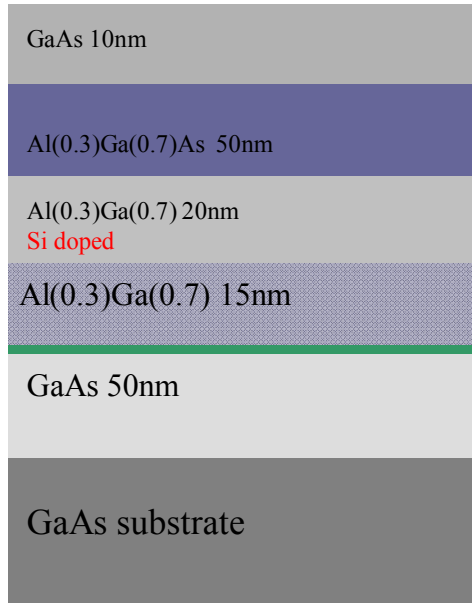
U 为两比特逻辑门

实用化量子计算要求：比特品质因子 $Q > 10000$

$Q = \text{Min}(\text{相干时间}) / \text{Max}(\text{基本逻辑门时间})$

- 核心任务：
1. **构造**高品质的量子比特；
 2. **调控**量子比特实现普适逻辑门
 3. **扩展**量子比特的量子数据总线架构

半导体基片材料的选取和设计

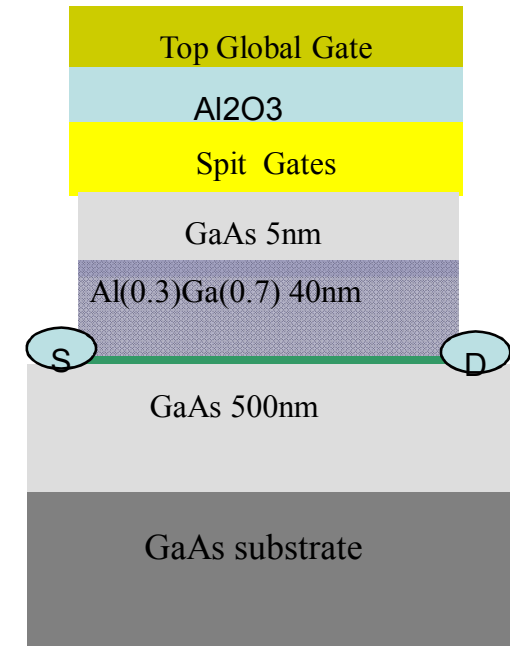


掺杂/非掺杂 GaAs量子点

#28 的 Si 是 1.0×10^{18}

#34 的 Si 应该是 9.0×10^{17}

掺杂越高浓度越大，量子点的制备和操控越简单，但本底噪声越大！



Customer Part: Custom GaAs 29

GaAs FET Structure Specification

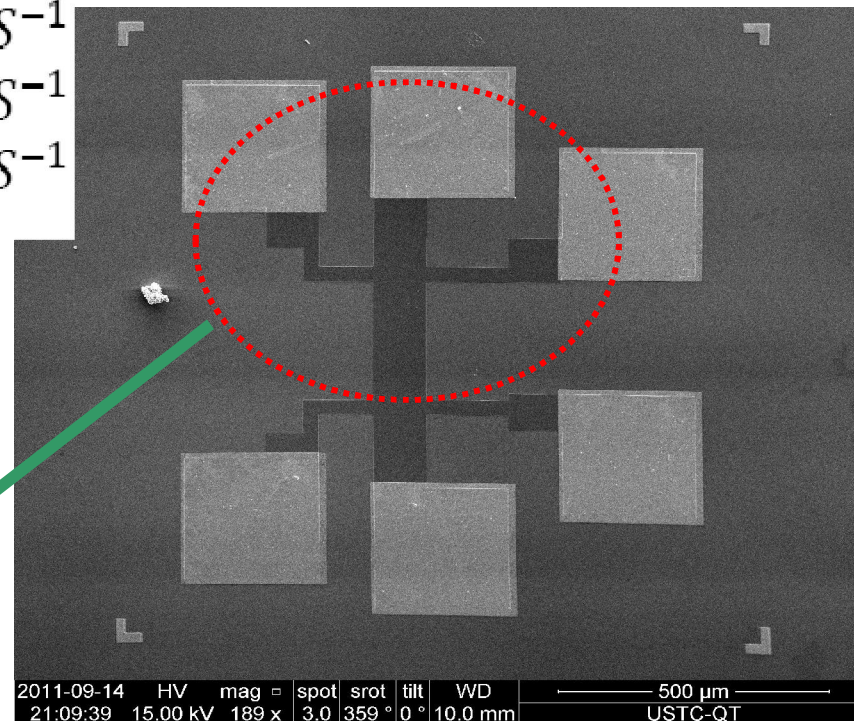
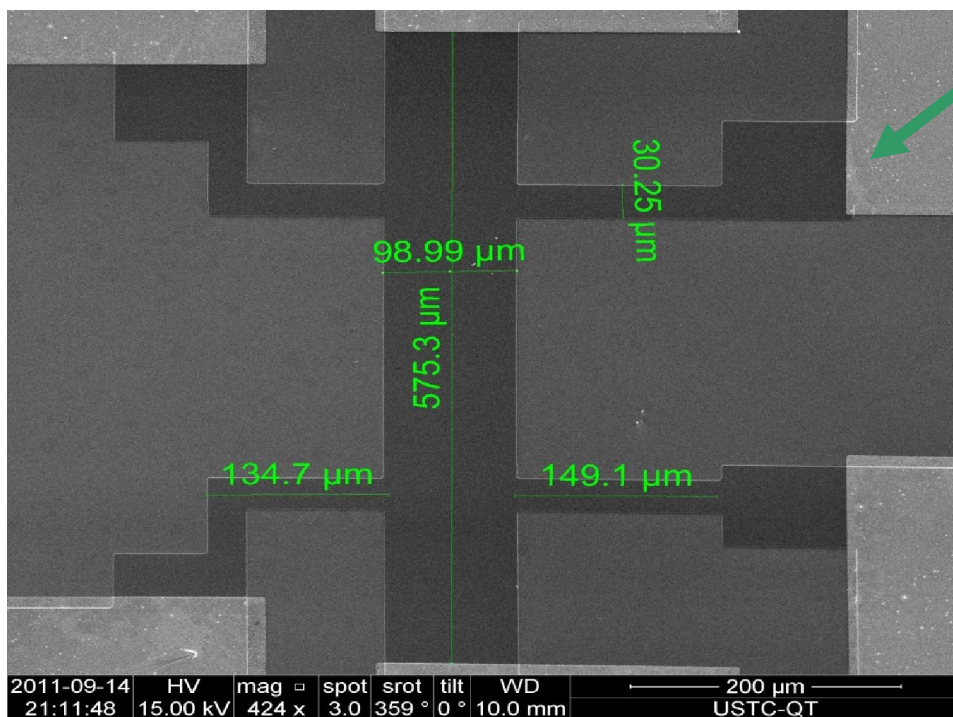
Layer	Type	Material	Group	Repeat	Mole Fraction (x)	Mole Fraction (y)	Strain (ppm)	PL (nm)	Thickness (Å)	Dopant	CV Level	Comments
5	i	GaAs							100			
4	i	Al(x)GaAs			30.0				500			
3	N	Al(x)GaAs			30.0				200	Si	7.0e17	
2	i	Al(x)GaAs			30.0				150			
1	i	GaAs							5000			
SUBSTRATE												

各种半导体基片参数的测定

#28: $n=3.2 \times 10^{11} \text{cm}^{-2}$ $\mu = 1.5 \times 10^5 \text{cm}^2 \text{V}^{-1} \text{S}^{-1}$

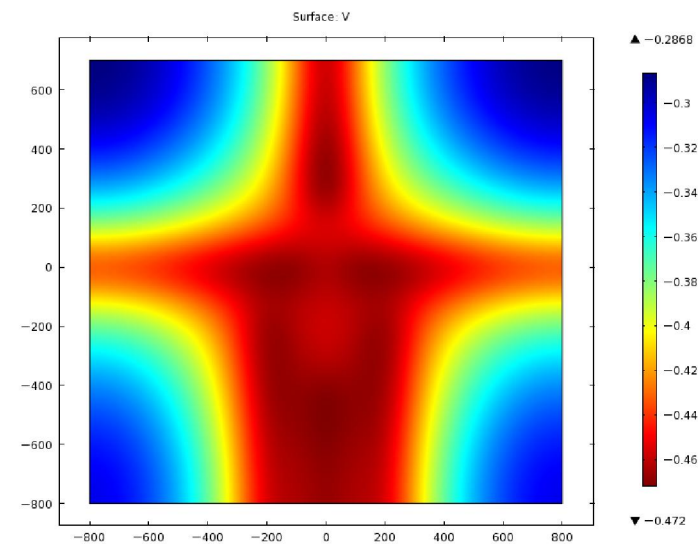
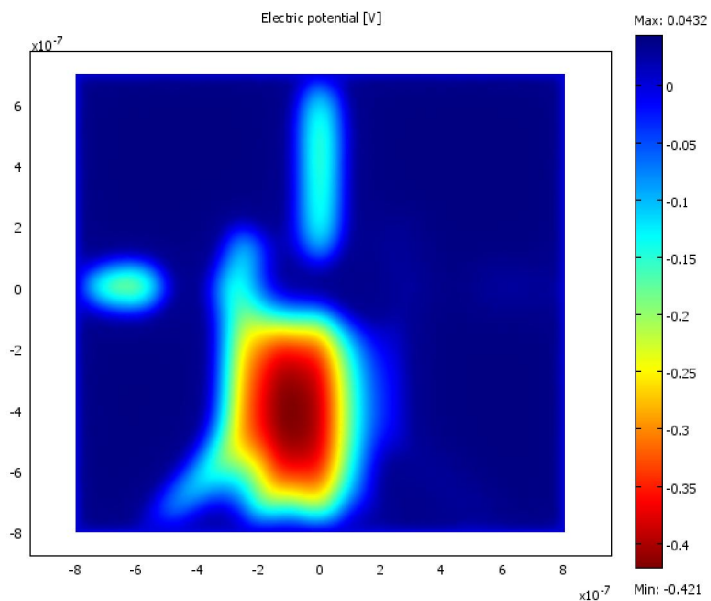
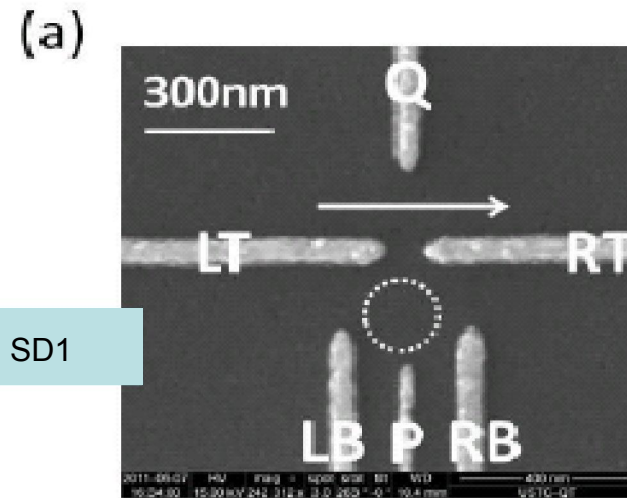
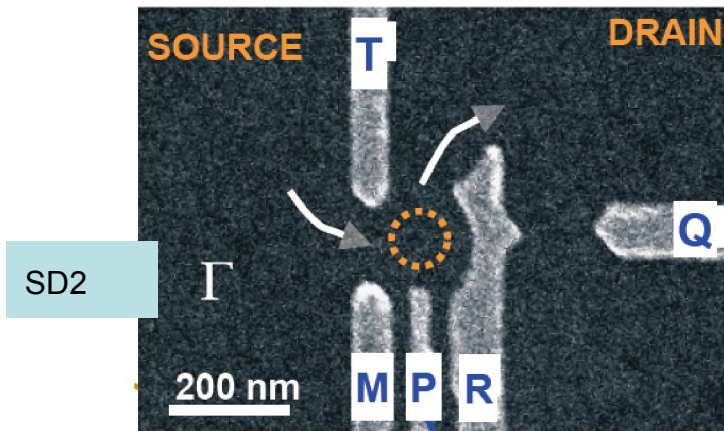
#34: $n=2.1 \times 10^{11} \text{cm}^{-2}$ $\mu = 0.6 \times 10^5 \text{cm}^2 \text{V}^{-1} \text{S}^{-1}$

#29: $n=1.3 \times 10^{11} \text{cm}^{-2}$ $\mu = 0.2 \times 10^5 \text{cm}^2 \text{V}^{-1} \text{S}^{-1}$

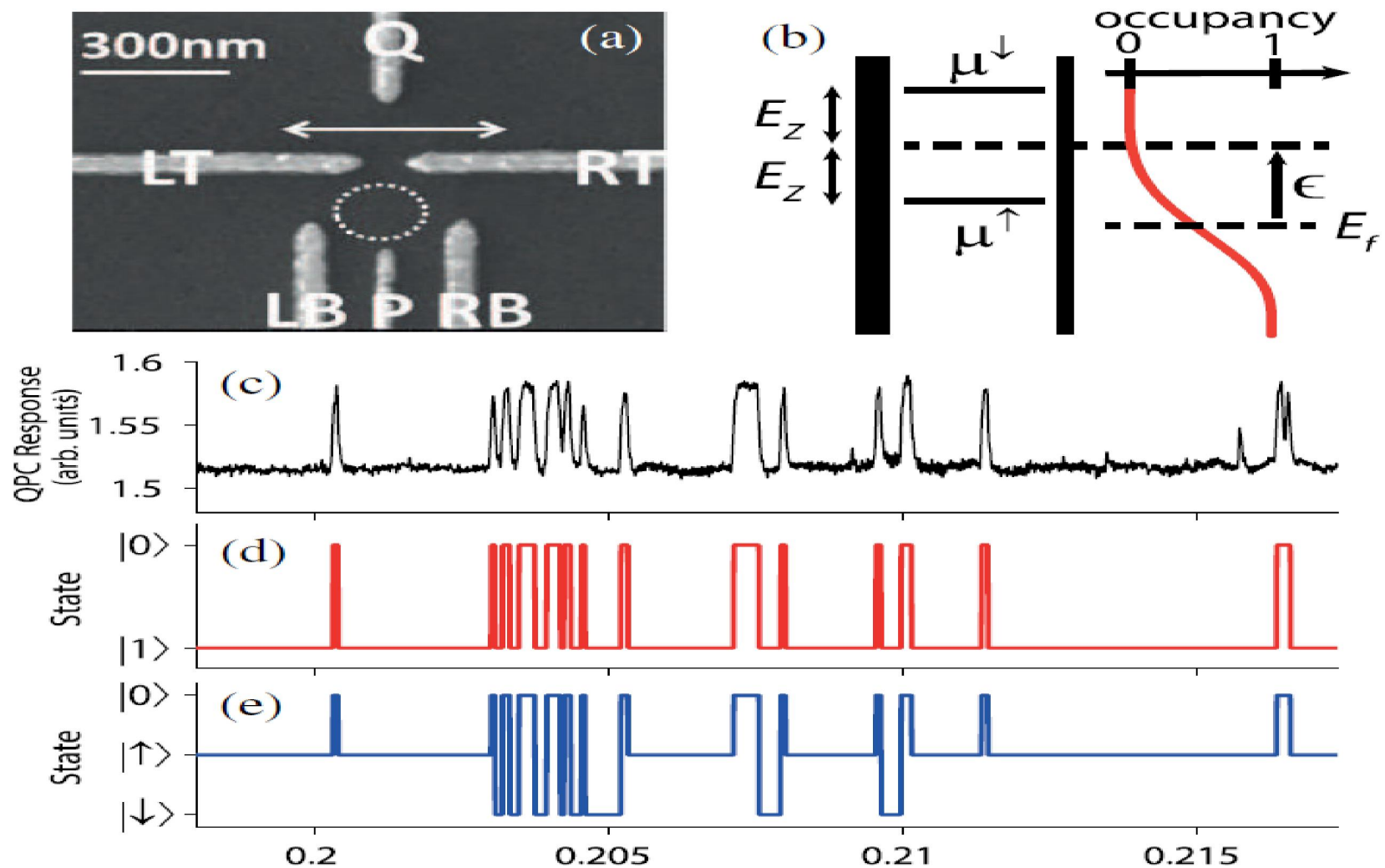


迁移率和电子气浓度对做量子点影响不大，但影响加工工艺和量子点尺寸以及欧姆接触的电导。

半导体栅型量子点的设计和制备



半导体栅型量子点的设计和制备



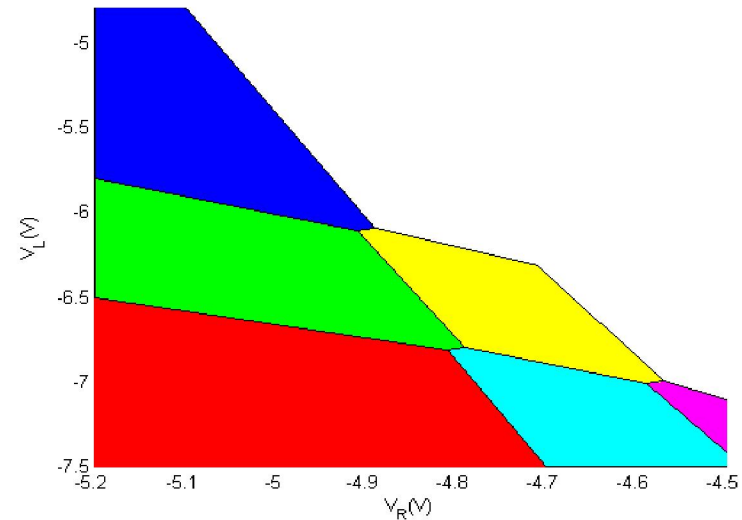
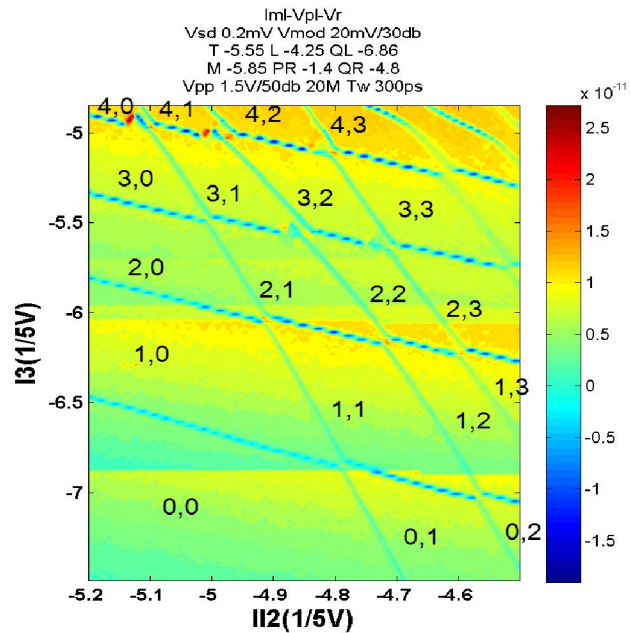
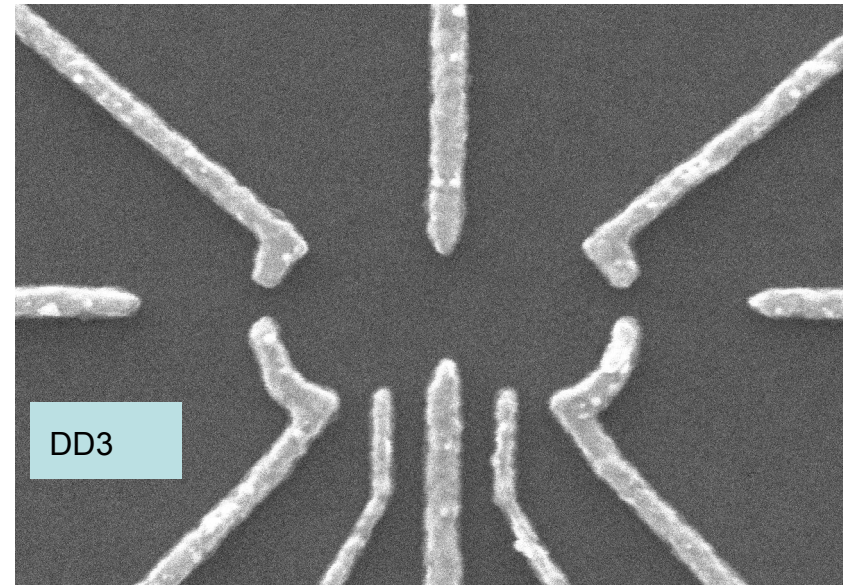
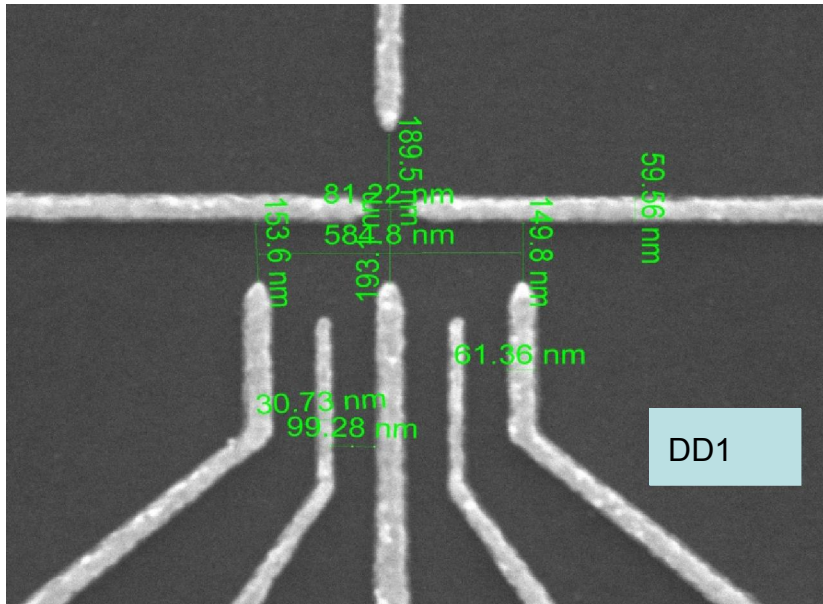
PRL 111, 126803 (2013)

PHYSICAL REVIEW LETTERS

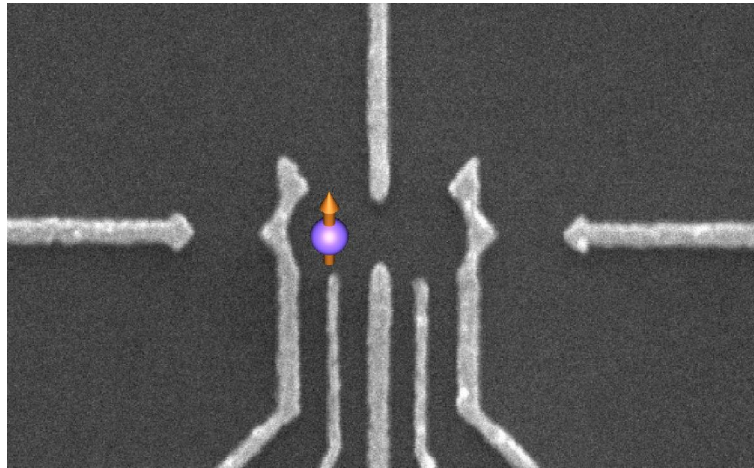
week ending
20 SEPTEMBER 2013

Detection and Measurement of Spin-Dependent Dynamics in Random Telegraph Signals

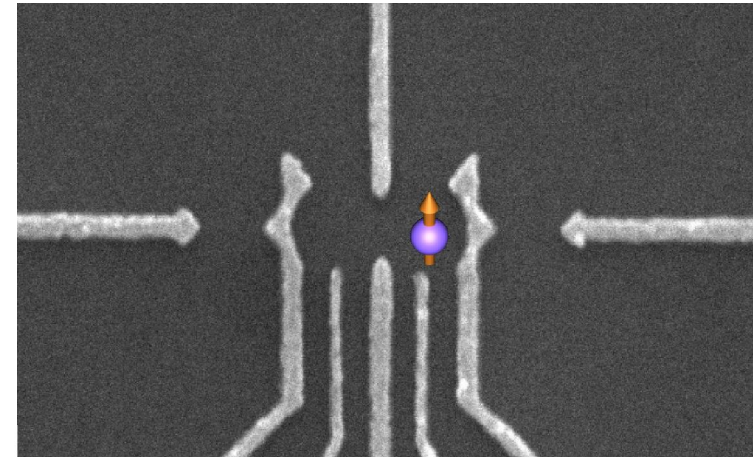
半导体栅型量子点的设计和制备



双量子点单电荷位置编码



$|L\rangle$

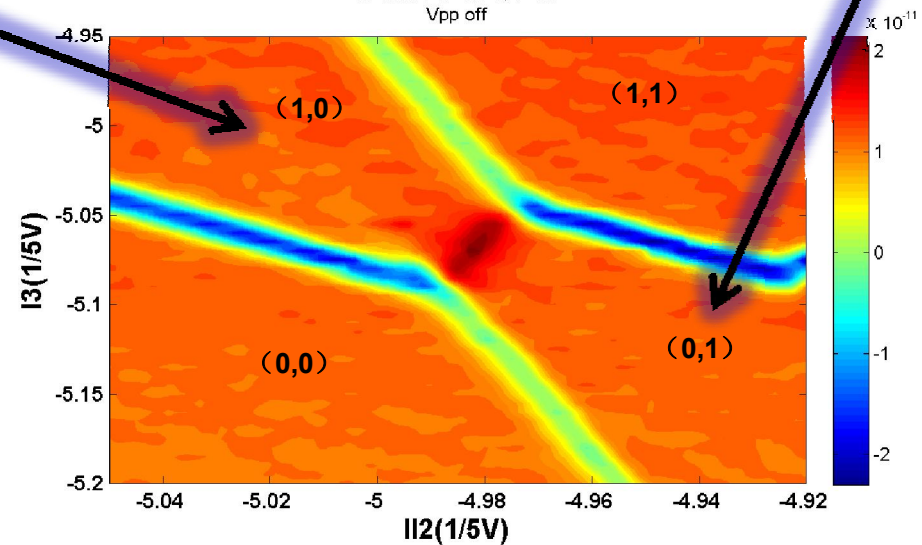


$|R\rangle$

the basis states $|L\rangle = (1, 0)$ and $|R\rangle = (0, 1)$

电荷量子比特的哈密顿量

$$H = \frac{1}{2} \epsilon \sigma_z + \Delta \sigma_x$$



超快单比特普适量子逻辑门

Ultrafast Universal Quantum Control of Quantum Dot Charge Qubit Using Landau-Zener-Stückelberg Interference

Gang Cao¹, HaiOu Li¹, Tao Tu¹, Li Wang¹, Chen Zhou¹, Ming Xiao¹, HongWei

Yan¹, GuoPing Guo¹, GuoPing Guo^{1*}

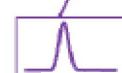
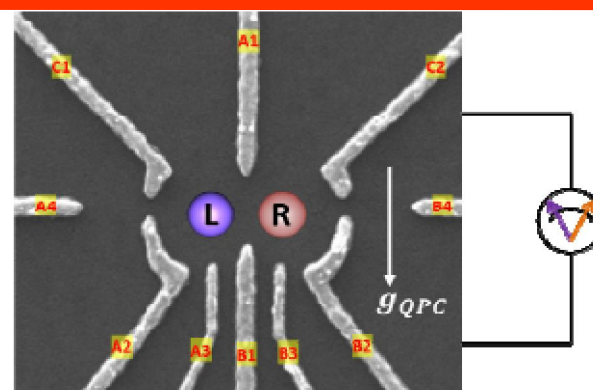
¹ Key Laboratory of Quantum Information and Quantum Optics, Chinese Academy of Sciences, University of Science and Technology of China, Hefei, Anhui, People's Republic of China
² Department of Physics, University of California at Los Angeles, 405 Hilgard Avenue, Los Angeles, California 90095, USA

利用LZ干涉实现超快普适的量子逻辑门操作

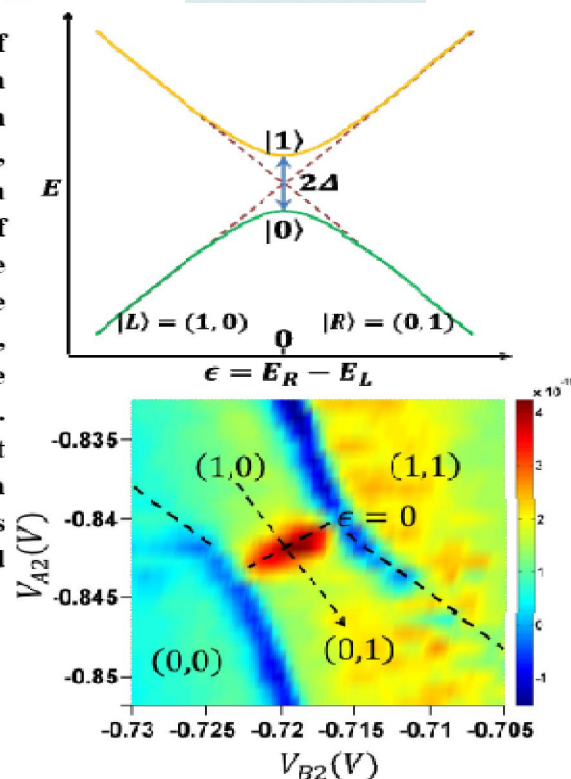
A basic requirement for quantum information processing is the ability of universal control the state of a single qubit¹⁻³ on timescales much shorter than the coherence time. Although ultrafast optical control of a single exciton (electron-hole pair)⁴ or a single spin⁵⁻⁷ in quantum dots have been achieved, scaling up such methods remains major challenging. Here we demonstrate a complete control of the quantum dot charge qubit with the time scales of picosecond, orders of magnitude faster than those previously measured in the charge^{8,9} or spin based qubits¹⁰⁻¹² also with electric fields in quantum dots. We observed tunable qubit dynamics in charge stability diagram, in time domain, and in pulse amplitude space of the driven pulse. Our analysis indicated that the results are well described by Landau-Zener-Stückelberg (LZS) interference¹³. These results establish the feasibility of a full set of all-electrical single qubit operations and enable potential multiple-qubit implementations^{14,15}. Although our experiment is carried out in a solid-state architecture, the technique is independent of the particular physical encoding of quantum information, and has the potential for wider application.

通过设计合适的高频脉冲序列，实现了电荷单量子比特超快量子逻辑门

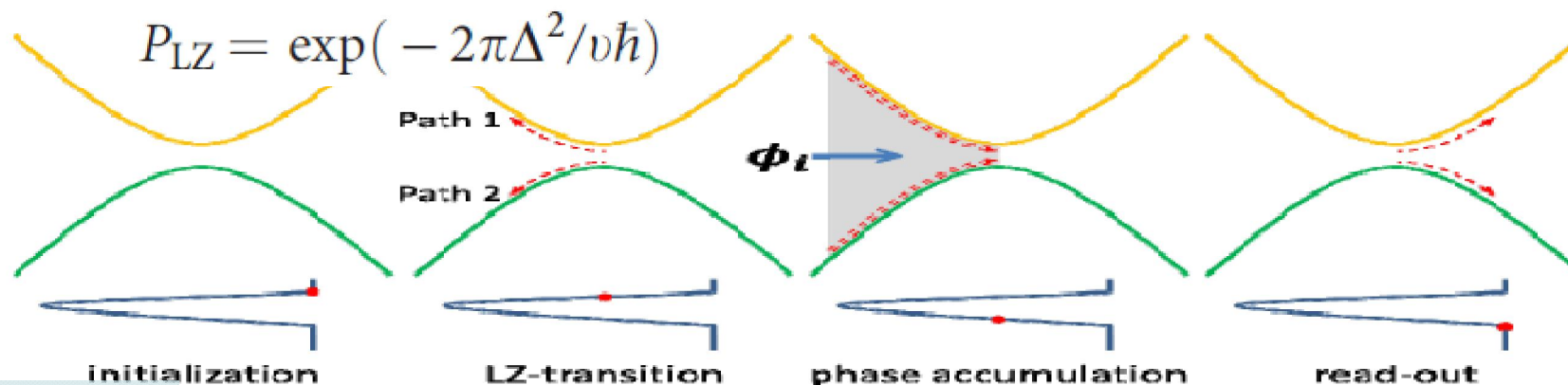
Nature Communications 4, 1401(2013)



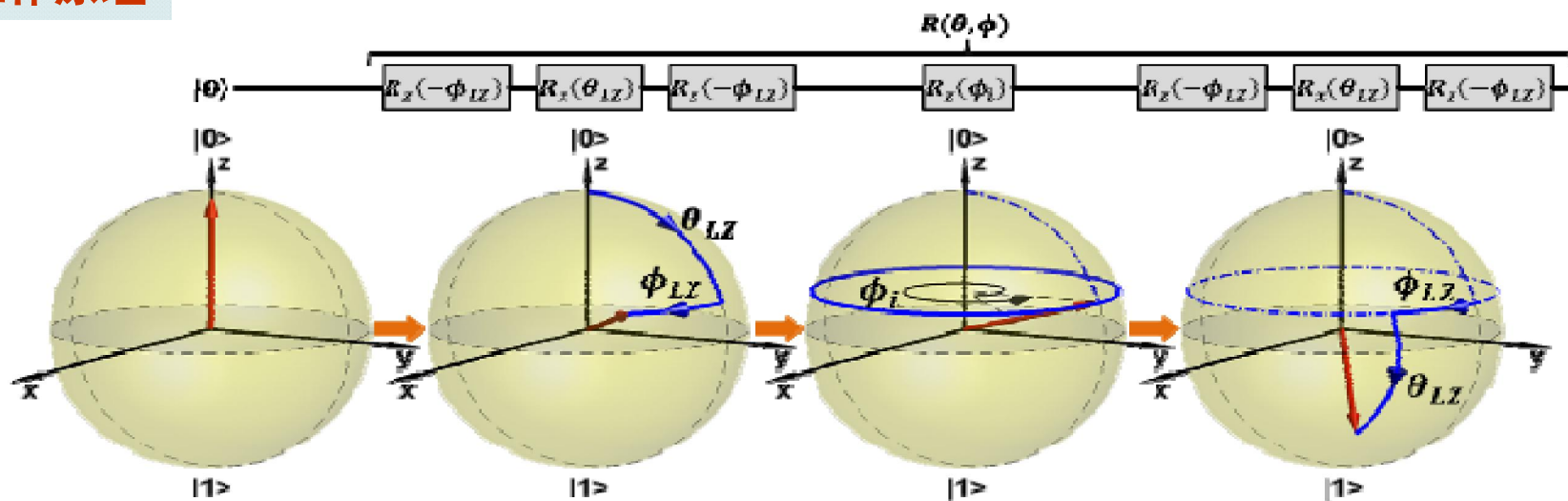
工作原理



超快单比特普适逻辑门操控



工作原理

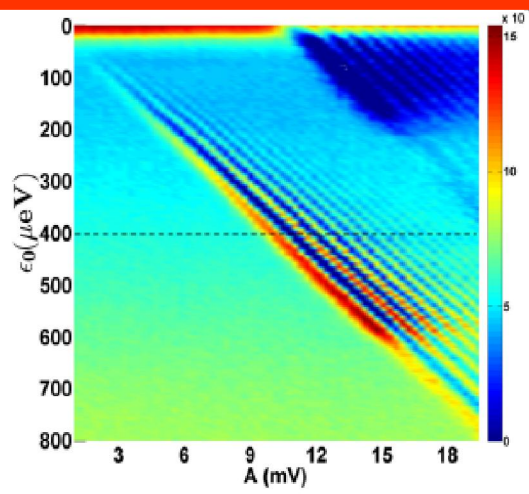


$$|\Psi_{out}\rangle = R_z(-\phi_{LZ})R_x(\theta_{LZ})R_z(-\phi_{LZ})R_z(\phi_i)R_z(-\phi_{LZ})$$

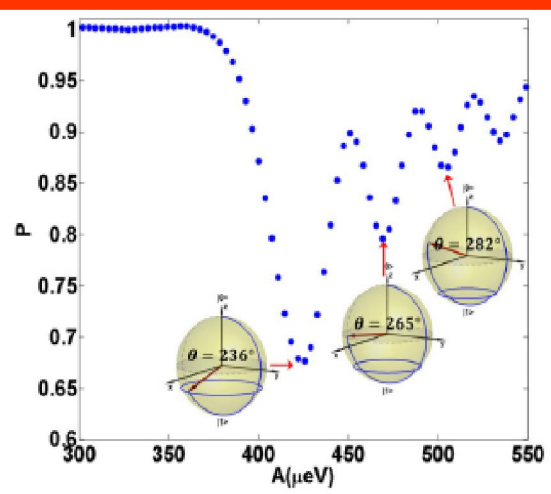
$$\times R_x(\theta_{LZ})R_z(-\phi_{LZ})|\Psi_{in}\rangle = R(\theta, \phi)|\Psi_{in}\rangle$$

◆控制脉冲高度和宽度在皮秒量级实现任意单量子比特逻辑门, 比自旋量子比特快了2个量级

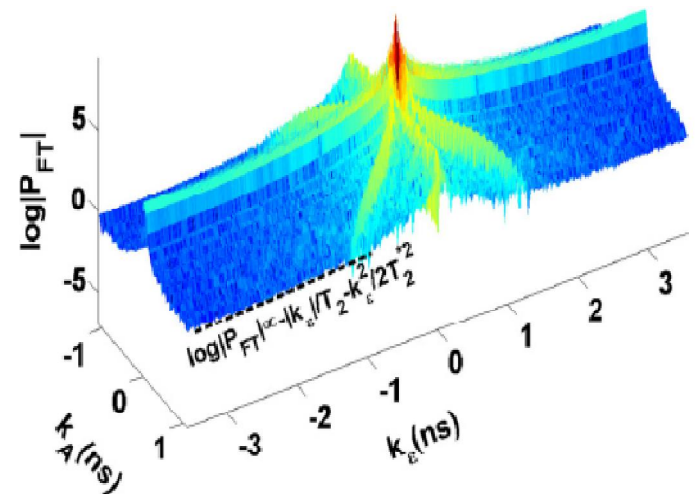
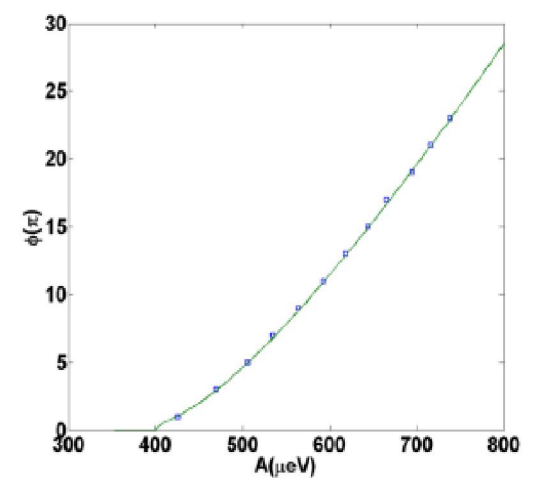
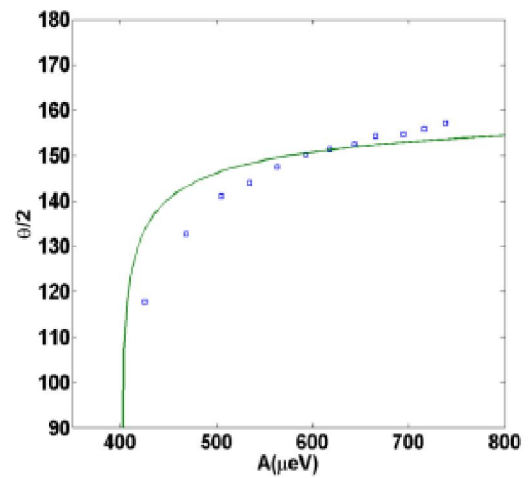
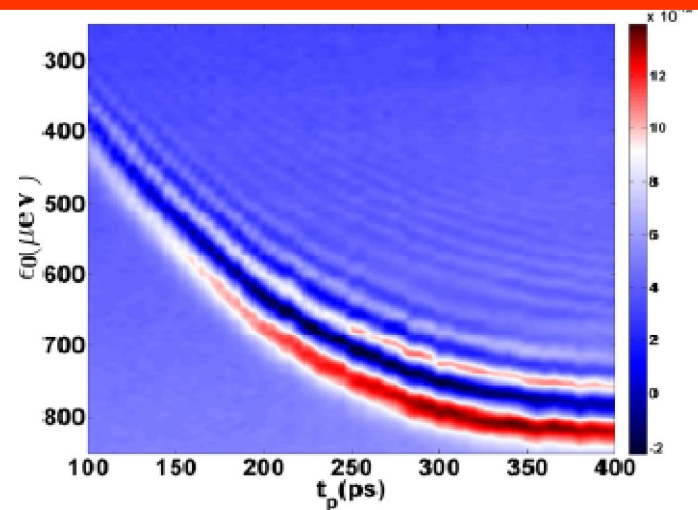
超快单比特普适逻辑门操控



(a)



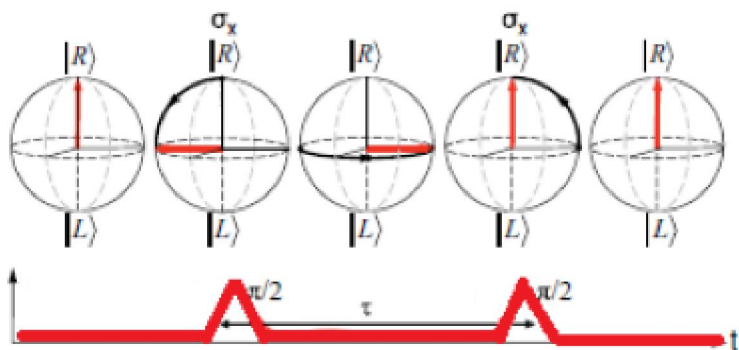
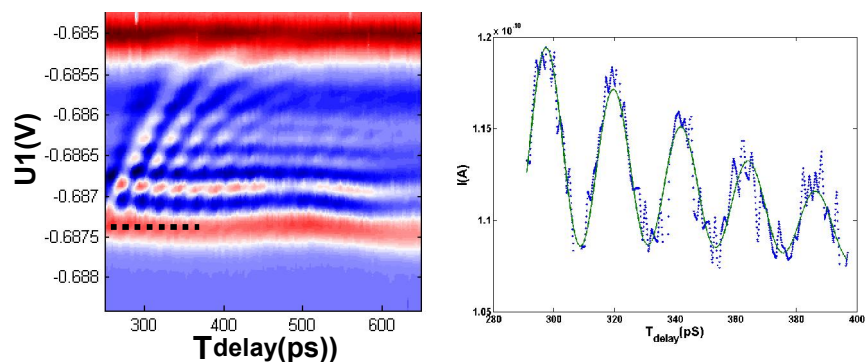
(b)



- ◆ 实现并观察到了单个电子的可控隧穿和路径干涉.
- ◆ 结合频域和时间域测量了 $T_2^* \sim 3\text{ns}$

高速脉冲序列调控量子态

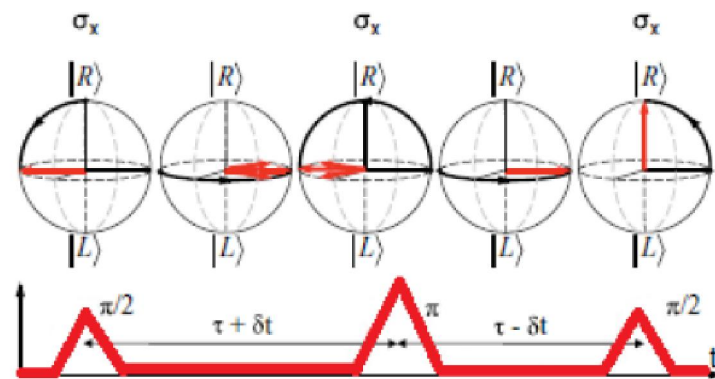
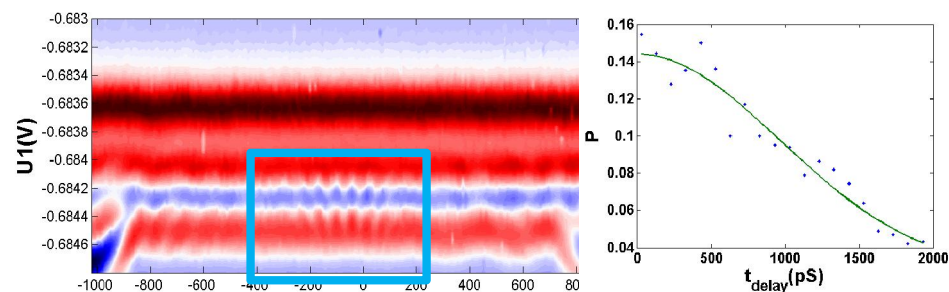
单量子比特操控



$$\Delta P_R = A_0 \times \exp \left[- \left(\frac{\tau - \tau_0}{T_2^*} \right)^2 \right] \times \cos [2\pi f (\tau - \tau_0)].$$

$$T_2^* \sim 230 \text{ ps}$$

单量子比特 Ramsey 干涉



$$P = y_0 + A \exp(- (t/T_2)^2)$$

$$T_2 \sim 1360 \text{ ps}$$

单量子比特 Echo 操控

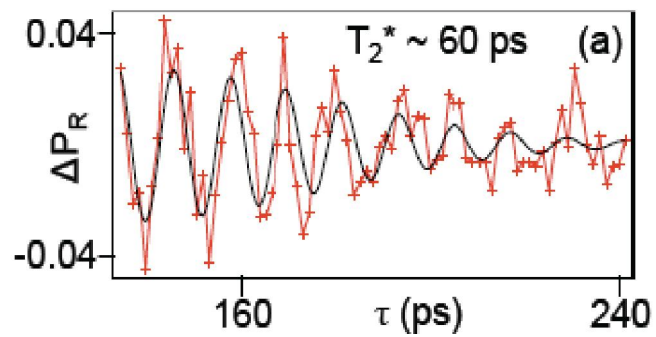
高速脉冲序列调控量子态

比较 Ramsey

Princeton, Petta组

掺杂GaAs

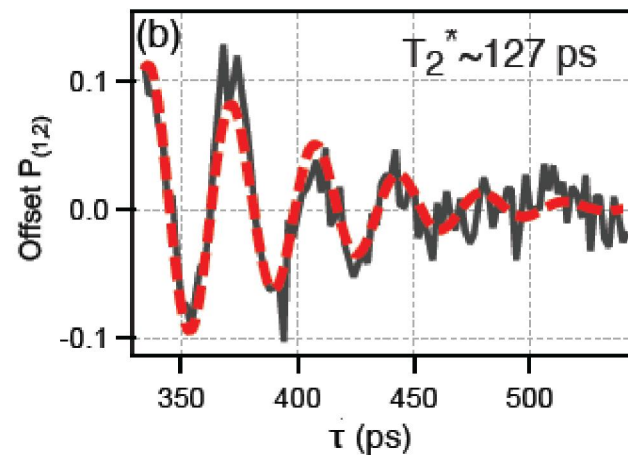
$T_2^* \sim 60 \text{ ps}$



Wisconsin, Eriksson组

掺杂GeSi

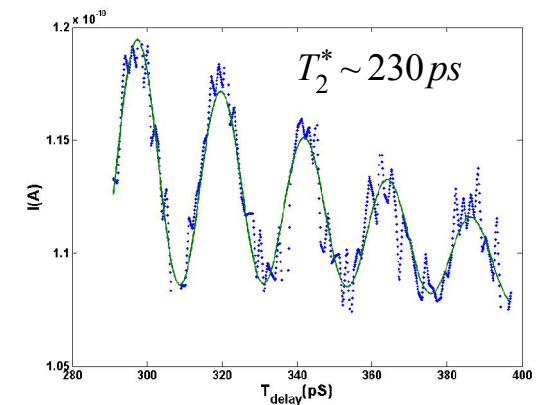
$T_2^* \sim 127 \text{ ps}$



USTC

掺杂GaAs

$T_2^* \sim 230 \text{ ps}$



高速脉冲序列操控量子态

比较 Echo

Princeton, Petta组

掺杂GaAs

没有成功

Wisconsin, Eriksson组

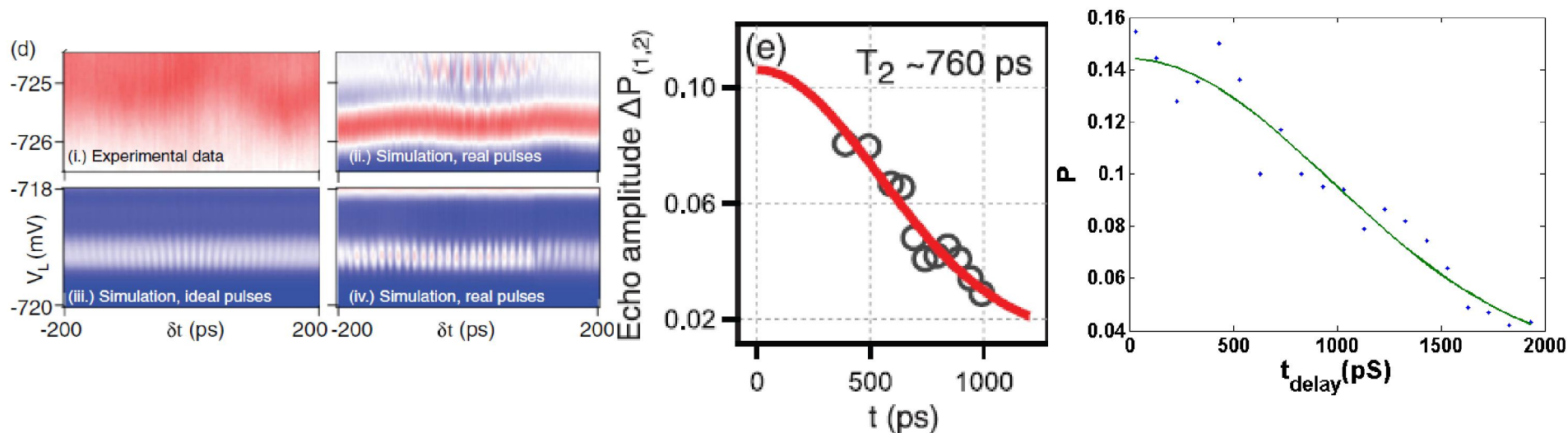
掺杂GeSi

$T_2 \sim 760$ ps

USTC

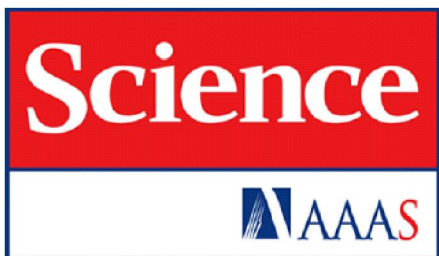
掺杂GaAs

$T_2 \sim 1360$ ps



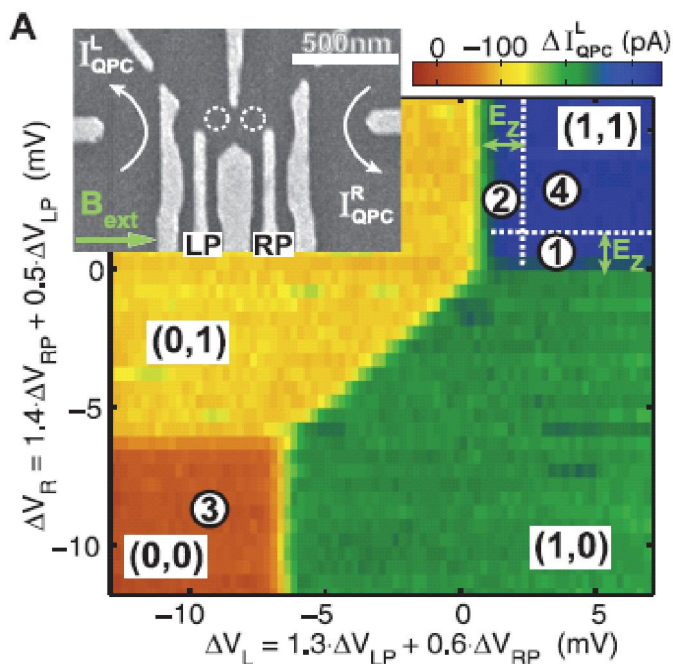
两比特量子控制非门

自旋编码双量子比特

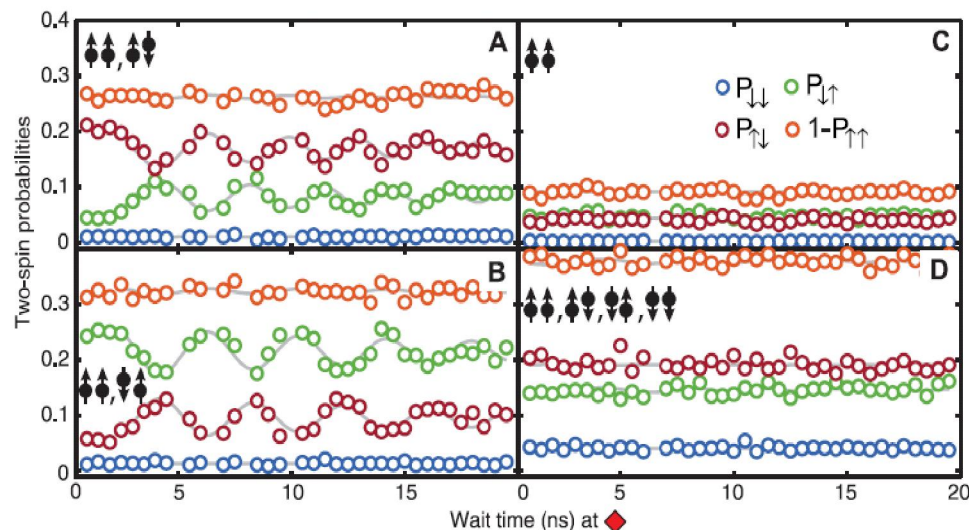
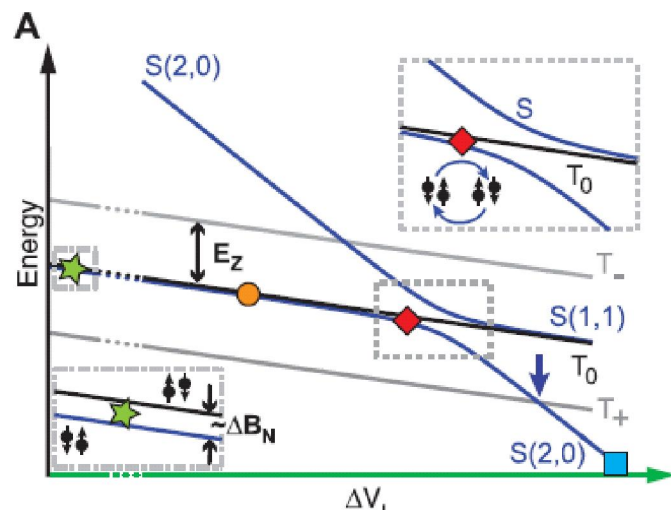


Single-Shot Correlations and Two-Qubit Gate of Solid-State Spins

K. C. Nowack et al.
Science **333**, 1269 (2011);
 DOI: 10.1126/science.1209524



单电子自旋编码



两比特量子控制非门

自旋编码双量子比特

双量子点双电子
自旋单三态编码



Demonstration of Entanglement of Electrostatically Coupled Singlet-Triplet Qubits

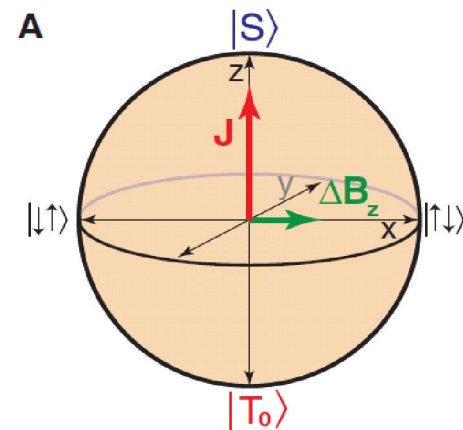
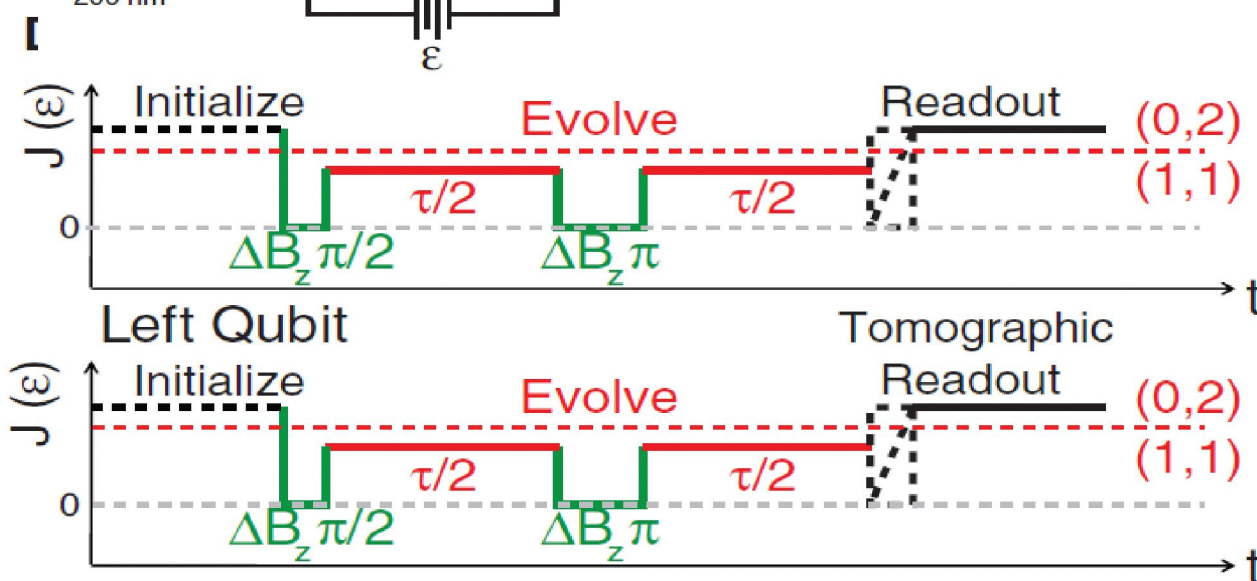
M. D. Shulman *et al.*

Science **336**, 202 (2012);

DOI: 10.1126/science.1217692



$$H_{2\text{-qubit}} = \frac{\hbar}{2} (J_1(\sigma_z \otimes I) + J_2(I \otimes \sigma_z) + \frac{J_{12}}{2} ((\sigma_z - I) \otimes (\sigma_z - I)) + \Delta B_{z,1}(\sigma_x \otimes I) + \Delta B_{z,2}(I \otimes \sigma_x))$$



两比特量子控制非门

自旋编码双量子比特



Demonstration of Entanglement of Electrostatically Coupled Singlet-Triplet Qubits

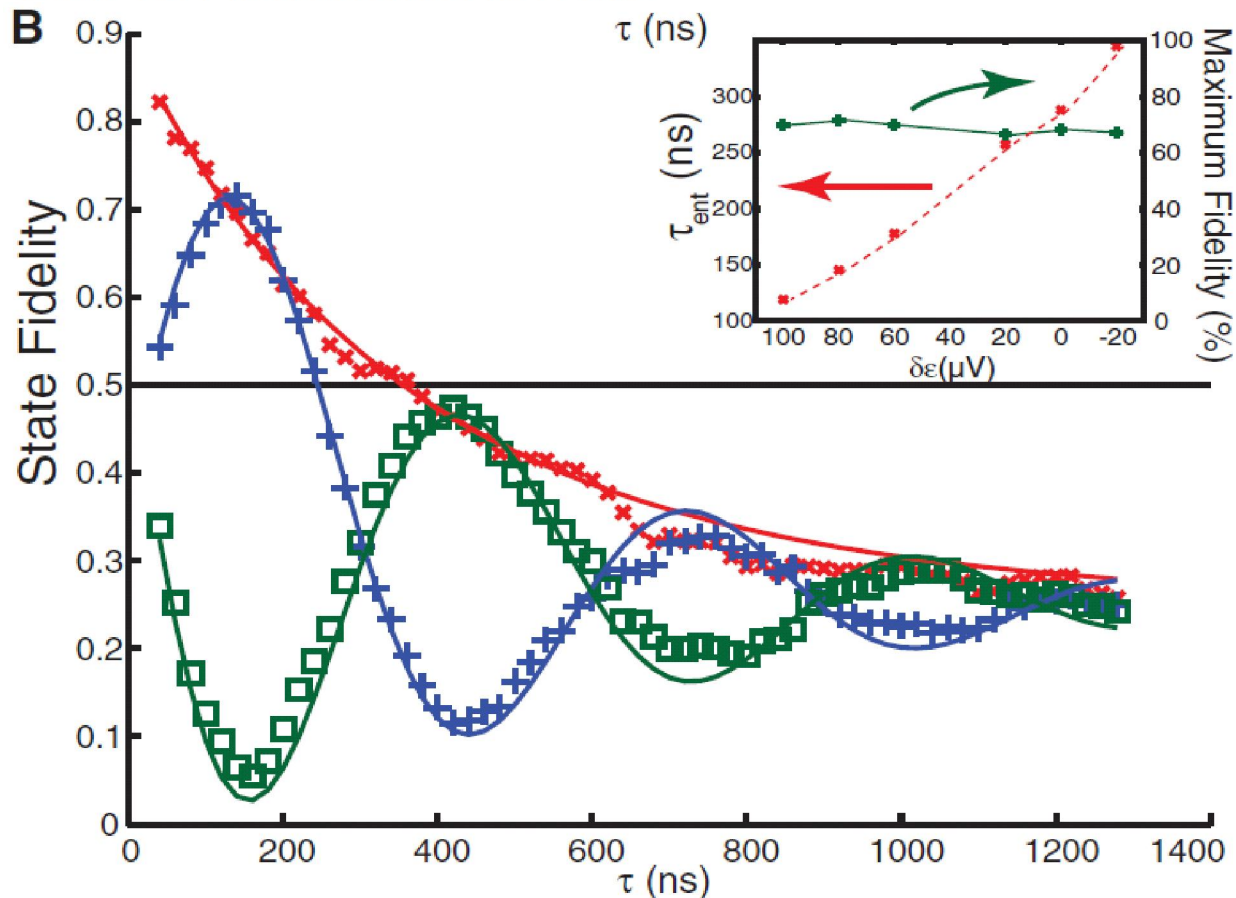
M. D. Shulman *et al.*

Science **336**, 202 (2012);

DOI: 10.1126/science.1217692

双量子点双电子
自旋单三态编码

操控时间长



两比特量子控制非门

ARTICLE

Received 24 Feb 2015 | Accepted 1 Jun 2015 | Published 17 Jul 2015

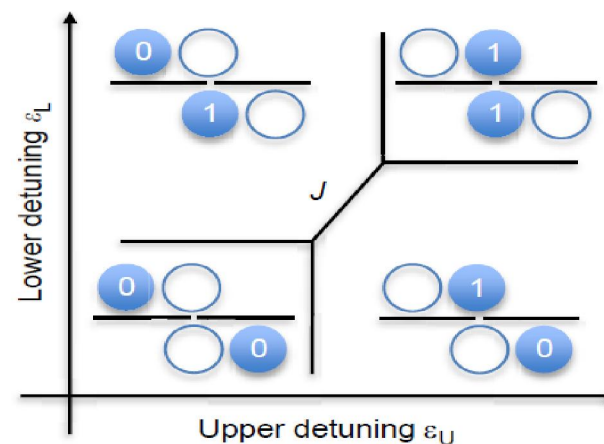
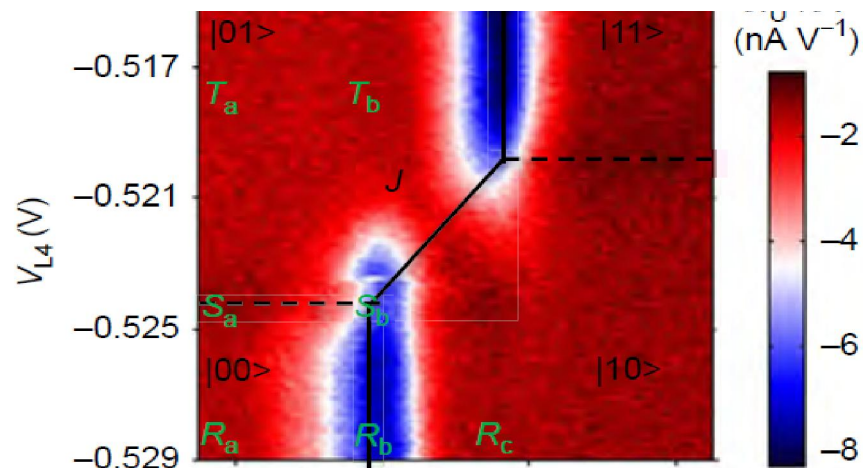
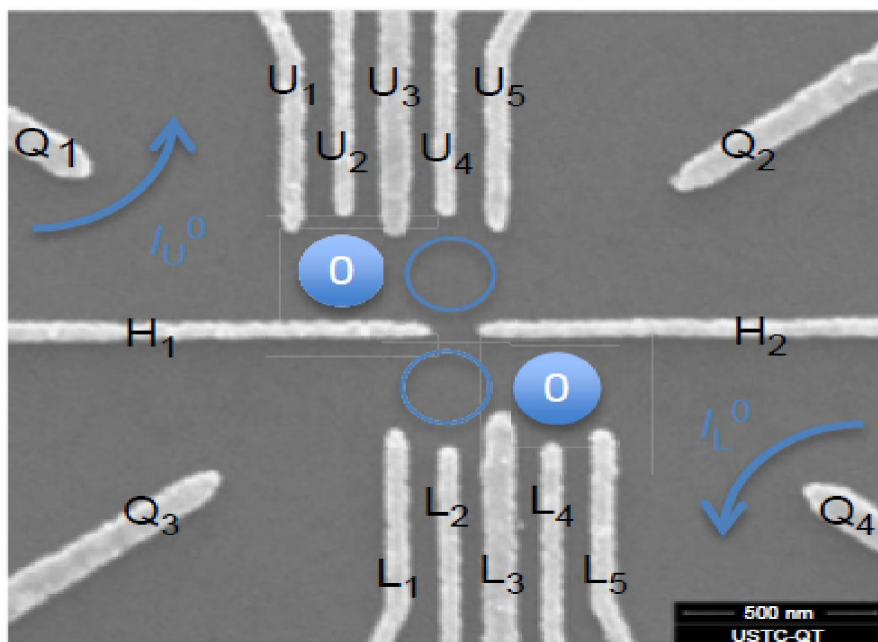
DOI: 10.1038/ncomms8681

OPEN

Conditional rotation of two strongly coupled semiconductor charge qubits

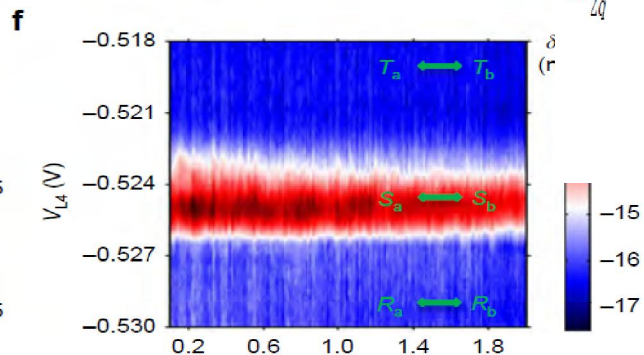
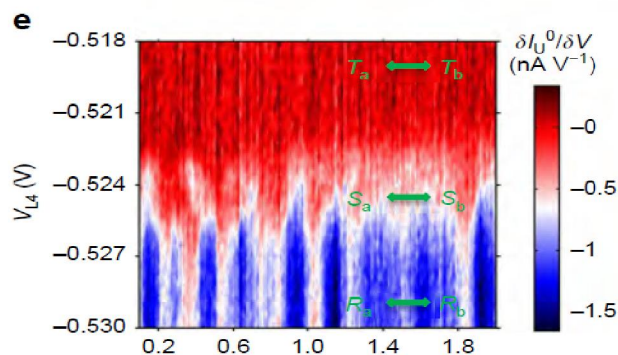
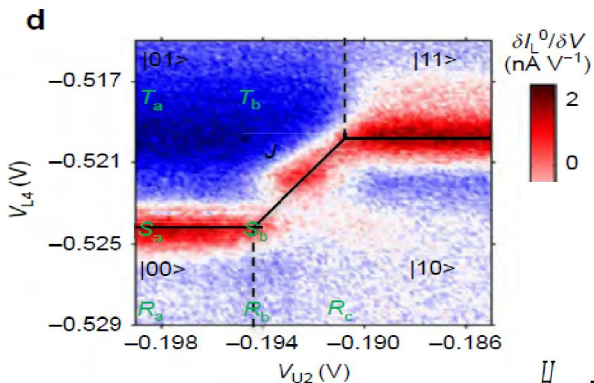
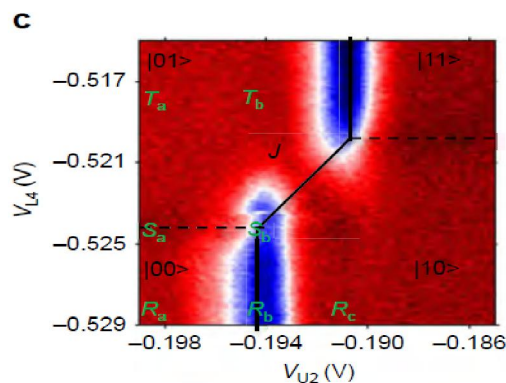
Hai-Ou Li^{1,2,*}, Gang Cao^{1,2,*}, Guo-Dong Yu^{2,3,*}, Ming Xiao^{2,3}, Guang-Hong Wen^{1,2}, Hong-Wen Jiang⁴ & Guo-Ping Guo^{1,2}

双量子点单电子
状态编码

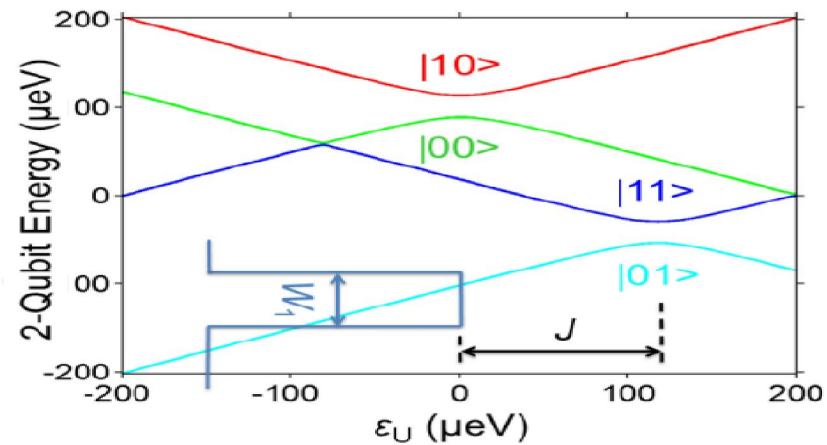
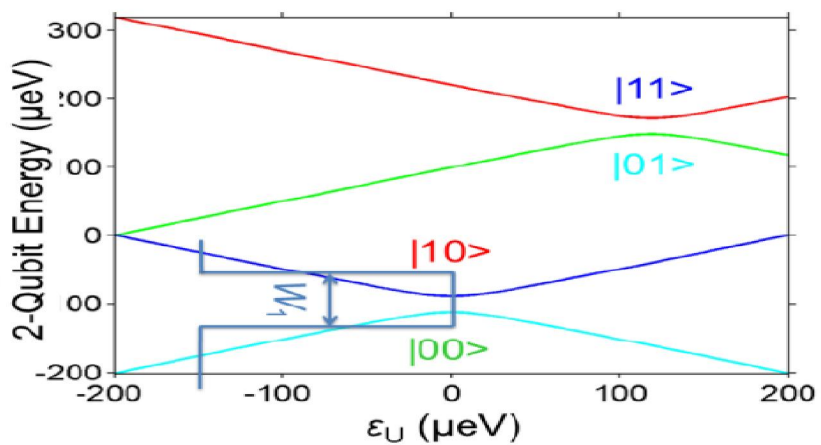


$$H_{2q} = \frac{1}{2} \sum_{i=1}^2 (\varepsilon_i \sigma_z^{(i)} + \Delta_i \sigma_x^{(i)}) + \frac{J}{4} (I - \sigma_z^{(1)}) \otimes (I - \sigma_z^{(2)})$$

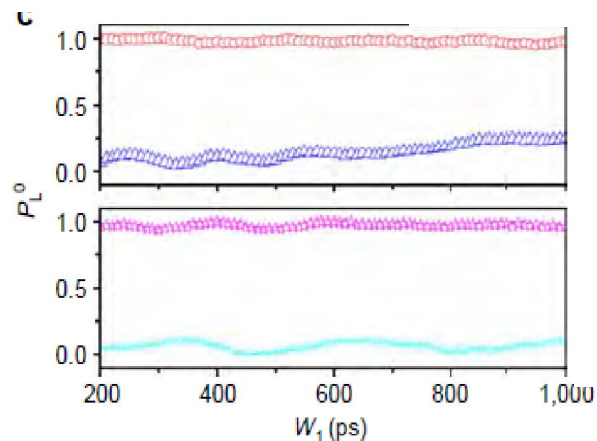
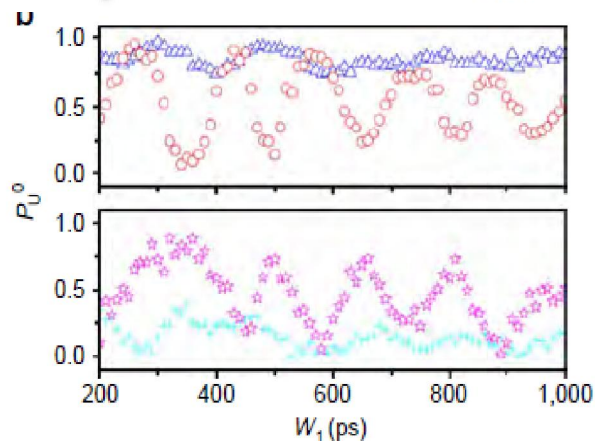
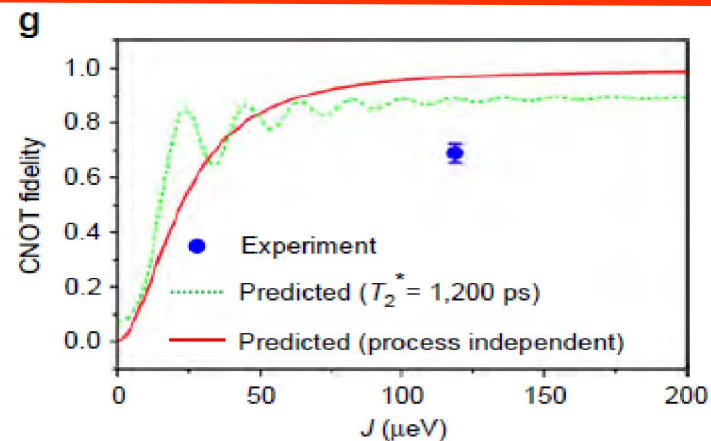
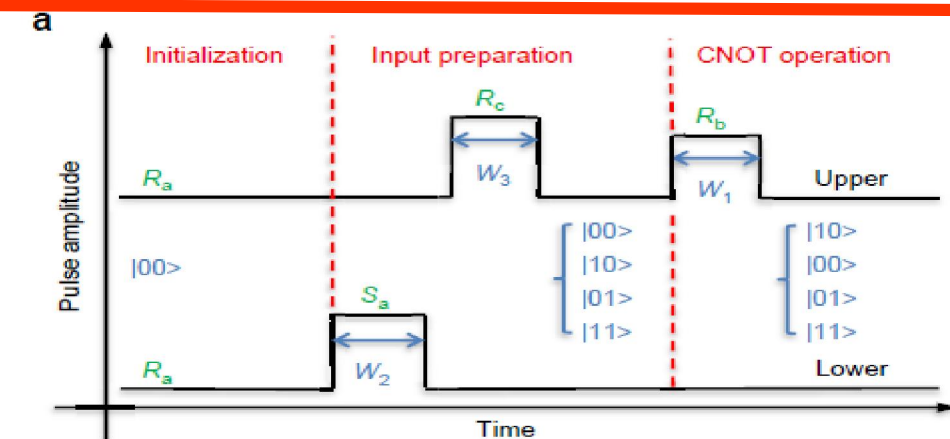
两比特量子控制非门



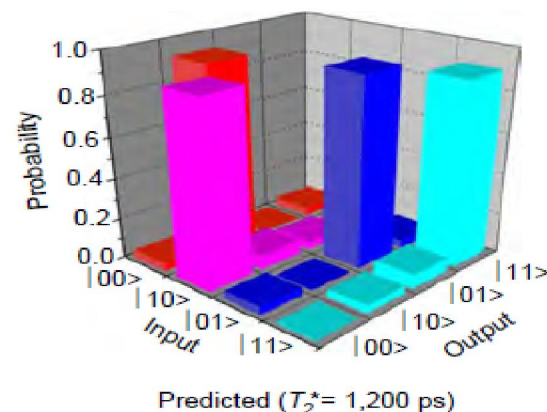
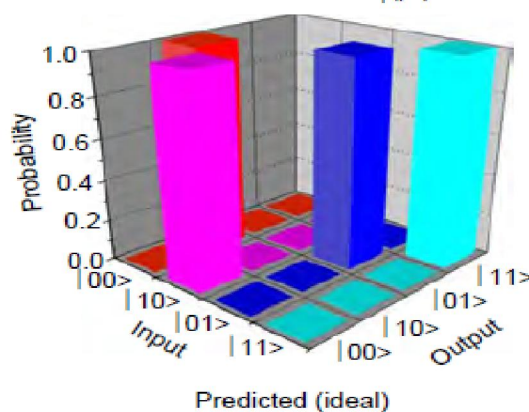
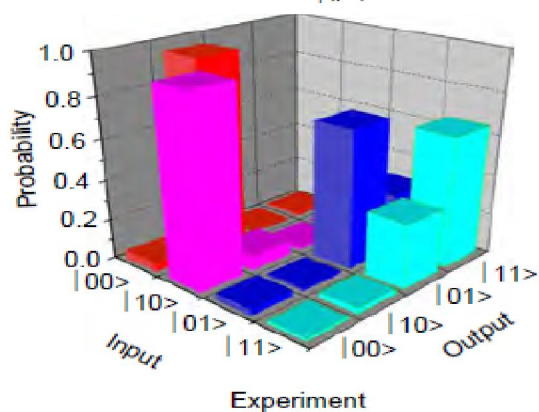
$$H_{2q} = \frac{1}{2} \begin{vmatrix} \epsilon_U + \epsilon_L & \Delta_U & \Delta_L & 0 \\ \Delta_U & -\epsilon_U + \epsilon_L & 0 & \Delta_L \\ \Delta_L & 0 & \epsilon_U - \epsilon_L & \Delta_U \\ 0 & \Delta_L & \Delta_U & -\epsilon_U - \epsilon_L + 2J \end{vmatrix}$$



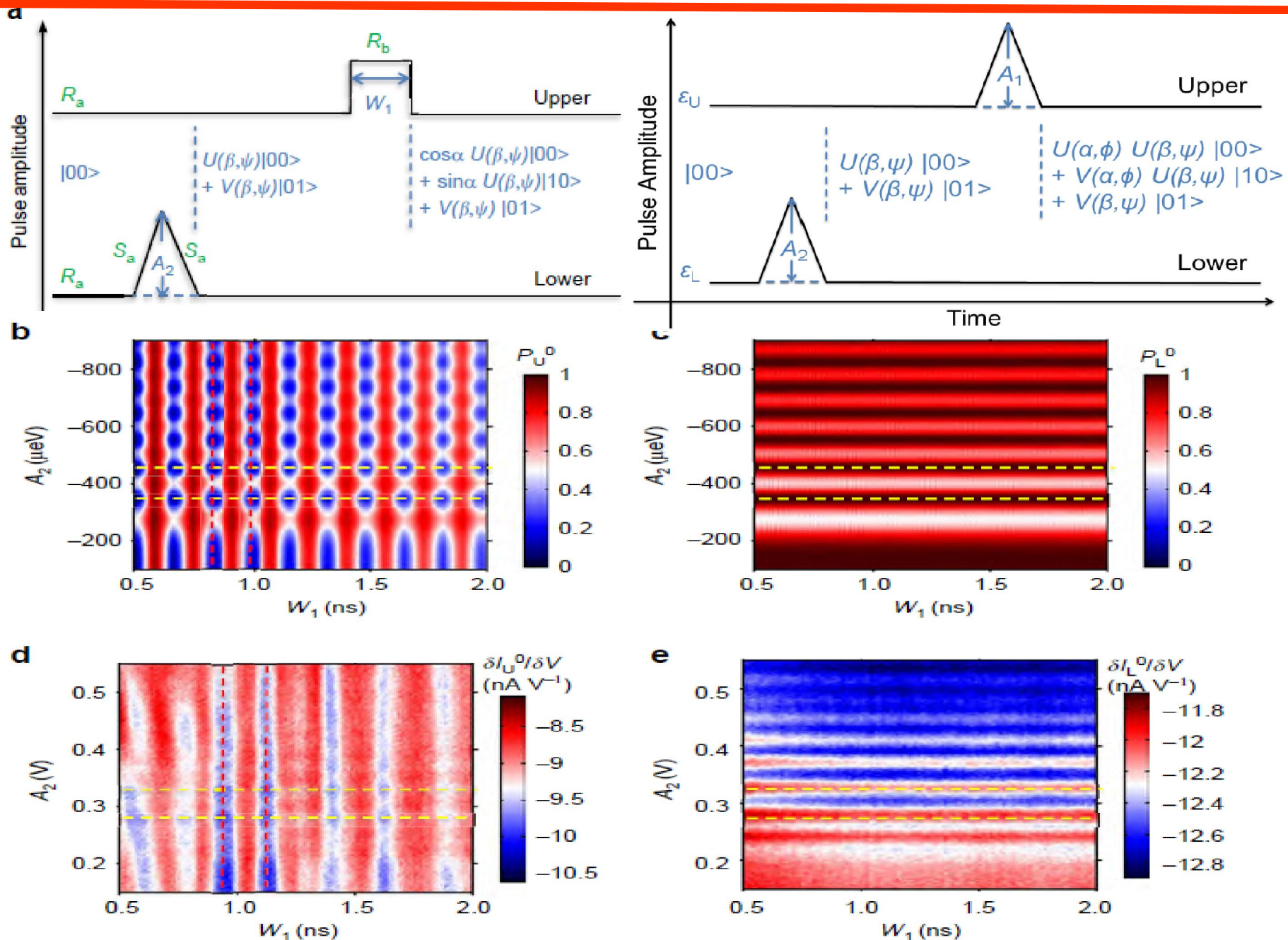
两比特量子控制非门



$$T_{\text{measured}} = \begin{pmatrix} 0.09 & 0.89 & 0.002 & 0.02 \\ 0.87 & 0.12 & 0.01 & 0.002 \\ 0.06 & 0.02 & 0.74 & 0.18 \\ 0.02 & 0.07 & 0.23 & 0.68 \end{pmatrix}$$



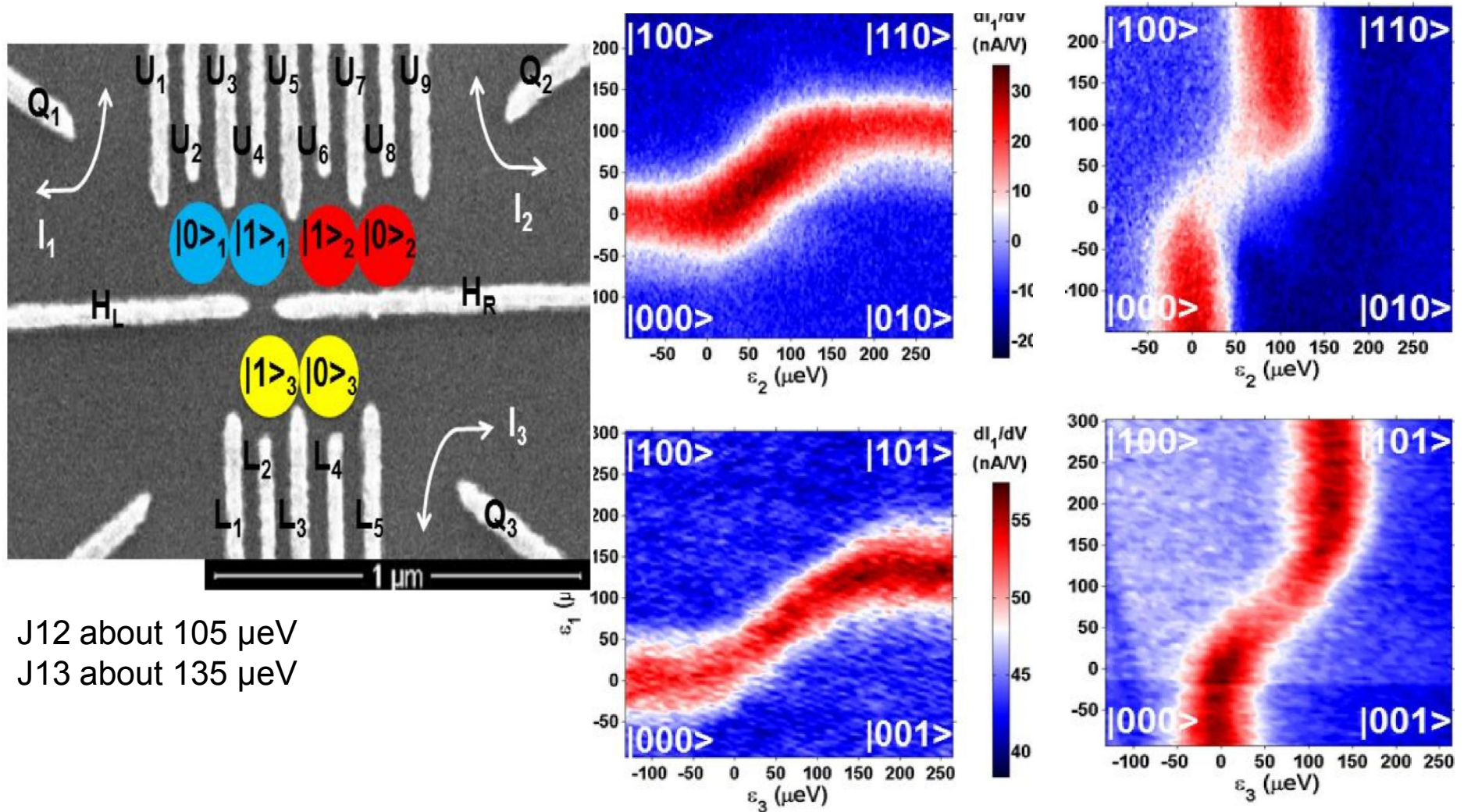
两比特量子控制非门



Basic Quantum gates of charge qubit

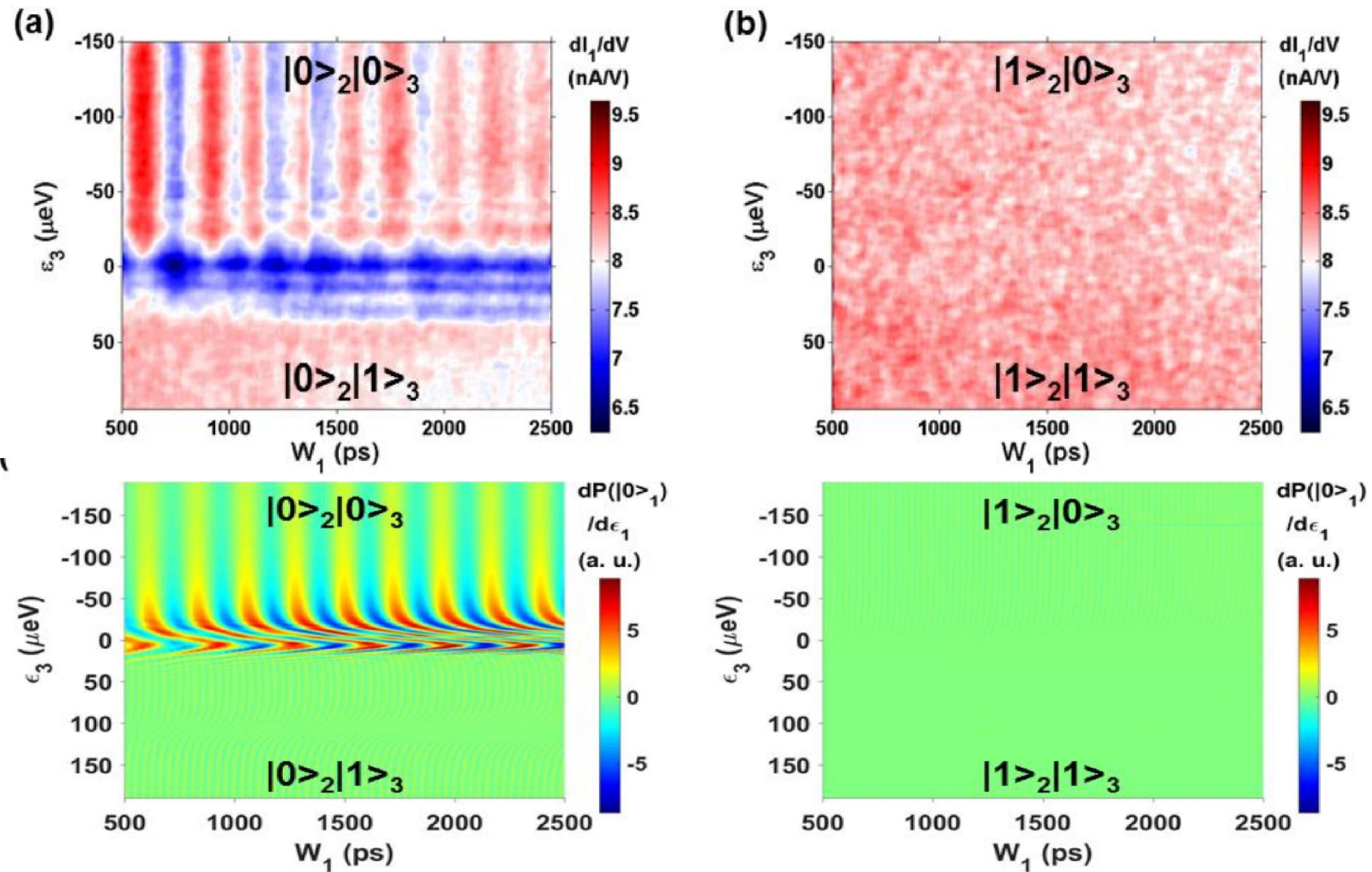
Controlled Quantum operations of a Semiconductor Three-Qubit System

arXiv 1610.06704



Basic Quantum gates of charge qubit

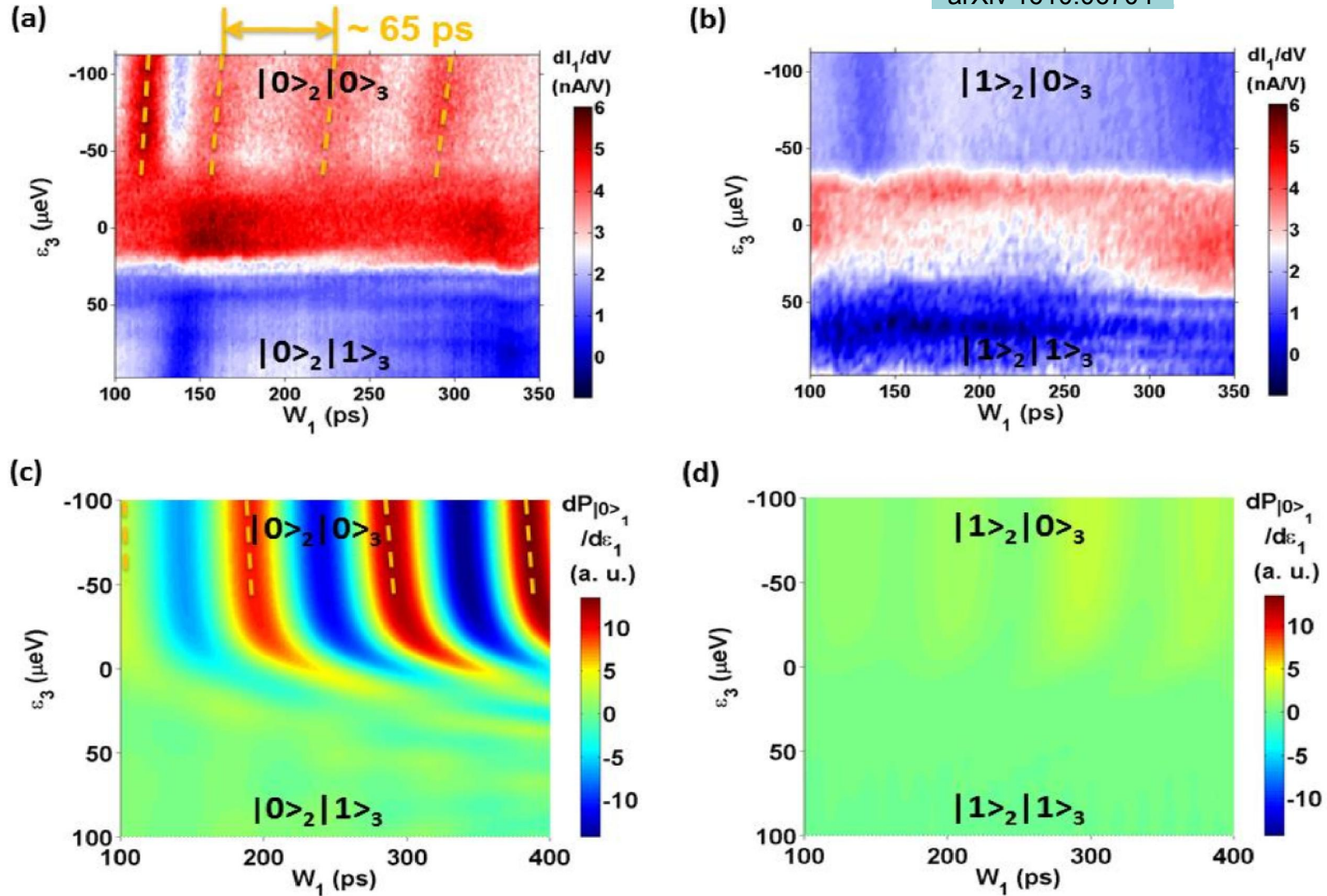
Controlled Quantum operations of a Semiconductor Three-Qubit System



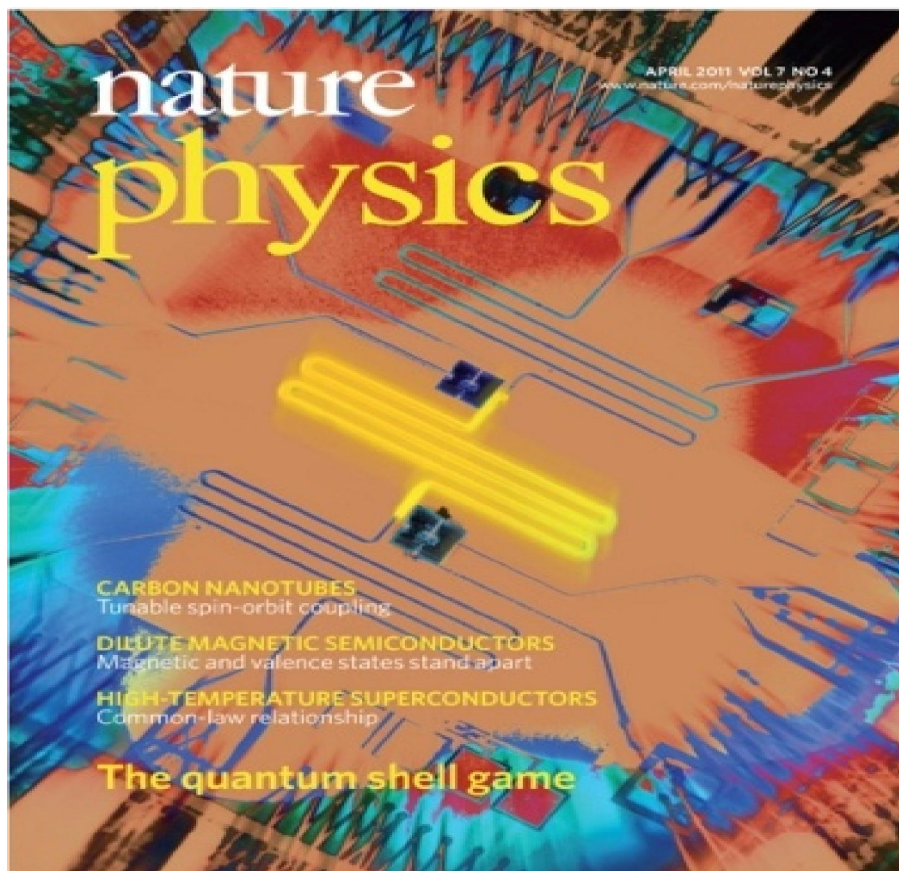
Basic Quantum gates of charge qubit

Controlled Quantum operations of a Semiconductor Three-Qubit System

arXiv 1610.06704

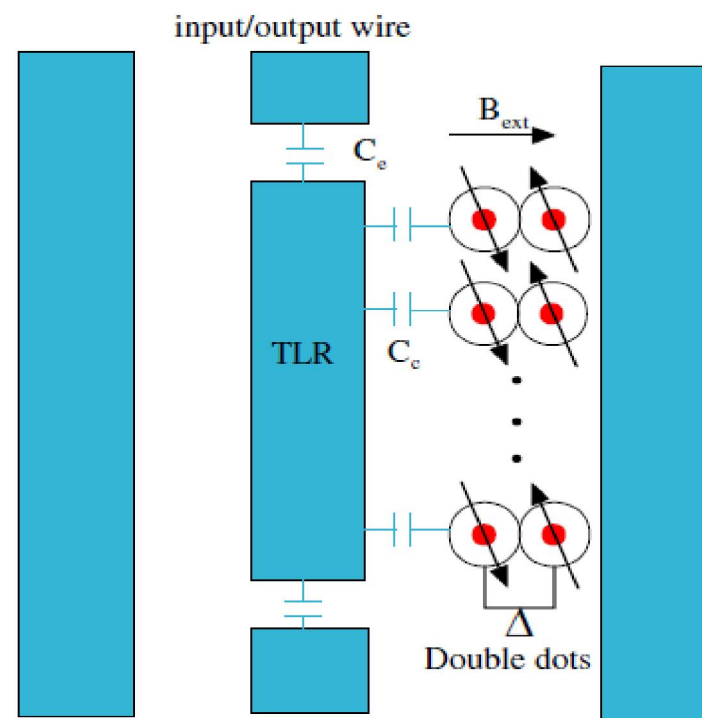
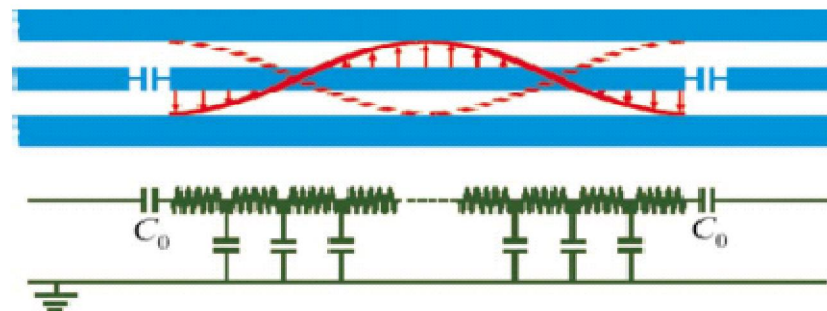


量子数据线和量子比特的长程耦合



J-C Model:

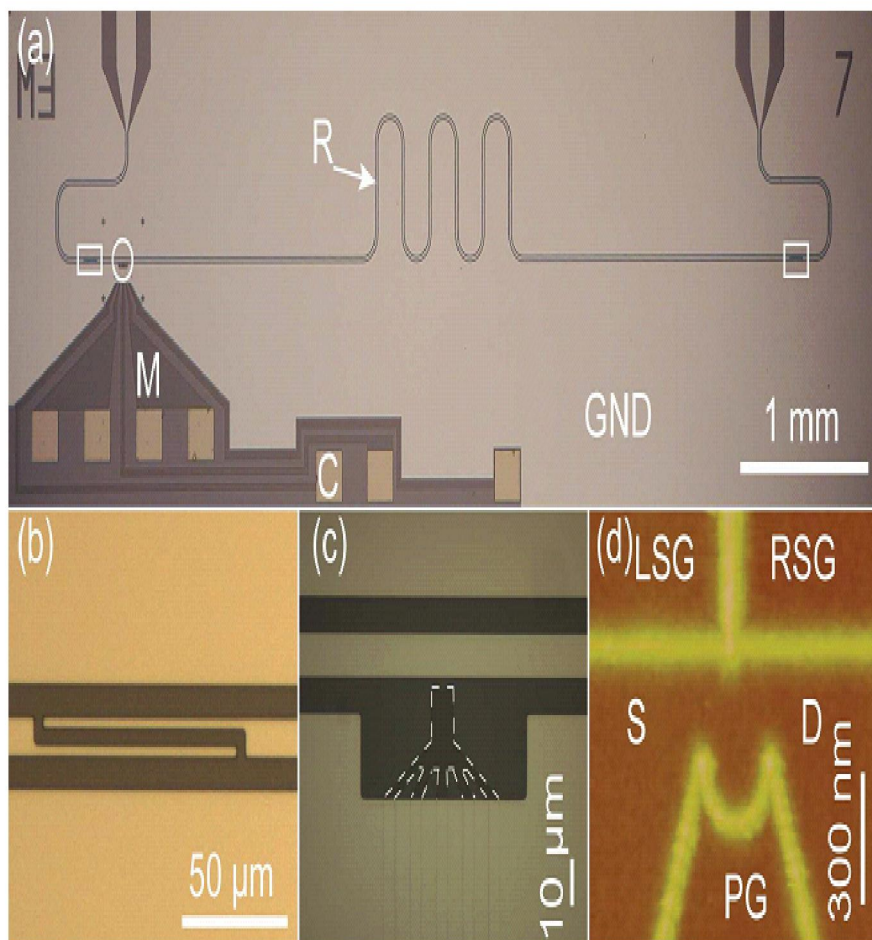
$$H = \hbar\omega_r(a^\dagger a + 1/2) + \hbar\frac{\omega_a}{2}\sigma_z + \hbar g(a^\dagger \sigma^- + a\sigma^+)$$



Phys. Rev. Lett. 101, 230501 (2008).
Phys. Rev. A 78, 020302(R) (2008).

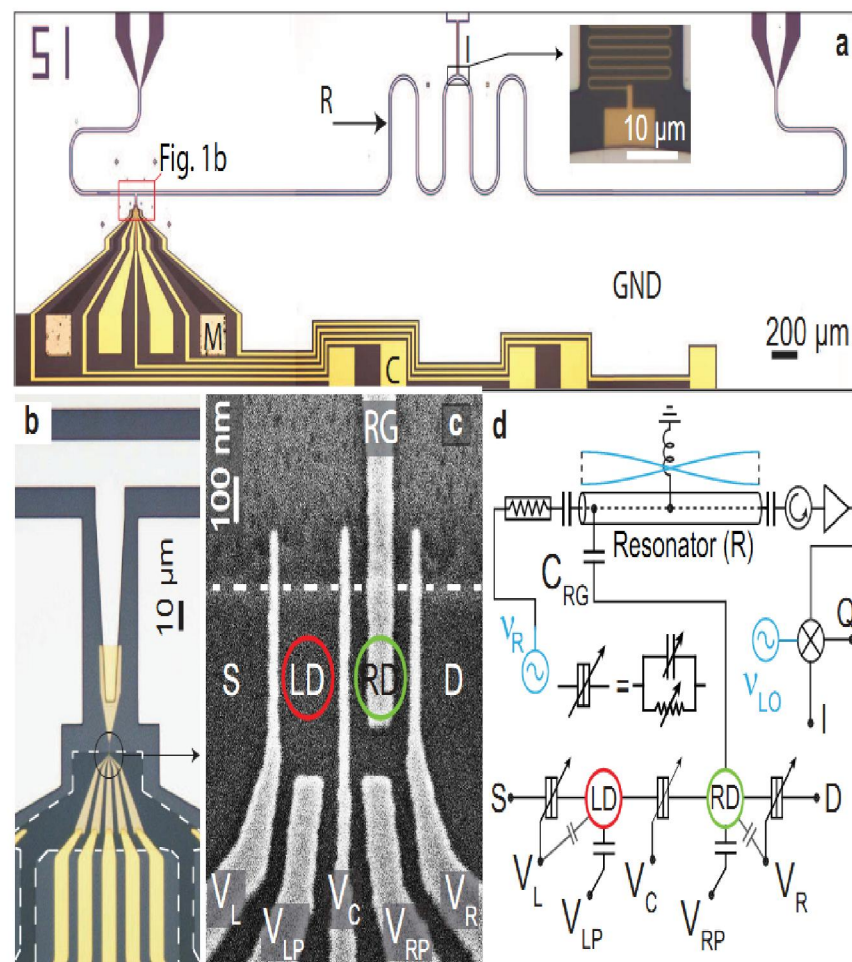
量子数据线和量子比特的长程耦合

Characterization of a microwave frequency resonator via a nearby quantum dot



Appl. Phys. Lett. 262105 98 (2011)

Dipole-coupling of a double quantum dot to a microwave resonator



Phys. Rev. Lett. 108, 046807 2012

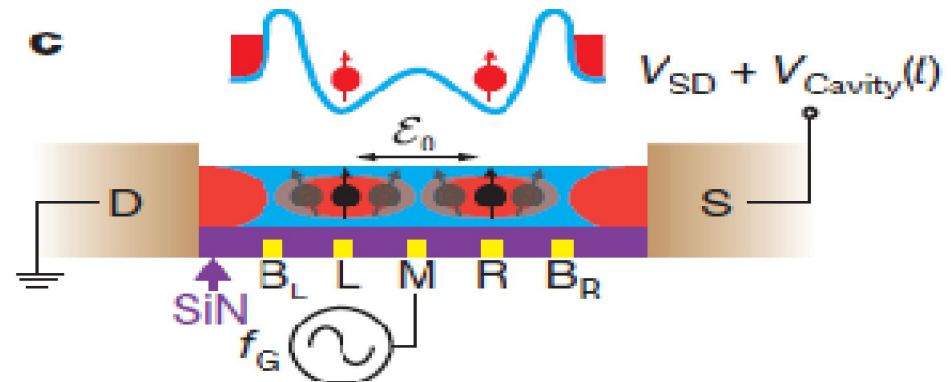
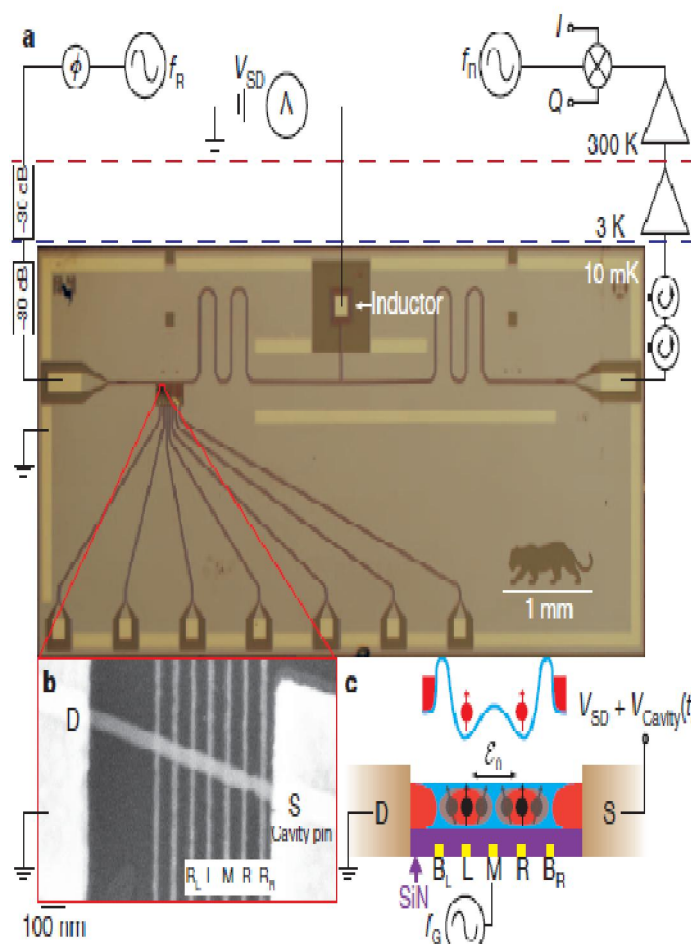
量子数据线和量子比特的长程耦合

LETTER

doi:10.1038/nature11559

Circuit quantum electrodynamics with a spin qubit

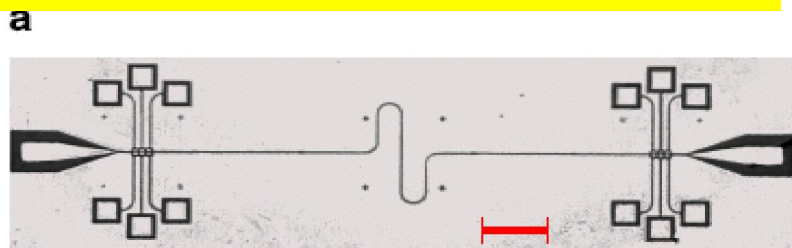
K. D. Petersson¹, L. W. McFaul¹, M. D. Schroer¹, M. Jung¹, J. M. Taylor², A. A. Houck³ & J. R. Petta^{1,4}



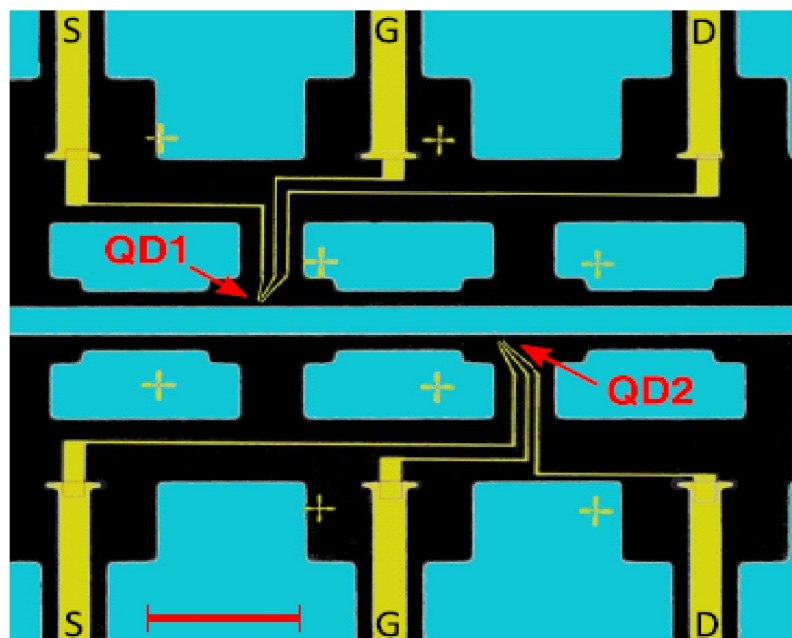
cavity^{8,9}. The architecture allows us to achieve a **charge-cavity coupling rate of about 30 megahertz**, consistent with coupling rates obtained in gallium arsenide quantum dots¹⁰. Furthermore, the **strong spin-orbit interaction** of indium arsenide allows us to drive spin rotations electrically with a local gate electrode, and the charge-cavity interaction provides a measurement of the resulting spin dynamics. Our results demonstrate how the cQED architecture can be used **as a sensitive probe of single-spin physics** and that a **spin-cavity coupling rate of about one megahertz** is feasible, presenting the possibility of long-range spin coupling via superconducting microwave cavities.

量子数据线和量子比特的长程耦合

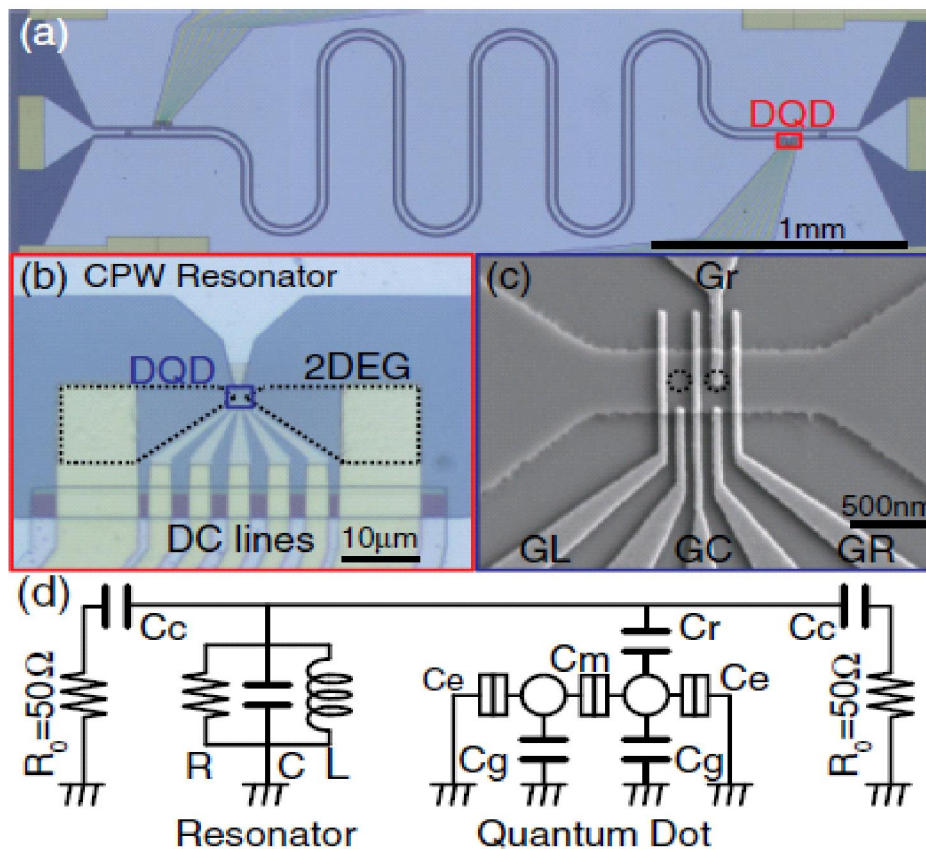
Photon mediated interaction between distant quantum dot circuits



Nature Communications 4, 1400 (2013)



Vacuum Rabi Splitting in a Semiconductor Circuit QED System ???



PRL 110, 066802 (2013) and its comment on arXiv 1304.3697

量子数据线和量子比特的长程耦合

J-C Model: $H = \hbar\omega_r (a^\dagger a + 1/2) + \hbar \frac{\omega_a}{2} \sigma_z + \hbar g (a^\dagger \sigma^- + a \sigma^+)$

Strong coupling: $2g \geq (\kappa + \gamma_{\text{dot}})/2$

$g_C/2\pi = 30 \text{ MHz}$ $\kappa/(2\pi) = 8 \text{ MHz}$; $(\gamma_1, \gamma_\phi, \gamma)/(2\pi) = (200, 200, 300) \text{ MHz}$

- Increasing **coherence time** of QD and **the coupling strength** is the Key to get strong coupling region.

How about New material: Graphene and TMDs QD ?

1. Fabrication of graphene QD and its manipulation ?
2. Design of superconductor resonator for graphene ?
3. Coherence time for graphene QD qubits ?

Theory for graphene quantum qubit

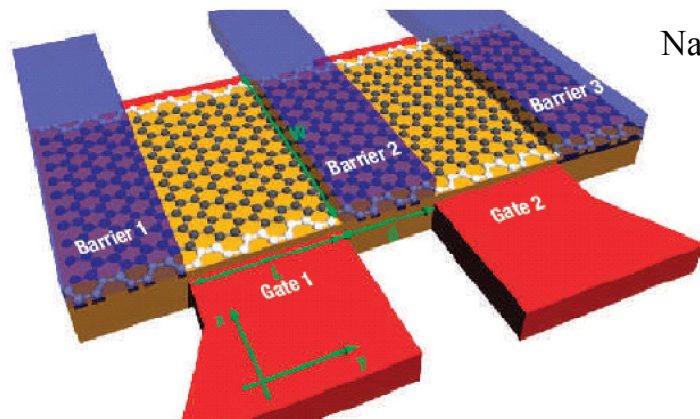
ARTICLES

Spin qubits in graphene quantum dots

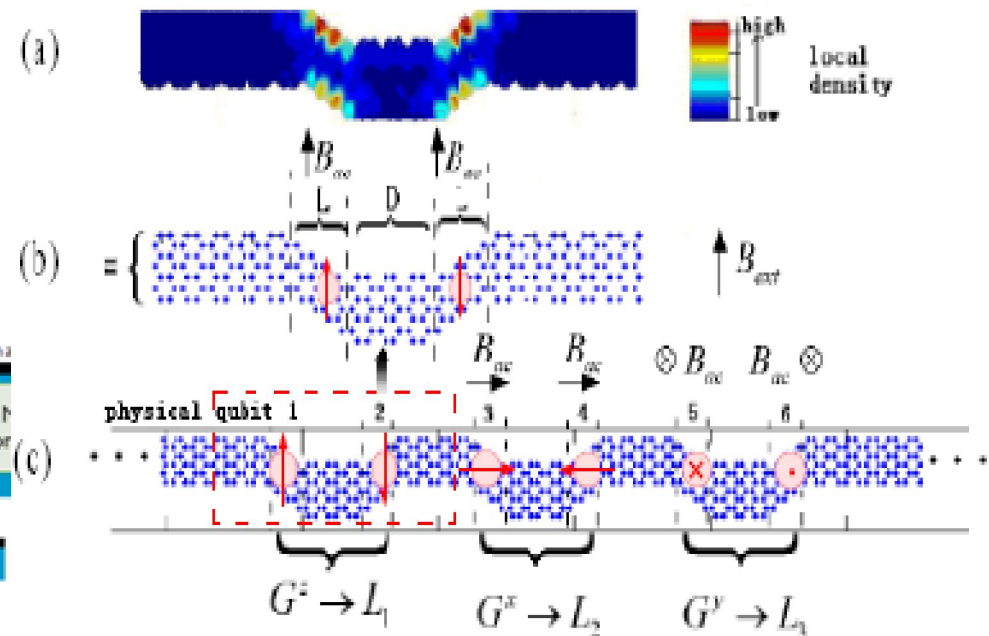
BJÖRN TRAUZETTEL, DENIS V. BULAEV, DANIEL LOSS AND GUIDO BURKARD*

Department of Physics and Astronomy, University of Basel, Klingelbergstrasse 82, CH-4056 Basel, Switzerland

*e-mail: Guido.Burkard@unibas.ch



Nature physics 3 192(2007)



G. P. Guo, et.al, [arXiv:0808.1618](https://arxiv.org/abs/0808.1618)
New J. Phys. 11 123005 (2009)

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LATEST NEWS ARTICLES

- Nano-lens goes cup-shaped
- Ripple control improves graphene electronics
- Flexible fabric that 'takes pictures'
- Quantum-dot arrays produced in nanotubes
- Photonic circuits move on

TECHNOLOGY UPDATE

Oct 24, 2008

Z-shaped result for quantum computing

Graphene nanoribbon could be used to make qubits for next-generation quantum computers. So say researchers at the Key Laboratory of Quantum Information in Hefei, China, who have shown that z-shaped ribbons can store electrons or holes in the corners of their Zs, where they can be then be written or read to.

Quantum coherence time T1 and T2 for graphene QD

ARTICLE

Received 18 Oct 2012 | Accepted 14 Mar 2013 | Published 23 Apr 2013

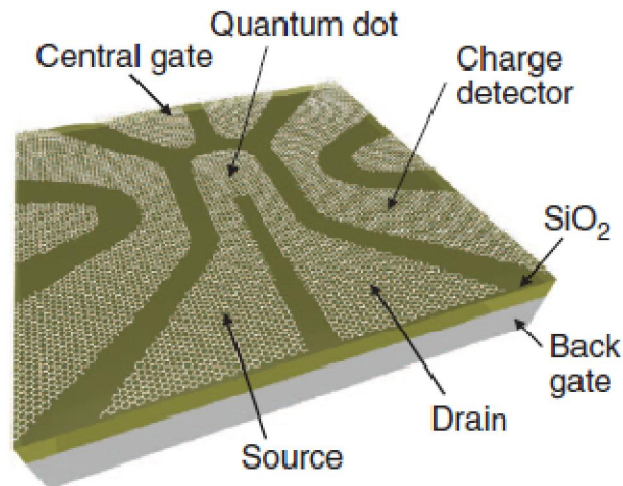
DOI: 10.1038/ncomms2738

OPEN

Probing relaxation times in graphene quantum dots

Christian Volk^{1,2,*}, Christoph Neumann^{1,2,*}, Sebastian Kazarski¹, Stefan Fringes¹, Stephan Engels^{1,2}, Federica Haupt³, André Müller^{1,2} & Christoph Stampfer^{1,2}

A lower bound for charge excited state relaxation times on the order of 60–100 ns.

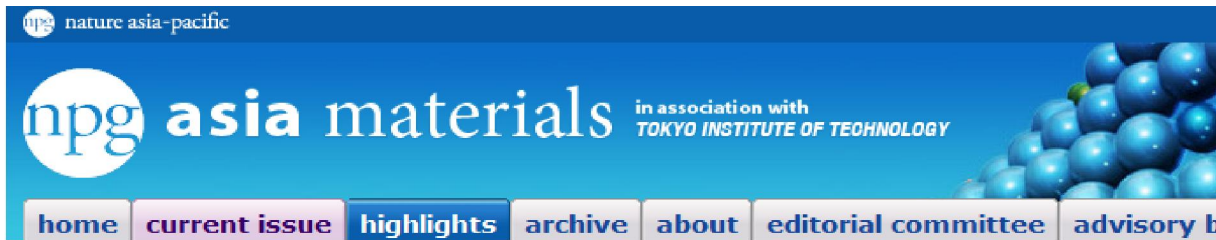


What is T2 for charge state ?

How about spin state ?

How to coupling distance graphene qubits ?

A graphene quantum dot with a single electron transistor as an integrated charge sensor



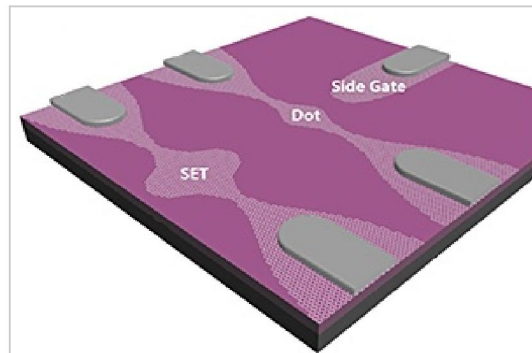
home » research highlight » Nanoelectronics: Quantum dots in charge

NPG Asia Materials research highlight | doi:10.1038/asiamat.2011.34
Published online 28 February 2011

Nanoelectronics: Quantum dots in charge

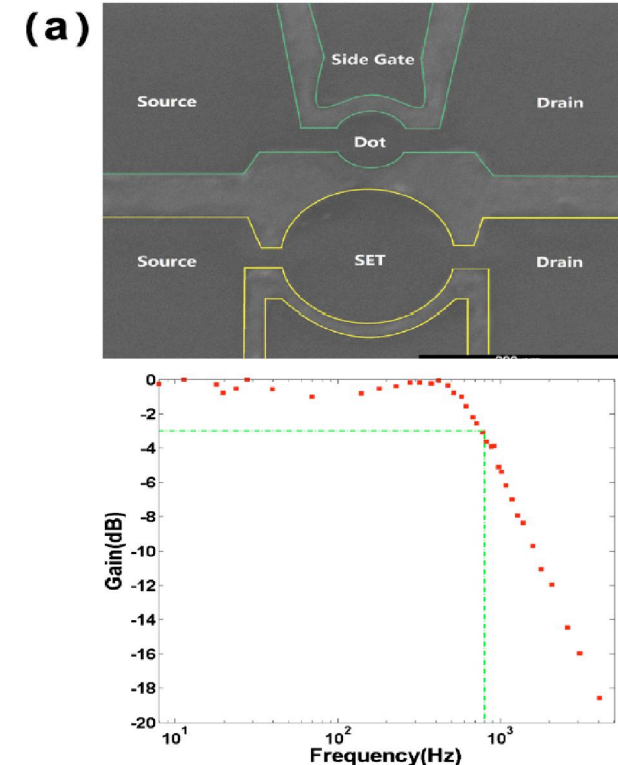
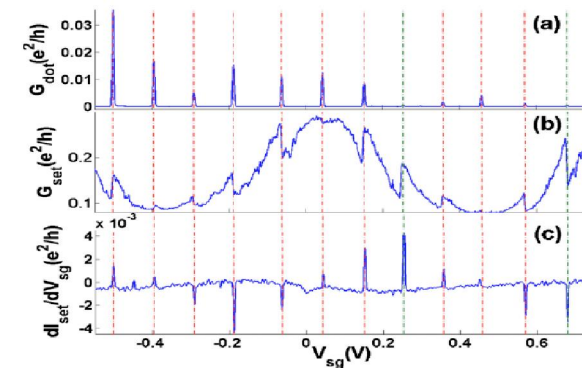
Paired graphene quantum dots could form the basis for an electron-counting information storage system.

Quantum dots could one day be used as solid-state quantum bits for quantum information processing. The information in a quantum dot would be encoded in the charge, spin or degree of freedom of electrons, and be read using ultrasensitive charge and spin detectors. A collaborative study between the University of Science and Technology of China and the University of California, Los Angeles, USA, has now demonstrated an integrated charge sensor system based on twin graphene dots¹, raising the hope that graphene's advantageous electronic and structural properties could serve for the future development of quantum information processing technologies.



Schematic illustration of an integrated charge-sensing system based on a pair of quantum dots

© 2011 AIP*



L. J. Wang, et.al. Appl. Phys. Lett. 97, 262113 (2010).

Graphene double quantum dot in series and in parallel

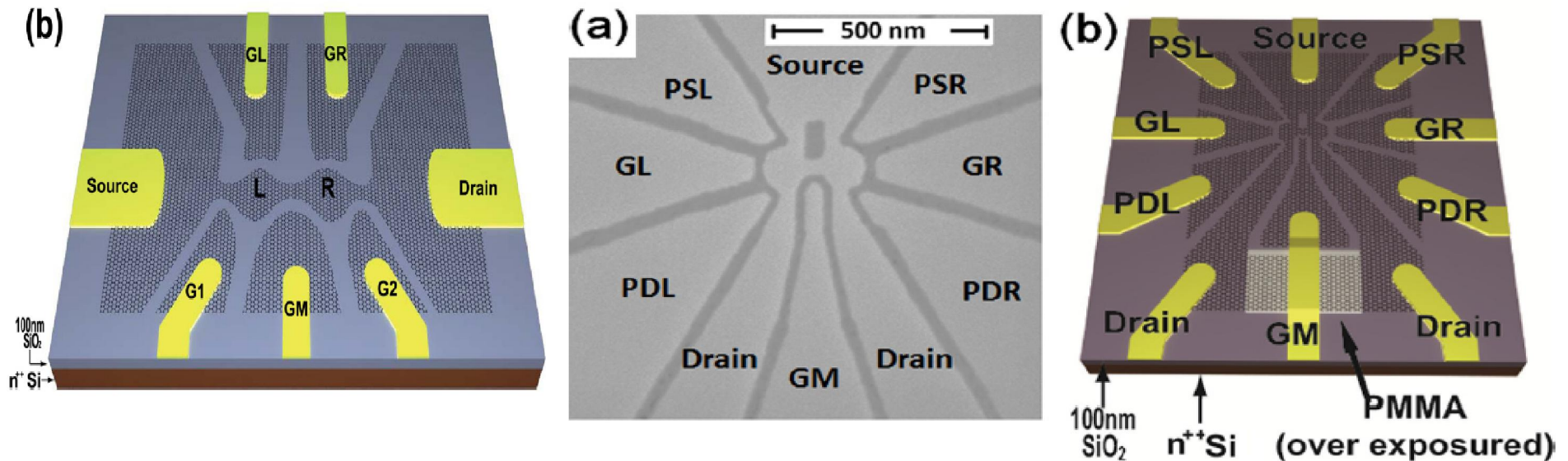
APPLIED PHYSICS LETTERS 99, 112117 (2011)

Gates controlled parallel-coupled double quantum dot on both single layer and bilayer graphene

Lin-Jun Wang,¹ Guo-Ping Guo,^{1,a)} Da Wei,¹ Gang Cao,¹ Tao Tu,^{1,b)} Ming Xiao,¹
Guang-Can Guo,¹ and A. M. Chang²

¹Key Laboratory of Quantum Information, University of Science and Technology of China, Chinese Academy of Sciences, Hefei 230026, People's Republic of China

²Department of Physics, Duke University, Durham, North Carolina 27708, USA



APPLIED PHYSICS LETTERS 100, 022106 (2012)

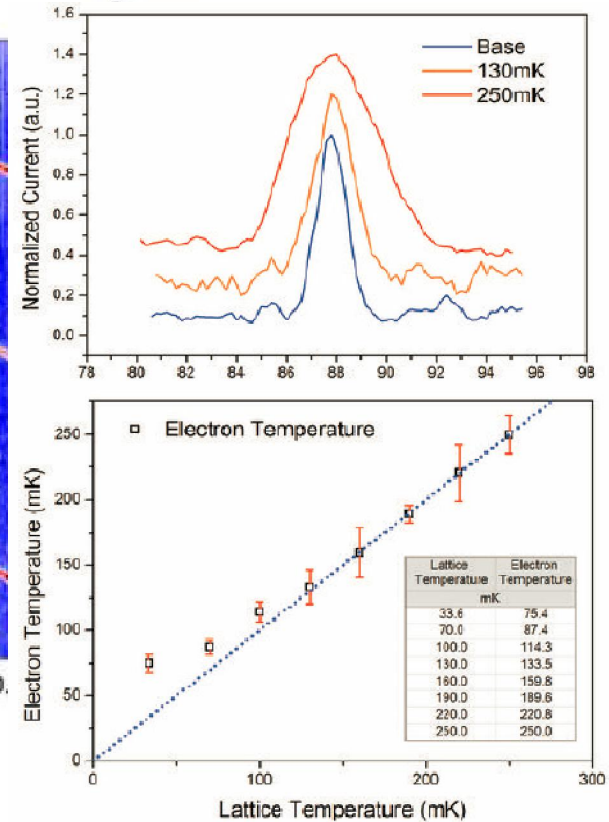
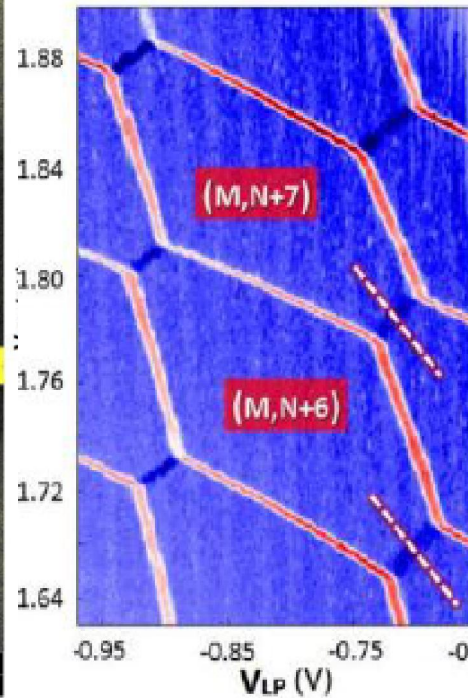
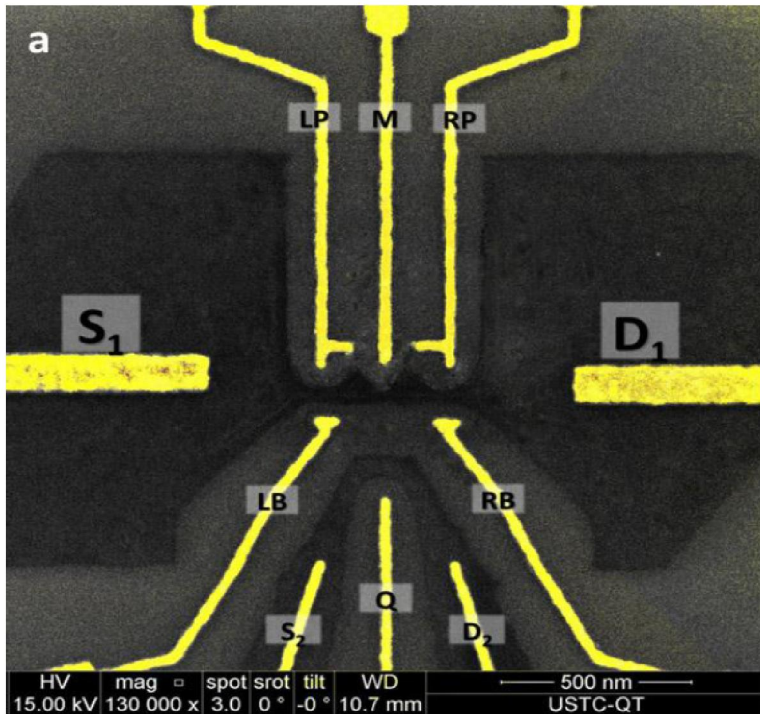
Controllable tunnel coupling and molecular states in a graphene double quantum dot

Lin-Jun Wang,¹ Hai-Ou Li,¹ Tao Tu,^{1,a)} Gang Cao,¹ Cheng Zhou,¹ Xiao-Jie Hao,¹ Zhan Su,¹
Ming Xiao,¹ Guang-Can Guo,¹ Albert M. Chang,² and Guo-Ping Guo^{1,b)}

¹Key Laboratory of Quantum Information, University of Science and Technology of China, Chinese Academy of Sciences, Hefei 230026, People's Republic of China

²Department of Physics, Duke University, Durham, North Carolina 27708, USA

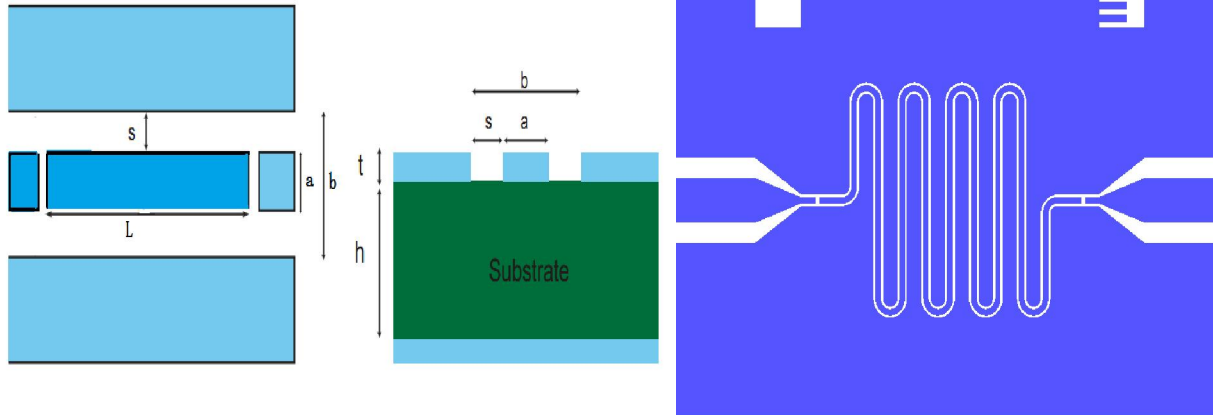
金属电极石墨烯双量子点



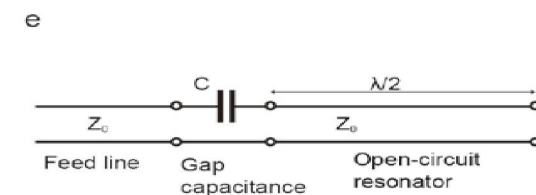
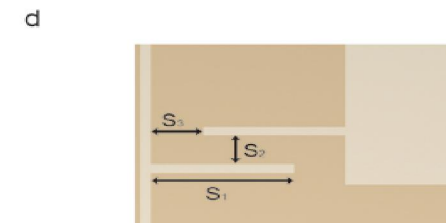
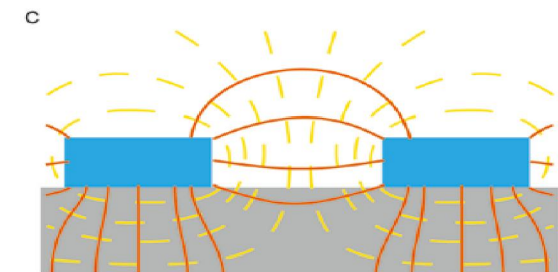
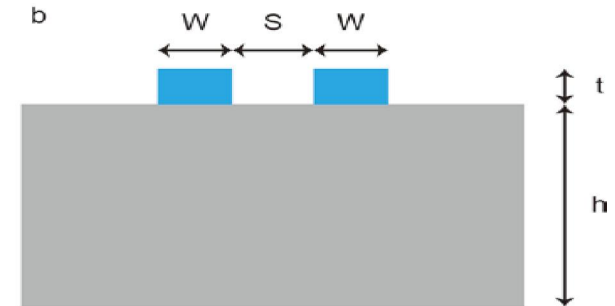
全金属电极：克服石墨烯电极边缘态的影响

Symmetric reflection line resonator (RLR)

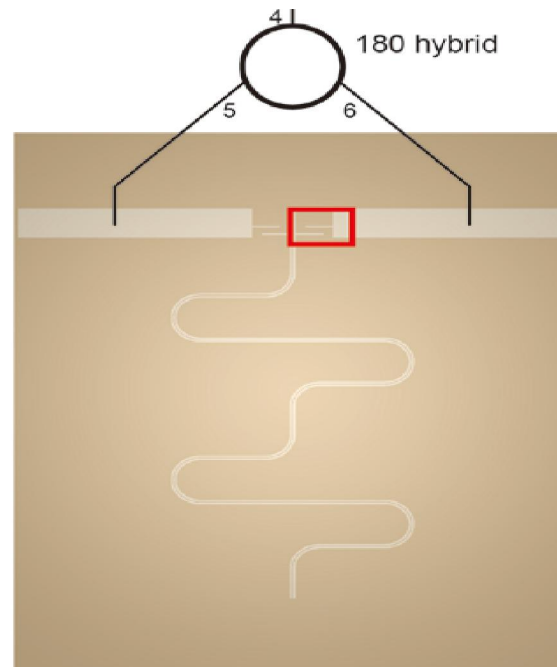
Transmission line resonator (TLR)



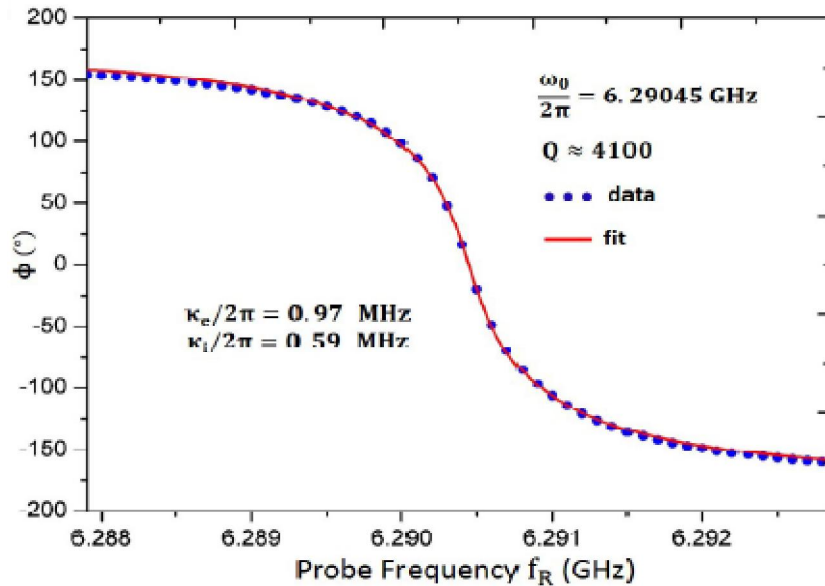
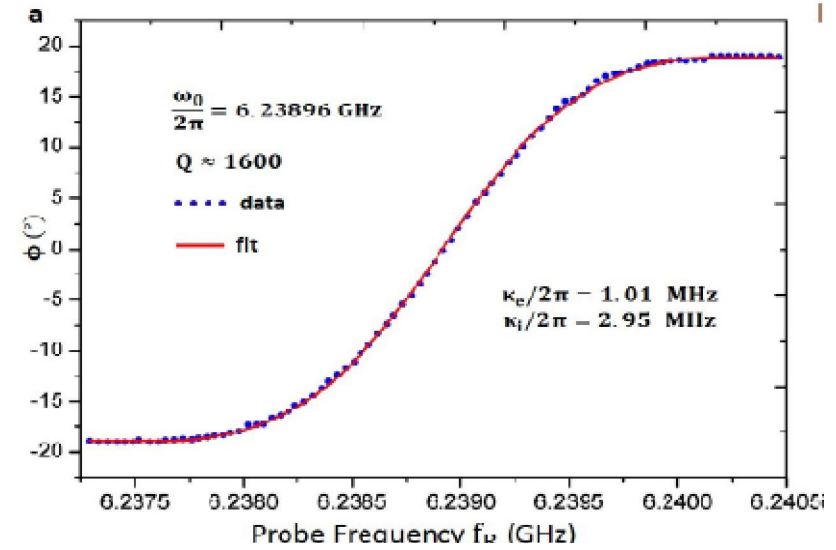
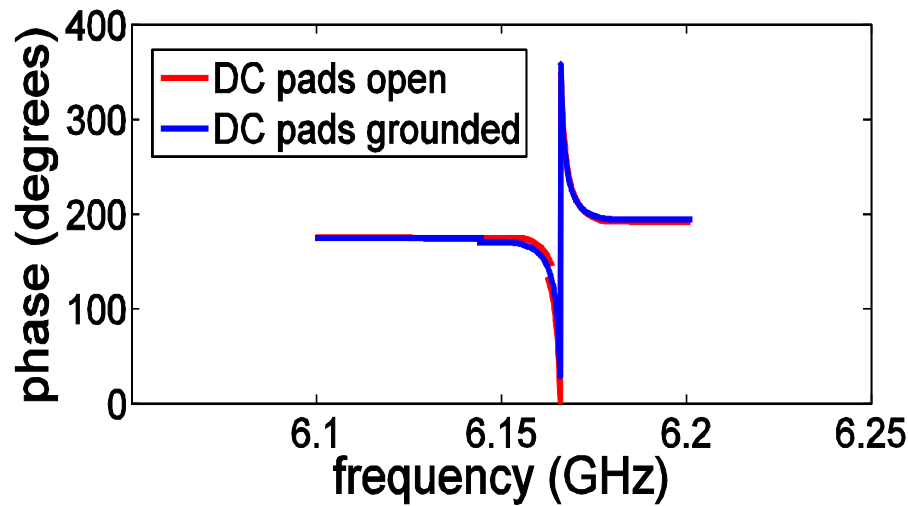
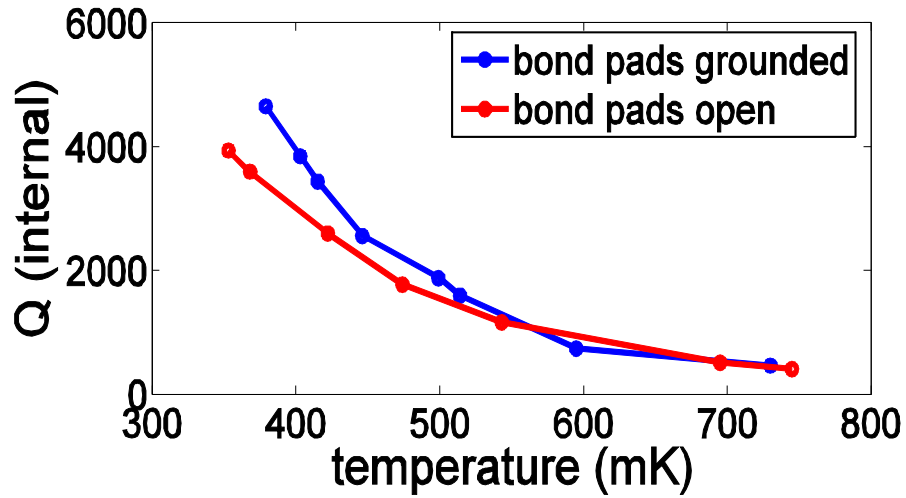
Our reflection line resonator (RLR)



- **1/2 or 1/4 wavelength reflection line resonator**
- **Two coplanar strip-lines geometry**
- **Coupling strength is expected to be double**
- **Bias T for independent DC bias of either side**

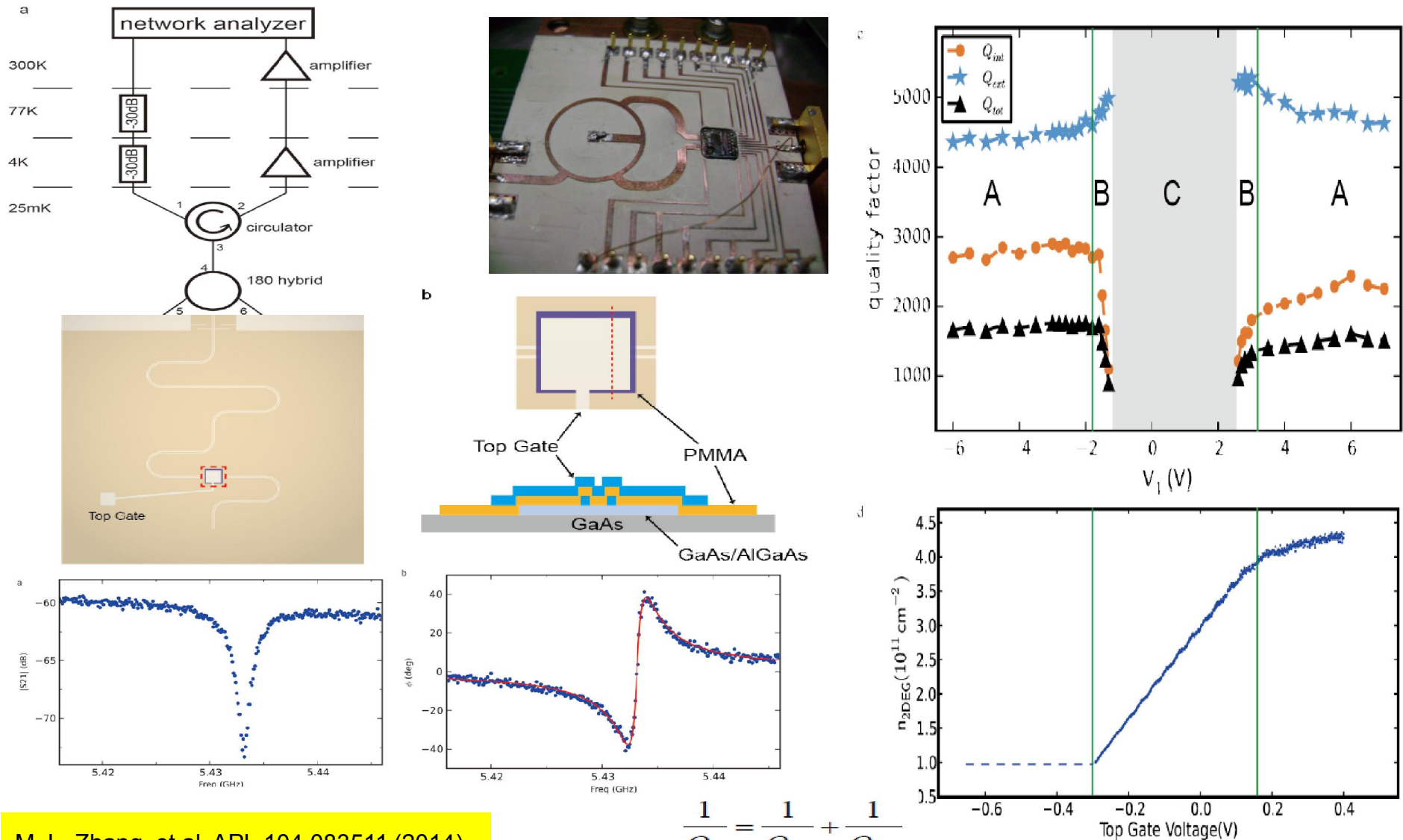


The property of reflection line resonator



The effect of DC pads, temperature, and dissipation to the cavity properties.

Tuning the Q factor of Symmetric reflection line resonator



M. L. Zhang, et.al, APL 104,083511 (2014).
M. L. Zhang, et.al, APL 105,073510 (2014).

$$\frac{1}{Q_L} = \frac{1}{Q_{int}} + \frac{1}{Q_{ext}}$$

石墨烯量子点与超导谐振腔复合系统

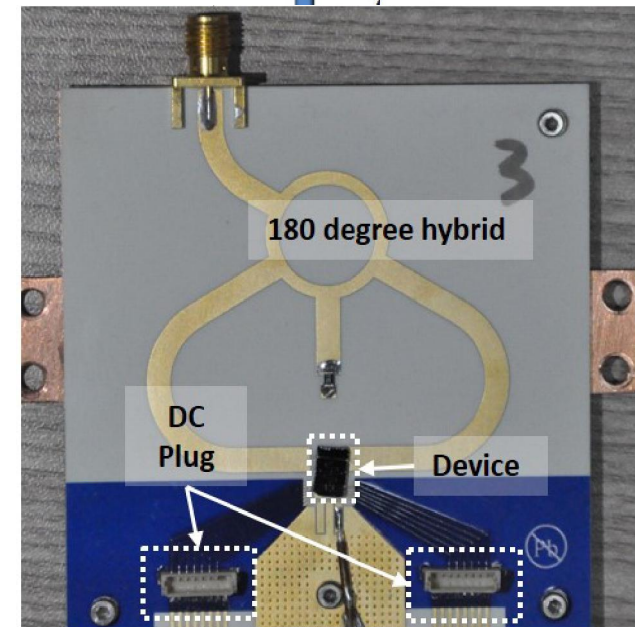
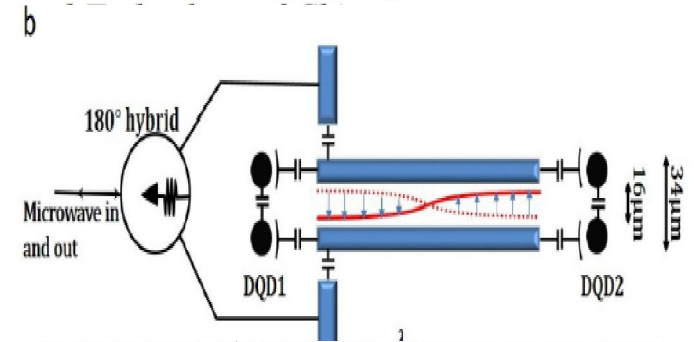
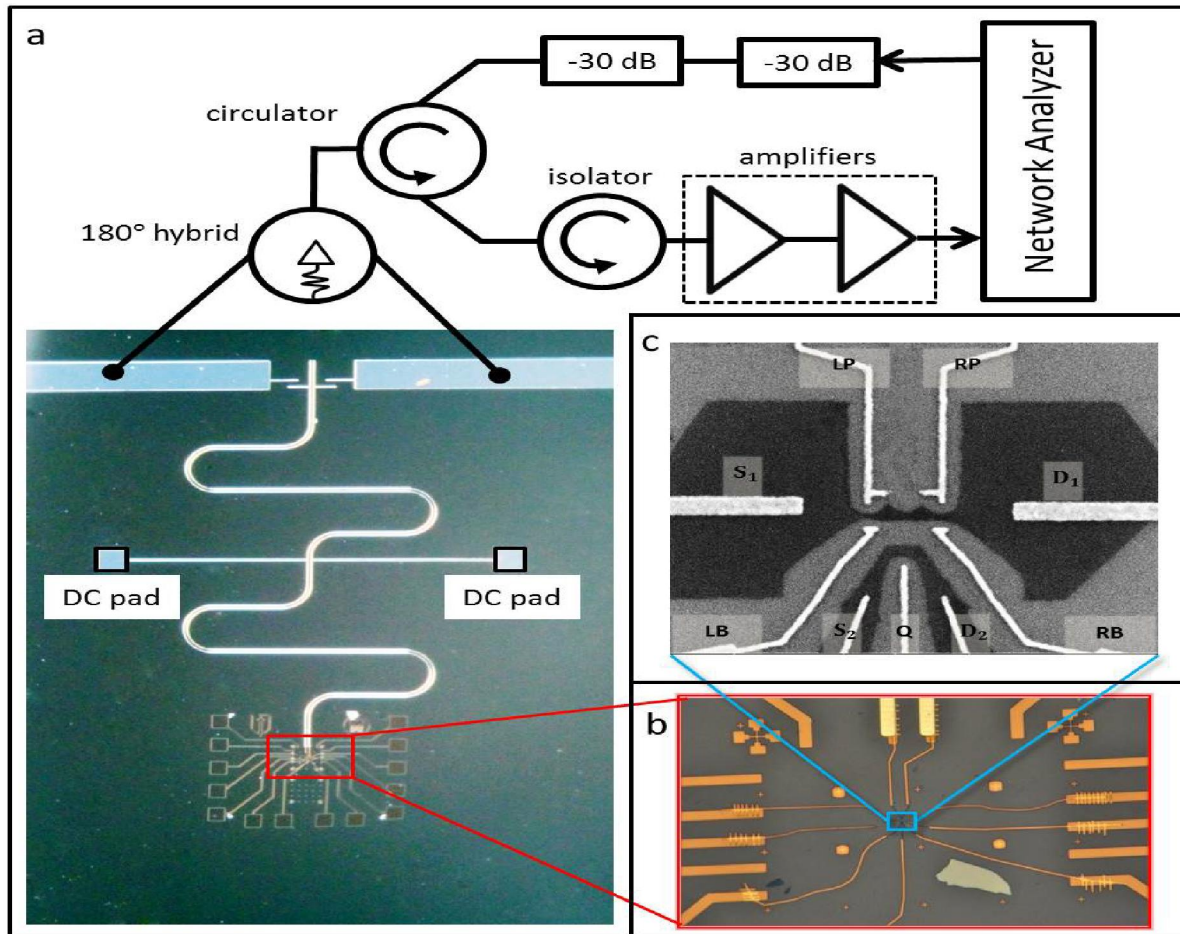
PRL 115, 126804 (2015)

PHYSICAL REVIEW LETTERS

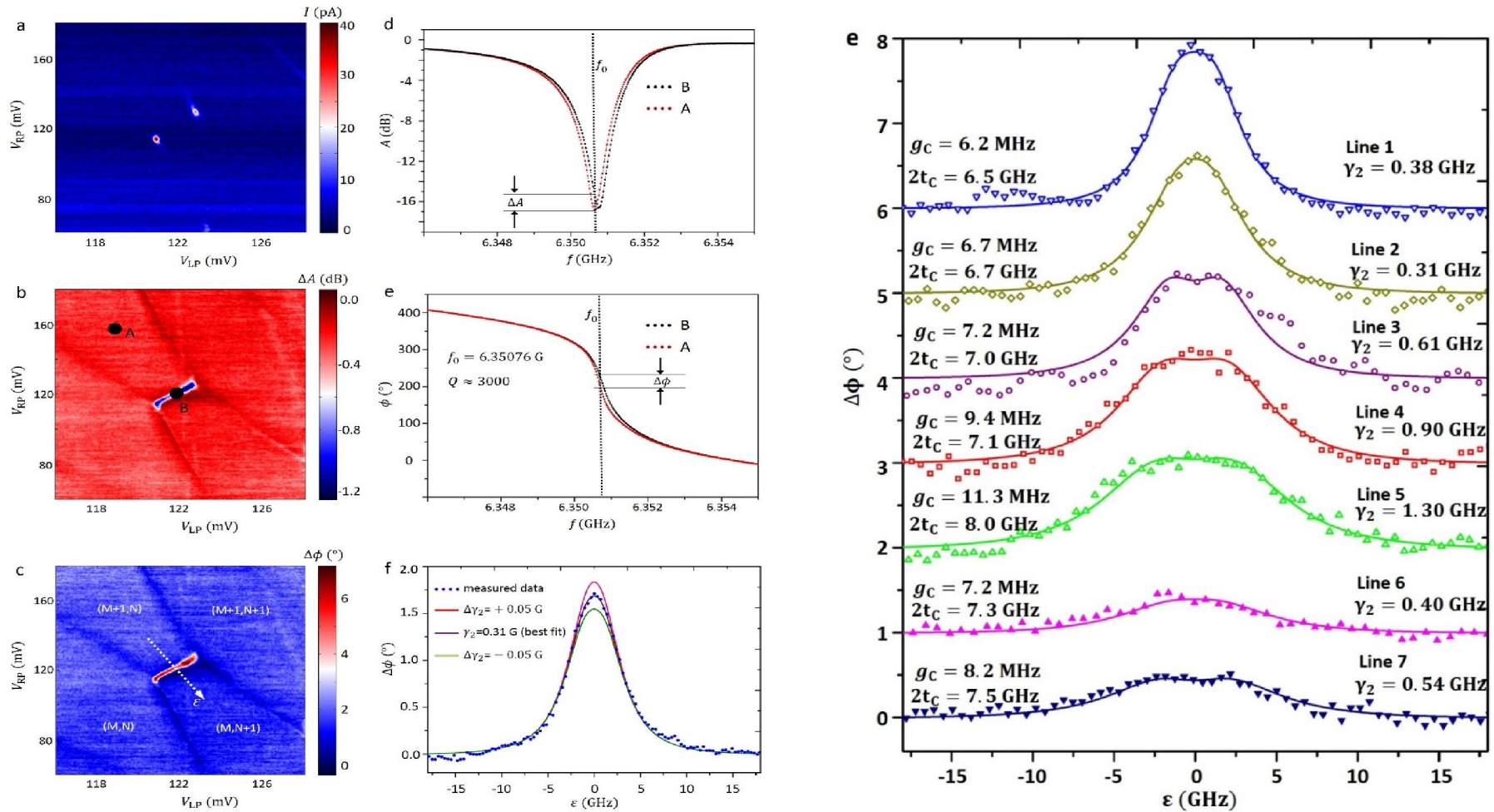
week ending
18 SEPTEMBER 2015

Charge Number Dependence of the Dephasing Rates of a Graphene Double Quantum Dot in a Circuit QED Architecture

Guang-Wei Deng,^{1,2} Da Wei,^{1,2} J. R. Johansson,³ Miao-Lei Zhang,^{1,2} Shu-Xiao Li,^{1,2} Hai-Ou Li,^{1,2} Gang Cao,^{1,2} Ming Xiao,^{1,2} Tao Tu,^{1,2} Guang-Can Guo,^{1,2} Hong-Wen Jiang,⁴ Franco Nori,^{5,6} and Guo-Ping Guo^{1,2,*}



石墨烯量子点与超导谐振腔复合系统

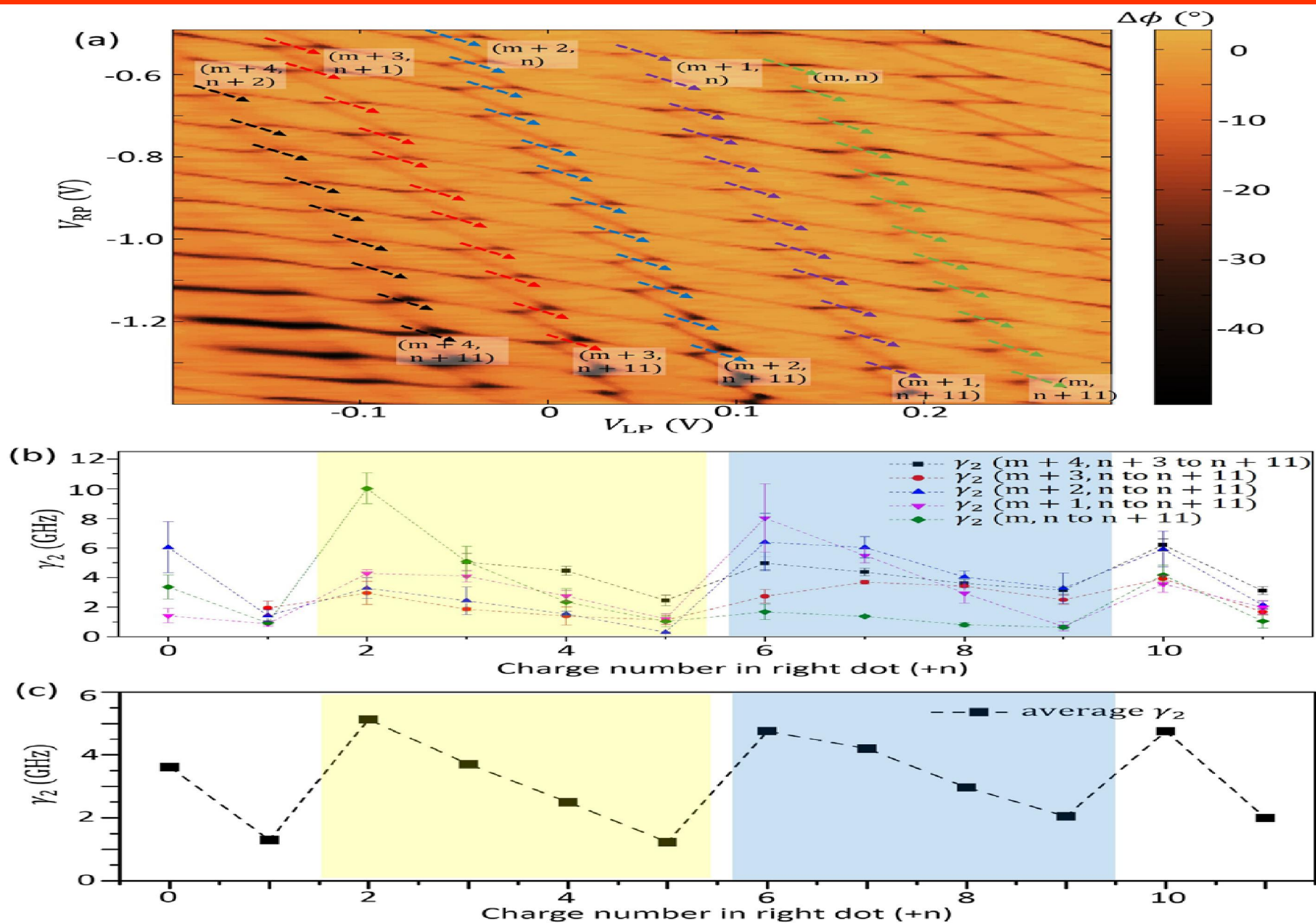


$$S_{11} = \frac{a_{out}}{a_{in}} = -\frac{i(\omega_0 - \omega) + g_{eff}\chi + \frac{\kappa_i - \kappa_e}{2}}{i(\omega_0 - \omega) + g_{eff}\chi + \frac{\kappa_i + \kappa_e}{2}}$$

$$g_{eff} = g_C \frac{2t_C}{\Omega}$$

$$\chi = \frac{g_{eff}}{i(\Omega - \omega) + \frac{1}{2}\gamma_1 + \gamma_2} \quad \Omega = \sqrt{(2t_C)^2 + \epsilon^2}$$

石墨烯量子点与超导谐振腔复合系统



石墨烯量子点与超导谐振腔复合系统

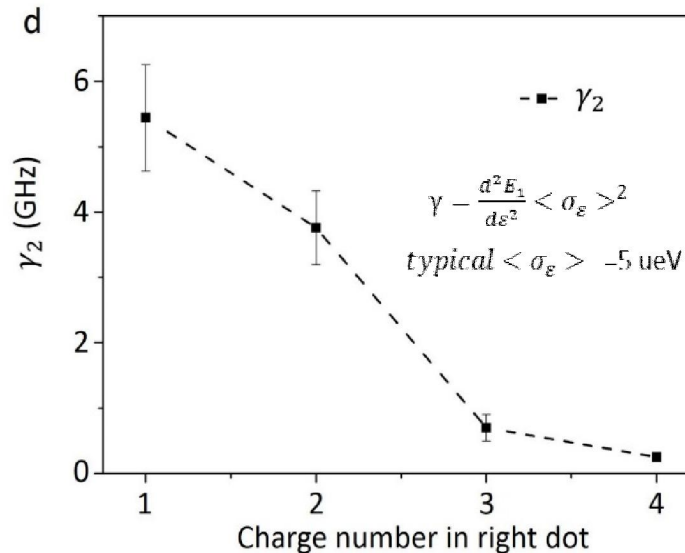
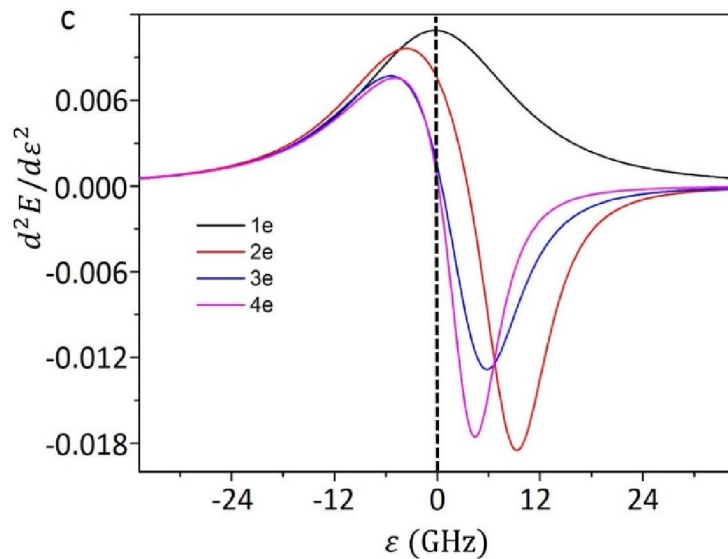
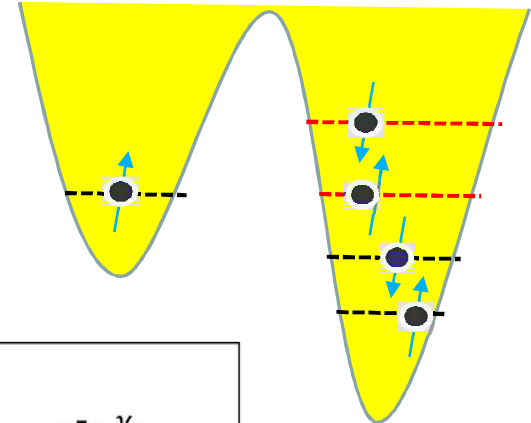
相干时间与能级差的偏倒数成正比：

PRL 108, 046807 (2012), PPL 105,063105 (2014), PRB 89, 165404 (2014)

$$\Omega(\delta, t) = \sqrt{\delta^2 + (2t)^2}$$

dephasing rate via the

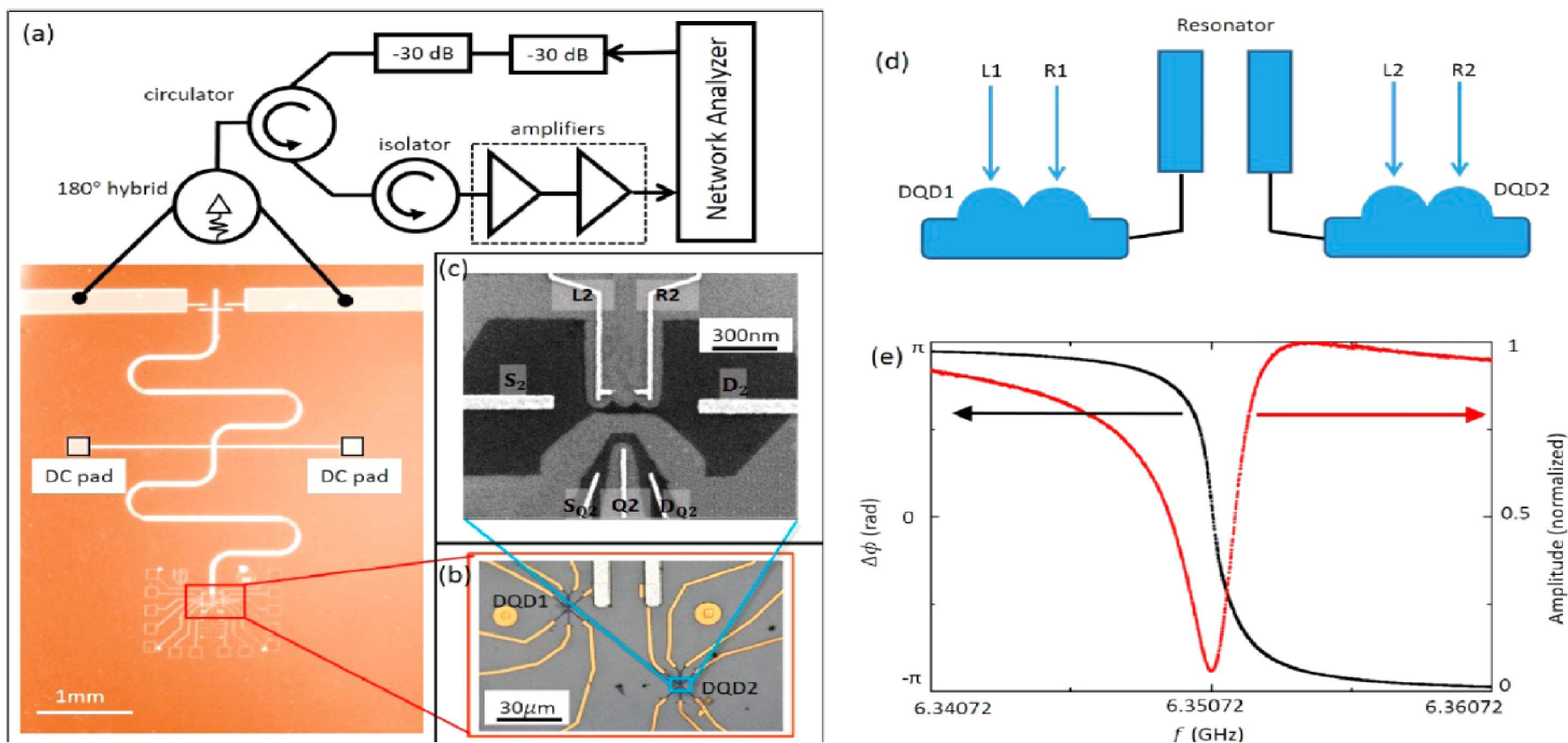
$$\gamma_\varphi/2\pi \approx \left. \frac{d^2\Omega}{d\delta^2} \right|_{\delta=0} S_\delta(\nu = 1 \text{ Hz}) = S_\delta(\nu = 1 \text{ Hz})/2t.$$



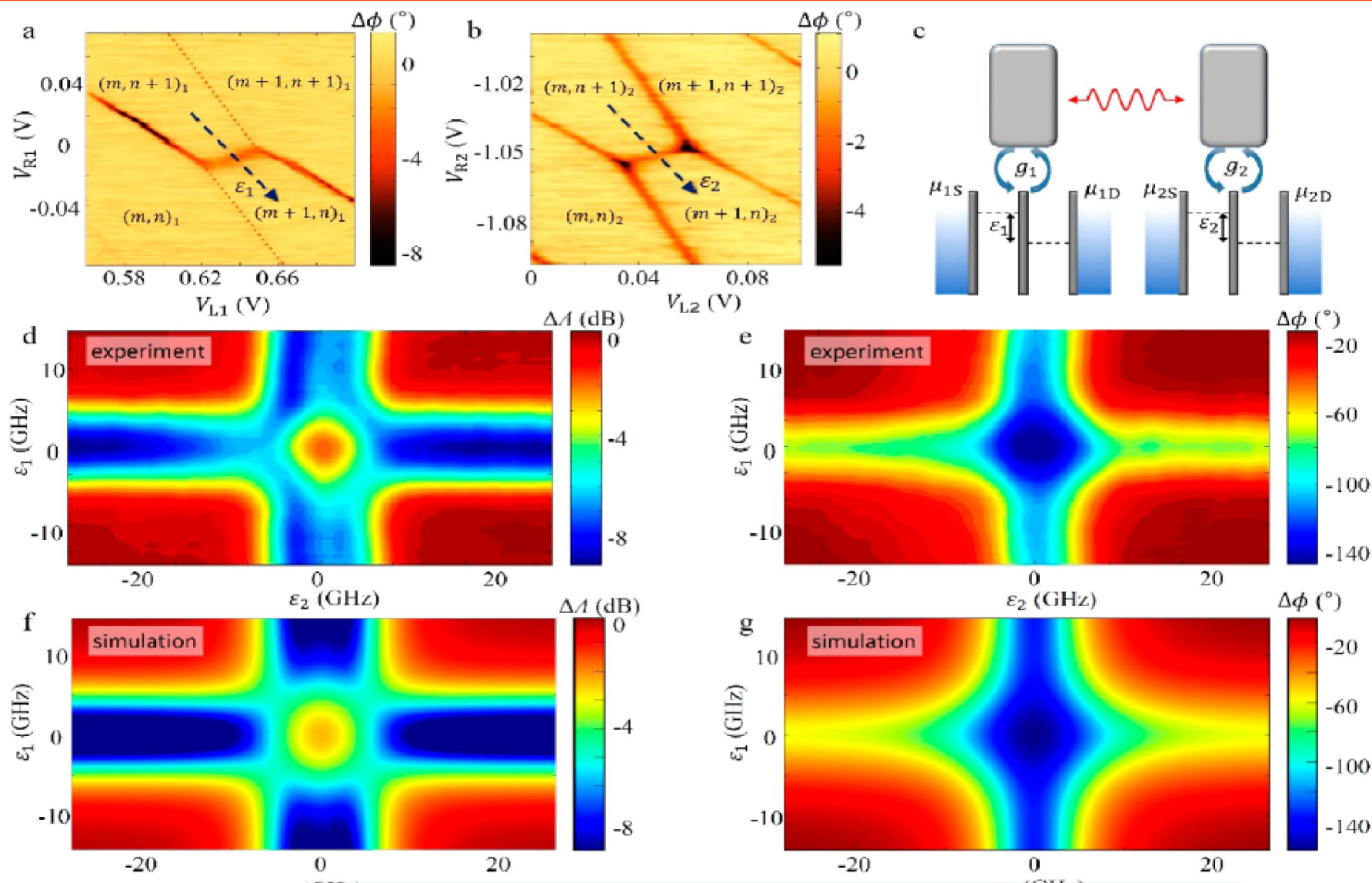
石墨烯量子比特的长程耦合

Coupling Two Distant Double Quantum Dots with a Microwave Resonator

Guang-Wei Deng,^{†,‡} Da Wei,^{†,‡} Shu-Xiao Li,^{†,‡} J. R. Johansson,[§] Wei-Cheng Kong,^{†,‡} Hai-Ou Li,^{†,‡} Gang Cao,^{†,‡} Ming Xiao,^{†,‡} Guang-Can Guo,^{†,‡} Franco Nori,^{||,⊥} Hong-Wen Jiang,[#] and Guo-Ping Guo^{*,†,‡}

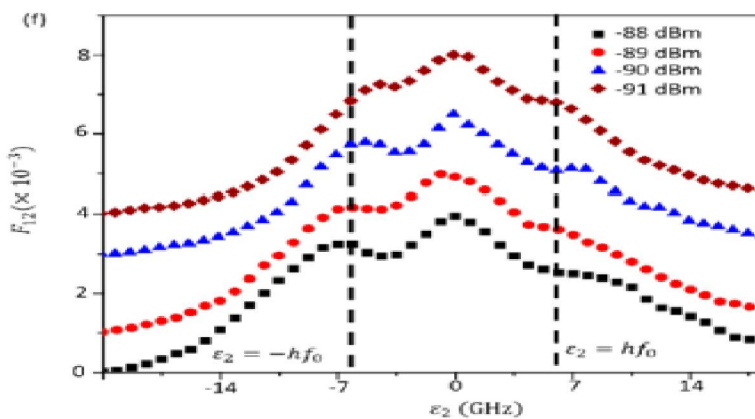
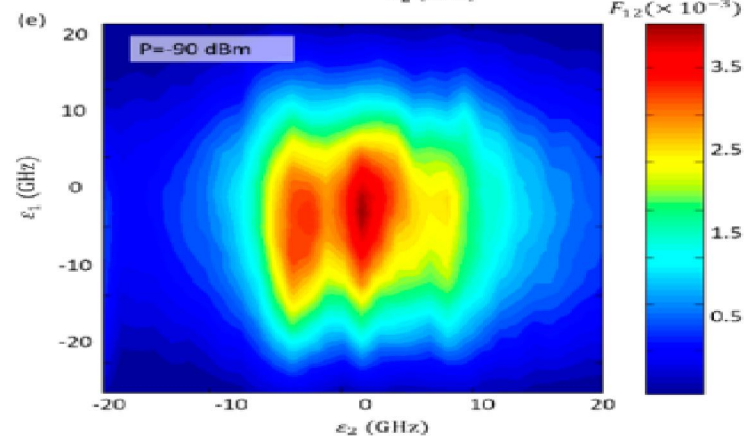
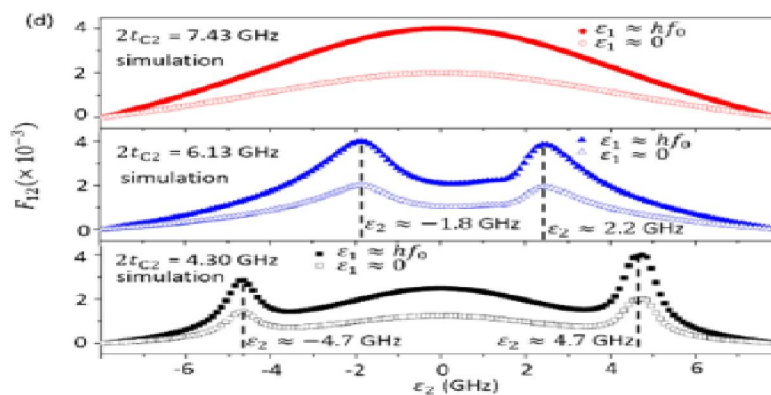
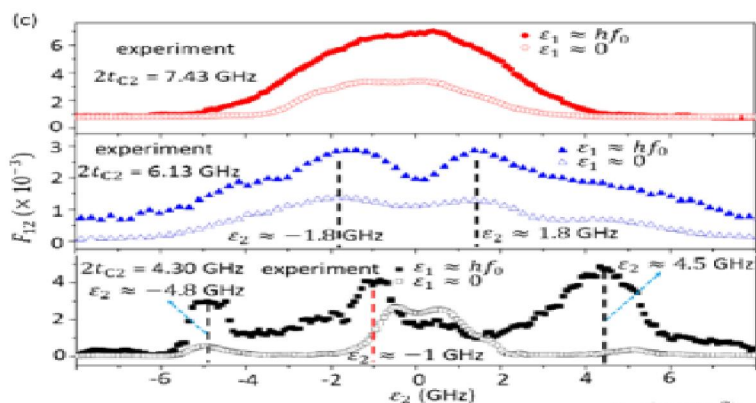
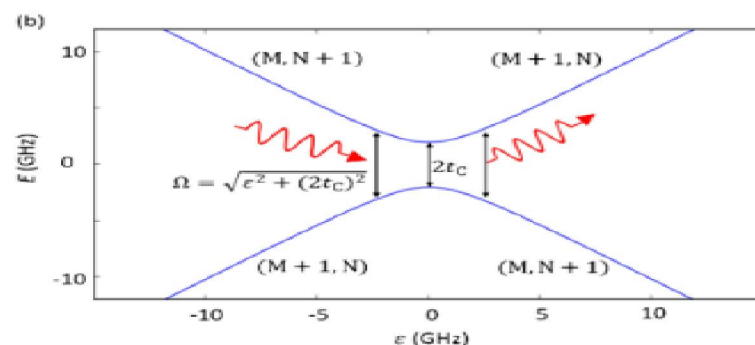
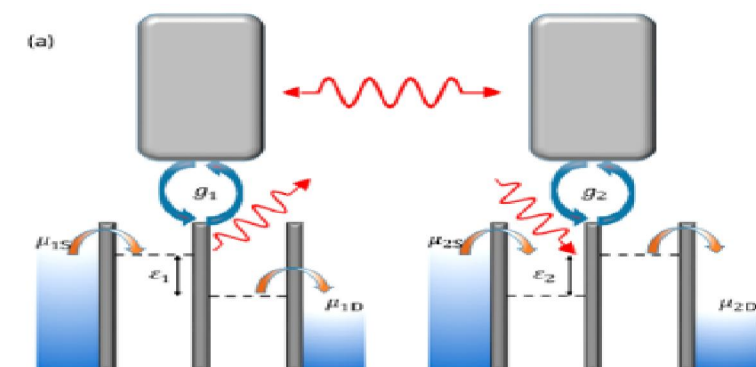


石墨烯量子比特的长程耦合

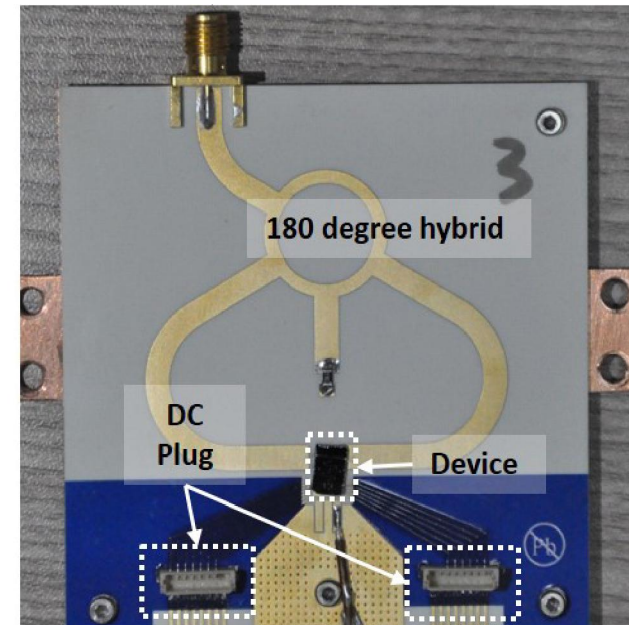
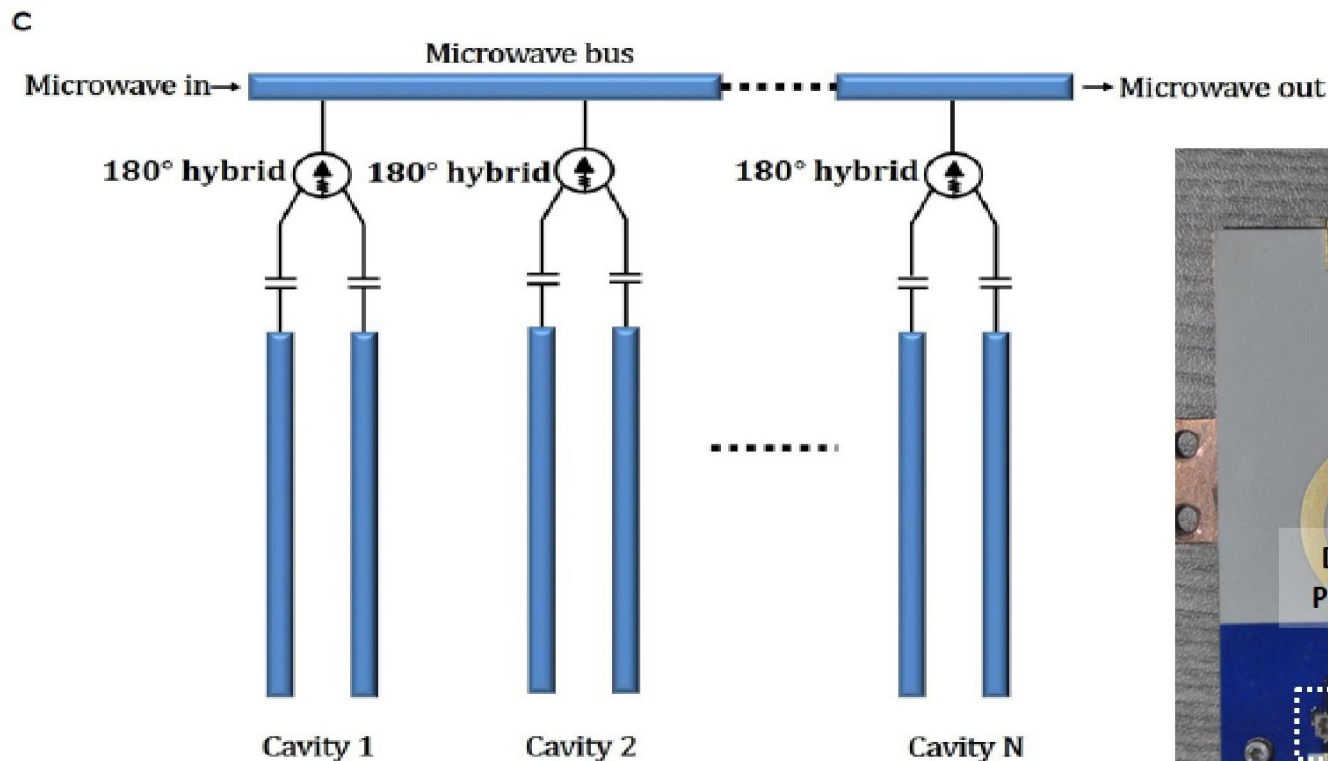
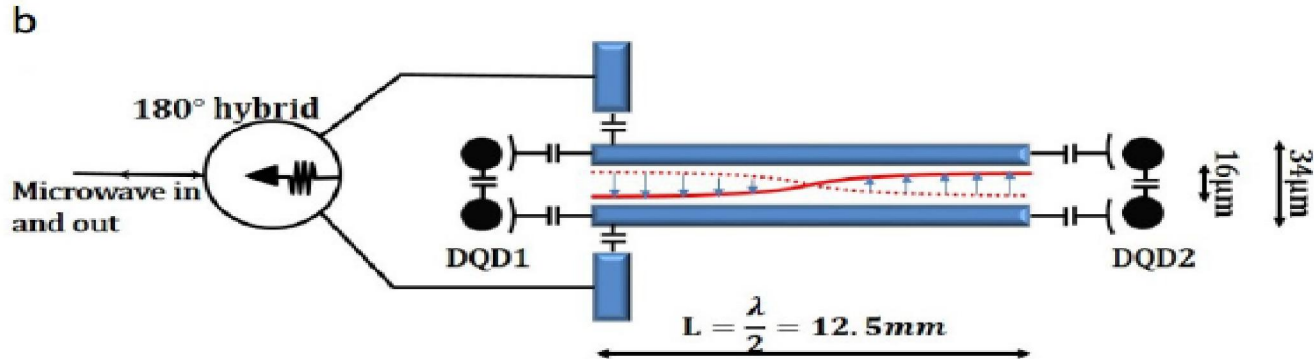


T-C Hamiltonian:
$$H = \omega_0 a^\dagger a + \sum_{i=1,2} \left[\frac{1}{2} \Omega_i \sigma_{zi} + g_i (\sigma_{+i} a + \sigma_{-i} a^\dagger) \right]$$

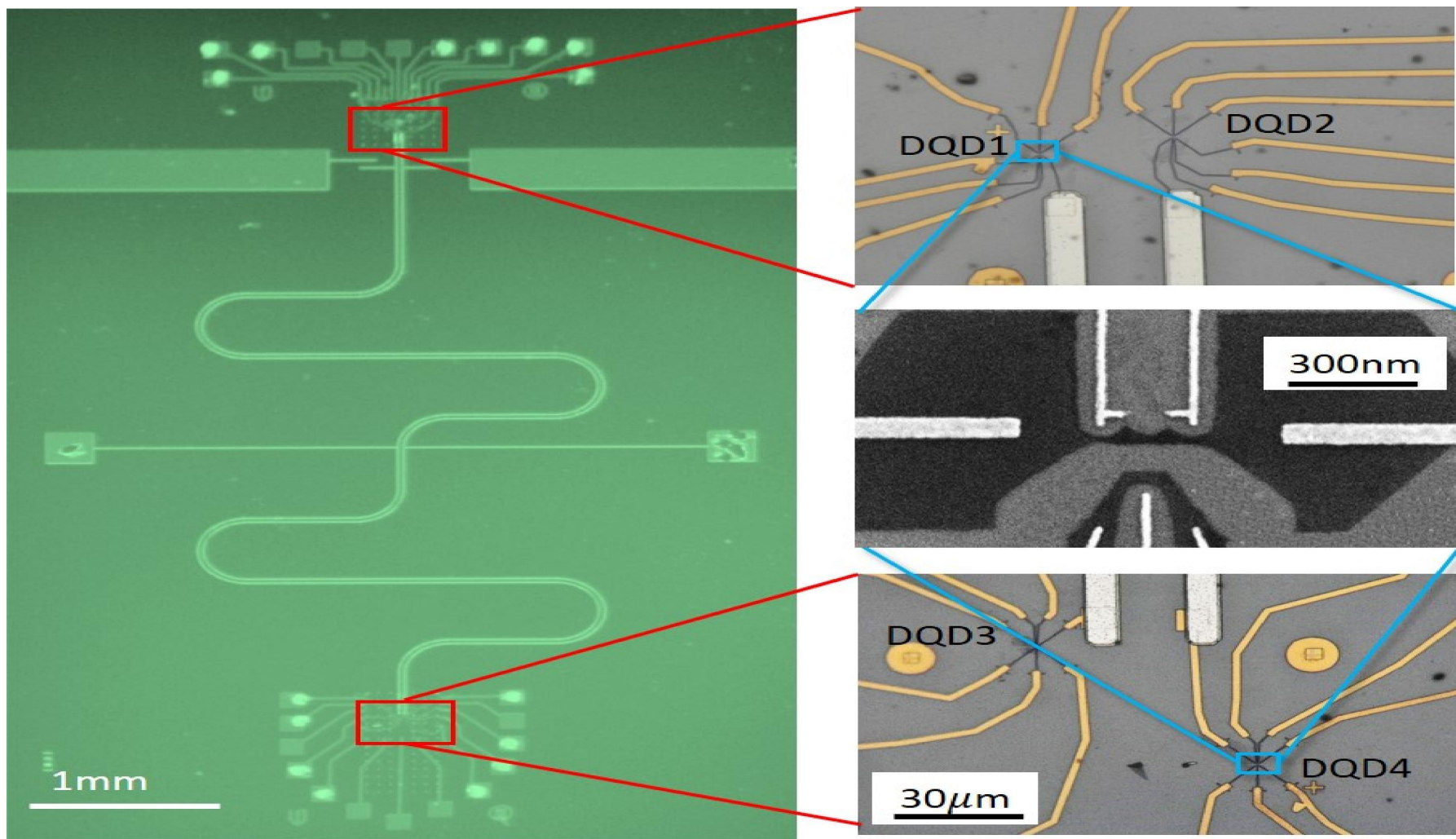
石墨烯量子比特的长程耦合



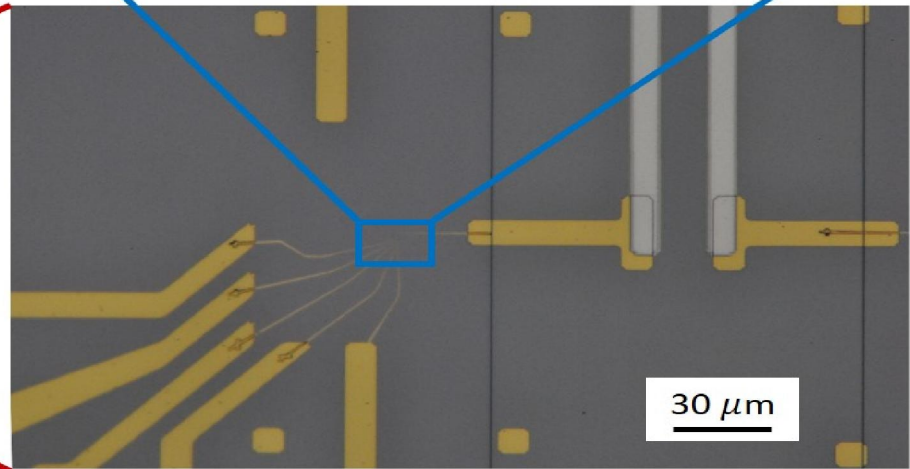
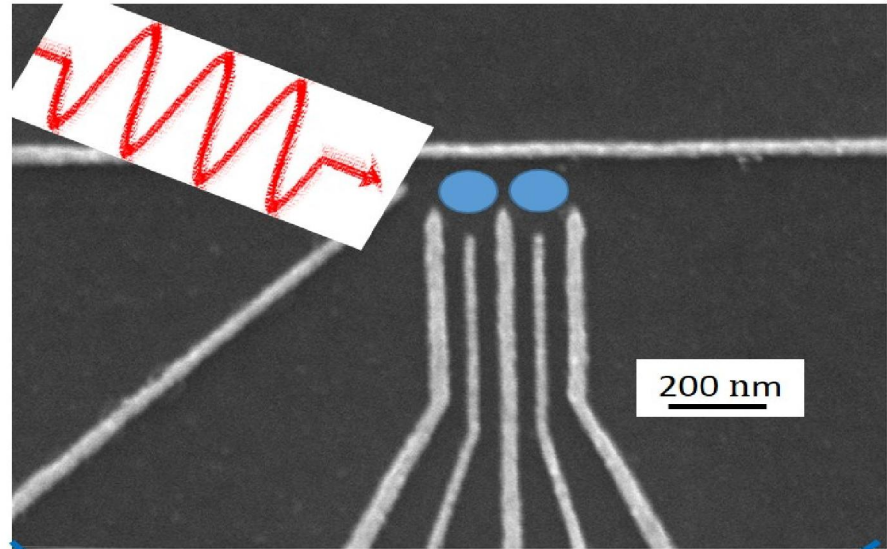
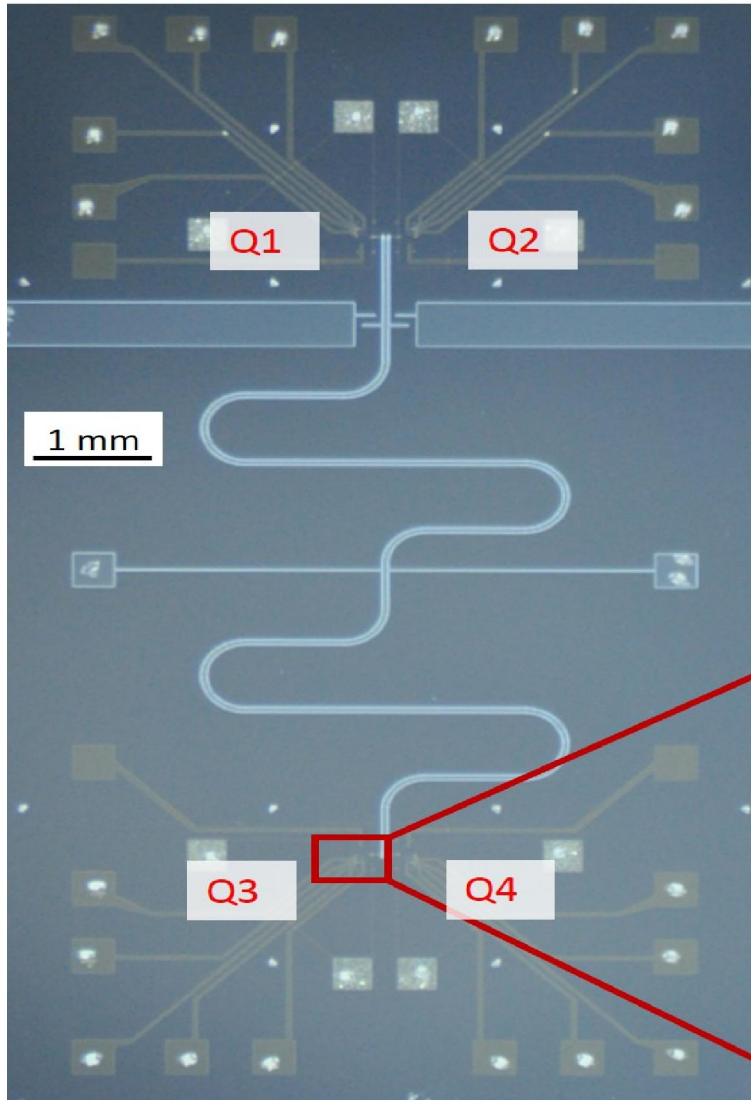
量子数据线和量子比特的长程耦合



量子数据线和量子比特的长程耦合

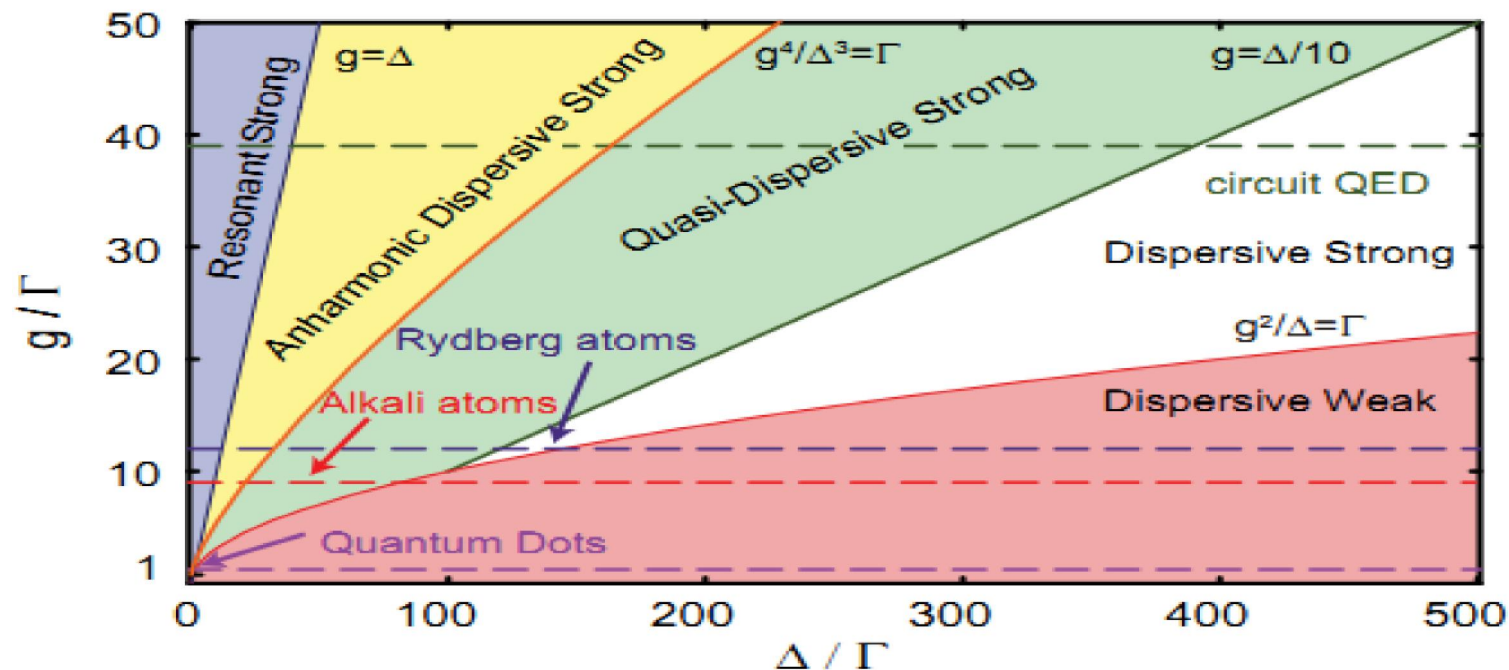


量子数据线和量子比特的长程耦合



量子数据线和量子比特的长程耦合

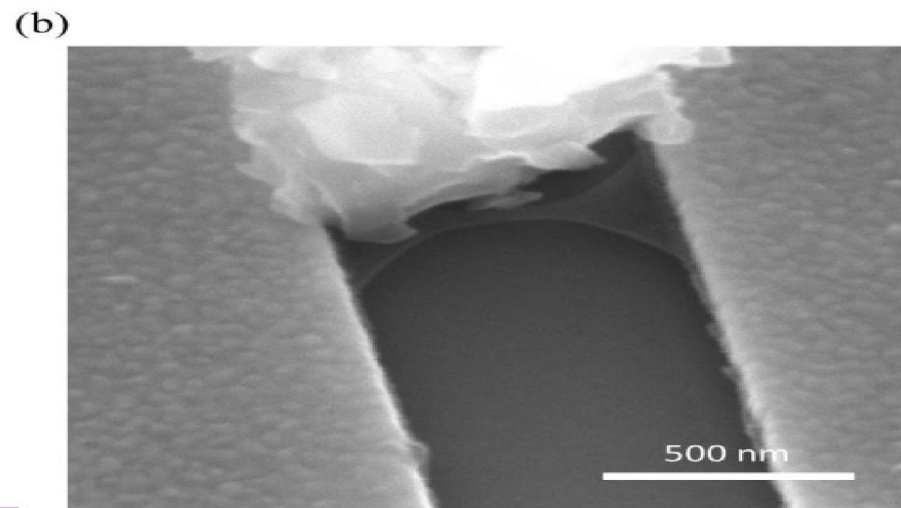
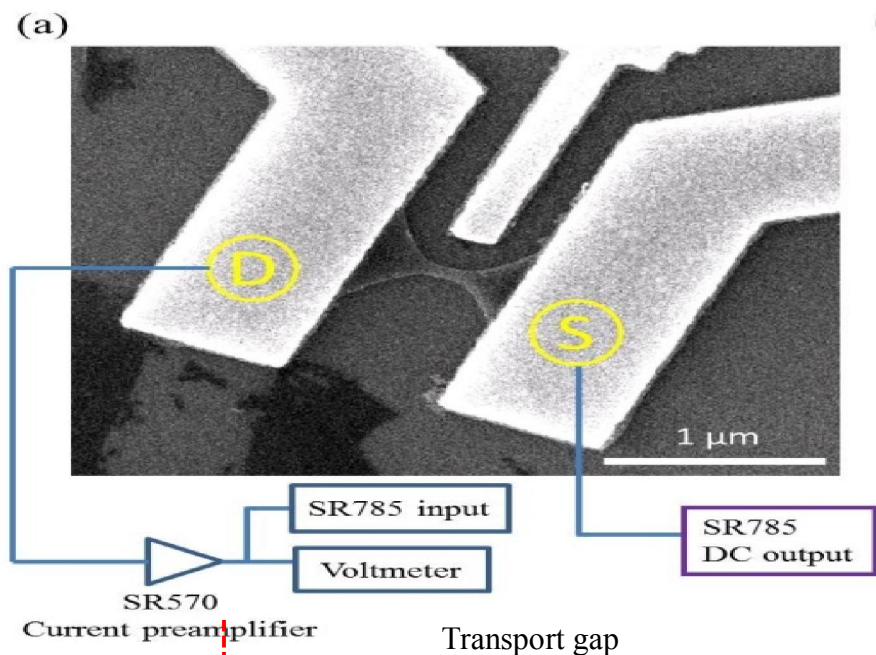
- Cavity decay $\kappa \sim \sim 3 \text{ MHz}$
- Quantum dot: $T_1 \sim 10 \text{ ns}$, $T_2 \sim 1 \text{ ns}$
- (Where is transmon mode for QD ?)
- Coupling strength $g \sim \sim 50 \text{ M}$



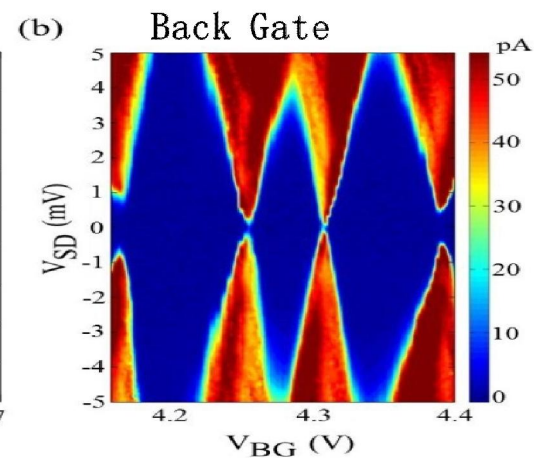
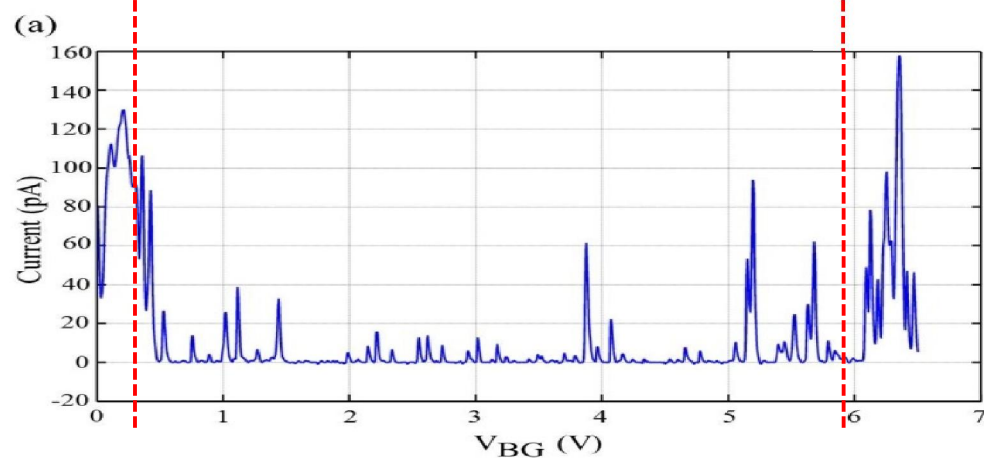
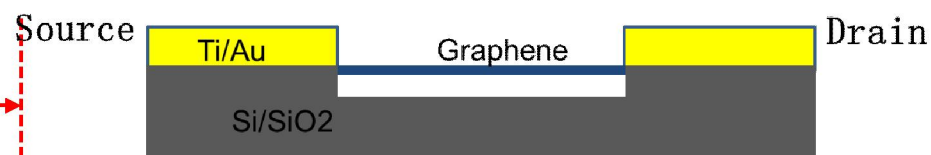
Quantum Information in dispersive region for QD CQED ?

G.-P. Guo, Phys. Rev. A 78, 020302 (2008). Phys. Rev. Lett. 101, 230501 (2008).

悬浮石墨烯量子点

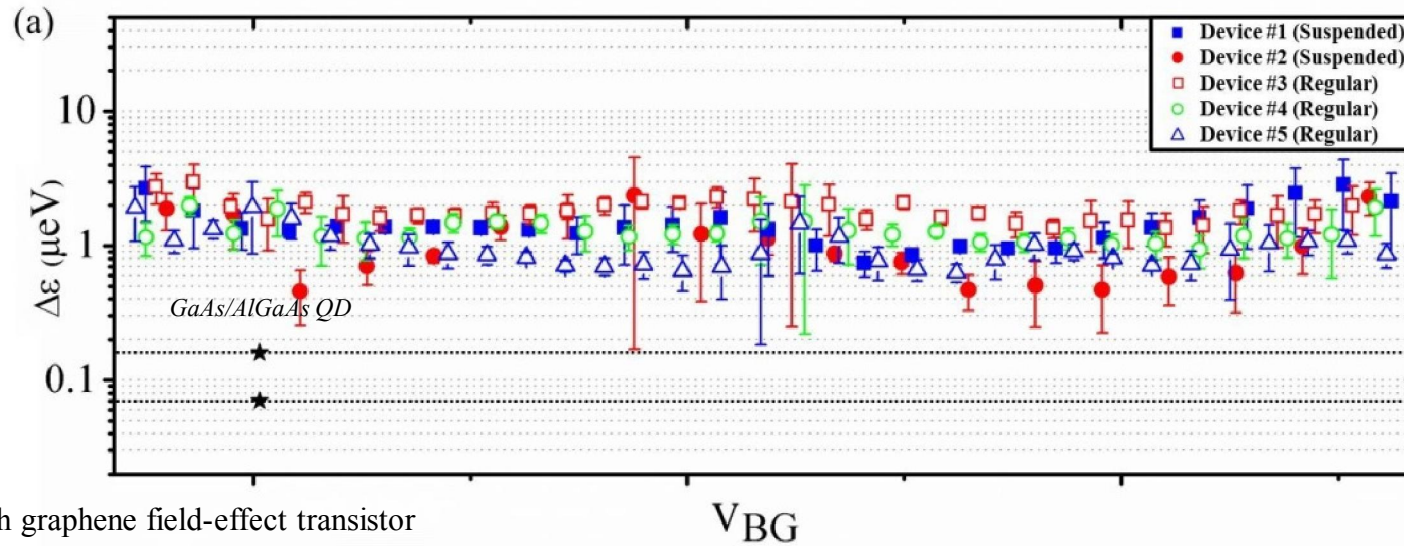


Floating graphene noise measurement

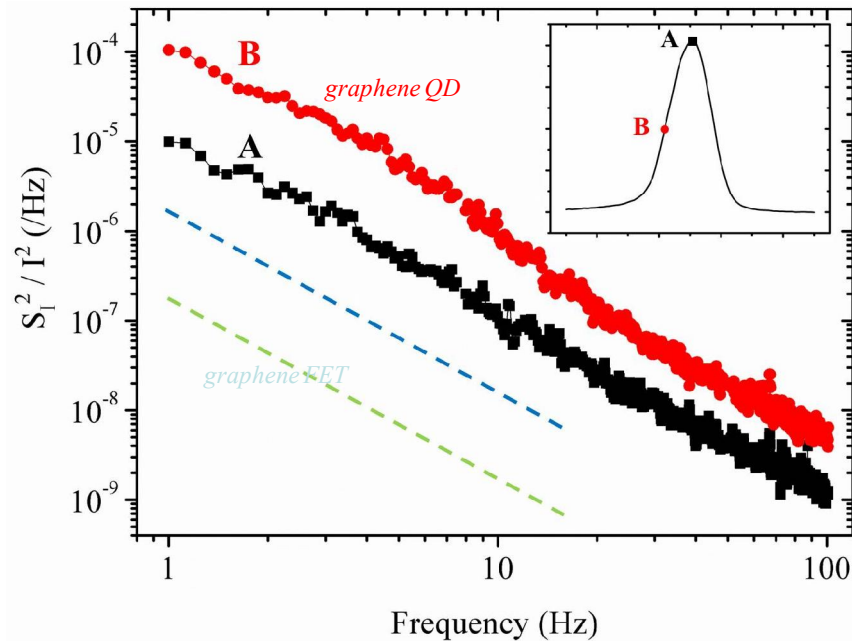


悬浮石墨烯量子点

Compare with GaAs/AlGaAs quantum dot

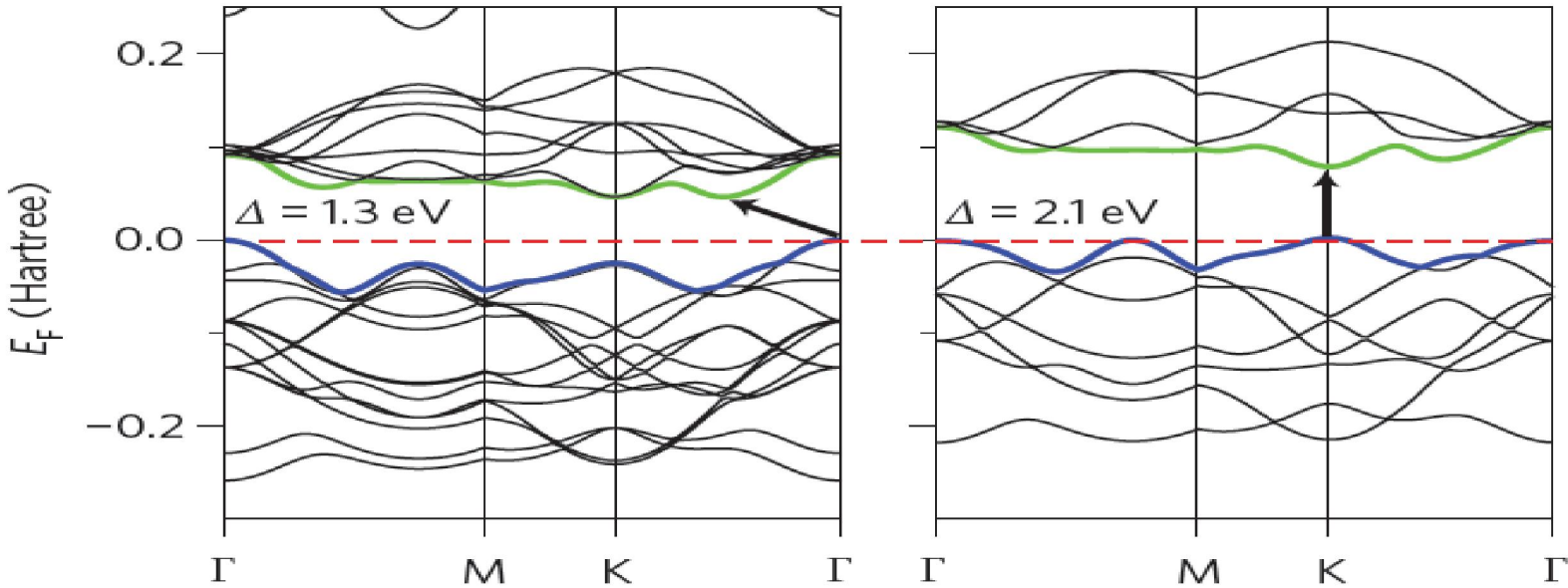
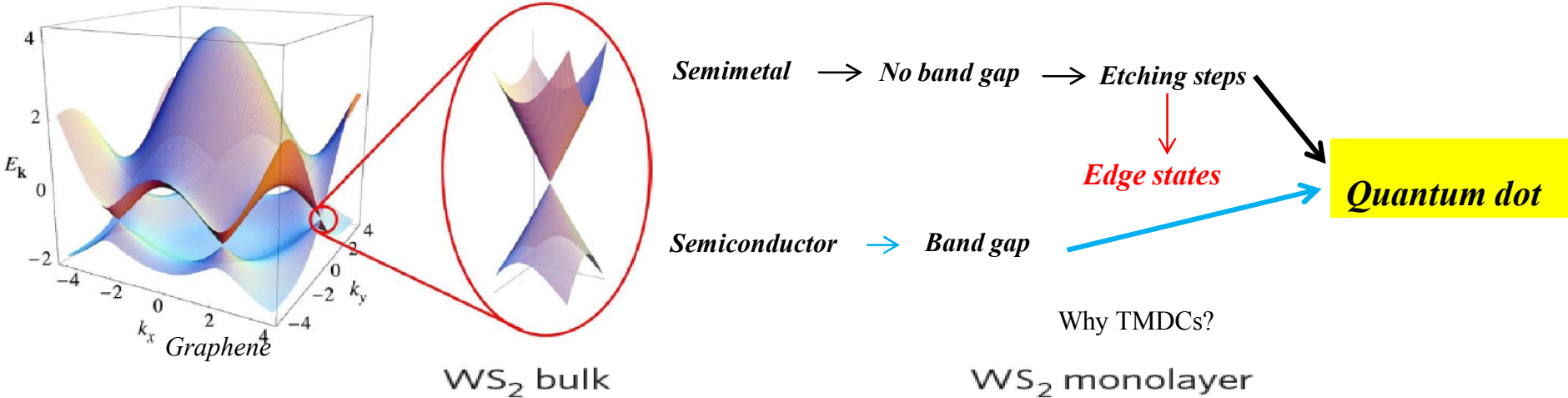


Compare with graphene field-effect transistor



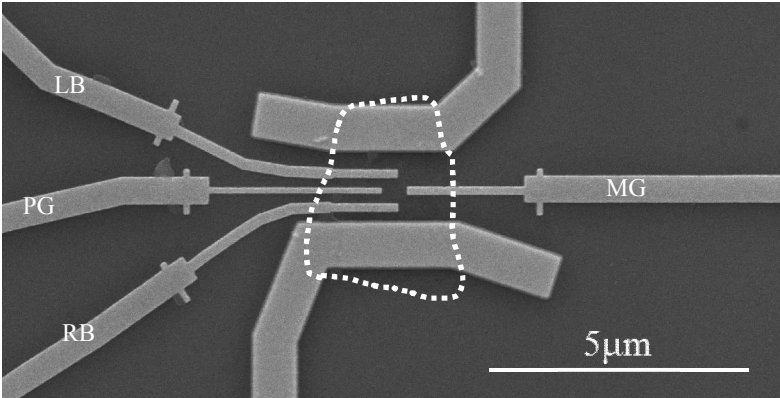
对于石墨烯纳米结构电噪声主要不是衬底应该是边缘态等缺陷

TMDs量子点

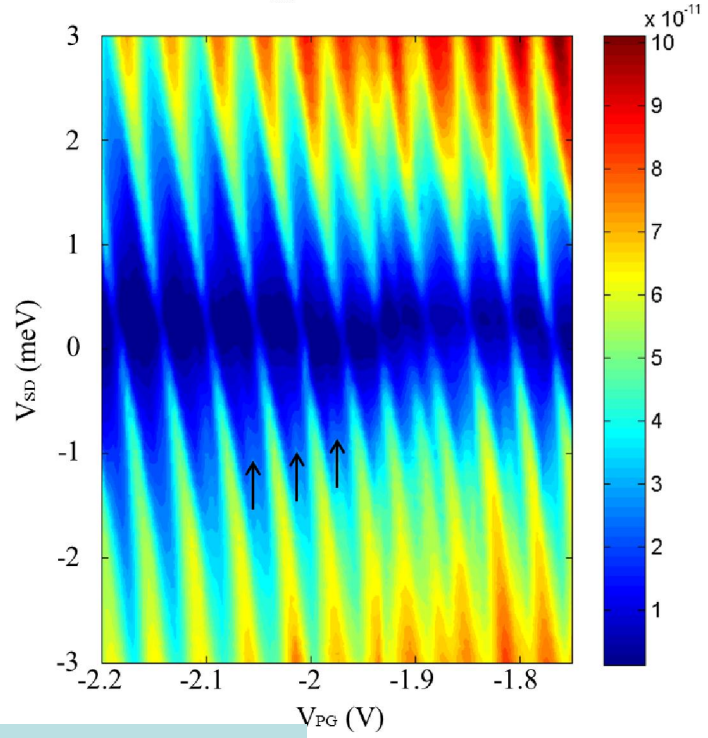
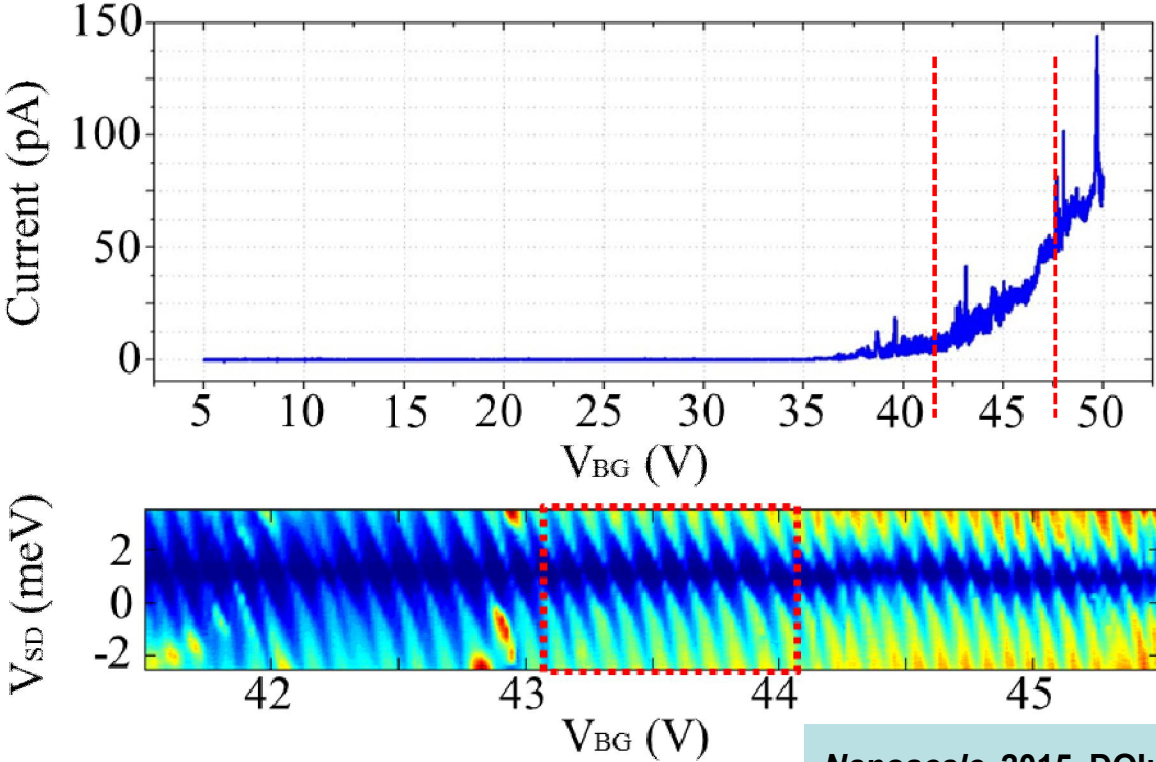
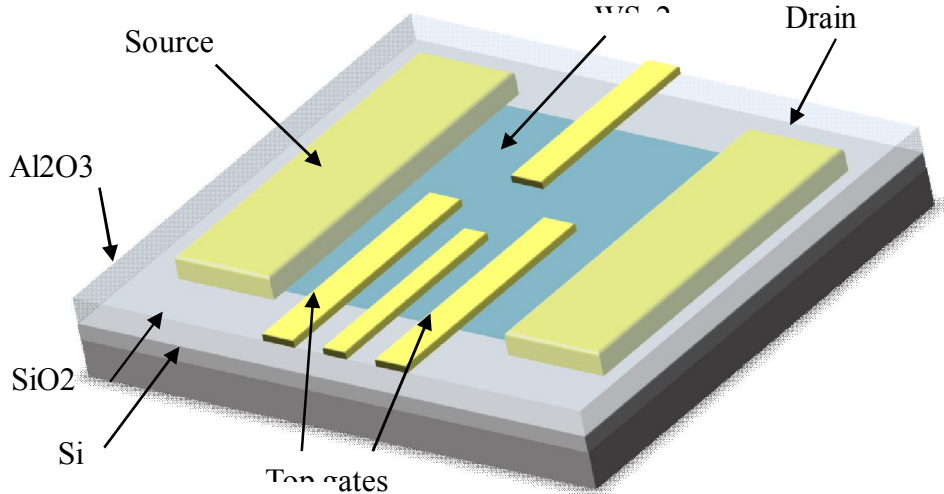


Transition metal dichalcogenides (TMDCs)

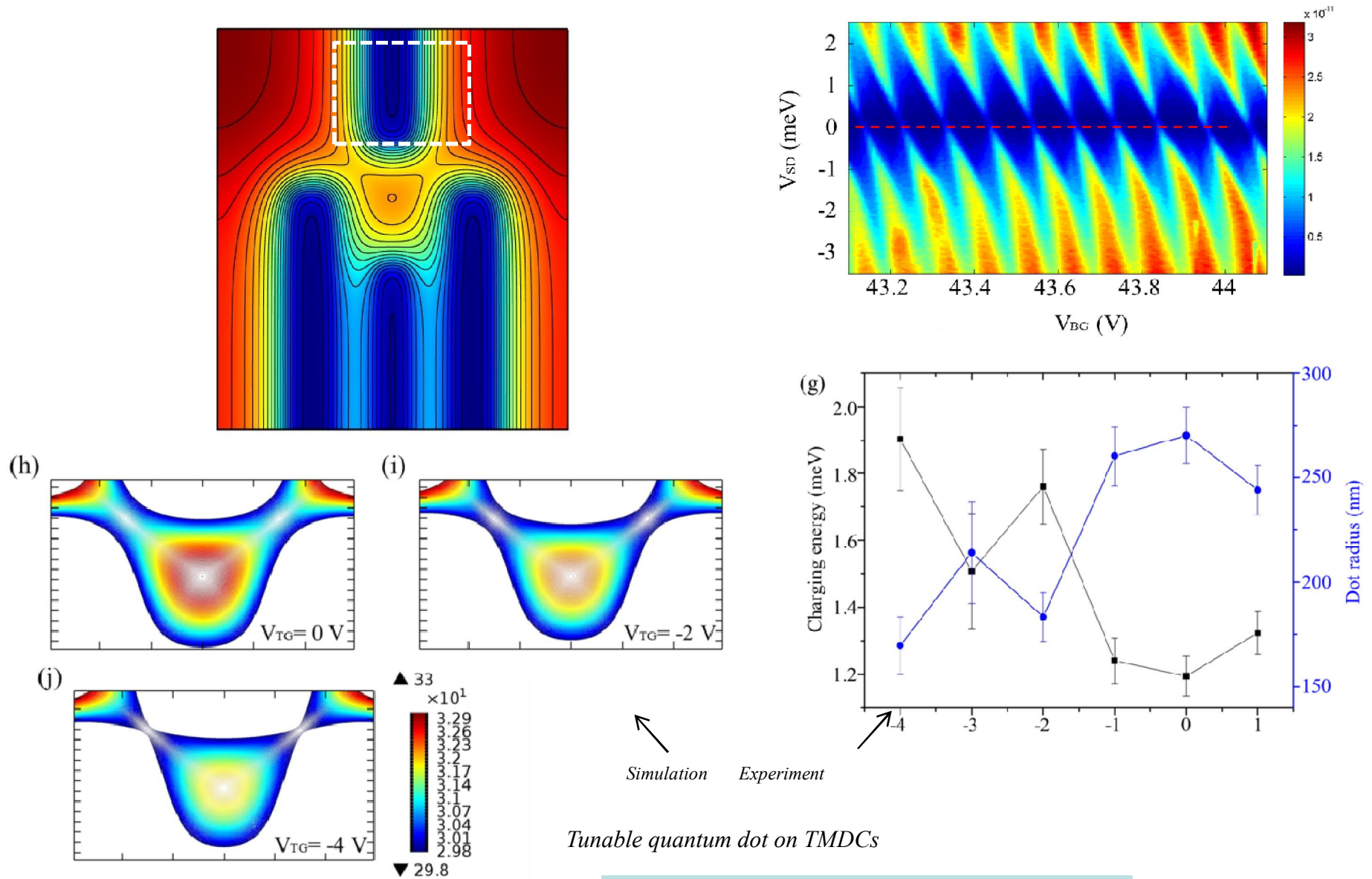
TMDs量子点



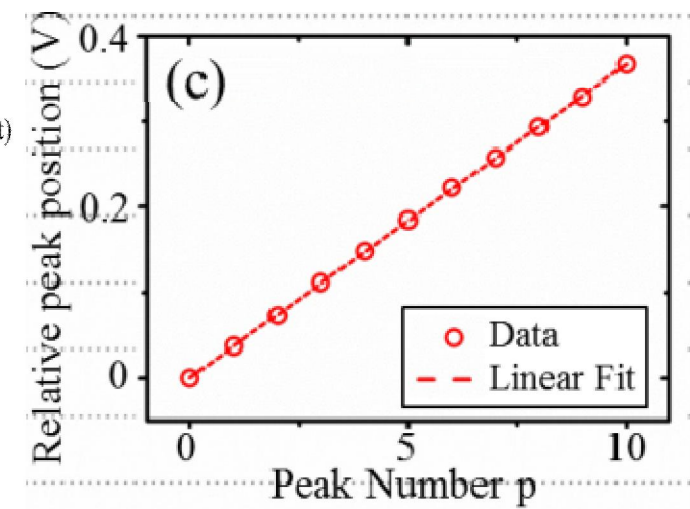
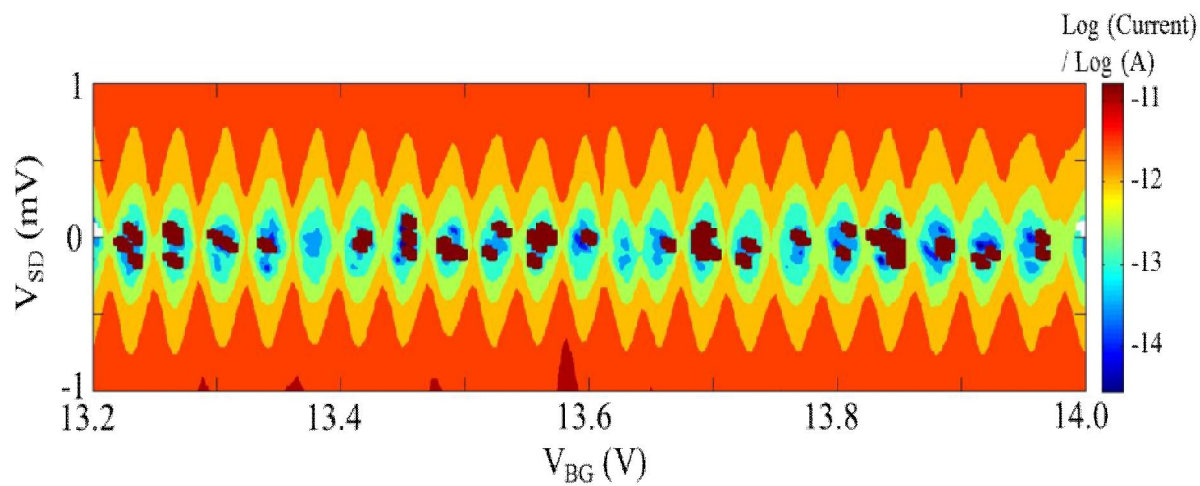
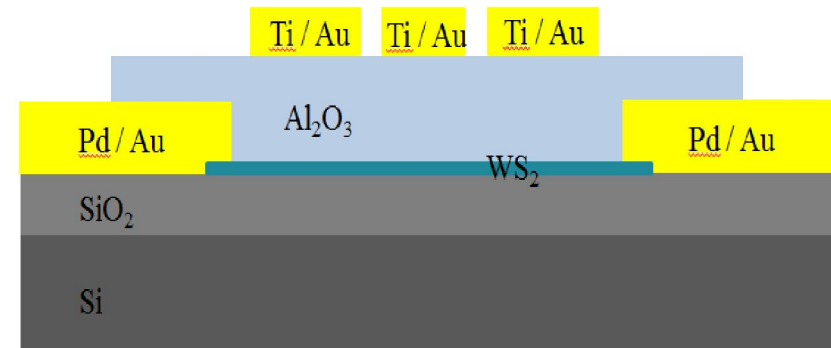
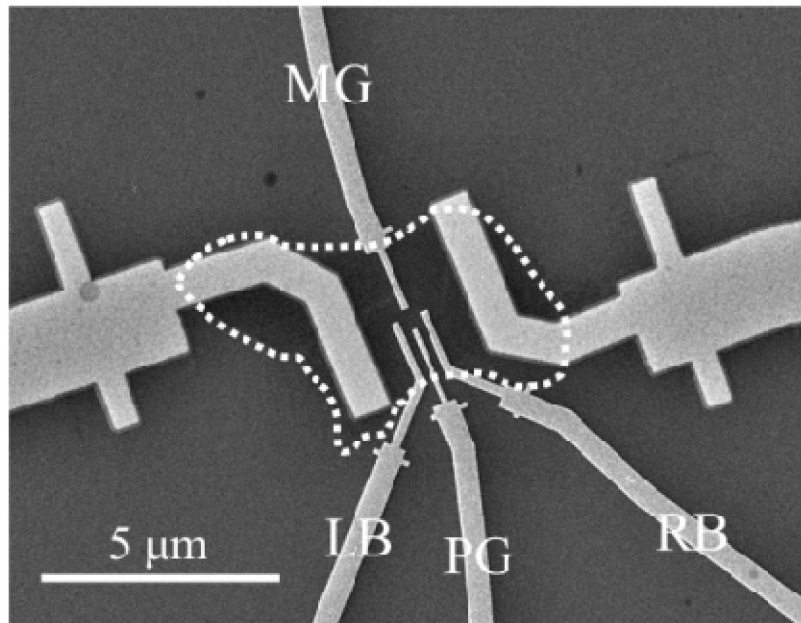
Quantum dot on WS₂



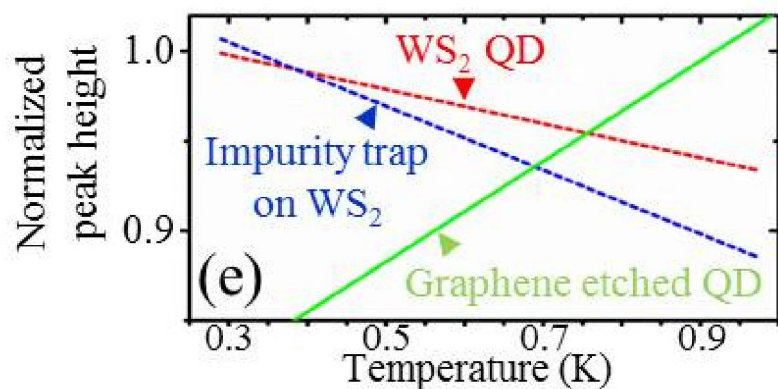
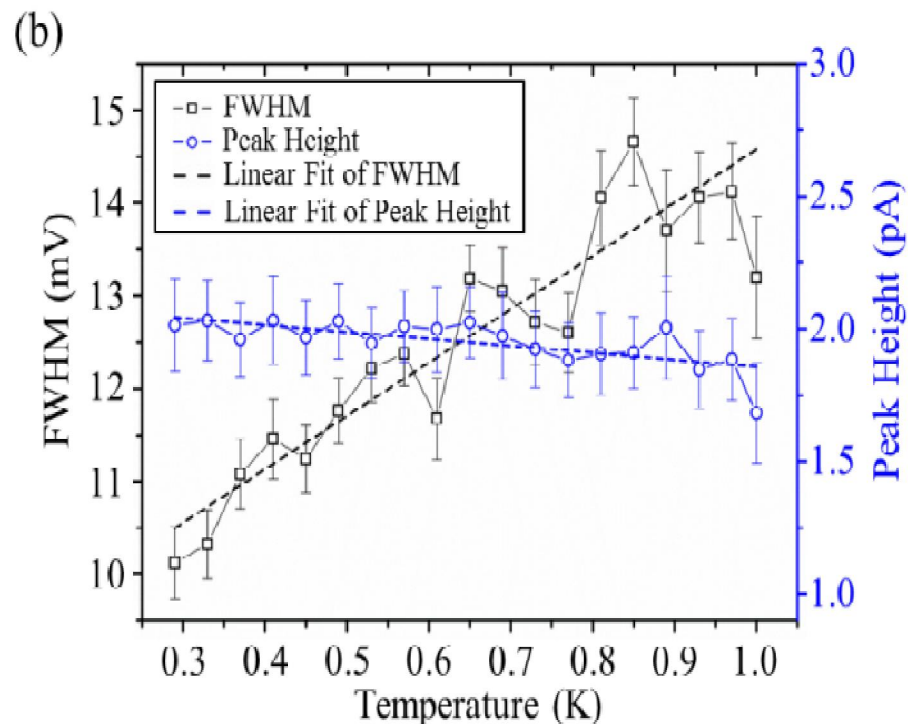
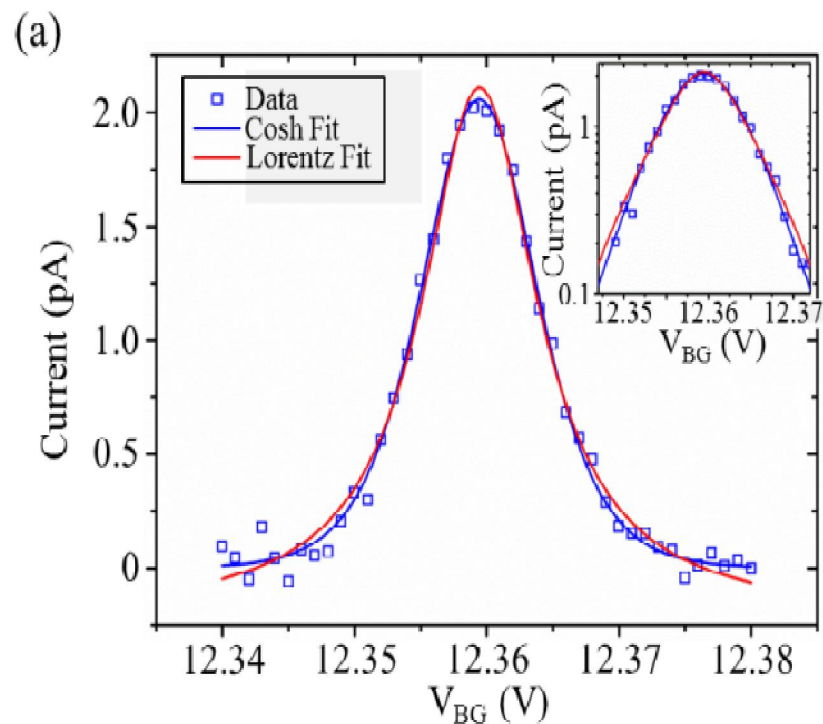
TMDs量子点



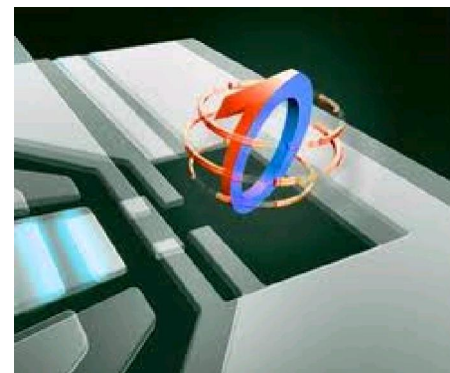
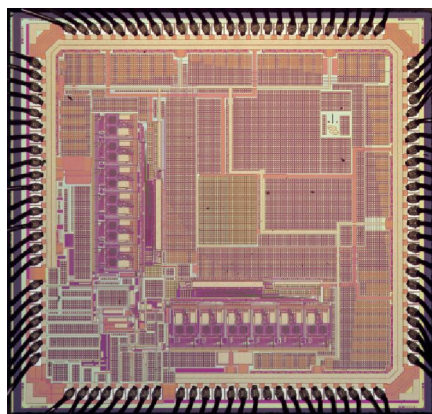
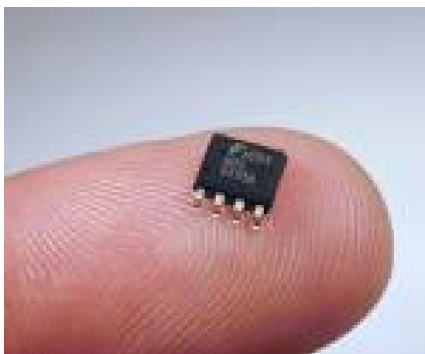
TMDs量子点



TMDs量子点



Disordered confining potential is absent in TMDCs quantum dot, which is different from graphene etched QD.



谢 谢!



中国科学院量子信息重点实验室

Key Lab of Quantum Information, CAS

量子计算与量子输运研究组