

Transport model study of conserved charge fluctuations in high temperature and high density QCD matter

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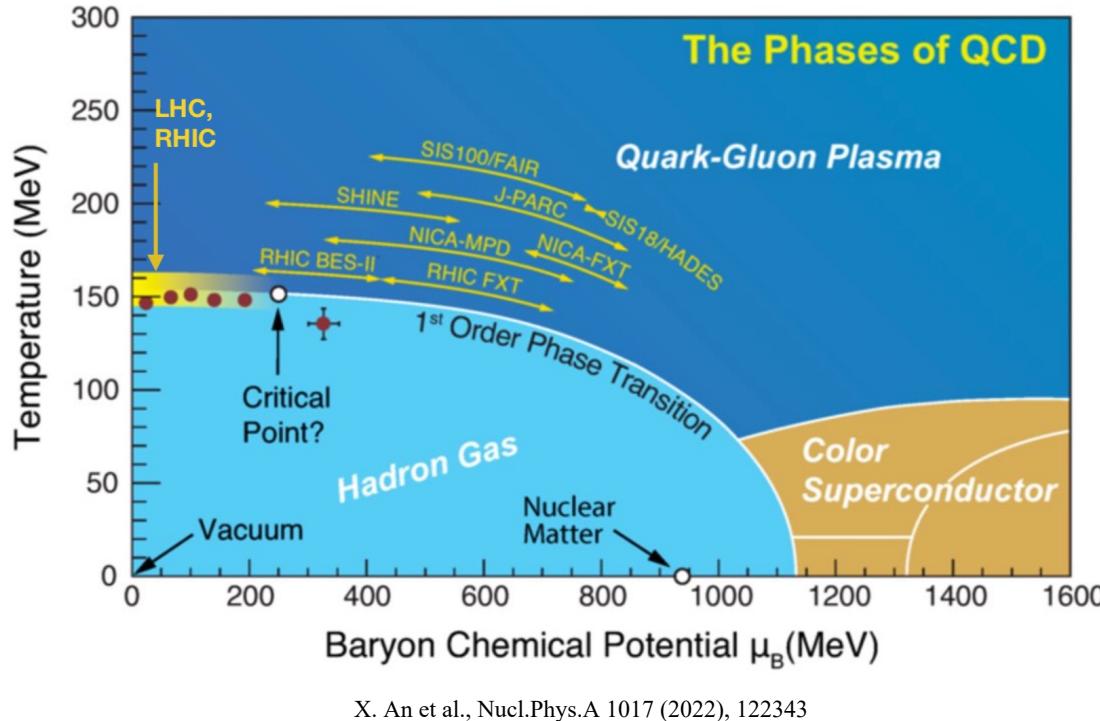
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QCD Phase Structure



- Small baryon chemical potential: Smooth Crossover Transition
- Large baryon chemical potential: First-order Phase Transition
- QCD Critical Endpoint: where the first-order phase transition ends(the key feature of QCD phase structure)
- Increase chemical potential by lowering the beam energy

Fluctuations of Conserved Charges (B, Q, S)

Conserved Charges : net-baryon (B)、 net-charge (Q)、 net-strangeness (S)

Moments and Cumulants: Variance (σ^2 、 C_2), skewness (S 、 C_3), kurtosis (κ 、 C_4)

Measured multiplicity N, $\langle \delta N \rangle = N - \langle N \rangle$

mean: $M = \langle N \rangle = C_1$

variance: $\sigma^2 = \langle (\delta N)^2 \rangle = C_2$

skewness: $S = \langle (\delta N)^3 \rangle / \sigma^3 = C_3 / C_2^{3/2}$

kurtosis: $\kappa = \langle (\delta N)^4 \rangle / \sigma^3 - 3 = C_4 / C_2^2$

Moments, cumulants and susceptibilities:

2nd order: $\sigma^2/M \equiv C_2/C_1 = \chi_2/\chi_1$

3rd order: $S\sigma \equiv C_3/C_2 = \chi_3/\chi_2$

4th order: $\kappa\sigma^2 \equiv C_4/C_2 = \chi_4/\chi_2$

The fluctuations of conserved charges are directly related to the system susceptibility (χ)

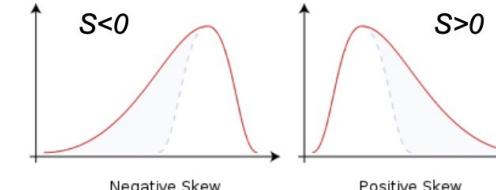
Cheng et al, PRD (2009) 074505. F. Karsch and K. Redlich, PLB 695, 136 (2011). B. Friman et al, EPJC 71 (2011) 1694. S. Gupta, X. Luo et al, Science, 332, 1525(2012). A. Bazavov et al, PRL109, 192302(12)/s. Borsanyi et al, PRL111, 062005(13)

$$\langle (\delta N)^4 \rangle \sim \xi^7$$

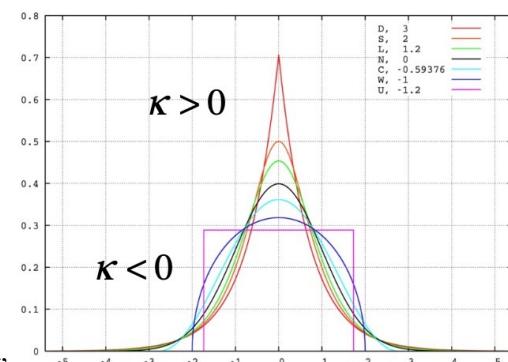
$$\langle (\delta N)^3 \rangle \sim \xi^{4.5}$$

The correlation length (ξ) of the critical point is divergent

Skewness (S) → asymmetry

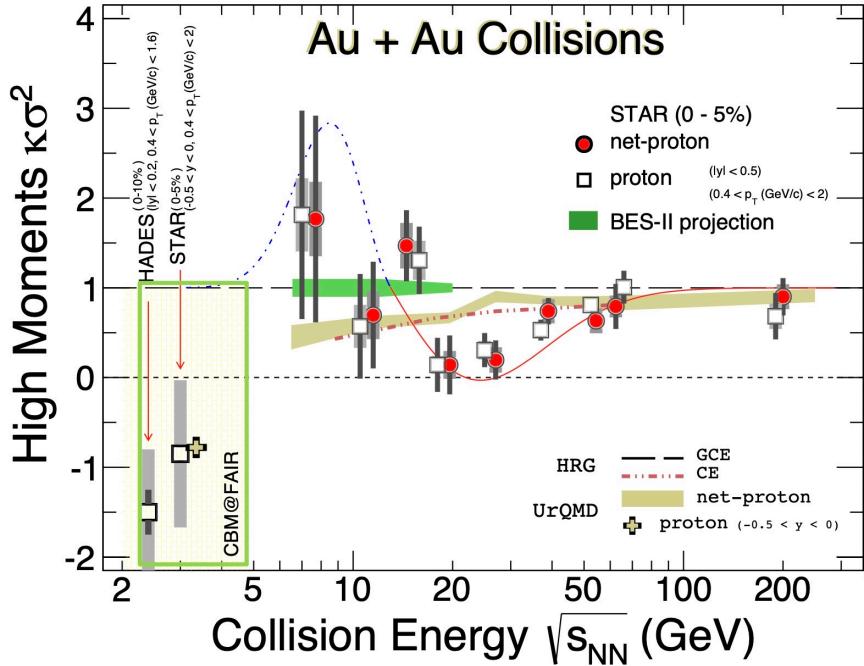
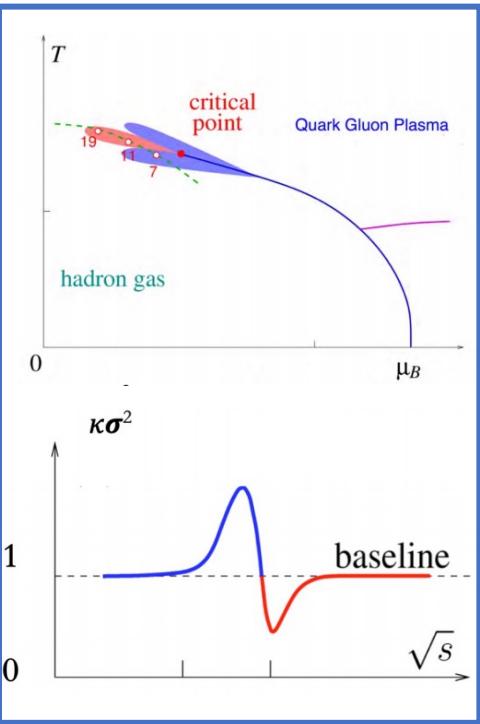


Kurtosis (κ) → Sharpness



Fluctuations of Conserved Charges (B)

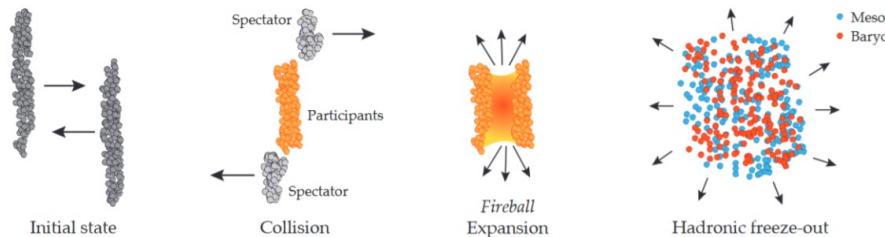
Fluctuations measured by STAR



- 1) Non-monotonic energy dependence:
hint of entering critical region.
- 2) 3 GeV proton high moments data:
Hadronic interaction dominant!
- 3) Energy gap between 3 and 7.7 GeV, important for
Critical Point search.

STAR:PRL126,92301(2021)PRL128,202303(2022)HADES: PRC102, 024914(2020)

Extended AMPT Model

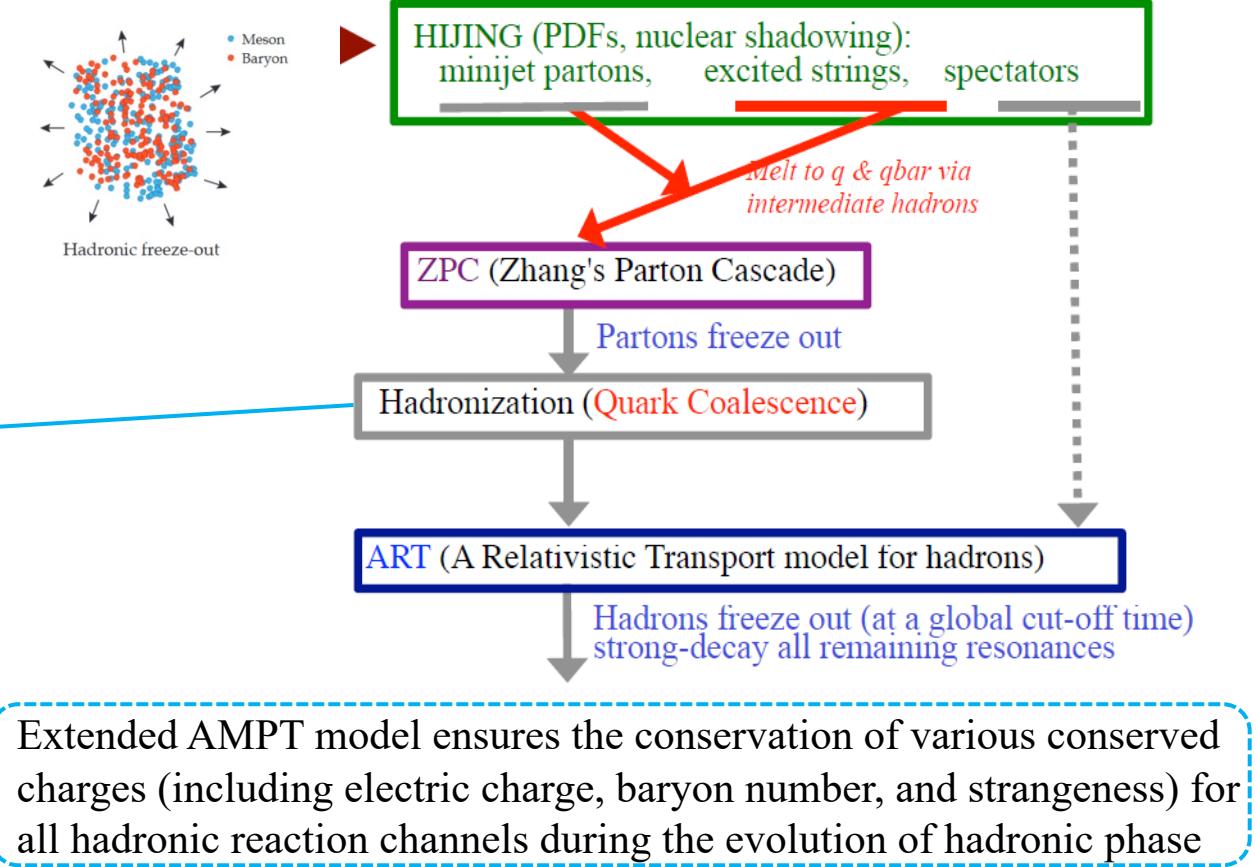


new quark coalescence:

quark to form either a meson or a baryon depending on the distance to its coalescence partner(s) (r_{BM})

$d_B < d_M * r_{BM}$: form a baryon
otherwise: form a meson

Y. He and Z.-W. Lin, Phys. Rev. C 96, 014910 (2017).



Fluctuations of Conserved Charges (B)

Cumulants:

$$C_1 = \kappa_1,$$

$$C_2 = \kappa_2 + \kappa_1,$$

$$C_3 = \kappa_3 + 3\kappa_2 + \kappa_1,$$

$$C_4 = \kappa_4 + 6\kappa_3 + 7\kappa_2 + \kappa_1.$$

Only for one particle! !



Correlation Functions:

$$\kappa_1 = C_1 = \langle N \rangle,$$

$$\kappa_2 = -C_1 + C_2,$$

$$\kappa_3 = 2C_1 - 3C_2 + C_3,$$

$$\kappa_4 = -6C_1 + 11C_2 - 6C_3 + C_4.$$

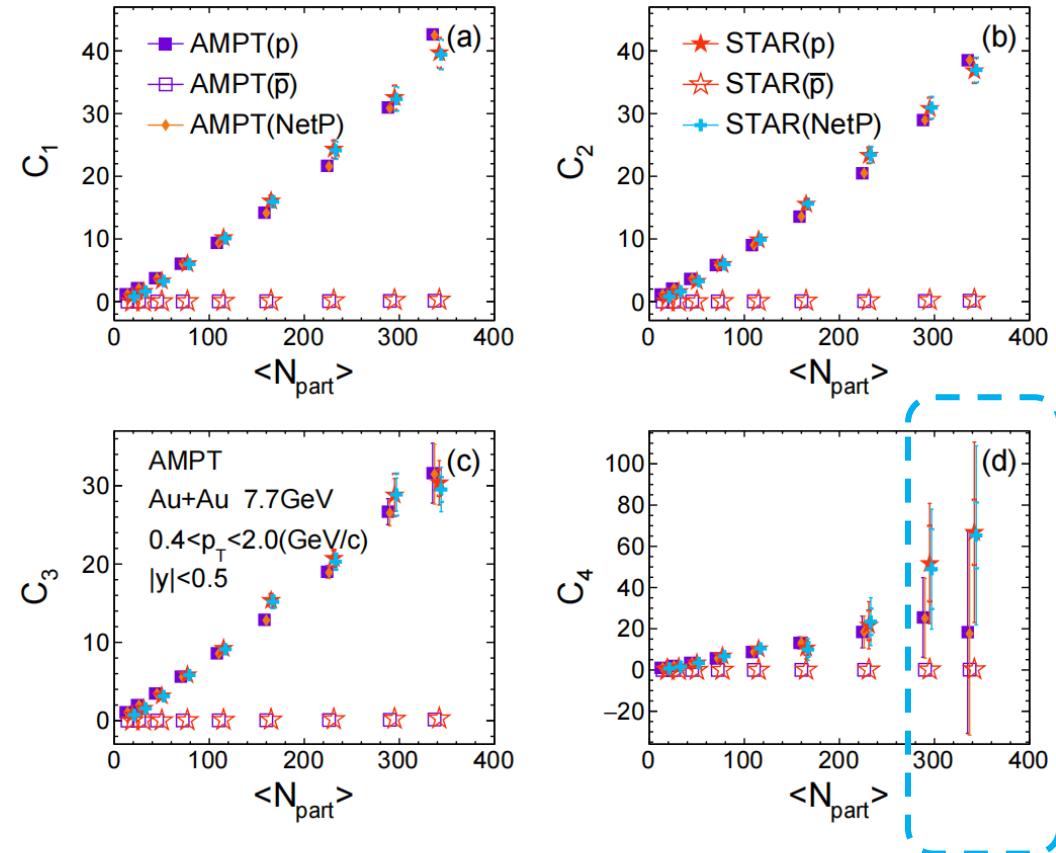
Factorial moments:

$$F_3 = \int dy_1 dy_2 dy_3 \rho_3(y_1, y_2, y_3) = F_1^3 + 3F_1 C_2 + C_3$$

$$\begin{aligned} \rho_3(y_1, y_2, y_3) = & \rho_1(y_1)\rho_1(y_2)\rho_1(y_3) + \rho_1(y_1)\underline{C_2(y_2, y_3)} \\ & + \rho_1(y_2)\underline{C_2(y_1, y_3)} + \rho_1(y_3)\underline{C_2(y_1, y_2)} \\ & + \underline{C_3(y_1, y_2, y_3)} \end{aligned}$$

Bzdak, Adam et al. Phys.Rev. C95 (2017) 5, 054906.

AMPT Results On Proton cumulants

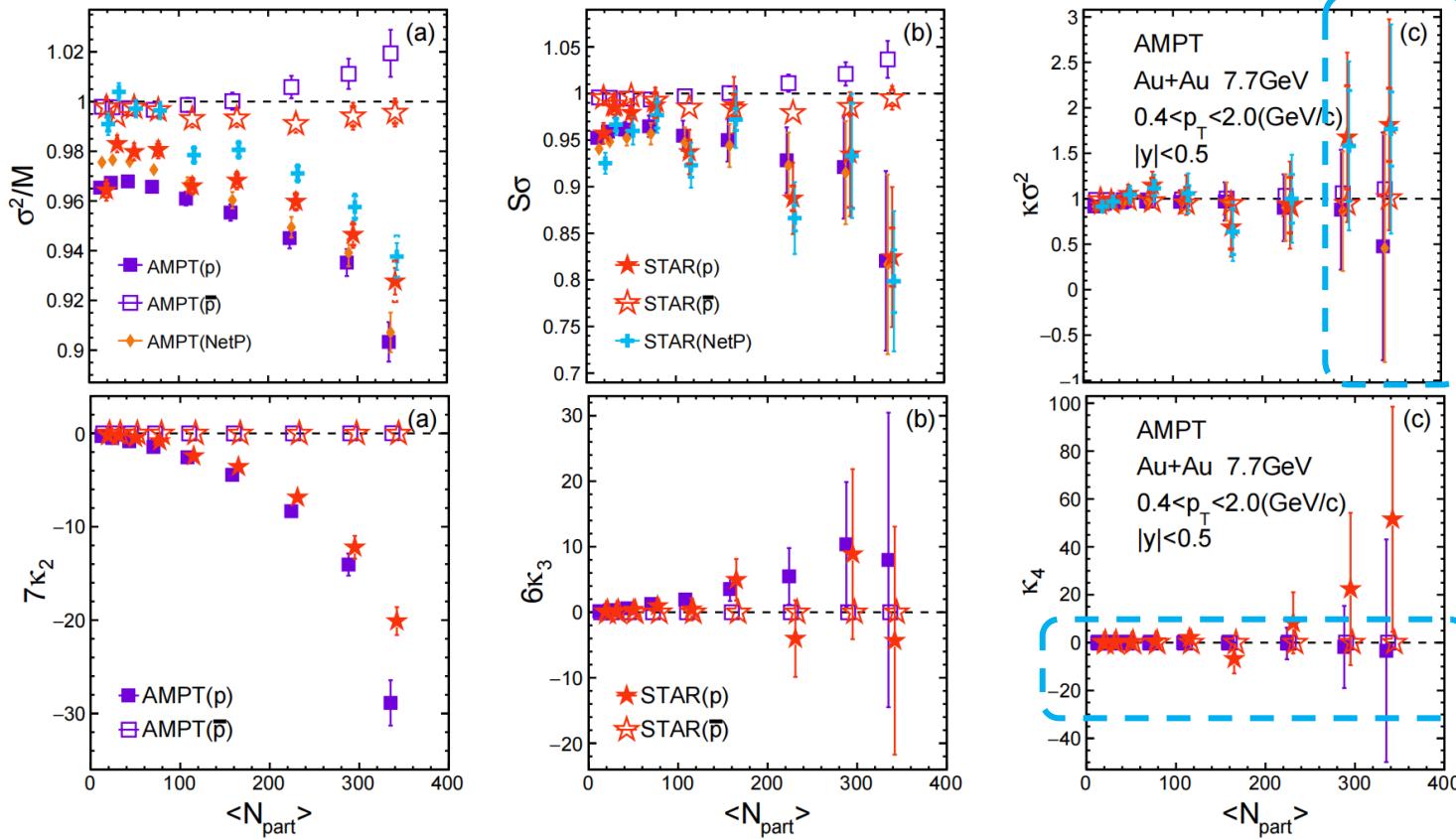


- The cumulants C_n for protons, antiprotons, and net-protons all show a similar increasing dependence on $\langle N_{part} \rangle$
- In the **0-5%** and **5-10%** centrality ranges, the fourth-order cumulant (C_4) in the AMPT model notably **underestimates** STAR's results

Qian Chen, Guo-Liang Ma, Phys.Rev.C 106 (2022) 014907

AMPT Results On Proton Cumulant Ratios And Correlations

Qian Chen, Guo-Liang Ma, Phys.Rev.C 106 (2022) 014907

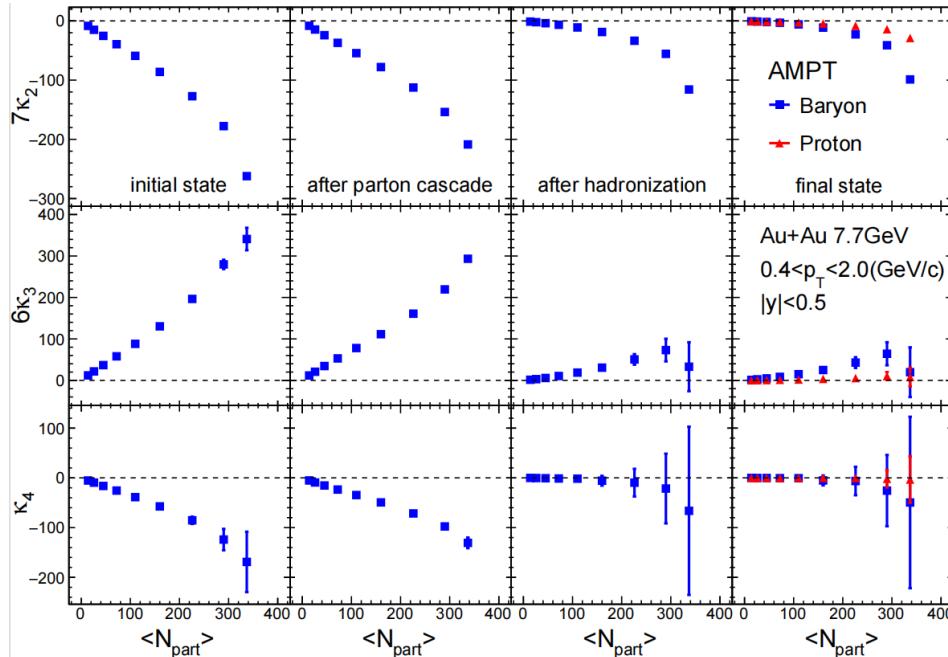


The four-proton correlation from AMPT is very small, consistent with zero.

Baryon Number Conservation

Expectation of baryon number conservation:

$$P(N) = \frac{B!}{N!(B-N)!} p^N (1-p)^{(B-N)}$$



Qian Chen, Guo-Liang Ma, Phys.Rev.C 106 (2022) 014907

n-baryon correlations:

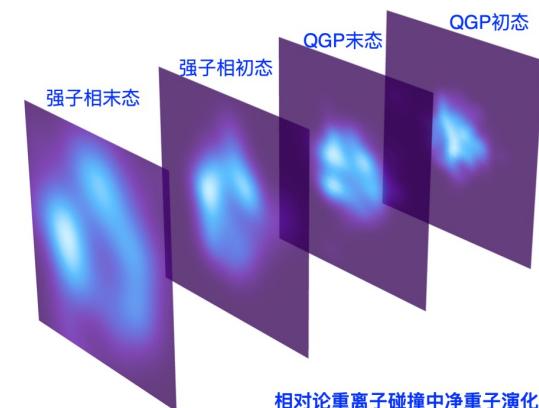
$$\kappa_1 = \langle N \rangle = pB$$

$$\kappa_2 = -\frac{\langle N \rangle^2}{B}$$

$$\kappa_3 = 2 \frac{\langle N \rangle^3}{B^2}$$

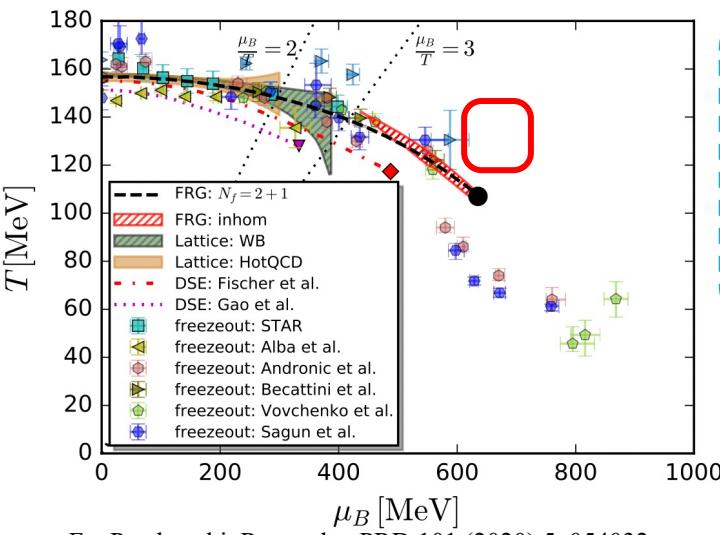
$$\kappa_4 = -6 \frac{\langle N \rangle^4}{B^3}$$

➤ Multi-baryon correlations are getting weaker with stage evolution of heavy-ion collisions



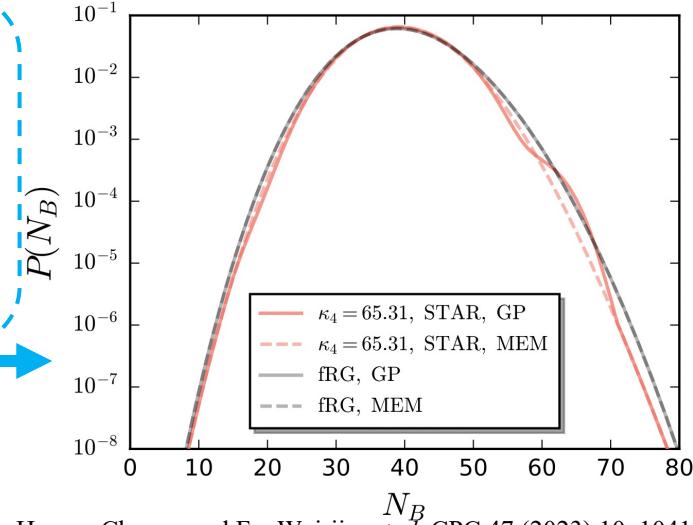
相对论重离子碰撞中净重子演化

Functional Renormalization Group



Fu, Pawłowski, Rennecke, PRD 101 (2020) 5, 054032

the baryon number distributions are reconstructed from the cumulants of different orders by means of the maximum entropy method



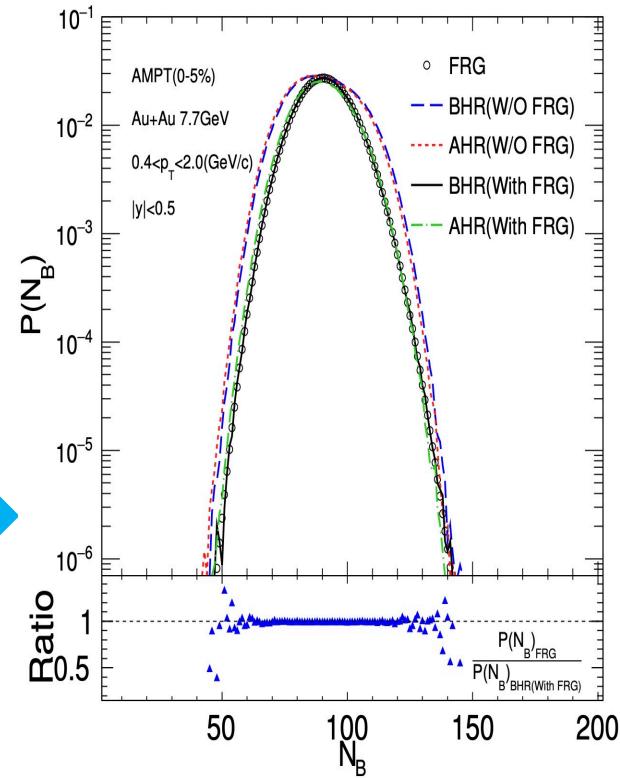
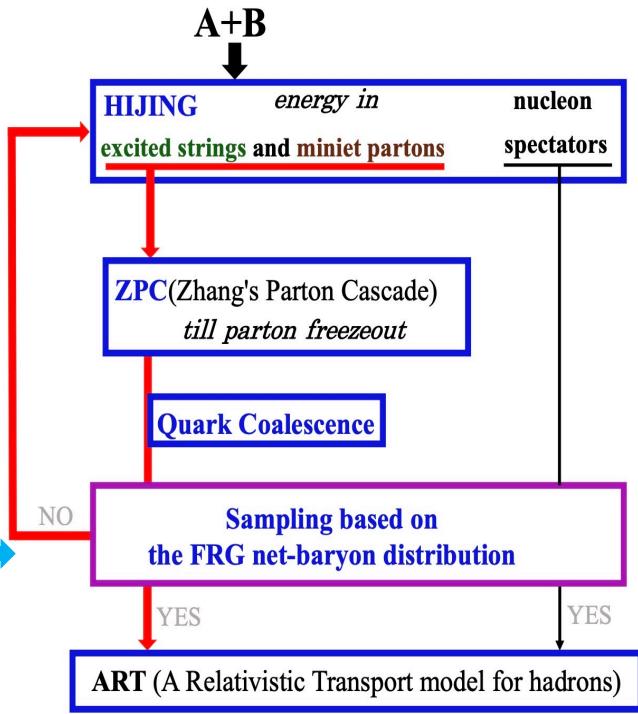
Huang, Chuang and Fu, Wei-jie, et al. CPC 47 (2023) 10, 104106

FRG enables the study of equations of state at both high and low baryon chemical potentials.

FRG with critical fluctuations mechanism without interactions between hadrons and decay processes

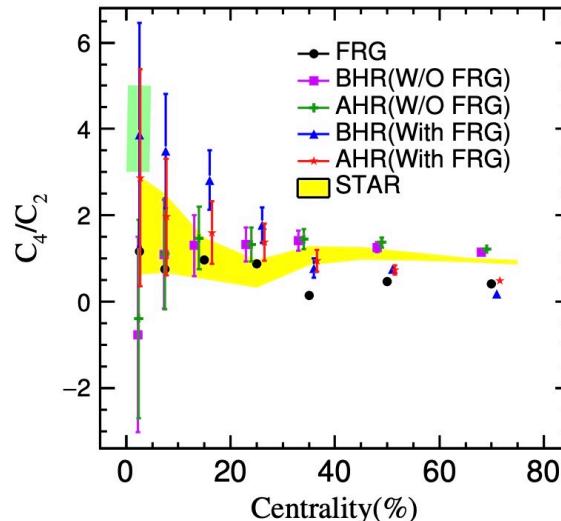
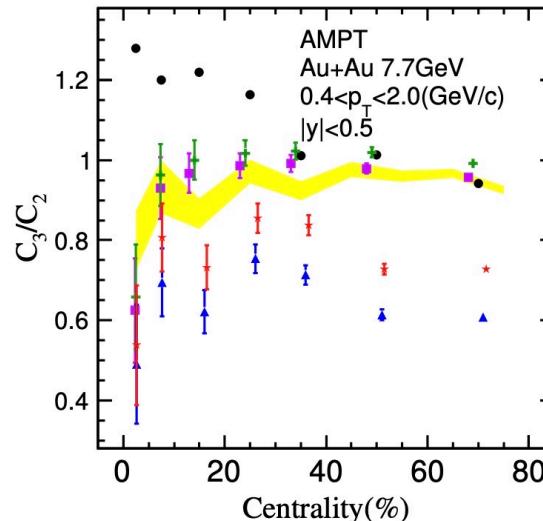
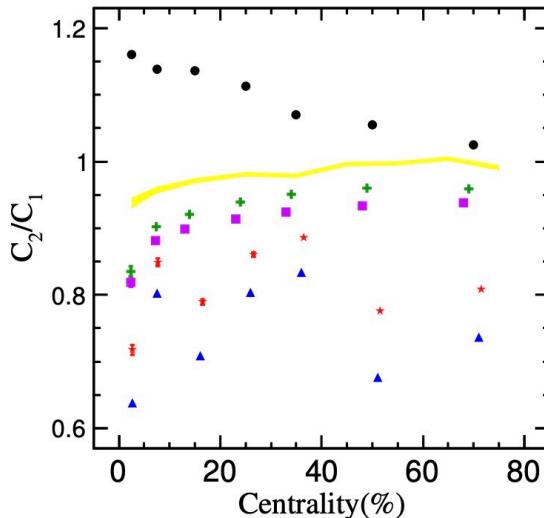
Incorporating FRG Into AMPT Model

FRG parameter input:
baryon chemical :
 $\mu_B = 399 \text{ MeV}$;
volume of the fire ball :
 $V = 980 \text{ fm}^3$;
pseudo-critical temperature :
 $T = 139 \text{ MeV}$;



Qian Chen, Rui Wen, Shi Yin, Wei-jie Fu, Zi-Wei Lin, and Guo-Liang Ma. arXiv:2402.12823.

Incorporating FRG Into AMPT Model



- The influences of hadronic rescatterings :increase the cumulant ratios C_2/C_1 and C_3/C_2 , but decrease C_4/C_2 in central collisions and increase C_4/C_2 in peripheral collisions.
- The effect of hadronic rescatterings is more significant for **critical fluctuations** than **dynamical fluctuations**.

Qian Chen, Rui Wen, Shi Yin, Wei-jie Fu, Zi-Wei Lin, and Guo-Liang Ma. arXiv:2402.12823.

Summary and Outlook

Summary:

- The AMPT results are consistent with the expectation from baryon number conservation.
- The incorporation of the FRG into the AMPT model reveals that the hadronic rescatterings process affects different orders of net-baryon cumulant ratios.

Outlook:

- ◆ Incorporation of critical fluctuation physics into AMPT : FRG、density fluctuations.
- ◆ nuclear thickness effects, coalescence mechanisms, different collision systems, effects of magnetic fields, ...

Thank you for your attentions !

Back Up

- In the old version, only K^+ and K^- were introduced in hadron rescatterings as explicit particles, but K^0 and \bar{K}^0 were omitted.



- In the old version, some isospin-averaged cross sections were used, and the charge of the final state particles is chosen randomly from all possible charges, independent of the total charge of the initial state.

For example:

- 1) $\pi^+ + \pi^+ \rightarrow \rho^+ + \rho^+ \quad \checkmark$
- 2) $\pi^+ + \pi^+ \rightarrow \rho^+ + \rho^- \quad \times$
- 3) $\pi^+ + \pi^+ \rightarrow \rho^- + \rho^- \quad \times$

Back Up

Cumulants:

$$C_2 = \langle N \rangle + \langle \bar{N} \rangle + \kappa_2^{(2,0)} + \kappa_2^{(0,2)} - 2\kappa_2^{(1,1)}$$

$$C_3 = \langle N \rangle - \langle \bar{N} \rangle + 3\kappa_2^{(2,0)} - 3\kappa_2^{(0,2)} + \kappa_3^{(3,0)} - \kappa_3^{(0,3)} - 3\kappa_3^{(2,1)} + 3\kappa_3^{(1,2)}$$

$$C_4 = \langle N \rangle + \langle \bar{N} \rangle + 7\kappa_2^{(2,0)} + 7\kappa_2^{(0,2)} - 2\kappa_2^{(1,1)} + 6\kappa_3^{(3,0)} + 6\kappa_3^{(0,3)} - 6\kappa_3^{(2,1)}$$

$$- 6\kappa_3^{(1,2)} + \kappa_4^{(4,0)} + \kappa_4^{(0,4)} - 4\kappa_4^{(3,1)} - 4\kappa_4^{(1,3)} + 6\kappa_4^{(2,2)}$$

Bzdak, Adam et al. Phys.Rev. C86 (2012) 044904

two or more kinds of particles ! ! !

Factorial moments:

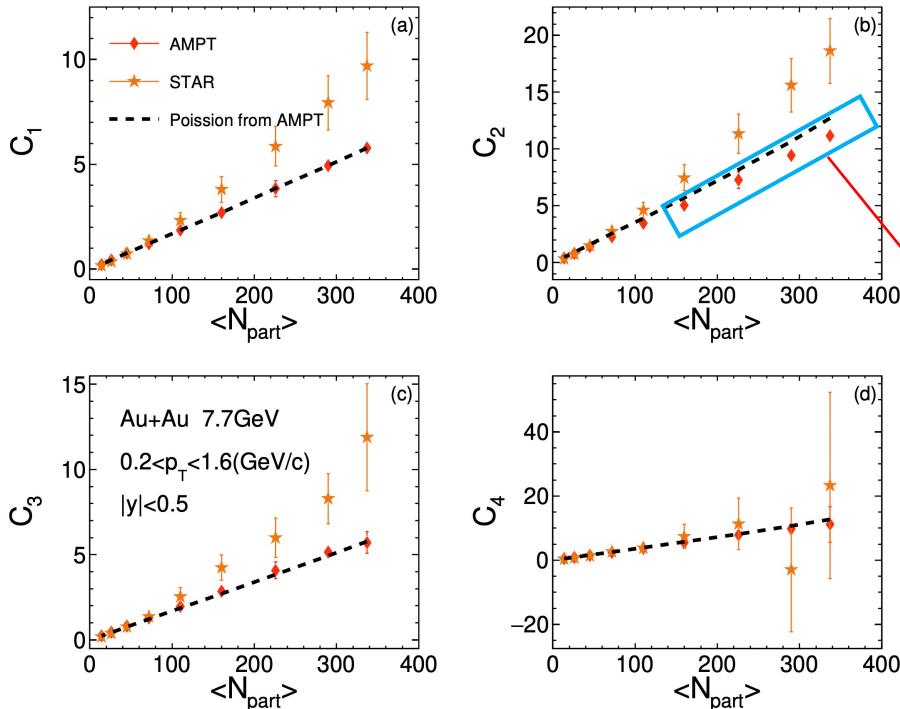
$$F_{i,k} = \left\langle \frac{N!}{(N-i)!} \frac{\bar{N}!}{(\bar{N}-k)!} \right\rangle = \frac{d^i}{dz^i} \frac{d^k}{d\bar{z}^k} H(z, \bar{z})|_{z=\bar{z}=1}$$

Correlation Functions:

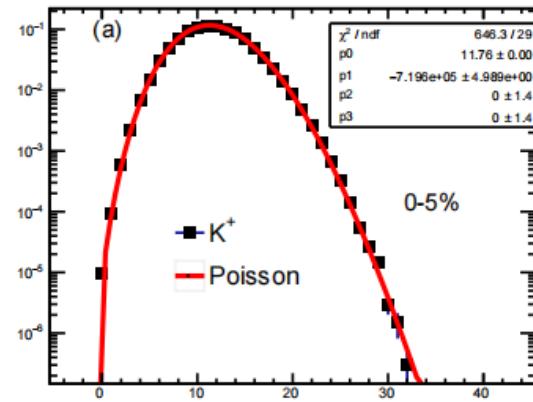
$$\kappa_2^{(2,0)} = -F_{1,0}^2 + F_{2,0},$$

$$\kappa_2^{(1,1)} = -F_{1,0}F_{0,1} + F_{1,1},$$

Back Up

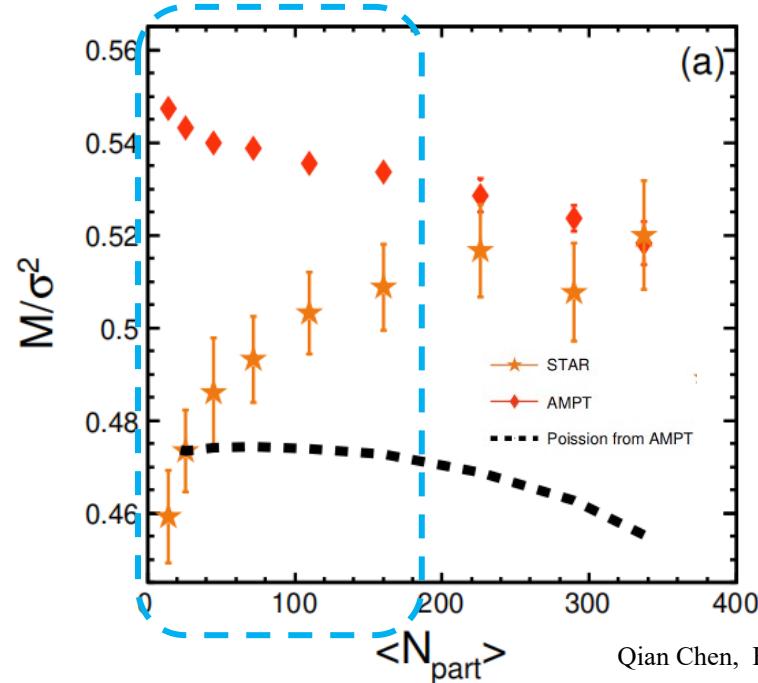


Qian Chen, Han-Sheng Wang, Guo-Liang Ma, Phys.Rev.C 107 (2023) 034910



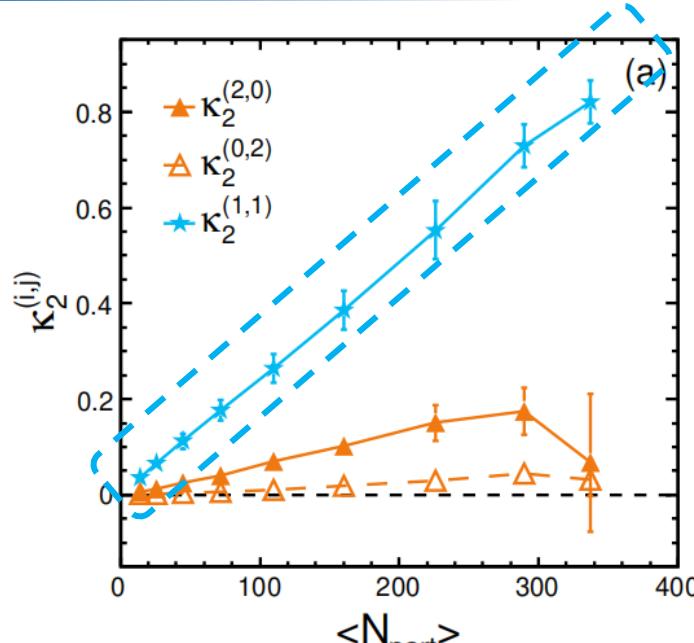
The C_2 for AMPT is slightly lower than Poisson baseline based on its mean multiplicity, suggesting a correlation between K^+ and K^-

Back Up



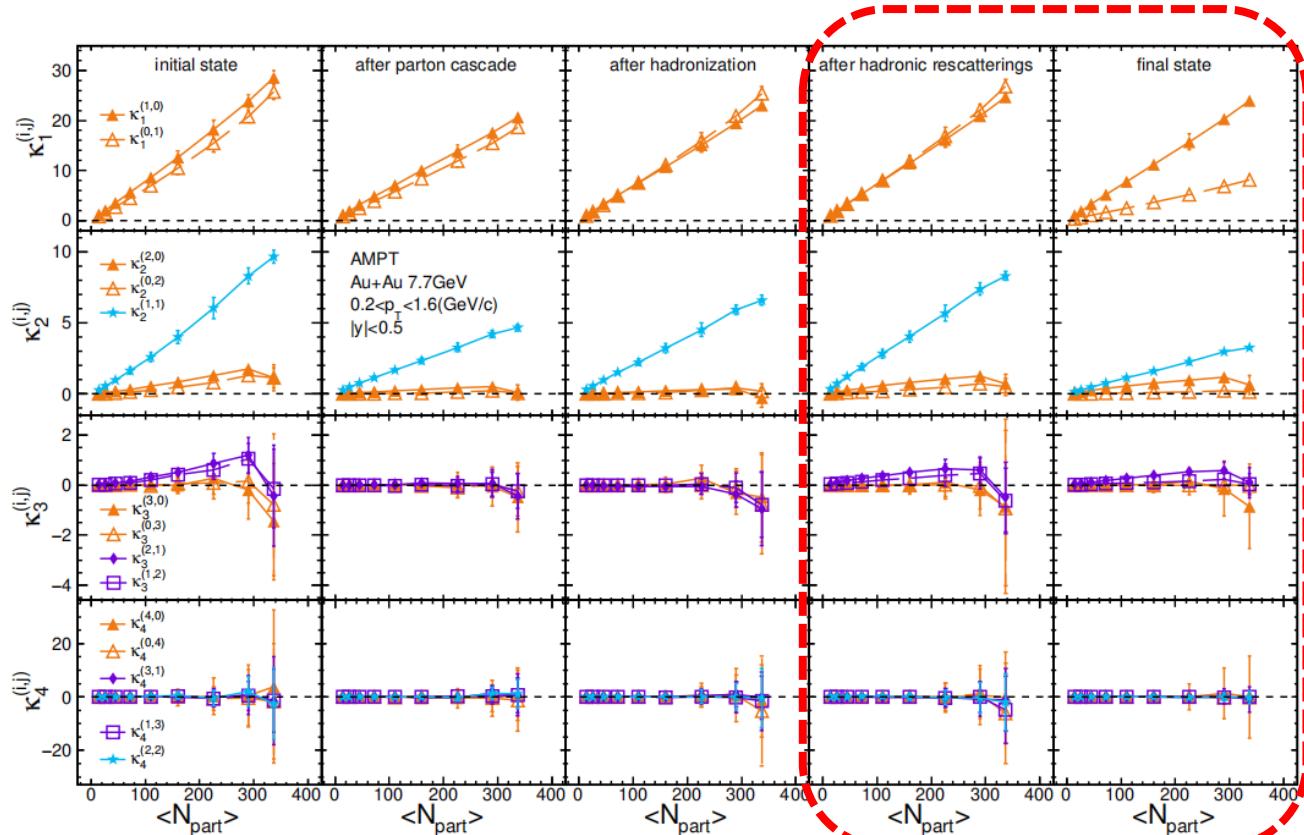
Caused by the new quarks coalescence mechanism

Qian Chen, Han-Sheng Wang, Guo-Liang Ma, Phys.Rev.C 107 (2023) 034910



Two-particle correlation function between the K^+ and K^- [$\kappa_2^{(1,1)}$] is dominants——pair production

Back Up



➤ The fluctuations of strangeness are notably influenced during the weak decay evolution stage