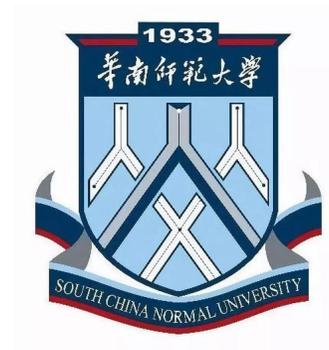


# Information field based global Bayesian inference of the jet transport coefficient

**Man Xie (解曼)**

Collaborate with Wei-Yao Ke, Han-Zhong Zhang and Xin-Nian Wang

*[PRC 108, L011901 (2023); arXiv: 2208.14419]*



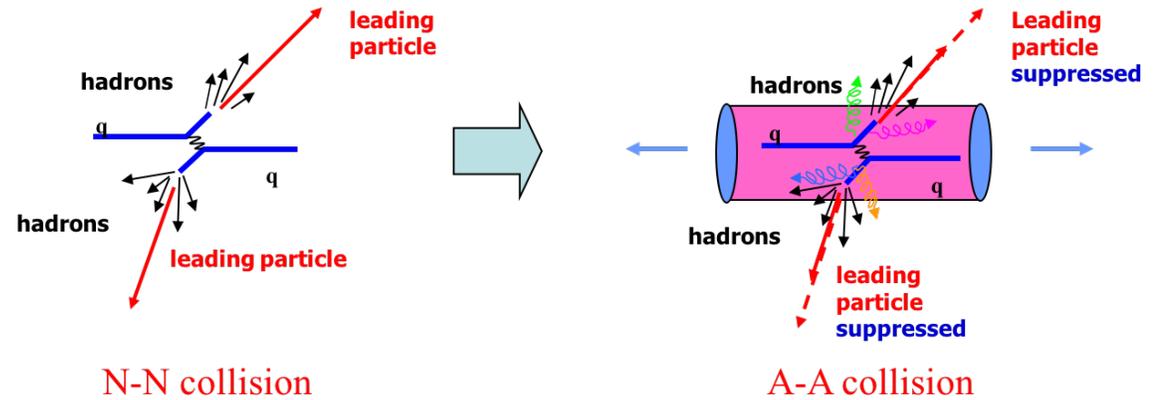
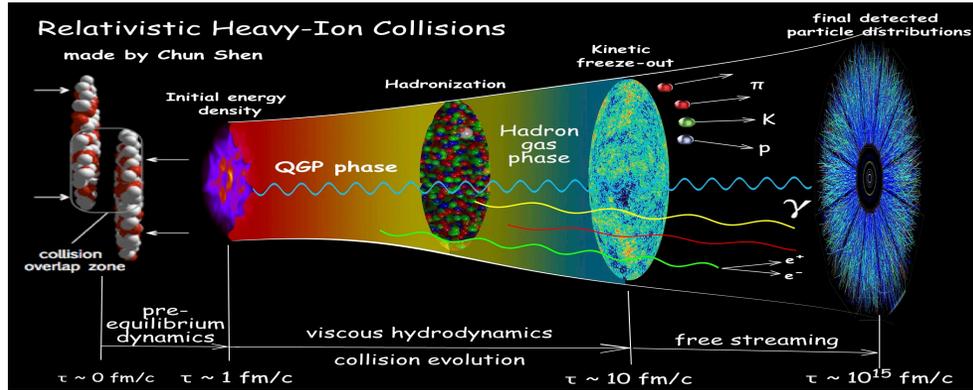
2024/5/17

# Outline

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- Motivation
- NLO parton model with energy loss
- Bayesian inference of  $\hat{q}(T)$  with information field (IF)
  - IF-assisted sensitivity analysis
  - Global Bayesian inference of  $\hat{q}(T)$
  - Validation and prediction
- Summary

# Motivation

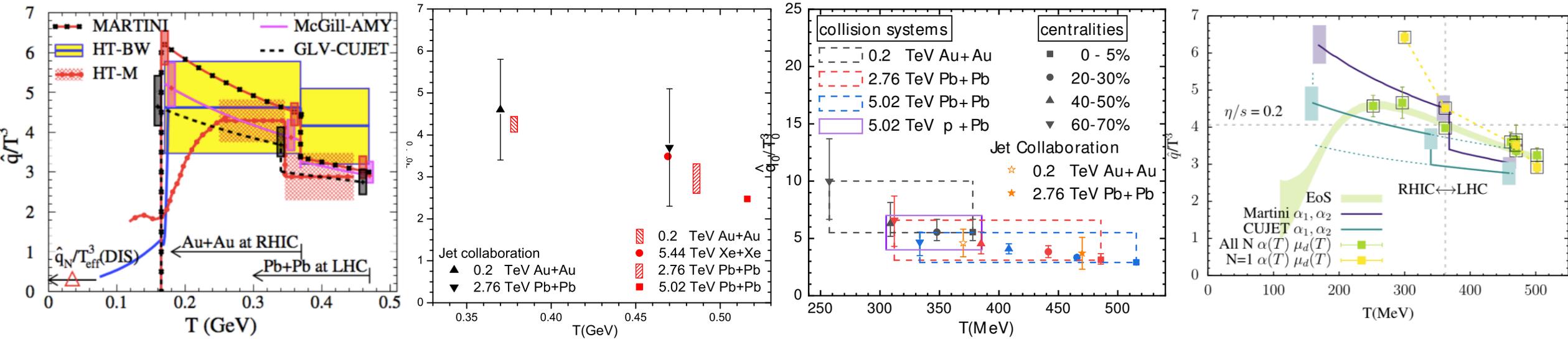


- Quark-gluon plasma (QGP) has been created in high-energy HIC.
- Jet quenching is an extremely useful tool to explore the properties of QGP. [X.-N. Wang and M. Gyulassy, PRL 68, 1480 (1992)]

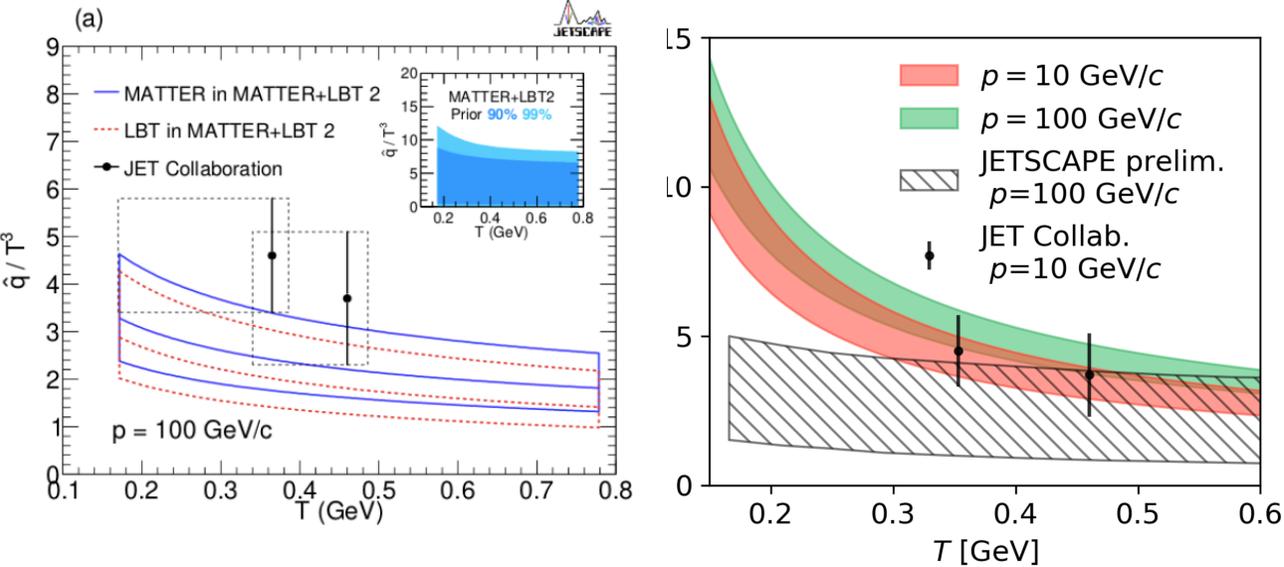
- The energetic jet loses a large amount of its energy via radiating gluon induced by multiple scatterings, leading to the suppression of the yield of high- $p_T$  jets and hadrons.
- Jet energy loss in QGP medium:  $\frac{dE}{dL} = \frac{\alpha_s N_c}{4} \hat{q} L$ .
- Jet transport parameter: hard parton's transverse momentum broadening squared per unit length.

$$\hat{q} = \rho \int dq_T^2 \frac{d\sigma}{dq_T^2} q_T^2. \text{ [BDMPS, NPB 483 (1997) 291]}$$

# Motivation: past efforts in the extraction of $\hat{q}/T^3$



[The JET colla., PRC 90. 014909] [M. Xie, H.Z. Zhang, et al., EPJC (2019)] [M. Xie, X.N. Wang, H.Z. Zhang, PRC (2021)] [X. Feal, et al., PLB (2021)]



- Form:  $\hat{q}/T^3 \propto C \longrightarrow \hat{q}/T^3(T)$
- Observable: single  $\longrightarrow$  multiple
- Method:  $\chi^2 \longrightarrow$  Bayesian inference
- Explicit parametrizations introduce unwarranted long-range correlations between different  $T$  regions.
- Uncertainty is large.

[S. Cao et al. (JETSCAPE), PRC (2021)] [WY Ke, X.N. Wang, JHEP (2021)]

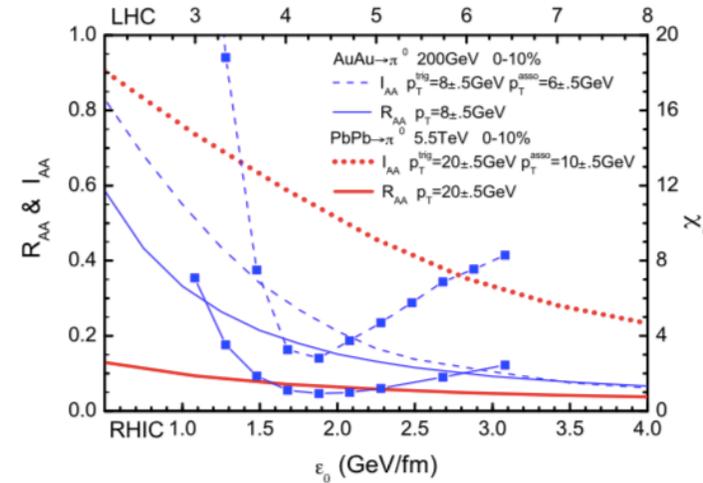
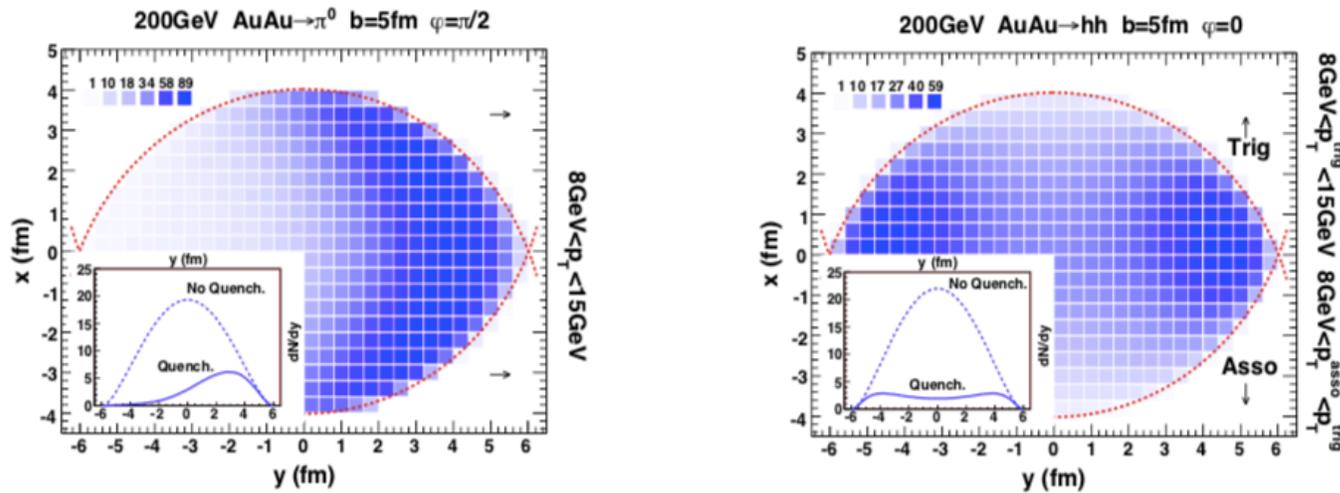
- Aim: the  $T$  dependence of  $\hat{q}/T^3$ .

- Data: all existing experimental data on  $R_{AA}^h, I_{AA}^{hh}, I_{AA}^{\gamma h}$  at RHIC and LHC with a wide range of centralities.

Single hadron,  $\gamma$ -hadron, di-hadron suppressions are direct consequences of JQ.

Dihadron and  $\gamma$ -hadron correlations are known to be more sensitive to the parton energy loss.

[H. Z. Zhang J.F. Owens, E. K. Wang and X.-N. Wang, PRL 98 (2007) 212301, PRL 103 (2009) 032302]



- Analysis method: Information field theory assisted Bayesian inference

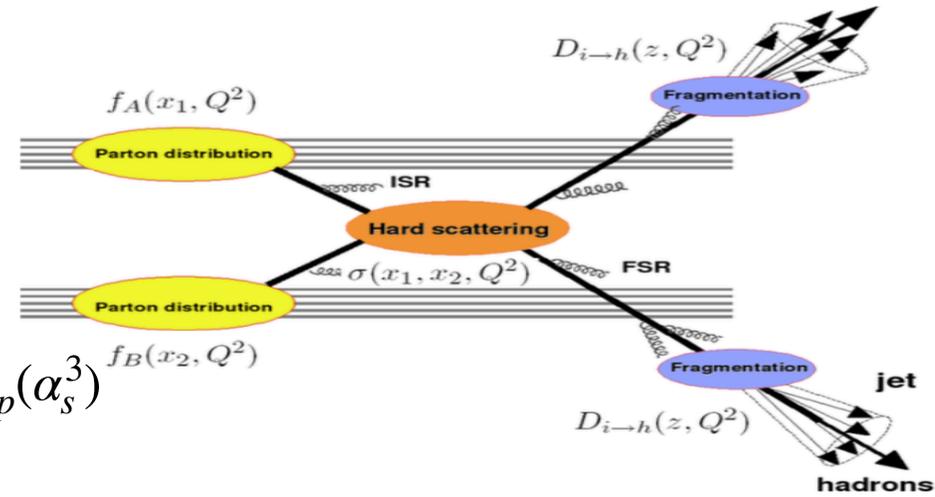
IF approach models the prior distribution of the unknown function that is free of long-range correlations

[Torsten A. Enßlin, Annalen Phys. 531 (2019) 3, 1800127], [J. C. Lemm, arXiv: physics/9912005]

# pQCD parton model In $p+p$ collisions:

Single hadron: [J.F.Owens, Rev. Mod. Phys 59,465(1987)]

$$\frac{d\sigma_{pp}^h}{dyd^2p_T} = \sum_{abcd} \int dx_a dx_b f_{a/p}(x_a, \mu^2) f_{b/p}(x_b, \mu^2) \frac{1}{\pi} \frac{d\sigma_{ab \rightarrow cd}}{d\hat{t}} \frac{D_{h/c}(z_c, \mu^2)}{z_c} + \Delta\sigma_{pp}^h(\alpha_s^3)$$



Di-hadron:

$$\frac{d\sigma_{pp}^{hh}}{dy_h^c d^2p_T^{h_c} dy_h^d d^2p_T^{h_d}} = \sum_{abcd} \int dz_c dz_d f_{a/p}(x_a, \mu^2) f_{b/p}(x_b, \mu^2) \frac{x_a x_b}{\pi z_c^2 z_d^2} \frac{d\sigma_{ab \rightarrow cd}}{d\hat{t}} D_{h/c}(z_c, \mu^2) D_{h/d}(z_d, \mu^2) \delta^2(\vec{p}_T^c + \vec{p}_T^d) + \Delta\sigma_{pp}^{hh}(\alpha_s^3)$$

$\gamma$ -hadron: only consider direct photon [M. Xie, X.N. Wang, H.Z. Zhang, PRC 103, 034911 (2021)]

$$\frac{d\sigma_{pp}^{\gamma h}}{dy_\gamma d^2p_T^\gamma dy_h d^2p_T^h} = \sum_{abd} \int dz_d f_{a/p}(x_a, \mu^2) f_{b/p}(x_b, \mu^2) \frac{x_a x_b}{\pi z_d^2} \frac{d\sigma_{ab \rightarrow \gamma d}}{d\hat{t}} D_{h/d}(z_d, \mu^2) \delta^2\left(\vec{p}_T^\gamma + \frac{\vec{p}_T^h}{z_d}\right) + \Delta\sigma_{pp}^{\gamma h}(\alpha_e \alpha_s^2)$$

Triggered fragmentation functions:

$$D_{pp}^{\text{trig}}(z_T) = p_T^{\text{trig}} \frac{d\sigma_{pp}/dy^{\text{trig}} dp_T^{\text{trig}} dy^{\text{assoc}} dp_T^{\text{assoc}}}{d\sigma_{pp}/dy^{\text{trig}} dp_T^{\text{trig}}}, z_T = p_T^{\text{assoc}}/p_T^{\text{trig}}.$$

$f_{a/p}(x_a, \mu^2)$  : CT14 PDF

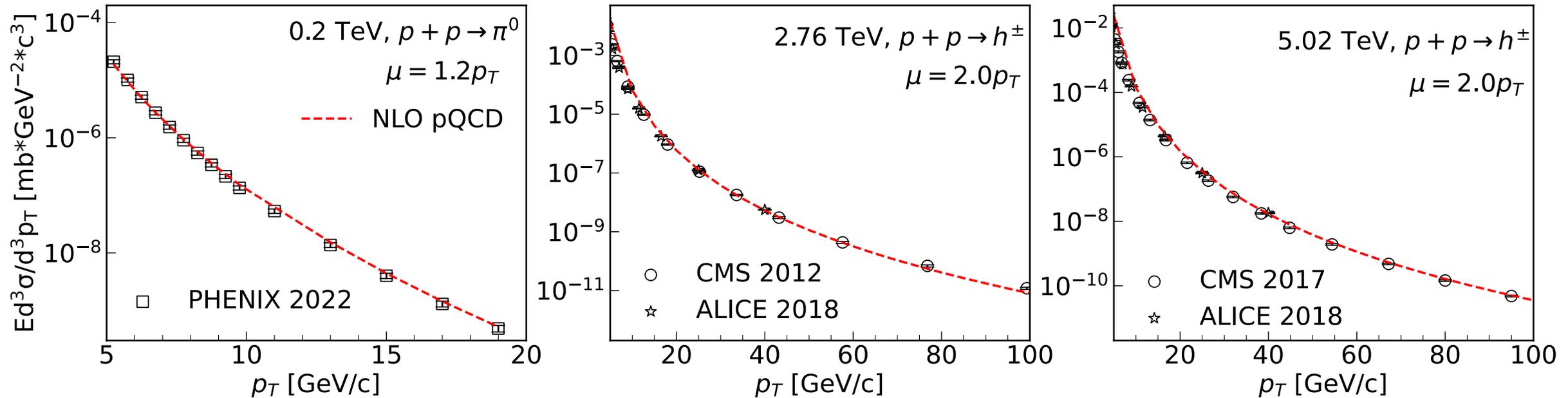
[CT14, Phys. Rev. D 95, 034003 (2017)]

$D_{h/d}(z_d, \mu^2)$  : KKP FFs

[KKP, Nucl. Phys. B 582,514(2000)]

# $pp$ baselines

## Single hadron:

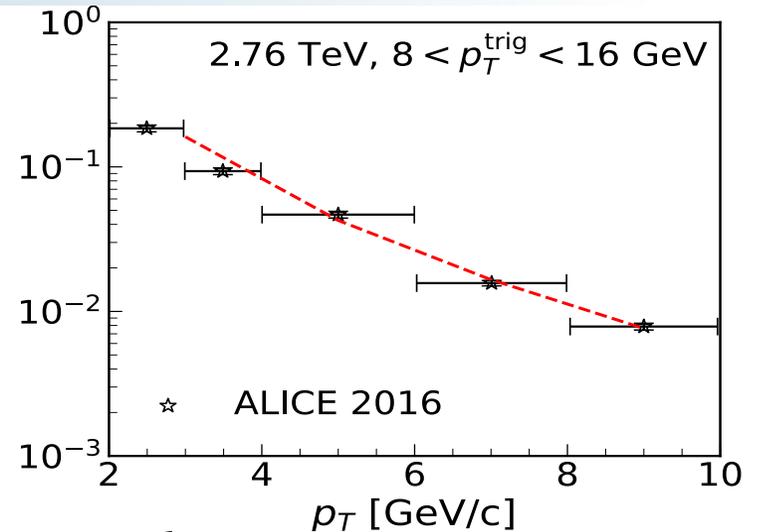
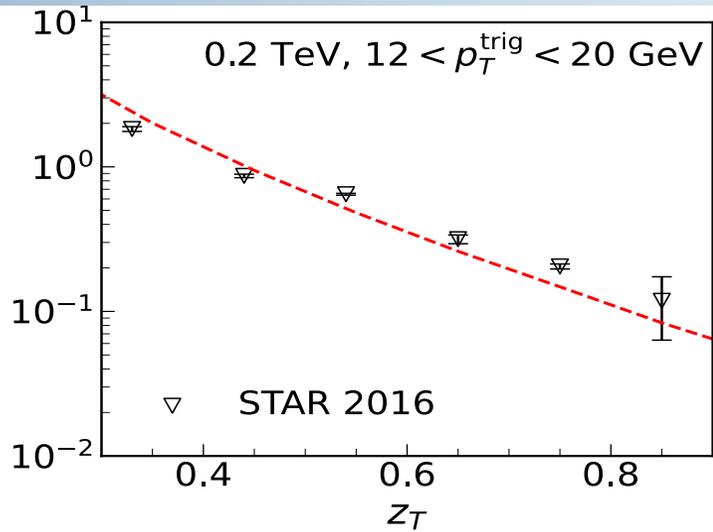
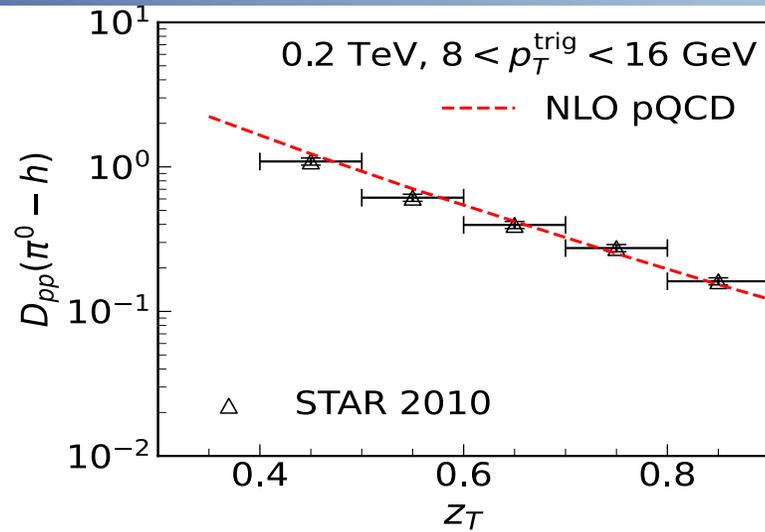


NLO pQCD describes data very well in  $p+p$  collisions.

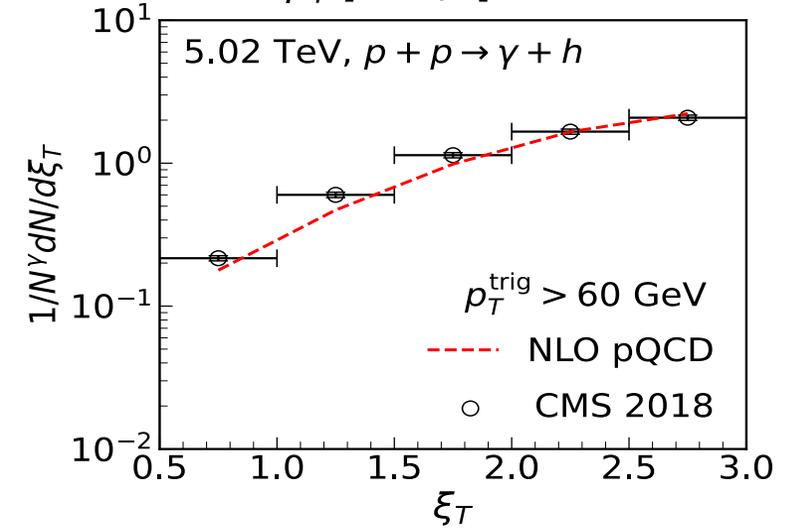
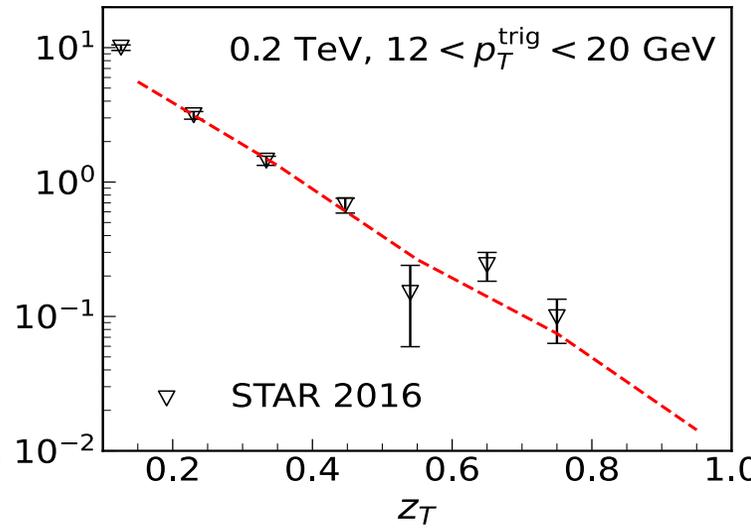
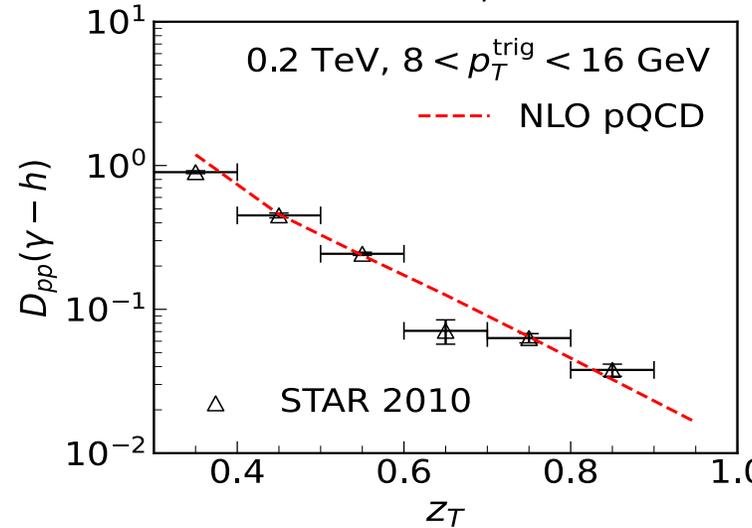
Data from [PHENIX, Phys. Rev. Lett. 101, 232301 (2008), Phys. Rev. C 105, 064902 (2022)]  
[CMS, Eur. Phys. J. C 72, 1945 (2012), JHEP 04, 039 (2017)], [ALICE, JHEP 11, 013 (2018)]

# $pp$ baselines

Di-hadron

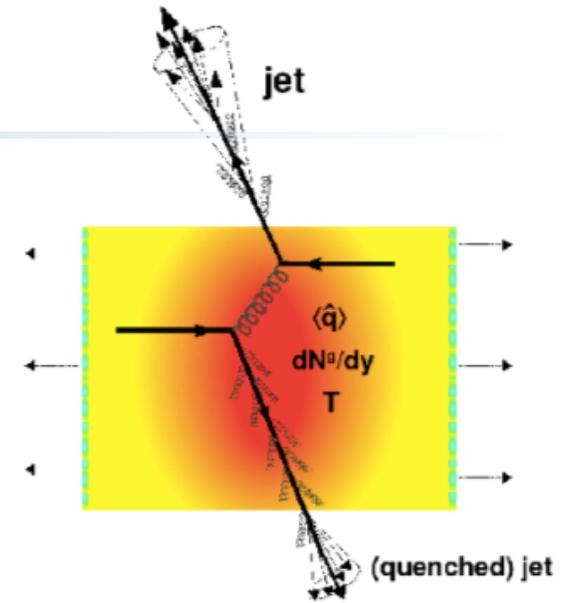
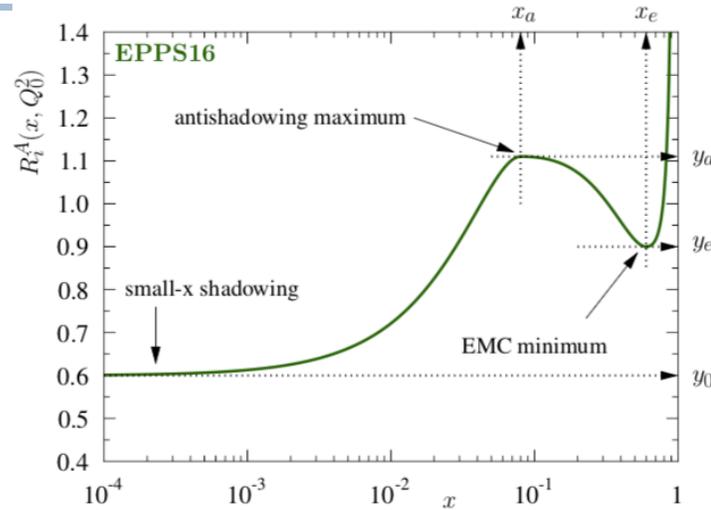
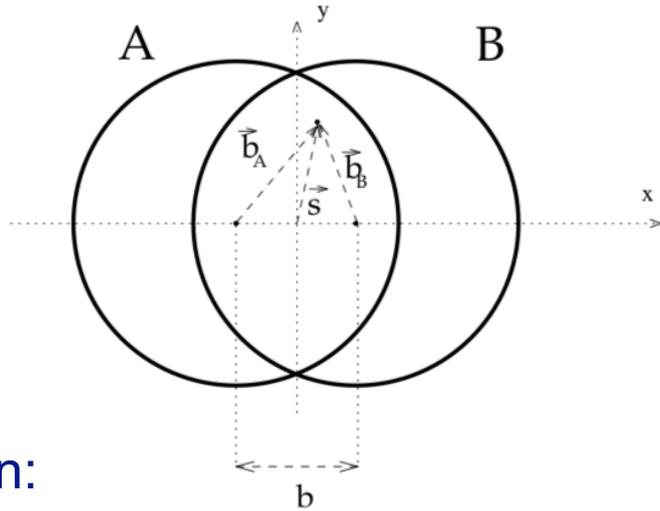


$\gamma$ -hadron



NLO pQCD parton model can give a good description of all the experimental data on single hadron,  $\gamma$ -hadron, and di-hadron production spectra at large  $p_T$  in  $p+p$  collisions.

# pQCD parton model In A+A collisions:



Di-hadron:

$$\frac{dN_{AB}^{hh}}{dy_h^c d^2p_T^{h_c} dy_h^d d^2p_T^{h_d}} = \sum_{abcd} \int d^2r dz_c dz_d t_A(\vec{r}) t_B(\vec{r} + \vec{b}) f_{a/A}(x_a, \mu^2, \vec{r}) f_{b/B}(x_b, \mu^2, \vec{r} + \vec{b}) \frac{x_a x_b}{\pi z_c^2 z_d^2} \frac{d\sigma_{ab \rightarrow cd}}{d\hat{t}} \times \tilde{D}_{h/c}(z_c, \mu^2, \Delta E_c) \tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d) \delta^2(\vec{p}_T^c + \vec{p}_T^d) + \mathcal{O}(\alpha_s^3)$$

$t_A(\vec{r})$  : nuclear thickness functions, integrating the Woods – Saxon nuclear density distribution;

$f_{a/A}(x_a, \mu^2, \vec{r})$  : EPPS16 nPDF [EPPS16, Eur. Phys. J. C 77, 163 (2017)]

$\tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d)$  : medium – modified FFs

# Medium-modified FFs

$$\tilde{D}_{h/d}(z_d, \mu^2, \Delta E_d) = (1 - e^{-\langle N_g \rangle}) \left[ \frac{z'_d}{z_d} D_{h/d}(z'_d, \mu^2) + \langle N_g \rangle \frac{z'_g}{z_d} D_{h/g}(z'_g, \mu^2) \right] + e^{-\langle N_g \rangle} D_{h/d}(z_d, \mu^2)$$

[X.-N. Wang, PRC70 (2004) 031901], [H. Z. Zhang, J.F. Owens, Enke Wang, X.-N. Wang, PRL 98.212301 (2007);103, 032302 (2009)]

## High-Twist approach:

$$\frac{\Delta E_d}{E} = \frac{2C_A \alpha_s}{\pi} \int d\tau \int \frac{dl_T^2}{l_T^4} \int dz [1 + (1 - z)^2] \hat{q}_d(\tau, \vec{r} + (\tau - \tau_0) \vec{n}) \sin^2\left(\frac{l_T^2(\tau - \tau_0)}{4z(1 - z)E}\right)$$

[W.T. Deng and X.-N. Wang, PRC 81,024902(2010), [E. Wang and X.-N. Wang, PRL 87, 142301 (2001); 89, 162301 (2002)]

$$\langle N_g^d \rangle = \frac{2C_A \alpha_s}{\pi} \int_{\tau_0}^{\infty} d\tau \int \frac{dl_T^2}{l_T^4} \int \frac{dz}{z} [1 + (1 - z)^2] \hat{q}_d(\tau, \vec{r} + (\tau - \tau_0) \vec{n}) \sin^2 \left[ \frac{l_T^2 (\tau - \tau_0)}{4z(1 - z)E} \right]$$

[N.B. Chang, W.T. Deng and X.-N. Wang, Phys. Rev. C 89, 03491 (2014)1

## Observables: nuclear modification factor

$$R_{AB} = \frac{dN_{AB}/dyd^2p_T}{T_{AB}(\vec{b})d\sigma_{pp}/dyd^2p_T}; \quad I_{AB}(z_T) = \frac{D_{AB}(z_T)}{D_{pp}(z_T)}.$$

# IF-assisted Bayesian inference of $\hat{q}(T)$

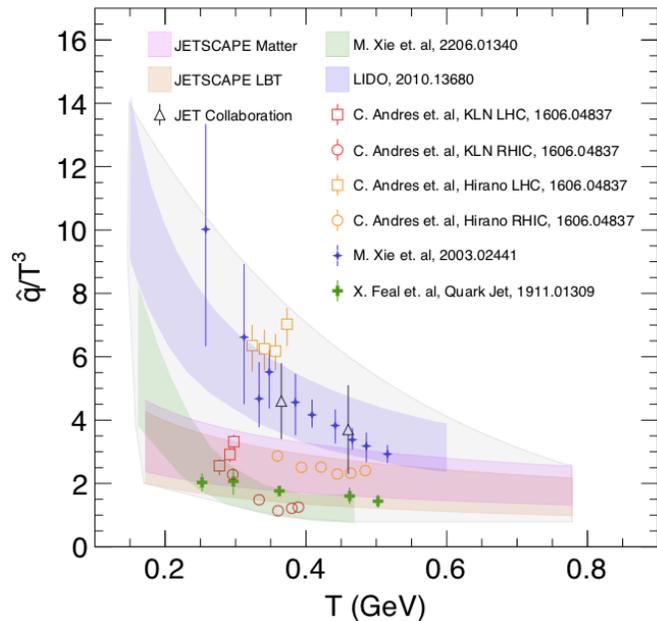
Construction of the functional prior  $\hat{q}(T)$ :

Gaussian random field ( $F$ ):  $\langle F(x) \rangle = \mu(x)$ ;  $\langle [F(x) - \mu(x)] [F(x') - \mu(x')] \rangle = C(x, x')$

[W. Bialek, C.G. Callan, et al, PRL 77, 4693 (1996)] [T. Ensslin, AIP Conf. Proc. 1553, 184 (2013)] [Y.Y. He, X.-N. Wang, et al, PRC 91,054908(2015)]

$$F(x) = \ln(\hat{q}/T^3); x = \ln(T/\text{GeV}); \quad \mu(x) \equiv C; C(x, x') = \sigma^2 \exp \left[ -(x - x')^2 / 2l^2 \right]$$

The  $l$  indicates that the high- $T$  prior is unaffected by data that are only sensitive to  $\hat{q}(T)$  at low  $T$ .



$$\mu = \langle \ln(\hat{q}/T^3) \rangle = 1.36, \sigma = 0.7;$$

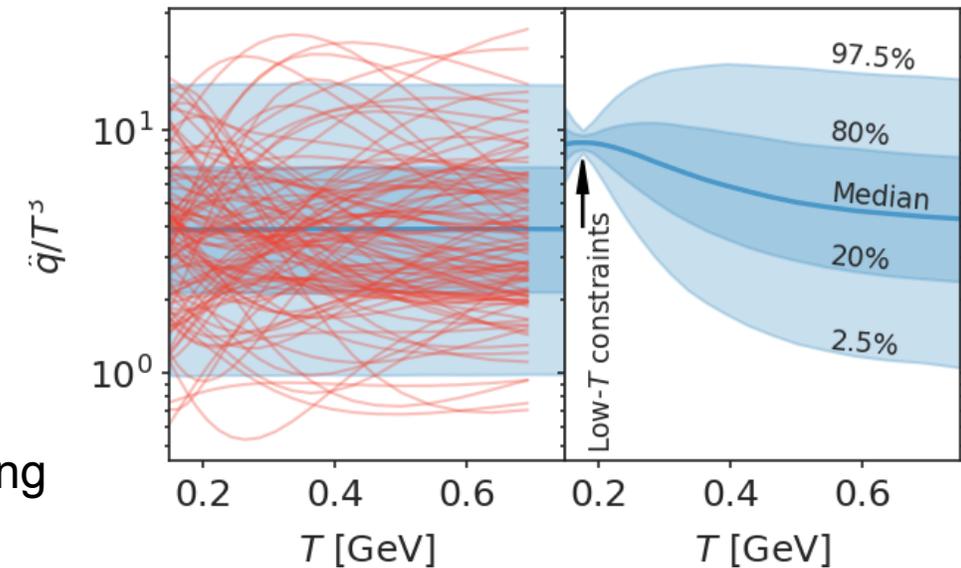
$$T_c < T < 4T_c : 0.8 < \hat{q}/T^3 < 15,$$

$$T_c = 0.165 \text{ GeV}.$$

$$l = \ln\left(\frac{T_{\max}}{T_{\min}}\right) \approx \ln 2;$$

The smallest  $\ln(T_{\max}/T_{\min})$  in the hydrodynamic simulations of all colliding systems and centralities.

Random realizations of prior  $\hat{q}/T^3$



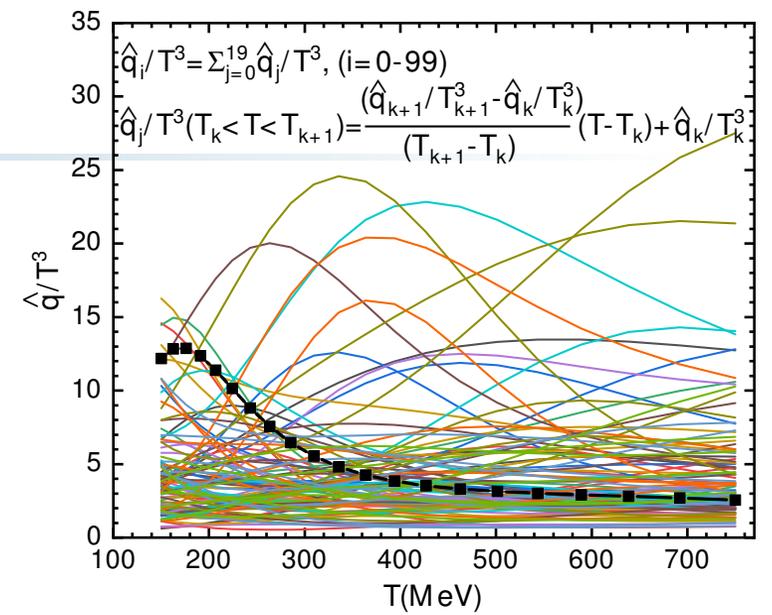
$\hat{q}$  individual variation in different  $T$  region.

# NLO pQCD model with 100 prior $\hat{q}(T)$

## Input parameter matrix

$$\left[ \frac{\hat{q}}{T^3} \right]_i = \sum_{j=1}^{20} \theta(T - T_j) \theta(T_{j+1} - T) \left( \left[ \frac{\hat{q}}{T^3} \right]_{i,j+1} \Delta_j + \left[ \frac{\hat{q}}{T^3} \right]_{i,j} (1 - \Delta_j) \right)$$

$$\Delta_j = \frac{T - T_j}{T_{j+1} - T_j}, \quad i = 1 - 100, \quad j = 1, 20.$$

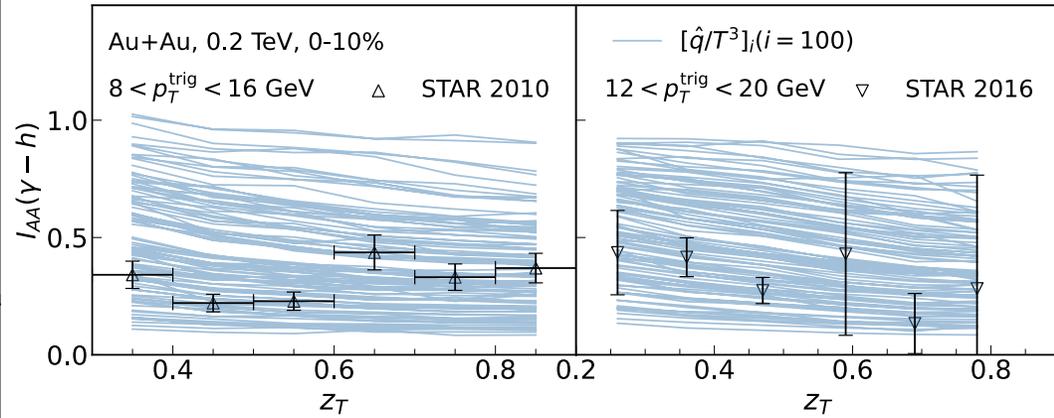
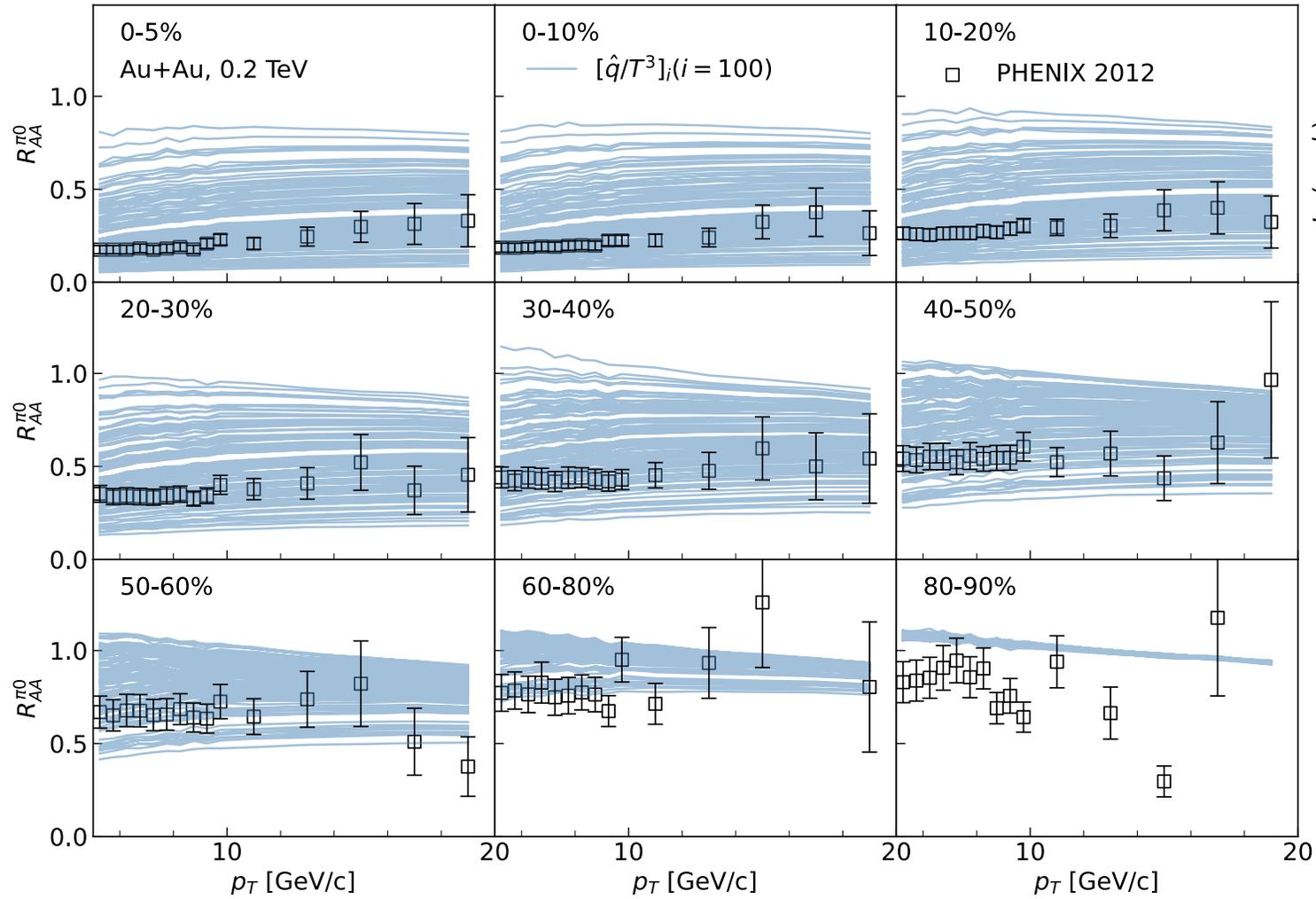


The hydrodynamic evolution of the QGP medium provided by the CLVisc (3+1) D hydrodynamic model.

[L. G. Pang, Q. Wang, X.-N. Wang, Phys.Rev.C 86 (2012) 024911], [L. G. Pang, H. Petersen, X.-N. Wang, Phys.Rev.C 97 (2018) 6, 064918]

- $R_{AA}^{\pi^0}$  in 0-5%, 0-10%, 10-20%, 20-30%, 30-40%, 40-50% Au-Au collisions at 0.2 TeV;
- $R_{AA}^{h^\pm}$  in 0-5%, 5-10%, 10-20%, 20-30%, 30-40%, 40-50% Pb-Pb collisions at 2.76 TeV;
- $R_{AA}^{h^\pm}$  in 0-5%, 5-10%, 10-20%, 10-30%, 20-30%, 30-50% Pb-Pb collisions at 5.02 TeV;
- $I_{AA}^{\gamma h^\pm}$  in 0-10% Au-Au collisions at 0.2 TeV, (3 trigger ranges);
- $I_{AA}^{\pi^0 h^\pm}$  in 0-10% Au-Au collisions at 0.2 TeV, (3 trigger ranges);
- $I_{AA}^{h^\pm h^\pm}$  in 0-5%, 0-10% Pb-Pb collisions at 2.76 TeV, (6 trigger ranges).

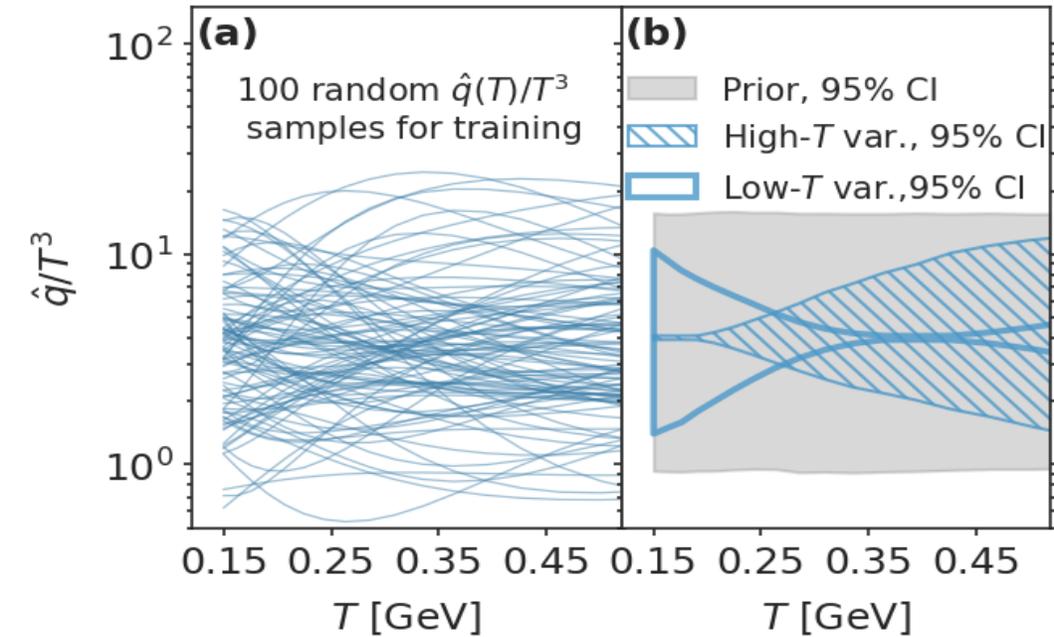
# Model calculations with 100 prior $\hat{q}(T)$



- 2500 curves of  $R_{AA}$
- 500 curves of  $I_{AA}^{\gamma h}$
- 1000 curves of  $I_{AA}^{hh}$

Data sets  $(\hat{q}_{ij}, \mathbf{y}_{ik})$  of the NLO pQCD calculations to train Gaussian Process emulators.

# IF-assisted sensitivity analysis

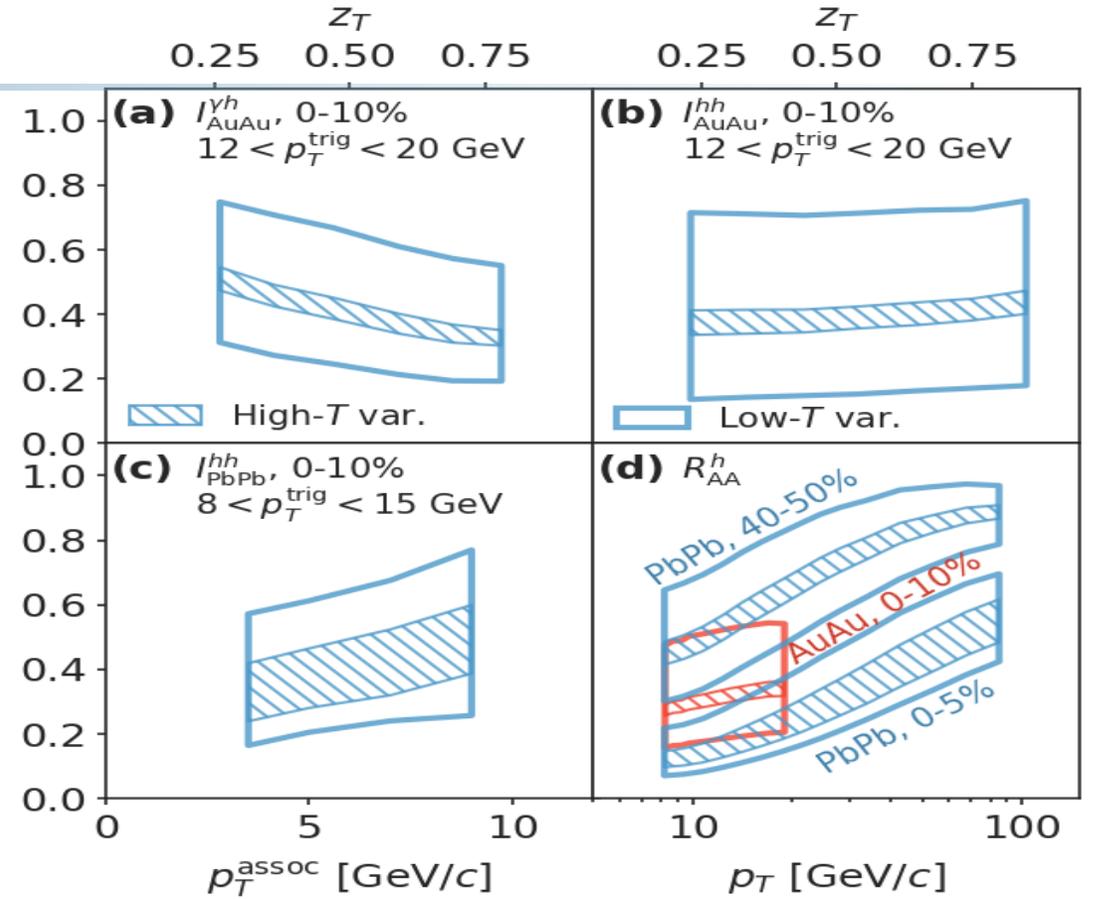


High- $T$  var:

$$0.15 < T < 0.2 \text{ GeV}, \hat{q}/T^3 = 4 \pm 0.1$$

Low- $T$  var:

$$0.3 < T < 0.4 \text{ GeV}, \hat{q}/T^3 = 4 \pm 0.1$$



- The sensitivities of  $R_{AA}$  and  $I_{AA}$  to high- $T$   $\hat{q}$  increase from RHIC to LHC energies and from peripheral to central collisions.

- The prior  $\hat{q}/T^3$  at high  $T$  and low  $T$  are uncorrelated.
- $I_{AA}$  are slightly more sensitive to high- $T$   $\hat{q}$  than  $R_{AA}$ .

# Global Bayesian inference of $\hat{q}(T)$

- Prior distribution of the random function:

[Jorsten A. Enßlin *Annalen Phys.* 531 (2019) 3, 1800127]

$$P_0[F] = \exp \left[ -\frac{1}{2} \int dx dx' \delta F(x) C^{-1}(x, x') \delta F(x') \right]$$

- Posterior distribution :

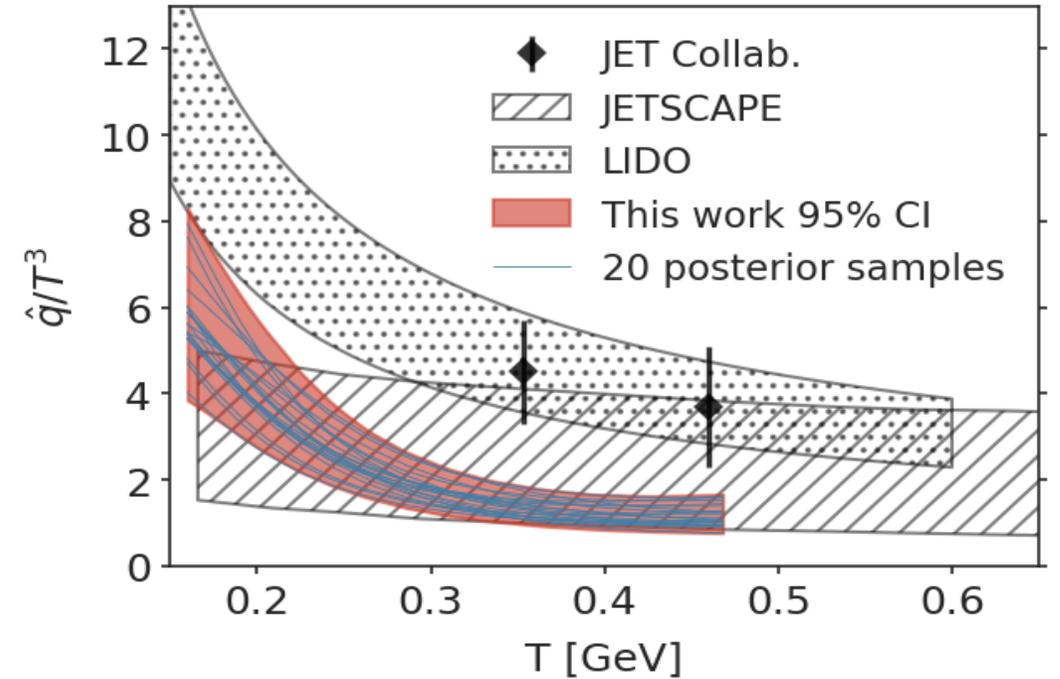
[python library emcee: multi-dimensional Gaussian]

$$\begin{aligned} P[F(x)] &\propto P_0[F(x)] \times \text{Likelihood} \\ &= \frac{\mathcal{N}}{\sqrt{\det(C)}} \exp \left\{ -\frac{1}{2} \int dx dx' \delta F(x) C^{-1}(x, x') \delta F(x') \right. \\ &\quad \left. -\frac{1}{2} \left( M'[F] - y_{\text{exp}} \right)^T \Sigma_{\text{error}}^{-1} \left( M'[F] - y_{\text{exp}} \right) \right\} \end{aligned}$$

$M'[F]$  : NLO pQCD calculations;  $y_{\text{exp}}$  : experimental data.

- To marginalize the values of  $F^*$  at a fixed input  $x^*$

$$p(F^*) = \int [\mathcal{D}F] \delta(F(x^*) - F^*) P[F(x)]$$



[JET Collaboration: PRC 90, 014909 (2014)]

[JETSCAPE: PRC 104, 024905 (2021)]

[LIDO, W.Y. Ke, X.-N. Wang, JHEP 05041]

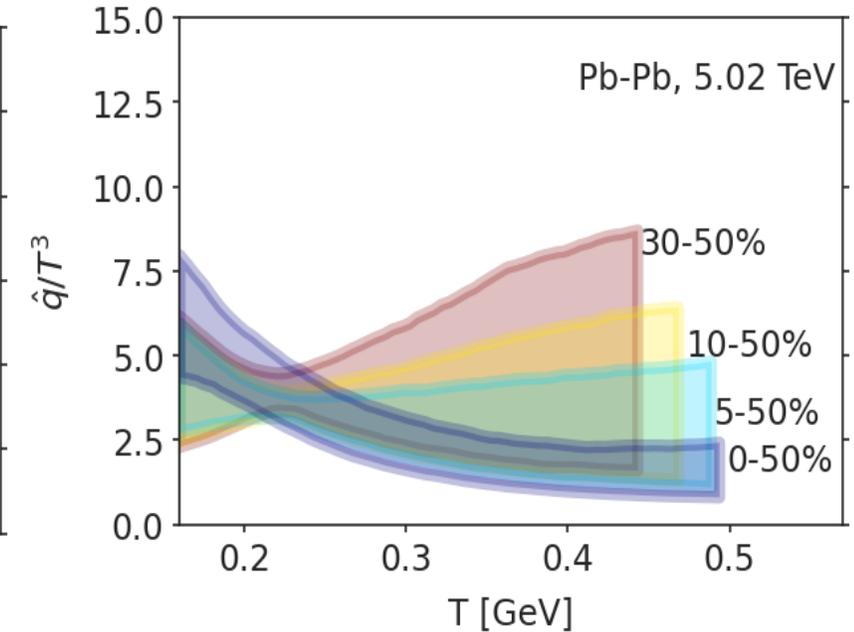
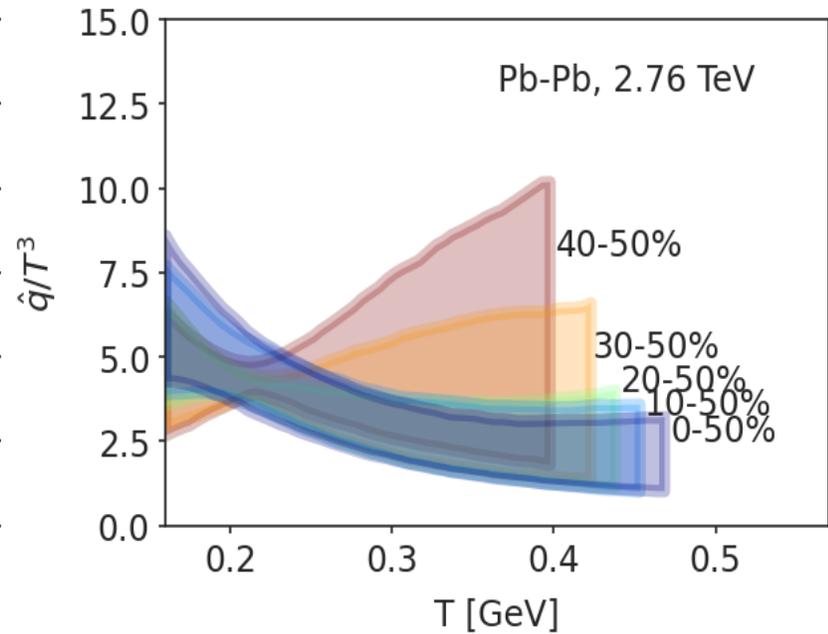
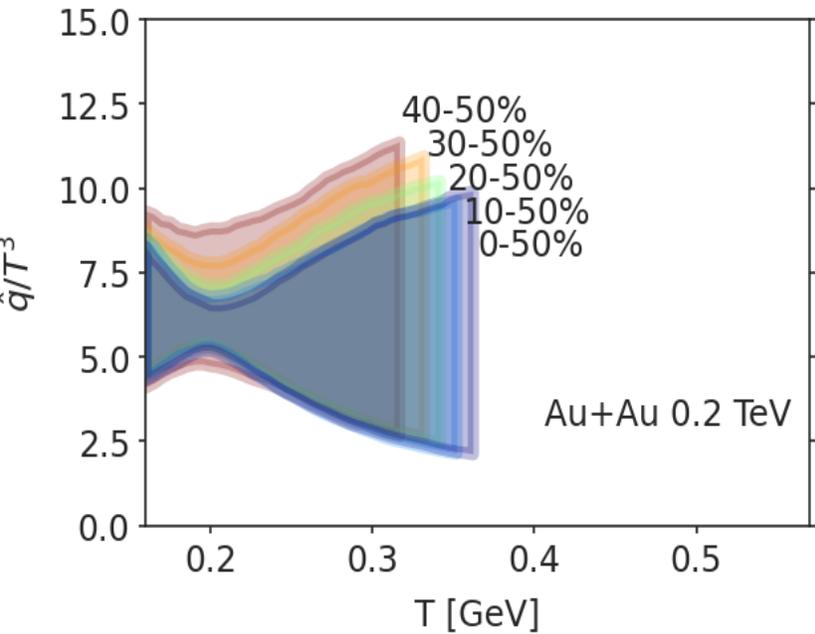
- The extracted  $\hat{q}/T^3$  is consistent with previous studies but exhibits a stronger temperature dependence.

# Progressive constraining power

From **low- $T$**   $\longrightarrow$  **high- $T$**  region;

**peripheral**  $\longrightarrow$  **central** collisions;

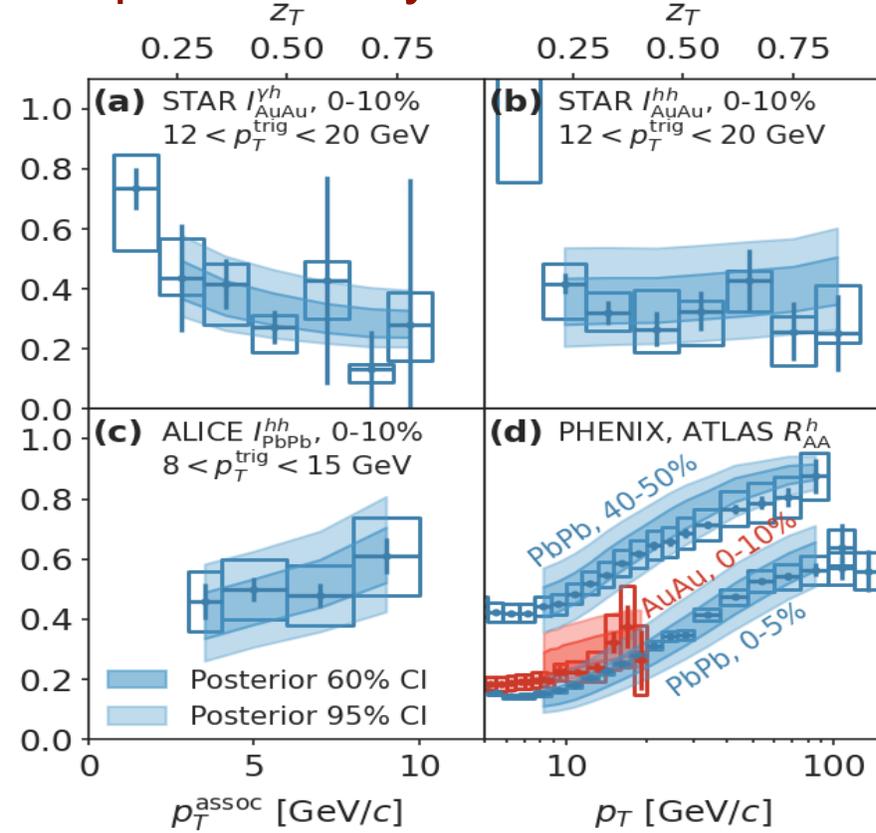
**RHIC**  $\longrightarrow$  **LHC** energies.



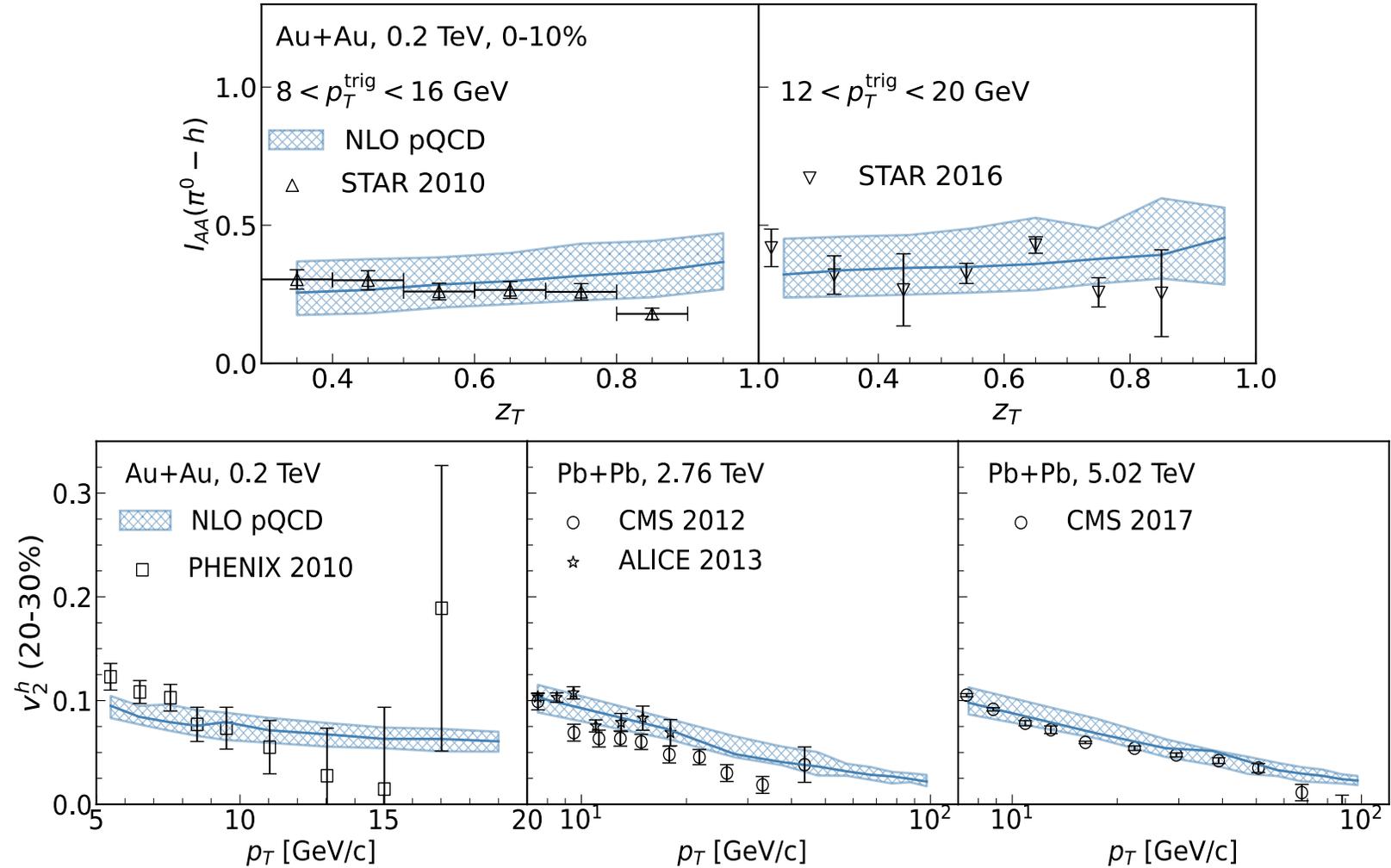
- $\hat{q}/T^3$  is progressively more constrained at higher  $T$  as one includes data from more central collisions and from higher colliding energies.
- The IF approach strongly suppresses correlations between high and low  $T$  prior values of  $\hat{q}/T^3$ .

# The posterior distribution of the observables

predicted by the emulator



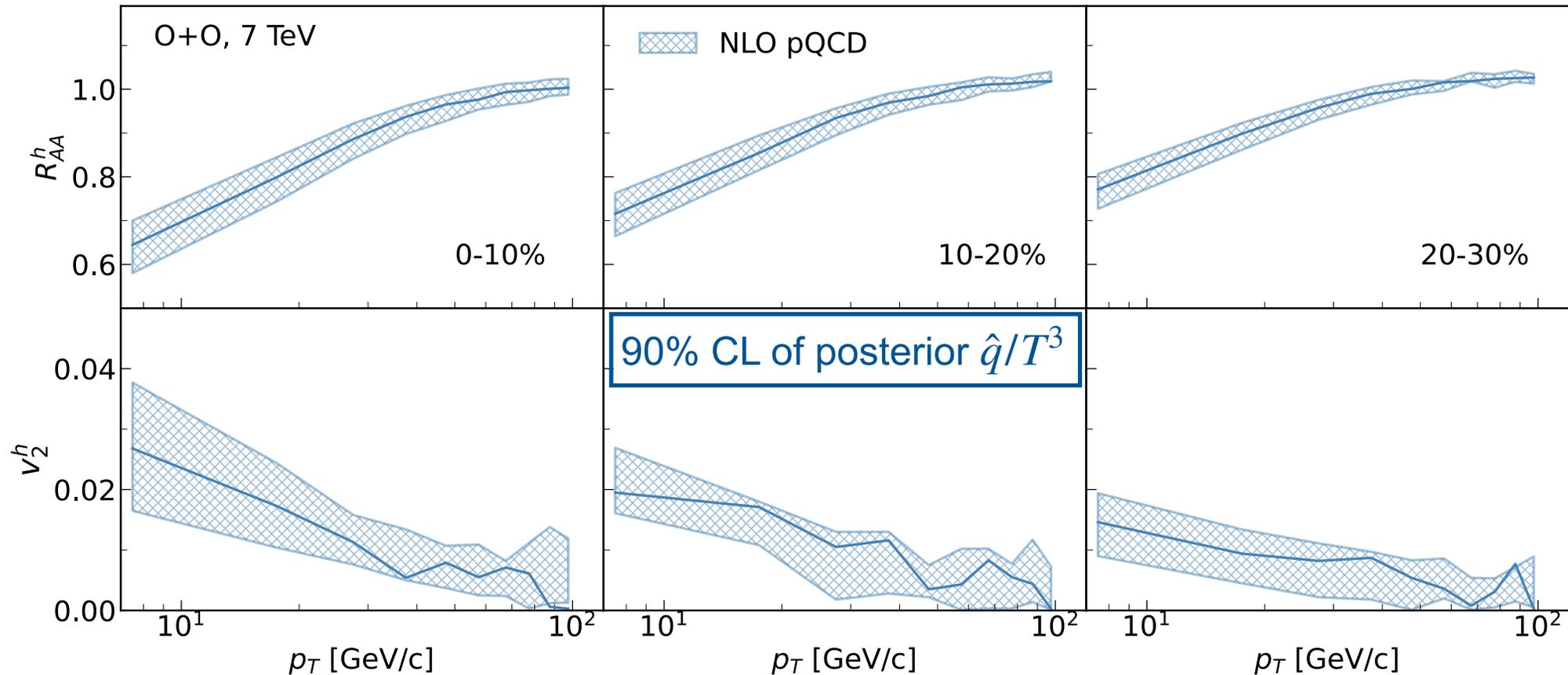
calculated by NLO pQCD model



- The experimental data over a large range of  $p_T$ , centralities and colliding energies are well reproduced.

# Prediction O-O collisions at 7 TeV Assume the QGP matter is formed.

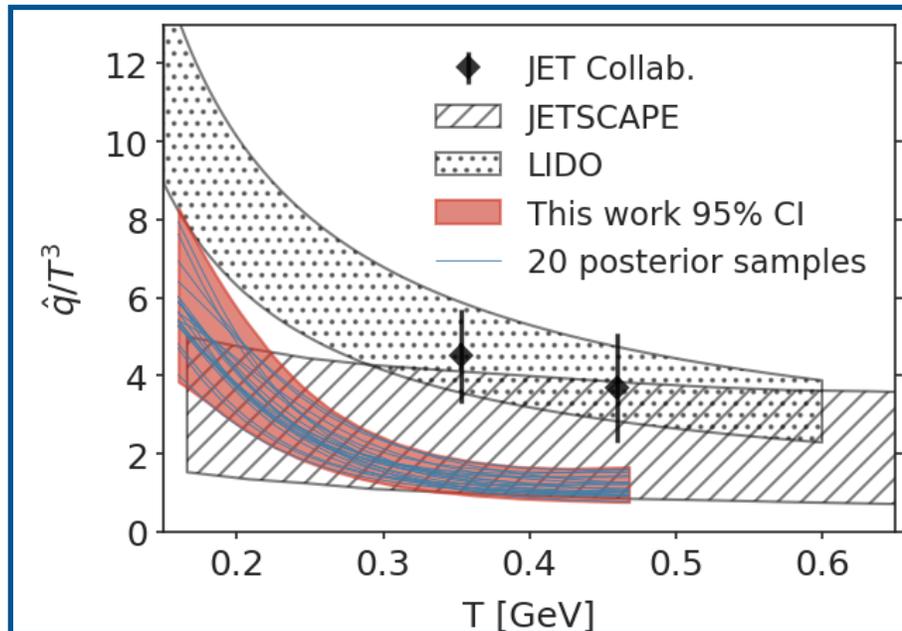
O-O collisions are proposed as an intermediate colliding system between heavy-ion and  $p$ -A collisions to investigate the properties of dense matter in small systems.



- In most central collisions, the suppression of single hadron is about 30~40% at  $p_T = 7.5 \sim 10$  GeV/c.
- The  $v_2$  is only 0.03 at  $p_T \sim 7.5$  GeV/c in 0-10% collisions.

# Summary

- We applied the IF approach to the first global Bayesian inference of the  $\hat{q}/T^3$  from combined experimental data on  $R_{AA}^h$ ,  $I_{AA}^{hh}$ , and  $I_{AA}^{\gamma h}$  in HIC at both RHIC and LHC energies.
- The IF approach can provide progressive constraints on  $\hat{q}$  from low to high  $T$  when experimental data in more central collisions and at higher colliding energies are included.
- The extracted  $\hat{q}/T^3$  is consistent with previous studies but exhibits a stronger  $T$  dependence.



***Thank you for your attention!***