



Singularity of density of states in the pi-flux lattice

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Outline

- 1 Motivation
 - Interactions vs disorder
 - π -flux model
 - Lattice deformations and ripples
- 2 Random hopping in π -flux model
 - Singularity of DOS
 - k -dependent self-energy and minimal conductivity
- 3 Summary

disorder in semimetals

Why disorder?

- ① inherently present in all solid-states systems
- ② determine/dominate properties
localization, symmetry broken, new phase...
- ③ adjustable
doping, impurities, vacancies, absorption, strain...

Interactions-disorder duality and critical phenomena

non-Anderson transition

$$H = \int \Psi^\dagger(r) \xi_p \Psi(r) d^d r - \frac{1}{2} \int \Psi^\dagger(r) \Psi^\dagger(r') \xi_p \Psi(r') \Psi(r) d^d r d^d r'$$

Interacting model

Disordered model

$$\overrightarrow{\text{---}}_{\omega, \mathbf{p}} = \frac{1}{i\omega - \xi_{\mathbf{p}}}$$

$$\overrightarrow{\text{---}}_{\mathbf{k}, \mathbf{p}} = \frac{1}{2} [\hat{G}^R(\mathbf{k}, \mathbf{p}) + \hat{G}^A(\mathbf{k}, \mathbf{p})]$$



$$= 1$$



$$= \hat{\sigma}_z$$

$$\overbrace{\text{~~~~~}}_{\omega, \mathbf{p}} = D(\omega, \mathbf{p})$$

$$\overbrace{\text{-----}}_{\mathbf{k}, \mathbf{p}} = -\tilde{D}(\mathbf{k}, \mathbf{p})$$

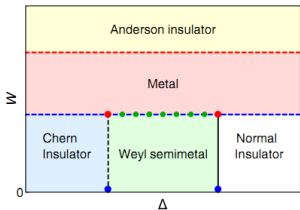
	Interacting model	Disordered model
Coordinates	(τ, \mathbf{r})	(τ_{d+1}, \mathbf{r})
Temperature/size	T	$1/\ell_{d+1}$
Coupling to interactions/disorder	$\hat{\Psi}^\dagger \hat{\Psi} \phi$	$\hat{\psi}^\dagger \hat{\sigma}_z \hat{\psi} u$
Observables	\hat{n} (density)	ρ_s , Eq. (5)

Syzranov, Sergey V., et al. Annu. Rev. Condens. Matter Phys. 9, 1 (2018); Sun, Shijun, and Sergey Syzranov. Phys. Rev. B 108 195132

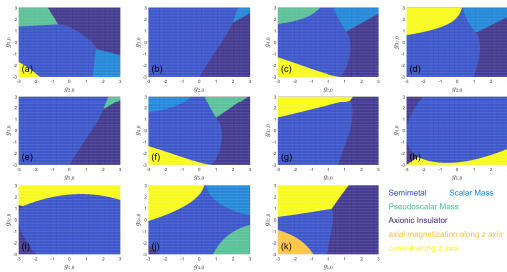
(2023)

disorder-driven transition in semimetals

global phase diagram of semi-Dirac semimetals



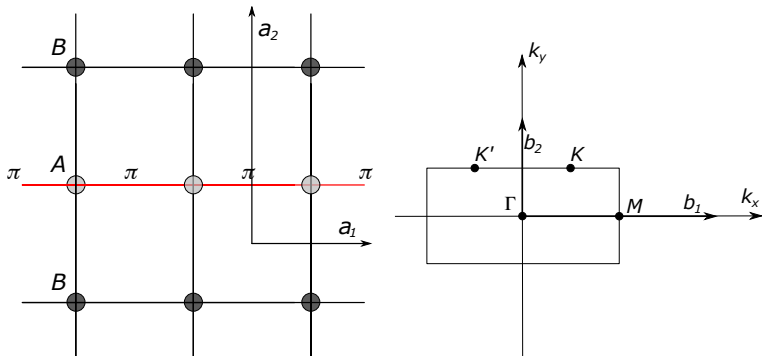
W : strength of disorder
 $\Delta = t_z - m_z$: QPT parameter of SM-IN transition



Physical Review B, 102, 085132 (2020); Journal of Physics: Condensed Matter, 33, 125601 (2021); Physical Review B, 107, 155125 (2023).

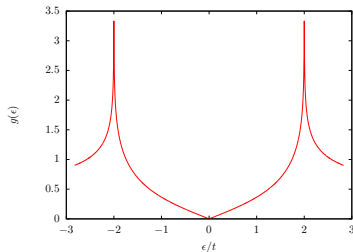
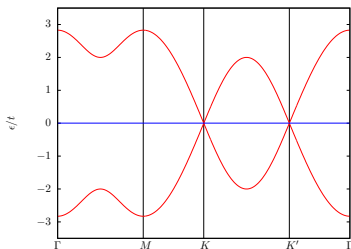
π -flux model

$$t_{j+\hat{x},j} = (-1)^{j_x}, \quad t_{j+\hat{y},j} = 1, \quad \epsilon_{\pm}(k) = \pm 2t\sqrt{\cos^2 k_x + \cos^2 k_y}$$



π -flux model

$$\epsilon_{\pm}(k) = \pm t \sqrt{1 + 4 \cos\left(\frac{k_x}{2}\right) \cos\left(\frac{\sqrt{3}k_y}{2}\right) + 4 \cos^2\left(\frac{k_x}{2}\right)}$$



Random hopping Dirac fermions on π -flux lattice

Gade's Model: $P[\delta t] \propto [-W, W]$,
 $\rho(E) = \frac{1}{E} \exp(-c \ln|E|^{1/x})$, $x = \frac{3}{2}$ or 2?

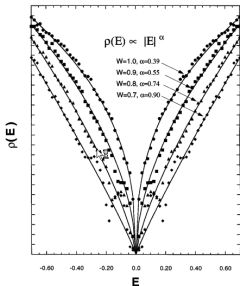


FIG. 1. The density of states $\rho(E)$ (with small imaginary part $\delta = 0.02$), where $L = 50$ and $W = 0.7, 0.8, 0.9$, and 1.0 . We have fitted the data by the power-law form $\rho(E) = CE^{\alpha(W)}$, where $\alpha(W) = 0.90, 0.74, 0.55$, and 0.39 for $W = 0.7, 0.8, 0.9$, and 1.0 , respectively.

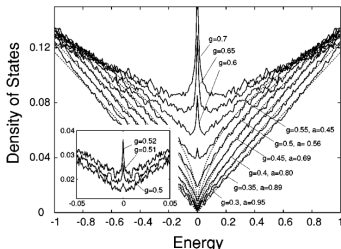


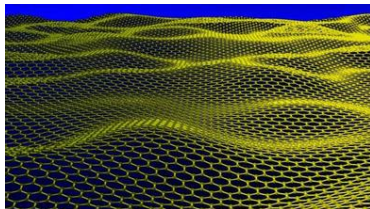
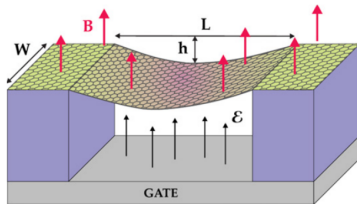
FIG. 4. Same as Fig. 3 on a 100×100 lattice for $g = 0.3 - 0.7$, and $\delta = 0.005$, averaged over 50 samples. Away from zero energy ($E > 0.1$) and for sufficiently weak randomness, data are fitted to a power law $\rho(E) \propto |E|^{\alpha}$. Inset: same as Fig. 3 on a 250×250 lattice for $g = 0.5, 0.51, 0.52$, and $\delta = 0.0005$, averaged over 70 samples.

NPB 398, 499 (1993), PRL, 79 (19) (1997), PRB, 58 (11) (1998), PRB 65, 064206(2002), PRB 67, 064202 (2003);

PRL 113, 5, 186802 (2014), PRL 117, 5, 116806 (2016)

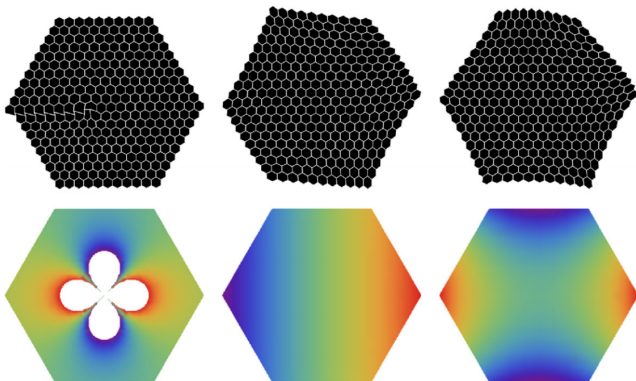
Lattice deformations

$$\Delta h \sim 1 - 4\text{\AA}, \quad u \sim 0.1 - 1\%, \quad L \sim 100\text{nm}$$

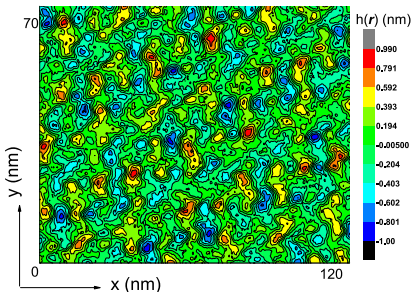


- ① hybridized orbitals and "topological" defects.
- ② modulate the on-site energies (ϵ_{ij}) and the hopping parameter (γ_{ij}).
- ③ pseudo magnetic field (gauge field)

effective magnetic fields induced by the deformations



Model of Ripples



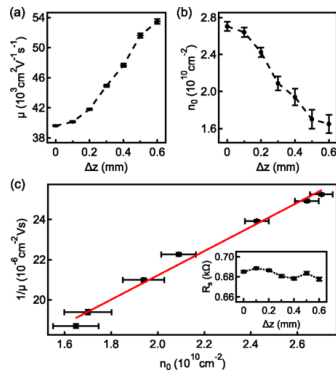
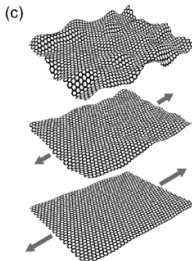
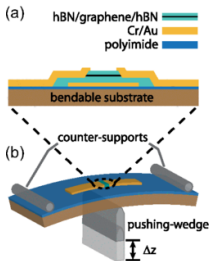
$$h(\mathbf{r}) = C \sum_i^{N_r} C_{\mathbf{q}_i} \sin(\mathbf{q}_i \cdot \mathbf{r} + \delta_i), \quad \delta_i \in [0, 2\pi], \quad \mathbf{q}_i \in [2\pi/L, 2\pi/3a]$$

$$C_{\mathbf{q}_i} = \begin{cases} \sqrt{2}/q_i^2, & q_i > q_{\text{ripple}}, \quad q_{\text{ripple}} = 2\pi/\lambda_{\text{ripple}} \\ \sqrt{2}/q_{\text{ripple}}^2, & \text{otherwise, } \lambda_{\text{ripple}} \approx 8\text{nm} \end{cases}$$

$$\delta t_{ij} = t\alpha \Delta a/a, \quad \Delta a = \sqrt{a^2 + (h(\mathbf{r}) - h(\mathbf{r}'))^2} - a, \quad \alpha = -2$$

Effects of Random Strain Fluctuations on mobility

Drude model: $G = \frac{1}{\frac{\alpha}{\epsilon\mu\sqrt{n^2+n_0^2}} + R_s}$, $\Delta z = 0.6 \text{ nm} \sim \bar{\sigma} \approx 0.2\%$



Wang, L. et al, Phys. Rev. Lett. 124, 157701 (2020); Wang, L., et al. Commun. Phys. 4, 147 (2021)

Random hopping model

A graphene or π -flux model with bond disorder belongs to the **chiral orthogonal symmetry** class since time-reversal symmetry is preserved

$$\mathcal{H} = \mathcal{H}_0 + V = \sum_{\langle ij \rangle} c_i^\dagger t_{ij} c_j + V + h.c.$$

$$V = \sum_{\langle ij \rangle} \delta t_{ij} |i\rangle \langle j| + h.c., \quad P[\delta t] \propto \exp[-\delta t^2 / 2g_b]$$

LMTB (Lanczos method on Tight-binding)

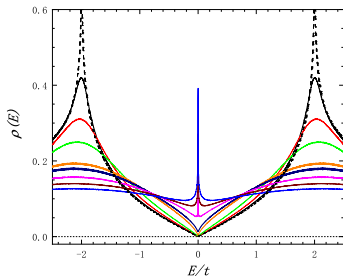
- large-scale accurate simulation in real/momentum space
- strong disorder (non-perturbative, beyond SCBA)

Physical Review B, 85 073407 (2012)

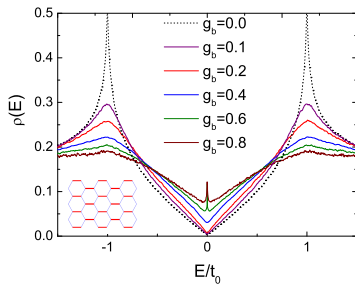
Average density of states

$$\rho(E) = -\frac{1}{\pi} \overline{\langle \psi | \frac{1}{E - H + i\gamma} | \psi \rangle}, \quad \gamma \sim 1/L$$

π -flux model



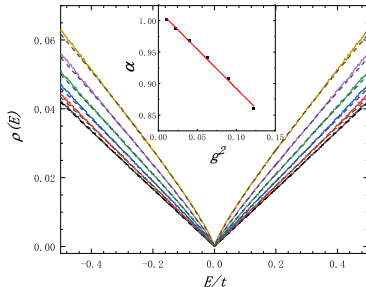
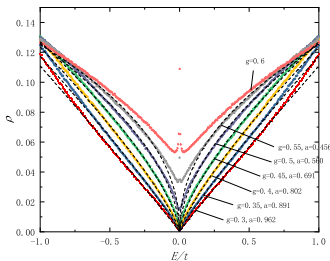
graphene



Singularity of DOS

Low energy DOS for weak disorder: Power Law

- $g_b = 0$, $\rho(E) = |E|$
- $g_b < g_b^C$, $\rho(E) = \rho_1 |E|^\alpha$, $\alpha = 1.05 - 1.24g_b$



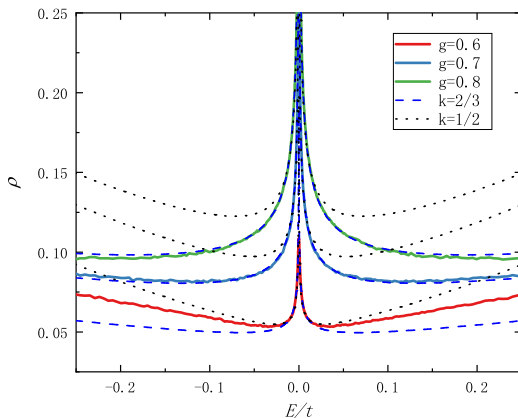
Born Approximation:

$$\rho(E) = -\frac{1}{\pi} \text{Im} \langle G_{rr} \rangle = \frac{A_c \lambda}{2\pi v^2} \left(\frac{E}{\lambda} \right)^{1-(2gA_c/\pi v^2)}$$

Singularity of DOS

Low energy DOS for strong disorder: Singularity

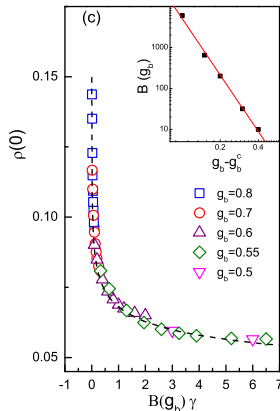
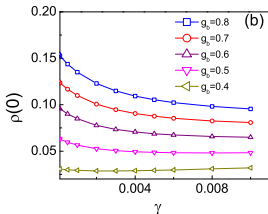
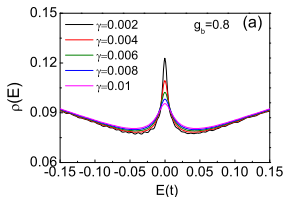
Gade's Model: $\rho(E) = \frac{1}{E} \exp(-c \ln|E|^{1/x})$, $x = \frac{3}{2}$ or 2? x = $\frac{3}{2}$



Gade, et al, Nucl. Phys. B 398, 499 (1993); Mudry, et al, Phys. Rev. B 67, 064202 (2003)

Low energy DOS for strong disorder

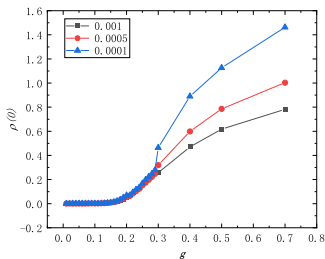
- $g_b > g_b^c$, $\rho(0) = \frac{A}{(B\gamma)^\nu}$, $B(g_b) = \frac{1}{(g_b - g_b^c)^\kappa}$
 with critical exponent $\nu = 0.12$ and $A = 0.07$



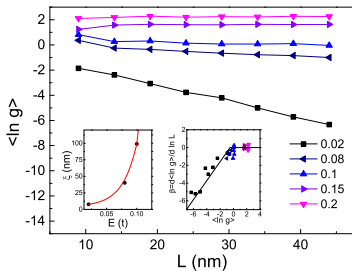
Crossover: semi-metal to diffusion metal

power-law behavior \rightarrow singular behavior

$\rho(0)$ vs g
 for different γ



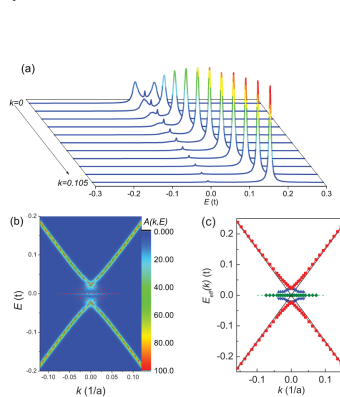
one-parameter scaling theory
 at $g = 0.7$



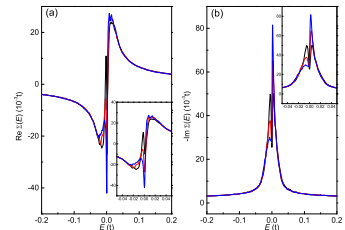
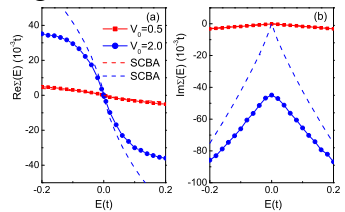
k -dependent self-energy and minimal conductivity

Accurate calculation of self-energy of Dirac fermions

spectral function and self-energy of strong disorder

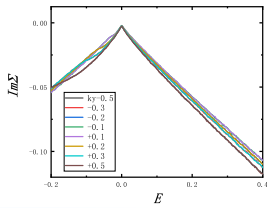
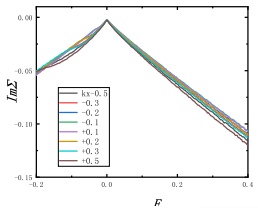
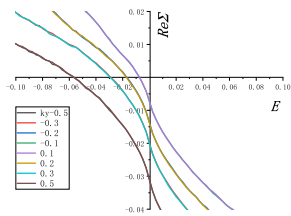
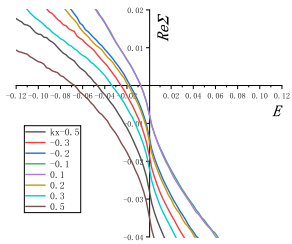


Zhu et al., Phys. Rev. B 85 .073407 (2012)



k -dependent self-energy and minimal conductivity

k -dependent of self-energy



k -dependent self-energy and minimal conductivity

Born approximation of self-energy

self-energy function $\Sigma(\mathbf{k}, E) = G_0^{-1}(\mathbf{k}, E) - G^{-1}(\mathbf{k}, E)$

Random hopping disorder

$$\Re\Sigma(k, E) = \gamma \left[-\frac{6E}{\pi} \ln \frac{\omega_c}{|E|} - \frac{3sa^2}{4\pi\hbar v_F} \left(\omega_c^2 + 2E^2 \ln \frac{\omega_c}{|E|} \right) k \right]$$

$$\Im\Sigma(k, E) = -\gamma \left(3|E| + \frac{3sa^2}{4\pi\hbar v_F} E^2 k \right)$$

$$\gamma = \frac{A_c n_i}{4\pi^2 \hbar v_F} \frac{\pi}{\hbar^2 v_F^2} \frac{w^2}{12}$$

lattice disorder

$$\Sigma(k, E) = \gamma \left[-\frac{2E}{\pi} \ln \frac{\omega_c}{|E|} - i|E| \right]$$

k -dependent self-energy and minimal conductivity

Minimal conductivity

Sample-dependent minimal conductivity $\sigma_{min} = \frac{1}{1-\alpha^2} \frac{4e^2}{\pi h}$

$$\Sigma(\mathbf{k}s, E) = \Sigma_1(E) - \alpha s \hbar v_F k + i \Sigma_2(E) = E - a - \alpha s \hbar v_F k + i \eta$$

$$\begin{aligned} \sigma_{xx}^0(E) &= \sigma_{xx}^{0,RA}(E) - \Re e[\sigma_{xx}^{0,RR}(E)] \\ &= \frac{2e^2}{\pi h} \frac{1}{1-\alpha^2} \left[1 + \left(\frac{a}{\eta} + \frac{\eta}{a} \right) \arctan \frac{a}{\eta} \right] = \frac{4e^2}{\pi h} \frac{1}{1-\alpha(g_b)^2} \Big|_{E \rightarrow 0} \end{aligned}$$

- scattering-independent value $\sigma_{min} = \frac{4e^2}{\pi h}$
- universal value $\sigma_{min} = \frac{4e^2}{h}$
- disorder-dependent value σ_{min}
- Experimental measurements : $\sigma_{min} \approx \frac{4e^2}{h} = C \times \frac{4e^2}{\pi h}$, $C = 1.7 \sim 10$

Nature 438,197 (2005), Nature 438,201 (2005), PNAS 102,10451(2005); PRL 99,216602(2007), PRL 98,076602(2007), PRL

98 256801(2007), PNAS 104 18392(2009)

Summary

Random hopping disorder in π -flux model

- 1 $g_b = 0, \quad \rho(E) = |E|$
- 2 $g_b < g_b^c, \quad \rho(E) = \rho_1 |E|^\alpha, \quad \alpha = 1.05 - 1.24g_b$
- 3 $g_b > g_b^c, \quad \rho(0) = \frac{A}{(B\gamma)^\nu}, \quad B(g_b) = \frac{1}{(g_b - g_b^c)^\kappa}$
- 4 $g_b > g_b^c, \quad \rho(E) = \frac{1}{E} \exp(-c \ln |E|^{2/3})$
- 5 Crossover: semi-metal to diffusion metal
- 6 k -dependent self-energy and sample-dependent minimal conductivity

Singularity of density of states in the pi-flux lattice, to be submitted

Criticism or suggestion welcome!