Exact solution of hydrodynamic and Boltzmann equations in longitudinally expanding systems

Shile Chen(陈诗乐) and SS, 2311.09575 SS, S. Jeon, and C. Gale, Phys.Rev.C(Letter) 105 (2022) 2, L021902

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hydrodynamics

Found (1+1)d analytical solutions

to rel. ideal hydro eqs. w/ simple EoS: $\mathcal{D}_{\mu}T^{\mu\nu} \equiv \partial_{\mu}T^{\mu\nu} + \Gamma^{\mu}_{\ \rho\mu}T^{\rho\nu} + \Gamma^{\nu}_{\ \rho\mu}T^{\rho\mu} = 0 \,,$ $T^{\mu\nu} = (\varepsilon + P) u^{\mu} u^{\nu} - P g^{\mu\nu},$ $P = c_s^2 \varepsilon \, .$

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hydrodynamics

Found

$$\begin{aligned} \frac{\varepsilon}{\varepsilon_{0}} &= \left(\frac{t_{0}}{\tau_{0}} + \frac{a\tau}{\tau_{0}}e^{\eta - \eta_{0}}\right)^{\frac{1 - c_{x}^{4}}{4c_{x}^{2}}\frac{1}{a^{2}} - \frac{(1 + c_{x}^{2})^{2}}{4c_{x}^{2}}}\left(\frac{t_{0}}{\tau_{0}} + \frac{\tau}{a\tau_{0}}e^{\eta_{0} - \eta}\right)^{\frac{1 - c_{x}^{4}}{4c_{x}^{2}}a^{2} - \frac{(1 + c_{x}^{2})^{2}}{4c_{x}^{2}}}, \\ u^{\tau} &= \frac{1}{2}\left(\sqrt{\frac{t_{0}e^{\eta_{0} - \eta} + \tau a}{t_{0}e^{\eta - \eta_{0}} + \tau / a}} + \sqrt{\frac{t_{0}e^{\eta - \eta_{0}} + \tau / a}{t_{0}e^{\eta_{0} - \eta} + \tau a}}\right), \\ u^{\eta} &= \frac{1}{2\tau}\left(\sqrt{\frac{t_{0}e^{\eta_{0} - \eta} + \tau a}{t_{0}e^{\eta - \eta_{0}} + \tau / a}} - \sqrt{\frac{t_{0}e^{\eta - \eta_{0}} + \tau / a}{t_{0}e^{\eta_{0} - \eta} + \tau a}}\right), \\ u^{x} &= u^{y} = 0. \end{aligned}$$

to rel. ideal hydro eqs. w/ simple EoS:

$$\mathcal{D}_{\mu}T^{\mu\nu} \equiv \partial_{\mu}T^{\mu\nu} + \Gamma^{\mu\nu}$$
$$T^{\mu\nu} = (\varepsilon + P) u^{\mu}$$
$$P = c_s^2 \varepsilon.$$

variables: τ : proper time η : spatial rapidity parameters (constants): C_{S} : speed of sound η_0 : Lorentz boost along \hat{z} τ_0 : scaling of time ε_0 : scaling of energy density *a*: [dimensionless] t_0 : [time unit]

 $\Gamma^{\mu}_{\ \rho\mu}T^{\rho\nu} + \Gamma^{\nu}_{\ \rho\mu}T^{\rho\mu} = 0,$ $^{\prime\prime}u^{\nu}-Pg^{\mu\nu},$

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proof

Coordinates, Metric and Connections

$$\begin{split} & \text{Im}[1] = \ \text{x1} = \{ \tau \operatorname{Cosh}[\eta], \tau \operatorname{Sinh}[\eta], \text{x}, \text{y} \}; \\ & \text{H} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}; \\ & \text{x2} = \{ \tau, \eta, \text{x}, \text{y} \}; \\ & \text{gDown} = \operatorname{Table} \Big[\operatorname{FullSimplify} \Big[\sum_{\alpha=1}^{4} \sum_{\beta=1}^{4} \operatorname{H}[\![\alpha, \beta]\!] \partial_{x2[\mu]} \text{x1}[\![\alpha]\!] \partial_{x2[\nu]} \text{x1}[\![\beta]\!] \Big], \\ & \{\mu, 4\}, \{\nu, 4\} \Big]; \\ & \text{gUp} = \operatorname{Inverse}[gDown]; \\ & \tau = \operatorname{Table} \Big[\operatorname{FullSimplify} \Big[\\ & \frac{1}{2} \sum_{\alpha=1}^{4} \operatorname{gUp}[\![\mu, \alpha]\!] \left(\partial_{x2[\nu]} \text{gDown}[\![\nu, \alpha]\!] + \partial_{x2[\nu]} \text{gDown}[\![\rho, \alpha]\!] - \\ & \partial_{x2[\nu]} \operatorname{gDown}[\![\nu, \rho]\!] \right) \Big], \{\mu, 4\}, \{\nu, 4\}, \{\rho, 4\} \Big]; \\ & \text{Print} \Big[\operatorname{"g}_{\mu\nu} = \operatorname{", MatrixForm}[gDown], \operatorname{", \t g}^{\mu\nu} = \operatorname{", } \\ & \operatorname{MatrixForm}[gUp], \operatorname{", \ nr=", r} \Big] \\ & g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -\tau^2 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}, \qquad g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -\frac{1}{\tau^2} & 0 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}, \\ & & \text{Figure int} \Big[\operatorname{"g}_{\mu\nu} = \operatorname{", \ nr=", r} \Big] \\ & \text{Figure int} \Big[\left\{ \{0, 0, 0, 0\}, \{0, \tau, 0, 0\}, \{0, 0, 0\}, \{0, 0,$$

Check the Solution

$$\begin{split} \mathsf{n}[\mathbf{0}] = \mathbf{p} = \mathbf{p} \Theta \left(\frac{\mathbf{t}\Theta}{\mathbf{t}\Theta} + \frac{\mathbf{a}\,\mathbf{\tau}}{\mathbf{t}\Theta} \,\mathbf{e}^{\eta-\eta\Theta} \right)^{\frac{1-c^4}{4c^2}} \frac{1}{a^2} - \frac{(1+c^2)^2}{4c^2}} \left(\frac{\mathbf{t}\Theta}{\mathbf{t}\Theta} + \frac{\mathbf{\tau}}{\mathbf{a}\,\mathbf{t}\Theta} \,\mathbf{e}^{\eta\Theta-\eta} \right)^{\frac{1-c^4}{4c^2}} \mathbf{a}^2 - \frac{(1+c^2)^2}{4c^2}}; \\ \mathbf{u}\tau = \frac{1}{2} \left(\sqrt{\frac{\mathbf{t}\Theta\,\mathbf{e}^{\eta\Theta-\eta} + \mathbf{\tau}\,\mathbf{a}}{\mathbf{t}\Theta\,\mathbf{e}^{\eta-\eta\Theta} + \mathbf{\tau}\,/\mathbf{a}}} + \sqrt{\frac{\mathbf{t}\Theta\,\mathbf{e}^{\eta-\eta\Theta} + \mathbf{\tau}\,/\mathbf{a}}{\mathbf{t}\Theta\,\mathbf{e}^{\eta\Theta-\eta} + \mathbf{t}\,\mathbf{a}}} \right]; \\ \mathbf{u}\eta = \frac{1}{2\,\tau} \left(\sqrt{\frac{\mathbf{t}\Theta\,\mathbf{e}^{\eta\Theta-\eta} + \mathbf{t}\,\mathbf{a}}{\mathbf{t}\Theta\,\mathbf{e}^{\eta-\eta\Theta} + \mathbf{t}\,/\mathbf{a}}} - \sqrt{\frac{\mathbf{t}\Theta\,\mathbf{e}^{\eta-\eta\Theta} + \mathbf{t}\,/\mathbf{a}}{\mathbf{t}\Theta\,\mathbf{e}^{\eta\Theta-\eta} + \mathbf{t}\,\mathbf{a}}}} \right]; \\ \mathbf{u} = \{\mathbf{u}\tau, \,\mathbf{u}\eta, \,\Theta, \,\Theta\}; \\ \mathsf{T}\mathsf{U}\mathsf{p} = \mathsf{Table}[(\mathbf{1} + \mathbf{c}^{-2})\,\mathsf{p}\,\mathsf{u}[\!\mu]\!]\,\mathsf{u}[\!\nu]\!] - \mathsf{p}\,\mathsf{g}\mathsf{U}\mathsf{p}[\!\mu], \,\nu]\!], \,\{\mu, \,4\}, \,\{\mathbf{v}, \,4\}]; \\ \mathsf{E}\mathsf{q} = \mathsf{Table}[\,\mathsf{FullSimplify}[\\ \mathsf{Sum}\left[\partial_{x2[\!\nu]}\,\mathsf{T}\mathsf{U}\mathsf{p}[\!\mu], \,\nu]\!] + \sum_{\alpha=1}^4 \left(\mathbf{r}[\!\mu], \,\alpha, \,\nu]\!]\,\mathsf{T}\mathsf{U}\mathsf{p}[\!\nu], \,\alpha]\!] + \mathbf{r}[\!\nu], \,\alpha, \,\nu]\!]\,\mathsf{T}\mathsf{U}\mathsf{p}[\!\mu], \,\alpha], \\ \{\mathbf{v}, \,\mathbf{1}, \,4\} \right], \\ \mathsf{Assumptions} \rightarrow \{\mathsf{p}\Theta > \Theta, \,\mathbf{a} > \Theta, \,\mathbf{\tau} > \Theta, \,\eta \in \mathsf{Reals}, \,\eta\Theta \in \mathsf{Reals}, \\ \tau\Theta > \Theta, \,\mathbf{t}\Theta > \Theta\} \right], \,\{\mu, \,4\} \right] \end{split}$$

Out[13]= $\{\Theta, \Theta, \Theta, \Theta\}$



parameter dependence

$$c_s = 1/\sqrt{3} \text{ (speed of sound)}$$

$$\eta_0 = 0 \text{ (Lorentz boost along } \hat{z}\text{)}$$

$$t_0 = 0.01\tau_0$$

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$$\sqrt{\frac{1 - c_s^2}{1 + c_s^2}} \le a \le \sqrt{\frac{1 + c_s^2}{1 - c_s^2}}$$

controls the degree of asymmetry

 $a \rightarrow 1/a$: parity reflection $\eta \rightarrow -\eta, \quad u^{\eta} \rightarrow -u^{\eta}$



 η



parameter dependence

$$c_s = 1/\sqrt{3}$$
 (speed of sound) $\frac{\tau/\tau_0}{1}$ $\eta_0 = 0$ (Lorentz boost along \hat{z}) $\frac{3}{5}$ $a = 1$ $\frac{9}{11}$

 $0 \le t_0$ controls the width of the plateau

$$\frac{\varepsilon}{\varepsilon_0} = \left(\frac{\tau^2 + 2t_0\tau\cosh\eta + t_0^2}{\tau_0^2}\right)^{-\frac{1+c_s^2}{2}} \\ = \left(\frac{(t+t_0)^2 - z^2}{\tau_0^2}\right)^{-\frac{1+c_s^2}{2}}$$

 t_0 : overlap time -- hydro starts at t = 0, collided at $t = -t_0$





phenomenology?

solid: *Qualitatively* consistent w/ exp. *Plateau Asymmetric* in pA

Quantitatively agreement if smeared



 η_p solid: Cooper-Frye from an isothermal surface dash: hadron cascade (gaussian convolution)



$$x^{\mu} = \{\tau, \eta, x, y\}$$

$$\begin{aligned} \frac{T_{\text{ideal}}}{T_i} &= \left(\frac{t_0}{\tau_i} + \frac{a\,\tau e^{\eta}}{\tau_i}\right)^{\frac{1-c_s^2}{4}\frac{1}{a^2} - \frac{1+c_s^2}{4}} \\ &\times \left(\frac{t_0}{\tau_i} + \frac{\tau e^{-\eta}}{a\,\tau_i}\right)^{\frac{1-c_s^2}{4}a^2 - \frac{1+c_s^2}{4}}, \\ u^{\tau} &= \frac{1}{2} \left(\sqrt{\frac{t_0 e^{-\eta} + \tau a}{t_0 e^{\eta} + \tau/a}} + \sqrt{\frac{t_0 e^{\eta} + \tau/a}{t_0 e^{-\eta} + \tau a}}\right), \\ u^{\eta} &= \frac{1}{2\tau} \left(\sqrt{\frac{t_0 e^{-\eta} + \tau a}{t_0 e^{\eta} + \tau/a}} - \sqrt{\frac{t_0 e^{\eta} + \tau/a}{t_0 e^{-\eta} + \tau a}}\right). \end{aligned}$$

$$\hat{x}^{\mu} = \{\hat{x}^0, \hat{x}^1, x, y\}$$



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$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{x}^{0} = \frac{2a\tau_{i}}{1+a^{2}} \left(\left(\frac{t_{0}+a\tau e^{\eta}}{\tau_{i}}\right)^{\frac{1}{a}} \left(\frac{t_{0}}{\tau_{i}}+\frac{\tau e^{-\eta}}{a\tau_{i}}\right)^{a} \right)^{\frac{1+a^{2}}{4a}},$$
$$\hat{x}^{1} = \frac{1+a^{2}}{4a} \ln \left(\left(\frac{t_{0}+a\tau e^{\eta}}{\tau_{i}}\right)^{\frac{1}{a}} / \left(\frac{t_{0}}{\tau_{i}}+\frac{\tau e^{-\eta}}{a\tau_{i}}\right)^{a} \right).$$





$$x^{\mu} = \{\tau, \eta, x, y\}$$

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\times \left(\frac{t_{0}}{\tau_{i}} + \frac{\tau e^{-\eta}}{a\,\tau_{i}}\right)^{\frac{1-c_{s}^{2}}{4}a^{2} - \frac{1+c_{s}^{2}}{4}}, \\
u^{\tau} = \frac{1}{2}\left(\sqrt{\frac{t_{0}e^{-\eta} + \tau a}{t_{0}e^{\eta} + \tau/a}} + \sqrt{\frac{t_{0}e^{\eta} + \tau/a}{t_{0}e^{-\eta} + \tau a}}\right), \\
u^{\eta} = \frac{1}{2\tau}\left(\sqrt{\frac{t_{0}e^{-\eta} + \tau a}{t_{0}e^{\eta} + \tau/a}} - \sqrt{\frac{t_{0}e^{\eta} + \tau/a}{t_{0}e^{-\eta} + \tau a}}\right).$$

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$$\hat{T}_{\text{ideal}} = \sqrt{\hat{g}^{00}} T_i \left(\frac{\tau_i}{\hat{x}^0}\right)^{c_s^2},$$
$$\hat{u}^0 = \sqrt{\hat{g}^{00}}, \quad \hat{u}^1 = \hat{u}^x = \hat{u}^y = 0.$$





$$x^{\mu} = \{\tau, \eta, x, y\}$$

$$\frac{T_{\text{ideal}}}{T_{i}} = \left(\frac{t_{0}}{\tau_{i}} + \frac{a\,\tau e^{\eta}}{\tau_{i}}\right)^{\frac{1-c_{s}^{2}}{4}\frac{1}{a^{2}} - \frac{1+c_{s}^{2}}{4}} \times \left(\frac{t_{0}}{\tau_{i}} + \frac{\tau e^{-\eta}}{a\,\tau_{i}}\right)^{\frac{1-c_{s}^{2}}{4}a^{2} - \frac{1+c_{s}^{2}}{4}}, \\
u^{\tau} = \frac{1}{2}\left(\sqrt{\frac{t_{0}e^{-\eta} + \tau a}{t_{0}e^{\eta} + \tau/a}} + \sqrt{\frac{t_{0}e^{\eta} + \tau/a}{t_{0}e^{-\eta} + \tau a}}\right), \\
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$$\hat{x}^{1} = \frac{1+a^{2}}{4a} \ln \left(\left(\frac{t_{0}+a\tau e^{\eta}}{\tau_{i}}\right)^{\frac{1}{a}} / \left(\frac{t_{0}}{\tau_{i}}+\frac{\tau e^{-\eta}}{a\tau_{i}}\right)^{a} \right).$$

$$\hat{g}^{\mu\nu} = \operatorname{diag}\left[e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -(\hat{x}^0)^{-2}e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -1, -1\right],$$

$$\hat{T}_{\text{ideal}} = \sqrt{\hat{g}^{00}} T_i \left(\frac{\tau_i}{\hat{x}^0}\right)^{c_s^2},$$
$$\hat{u}^0 = \sqrt{\hat{g}^{00}}, \quad \hat{u}^1 = \hat{u}^x = \hat{u}^y = 0.$$







$$\hat{x}^{\mu} = \{\hat{x}^0, \hat{x}^1, x, y\}$$

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"ture" 1+1D system: — no transverse coordinates

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$

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$$\hat{x}^{\mu} = \{\hat{x}^0, \hat{x}^1, x, y\}$$

Milne metric x global re-scaling

$$\hat{g}^{\mu\nu} = \text{diag}\left[1, -(\hat{x}^0)^{-2}\right] \times e^{2\frac{1-a^2}{1+a^2}\hat{x}^1},$$

$$\hat{g}^{\mu\nu} = \operatorname{diag}\left[e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -(\hat{x}^0)^{-2}e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -\frac{1}{1, -1}\right],$$

$$\hat{T}_{\text{ideal}} = \sqrt{\hat{g}^{00}} T_i \left(\frac{\tau_i}{\hat{x}^0}\right)^{c_s^2},$$
$$\hat{u}^0 = \sqrt{\hat{g}^{00}}, \quad \hat{u}^1 = \hat{u}^x = \hat{u}^y = 0.$$









"ture" 1+1D system: — no transverse coordinates

"false" 1+1D system: --- with transverse coordinates, but homogeneous...

$$\hat{x}^{\mu} = \{\hat{x}^0, \hat{x}^1, x, y\}$$

Milne metric x global re-scalir $-(\hat{x}^{0})^{-2}$ <u>- diag 1,</u> <u> ^ //1/</u>

$$\hat{g}^{\mu\nu} = \operatorname{diag}\left[e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -(\hat{x}^0)^{-2}e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -1, -1\right],$$

$$\begin{split} \hat{T}_{\text{ideal}} &= \sqrt{\hat{g}^{00}} \ T_i \left(\frac{\tau_i}{\hat{x}^0}\right) c_s^2, \\ \hat{u}^0 &= \sqrt{\hat{g}^{00}}, \quad \hat{u}^1 = \hat{u}^x = \hat{u}^y = 0. \end{split}$$









"ture" 1+1D system: — no transverse coordinates

"false" 1+1D system: --- with transverse coordinates, but homogeneous...

better properties in "true" 1+1D

$$\hat{x}^{\mu} = \{\hat{x}^0, \hat{x}^1, x, y\}$$

Milne metric x global re-scaling $\hat{g}^{\mu\nu} = \text{diag}\left[1, -(\hat{x}^0)^{-2}\right]$

$$\hat{g}^{\mu\nu} = \operatorname{diag}\left[e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -(\hat{x}^0)^{-2}e^{2\frac{1-a^2}{1+a^2}\hat{x}^1}, -1, -1\right],$$

$$\hat{T}_{\text{ideal}} = \sqrt{\hat{g}^{00}} T_i \left(\frac{\tau_i}{\hat{x}^0}\right) c_s^2,$$
$$\hat{u}^0 = \sqrt{\hat{g}^{00}}, \quad \hat{u}^1 = \hat{u}^x = \hat{u}^y = 0.$$









Boltzmann eq. w/ RTA

$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu} + \hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta}) = \mathscr{C}[f]\,,\\ & \mathscr{C}[f] = \frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta}) - f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{g}^{\mu\nu} = \text{diag}\left[e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -(\hat{x}^{0})^{-2}e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -1, -1\right],$$



Boltzmann eq. w/ RTA

$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu}+\hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta})=\mathscr{C}[f]\,,\\ & \mathscr{C}[f]=\frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta})-f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

formal analytical solutions:

- Milne coordinates
- Gubser coordinates
- our new coordinates? —

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{g}^{\mu\nu} = \text{diag}\left[e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -(\hat{x}^{0})^{-2}e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -1, -1\right],$$

[Florkowski, Ryblewski, Strickland PRC 88 024903]

[Denicol, Heinz, Martinez, Noronha, Strickland PRL 113 202301]



$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu} + \hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta}) = \mathscr{C}[f]\,,\\ & \mathscr{C}[f] = \frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta}) - f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

"ture" 1+1D system

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
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$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu} + \hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta}) = \mathcal{C}[f]\,, \\ & \mathcal{C}[f] = \frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta}) - f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

"ture" 1+1D system

$$\begin{split} \hat{T}_{\text{eff}}^2(\hat{x}^0) &= \frac{(\hat{x}^0)^2}{(\hat{x}^0_i)^2} e^{-\frac{1}{5\bar{\eta}} \int_{\hat{x}^0_i}^{\hat{x}^0_i} \hat{T}_{\text{eff}}(x') dx'} \hat{T}_0^2 \\ &+ \frac{1}{5\bar{\eta}} \int_{\hat{x}^0_i}^{\hat{x}^0} dx' \frac{(\hat{x}^0)^2}{(\hat{x}')^2} e^{-\frac{1}{5\bar{\eta}} \int_{\hat{x}'}^{\hat{x}^0} \hat{T}_{\text{eff}}(x'') dx''} \hat{T}_{\text{eff}}^3(\hat{x}') \,. \end{split}$$

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{g}^{\mu\nu} = \text{diag}\left[e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -(\hat{x}^{0})^{-2}e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -\frac{1}{1, -1}\right],$$

"false" 1+1D system



$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu} + \hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta}) = \mathcal{C}[f]\,, \\ & \mathcal{C}[f] = \frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta}) - f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

"ture" 1+1D system

$$\hat{T}_{\text{eff}}^{2}(\hat{x}^{0}) = \frac{(\hat{x}^{0})^{2}}{(\hat{x}_{i}^{0})^{2}} e^{-\frac{1}{5\bar{\eta}} \int_{\hat{x}_{i}^{0}}^{\hat{x}^{0}} \hat{T}_{\text{eff}}(x')dx'} \hat{T}_{0}^{2} \text{ decay}$$

$$+ \frac{1}{5\bar{\eta}} \int_{\hat{x}_{i}^{0}}^{\hat{x}^{0}} dx' \frac{(\hat{x}^{0})^{2}}{(\hat{x}')^{2}} e^{-\frac{1}{5\bar{\eta}} \int_{\hat{x}'}^{\hat{x}^{0}} \hat{T}_{\text{eff}}(x'')dx''} \hat{T}_{0}^{2}$$

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{g}^{\mu\nu} = \text{diag}\left[e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -(\hat{x}^{0})^{-2}e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -\frac{1}{1, -1}\right],$$

"false" 1+1D system

of initial state

 $\hat{T}_{eff}^{3}(\hat{x}')$ approach to equilibrium



$$\begin{split} & \left(\hat{p}^{\mu}\hat{\partial}_{\mu}+\hat{\Gamma}^{\rho}{}_{\mu\nu}\hat{p}^{\mu}\hat{p}_{\rho}\frac{\partial}{\partial\hat{p}_{\nu}}\right)f(\hat{x}^{\alpha},\hat{p}_{\beta})=\mathscr{C}[f]\,,\\ & \mathscr{C}[f]=\frac{\hat{p}_{\mu}\hat{u}^{\mu}(x)}{\tau_{r}(\hat{x})}\left(f(\hat{x}^{\alpha},\hat{p}_{\beta})-f_{eq}(\hat{x}^{\alpha},\hat{p}_{\beta})\right), \end{split}$$

"ture" 1+1D system

 $f(\hat{x}^0, \hat{p}_1, \hat{p}_T) = D(\hat{x}^0, \hat{x}_i^0)$

physics: non-trivial because rapidity dependence is introduced!

$$\hat{x}^{\mu} = \{\hat{x}^{0}, \hat{x}^{1}, x, y\}$$
$$\hat{g}^{\mu\nu} = \text{diag}\left[e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -(\hat{x}^{0})^{-2}e^{2\frac{1-a^{2}}{1+a^{2}}\hat{x}^{1}}, -1, -1\right],$$

"false" 1+1D system only for *a*=0 (symmetric)

$$f_0(\hat{x}_i^0, \hat{p}_1, \hat{p}_T) + \frac{1}{5\bar{\eta}} \int_{\hat{x}_i^0}^{\hat{x}^0} d\hat{x}' D(\hat{x}^0, \hat{x}') \hat{T}(\hat{x}') e^{-\frac{\hat{p}_0(\hat{x}')}{\hat{T}(\hat{x}')}}$$

mathematics: simply replacing variables $\{\tau \rightarrow \hat{x}^0, \eta \rightarrow \hat{x}^1\}$ in the Bjorken solution







physics: non-trivial because rapidity dependence is introduced!



solid:[exact solution] dash:[NS]





left: [exact solution] - [NS]

physics: non-trivial because rapidity dependence is introduced!

solid:[exact solution] dash:[NS]





physics: non-trivial because rapidity dependence is introduced!



summary and outlook

found analytical solutions in 1+1D expanding systems:

Hydrodynamics RTA Boltzmann "true" (1+1) RTA Boltzmann "false" (1+1)

- anisotropic hydro?
- analytical solutions in magnetohydro?

rapidity dependence

could it be asymmetric?

