

Weak magnetic effect in heavy-ion collision



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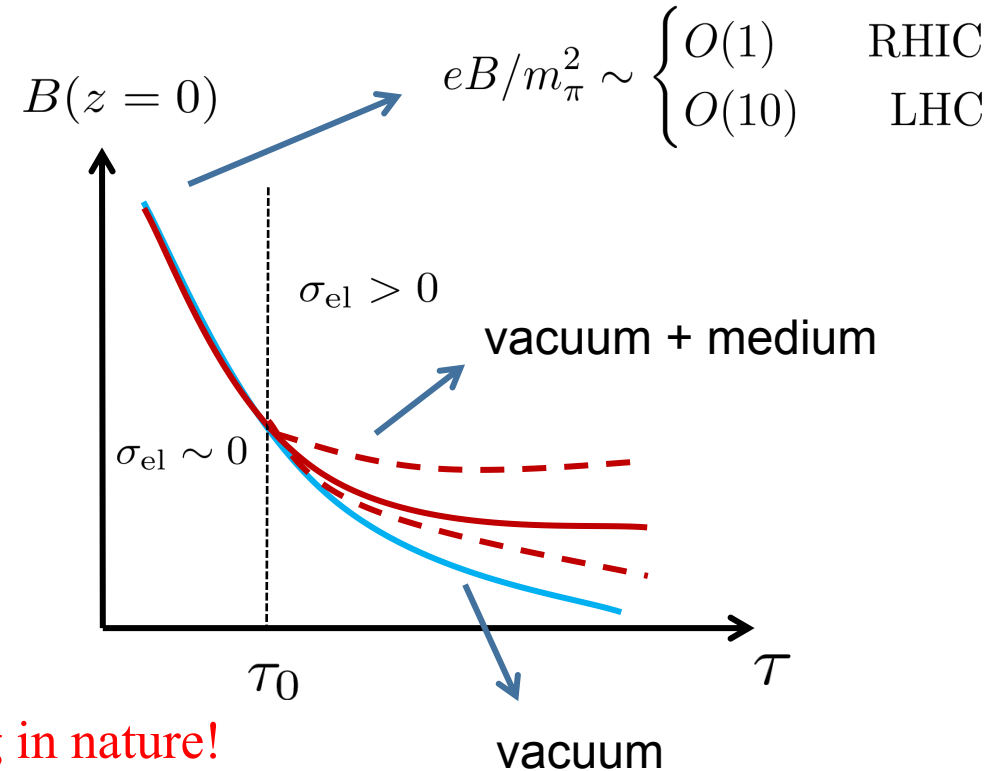
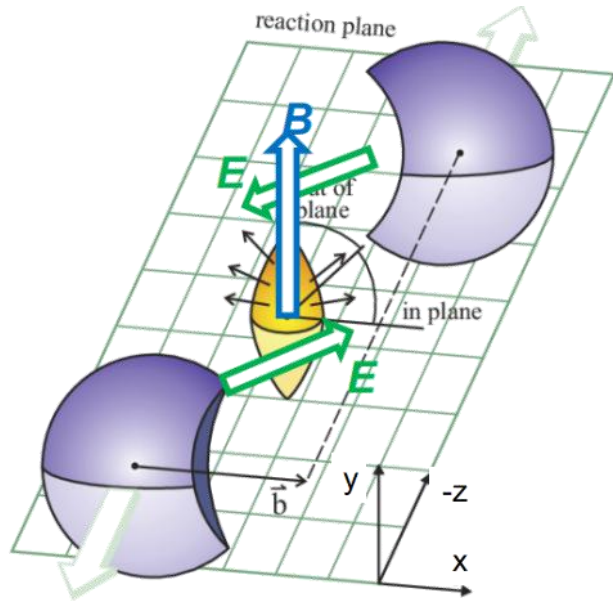
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Strong EM fields, UPC and EIC/EicC , USTC, 2024, 04/13

Electromagnetic field in heavy-ion collisions

- Generated from the relativistic motion of nucleus.
- B field is extremely strong initially, but becomes weak as QGP evolves hydrodynamically.

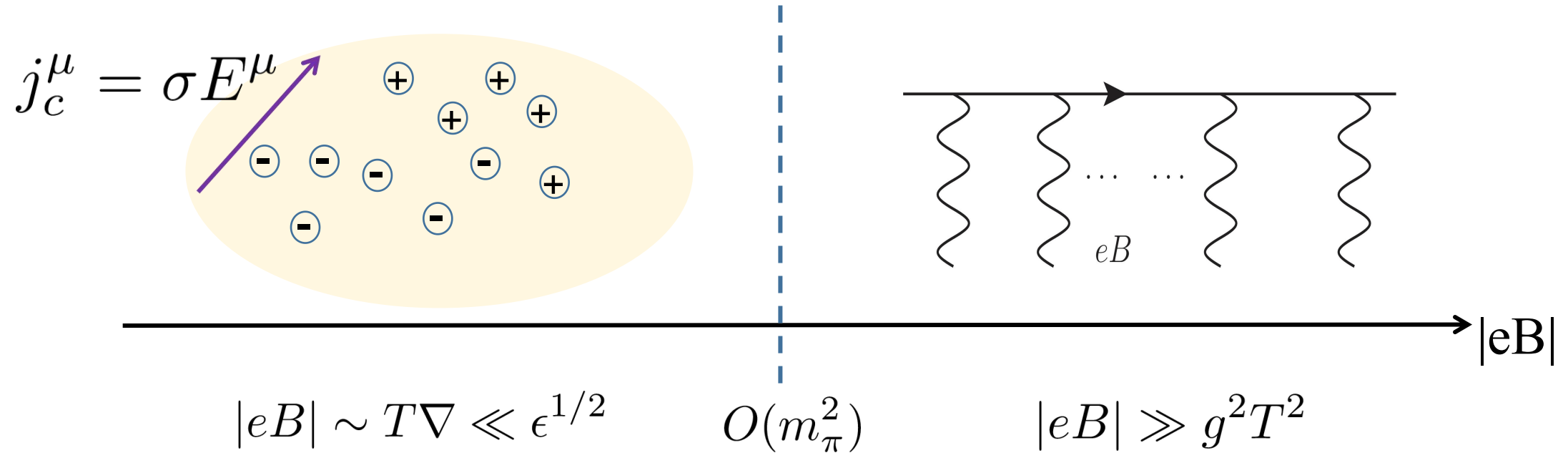
[Skokov and Bzdak 2012, Deng and Huang 2012, Kharzeev 2008, Tuchin 2010, Skokov and McLerran 2013, ...]



Weak field in terms of QCD scales, but still quite strong in nature!

Magnetic field in QGP: strong vs. weak

- effect of eB in a charged medium, e.g., QGP



Weak magnetic field and quark distribution

Solving transport eq. using Chapman-Enskog method,

$$p^\mu \partial_\mu f + q F^{\mu\nu} p_\mu \frac{\partial}{\partial p^\nu} f = C[f] \sim \frac{f - n_{\text{eq}}}{\tau_R}$$

at leading order in eB/T^2 one finds solution (scalar form solution),

$$\delta f_{\text{EM}} \sim \tau_R q F^{\mu\nu} p_\mu \frac{\partial}{\partial p^\nu} n_{\text{eq}} \quad \leftrightarrow \quad \Delta J_{\text{el}}^\mu = \sigma_{\text{el}} F^{\mu\nu} u_\nu$$

Concerning quark spin dof, it is similar but one needs to generalize n_{eq} to \mathcal{F}_{eq} .

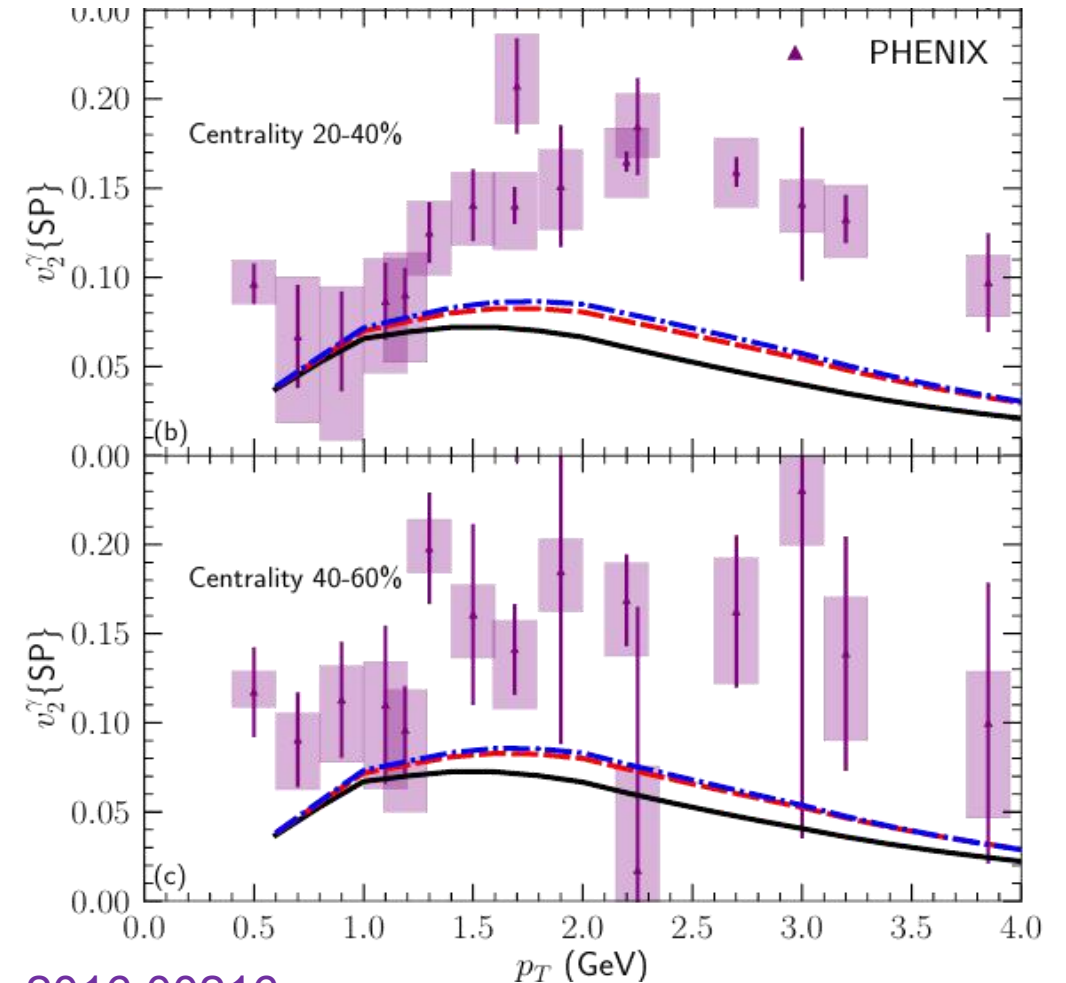
$$\delta \mathcal{F}_{\text{EM}} = -\frac{\bar{\tau}}{T} Q F^{\mu\nu} p_\mu \frac{\partial}{\partial p^\nu} \mathcal{F}_{\text{eq}}$$

Weak magnetic effects in

1. Direct photon production
2. Lambda hyperon local polarization
3. Virtual photon polarization

Up-to-date realistic eBe hydro predictions

- QGP with elastic and inelastic processes
- NNLO pQCD for prompt photons
- LO thermal AMY rate
- Hadron gas photon productions
- Dissipative corrections from shear and bulk
- Chemical equilibration in QGP
- ...
- **There is no EM field!**



[C. Gale et al, 2016.00216]

Weak magnetic photon emission from QGP

- Origin of momentum anisotropy:

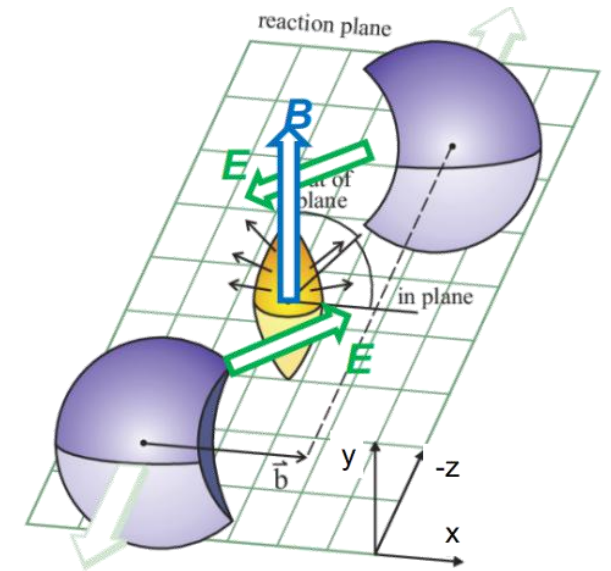
$$R_{\text{EM}}^\gamma \sim \delta f_{\text{EM}} \sim \frac{n'_{\text{eq}}}{\chi_n} p^\mu N_\mu = -\sigma_{\text{el}} \frac{n'_{\text{eq}}}{\chi_n} p^\mu u^\nu F_{\mu\nu}$$

properties of background QGP

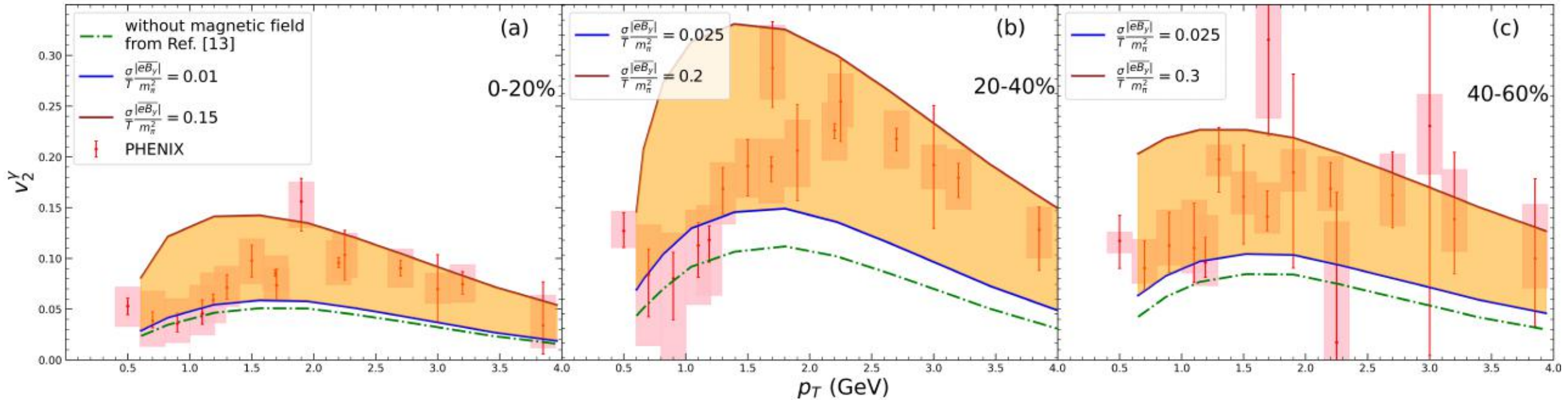
$$\ni \cos n\phi_p$$

properties of external E&B field, e.g.,

$$B_y p_T u^z \cos \phi_p$$



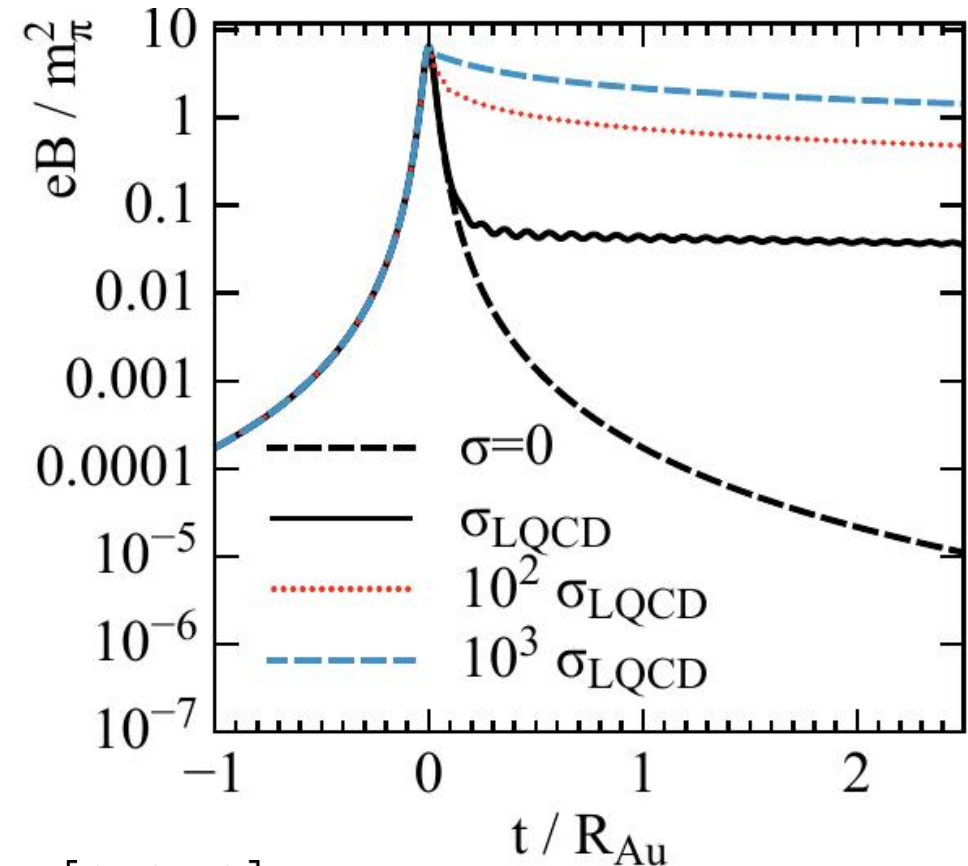
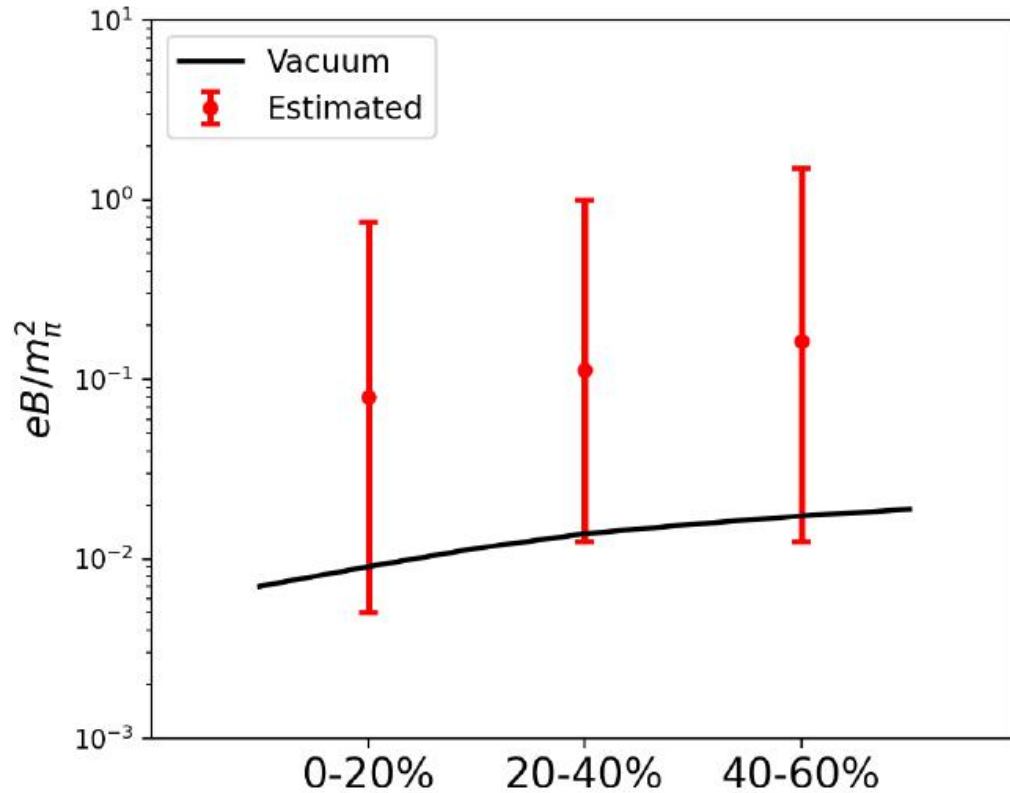
EbE hydro results: RHIC direct photon v2



- Experimental data can be reproduced, with a proper parameter $\rho \equiv \frac{\sigma_{\text{el}} \overline{eB_y}}{T m_\pi^2}$
- There is a systematic increase in the parameter.

Estimate of the time averaged eB

[Skokov and McLerran, 1305.0774]

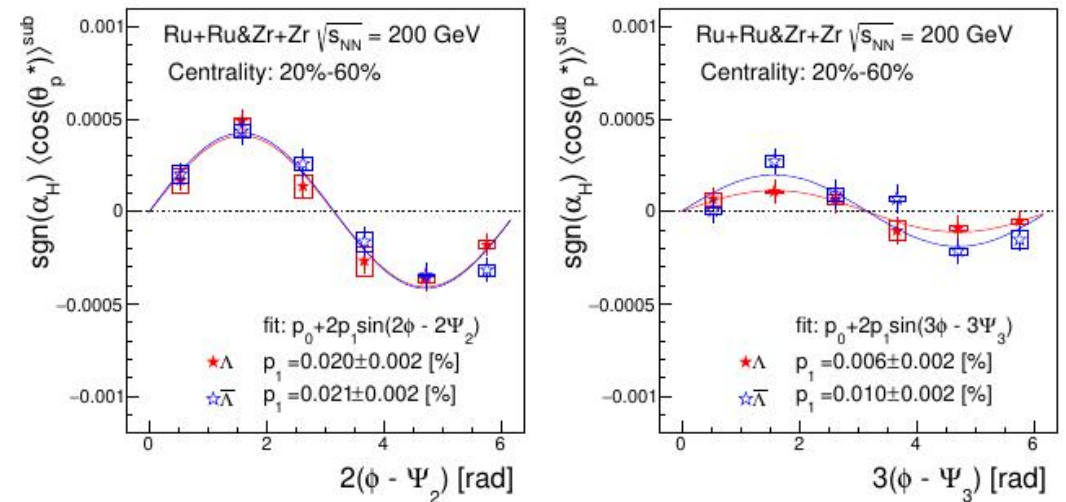
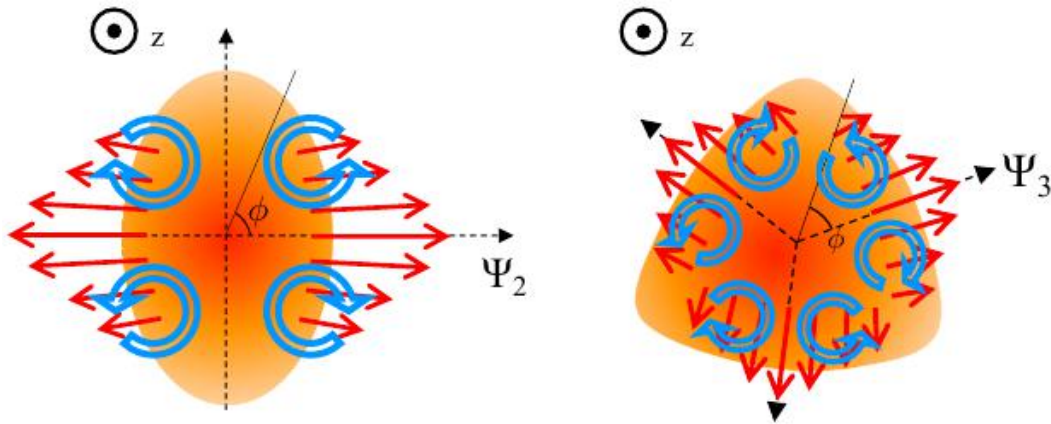


- Extract eB by direct photon v2, with $\sigma_{el}/T \in [0.2, 2]$ [S Floerchinger et al., 2112.12497]
- This is only a rough estimate of the upper bound of realistic eB.

Lambda hyperon polarization

[Liang and Wang, PRL 2006, The STAR collaboration, 2303.09074]

Local polarization along beam axis: spin of particles records rotating properties of QGP



- In thermal equilibrium: spin \sim thermal vorticity generated owing to initial geometry.
- Sign observed against naive hydro expectation using thermal vorticity: **sign problem**

Polarized particle from QGP

[F. Becattini et al., Annals of Physics 338 (2013) 32–49]

Spin of particles emitted from fluid: converted from thermal vorticity, SIP, etc.

- In equilibrium solution:

$$P^\mu(p) = \frac{p_\tau}{2m} \epsilon^{\mu\nu\rho\sigma} \frac{\int d\Sigma \cdot p n_{\text{eq}} \bar{\omega}^{\rho\sigma}}{\int d\Sigma \cdot p n_{\text{eq}}}$$

- With also SIP contribution: (LY and BBP)

[S. Y. F. Liu and Y. Yin, JHEP 07, 188 (2021)]

F. Becattini, et al, Pib 820, 136519 (2021), arXiv:2103.10917]

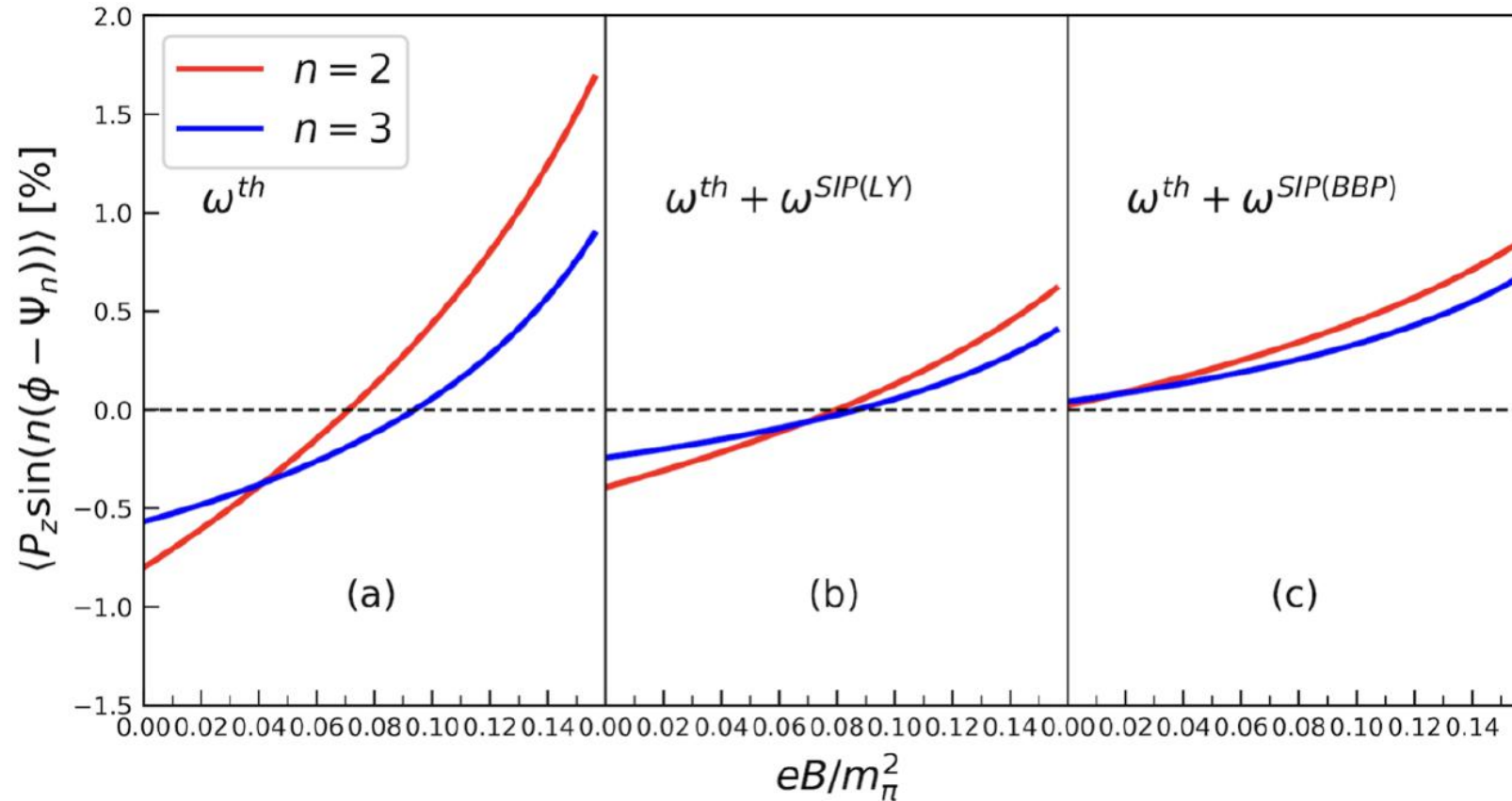
$$\bar{\omega}^{\rho\sigma} \rightarrow \bar{\omega}^{\rho\sigma} + \text{SIP}$$

- With weak EM dissipative contribution:

$$P^\mu(\mathbf{p}) = -\frac{1}{8m} \epsilon^{\mu\alpha\beta\sigma} p_\sigma \frac{\int d\Sigma \cdot p [n_F(1 - n_F) + (1 - 2n_F)\delta f_{\text{EM}}] \omega_{\alpha\beta}}{\int d\Sigma \cdot p (n_F + \delta f_{\text{EM}})}$$

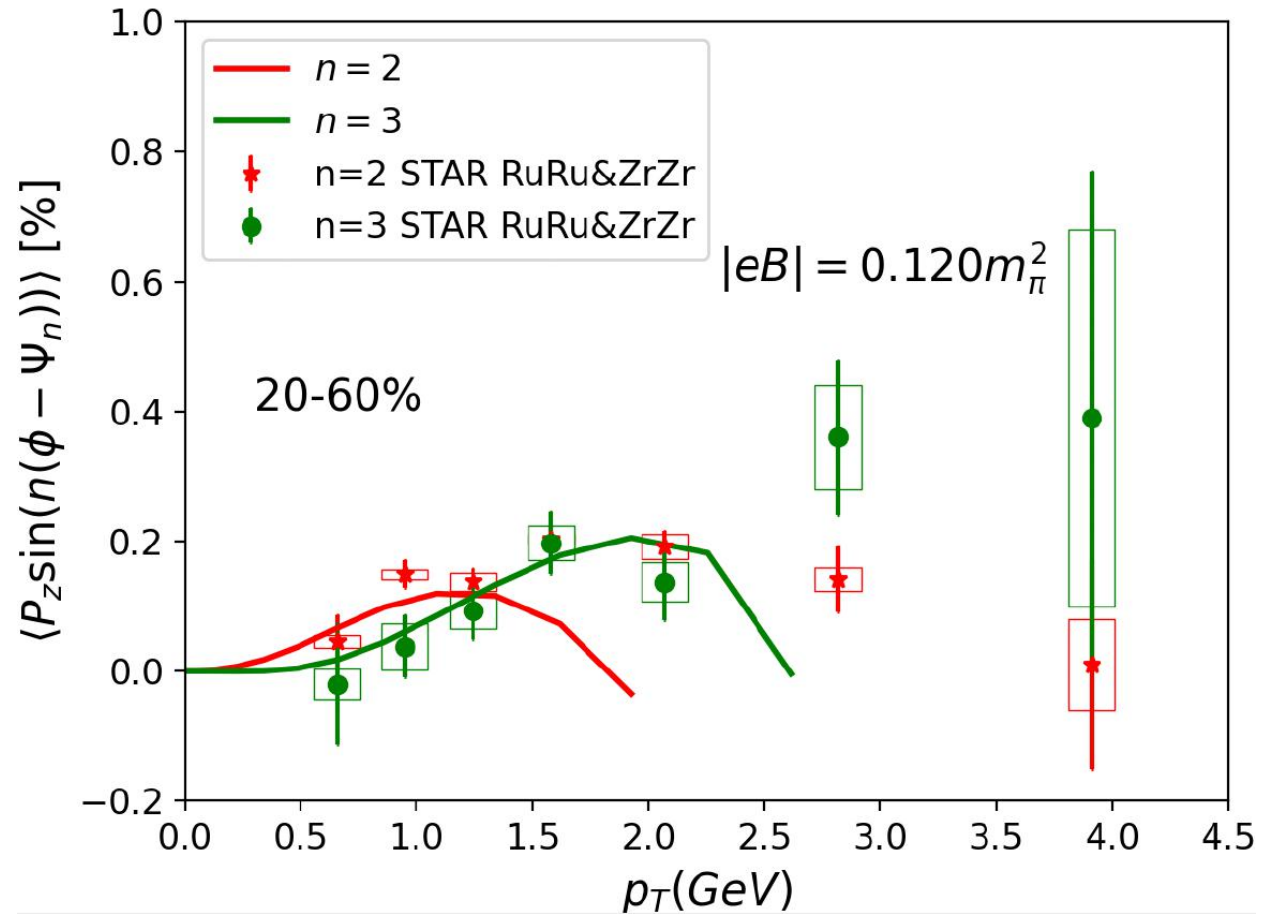
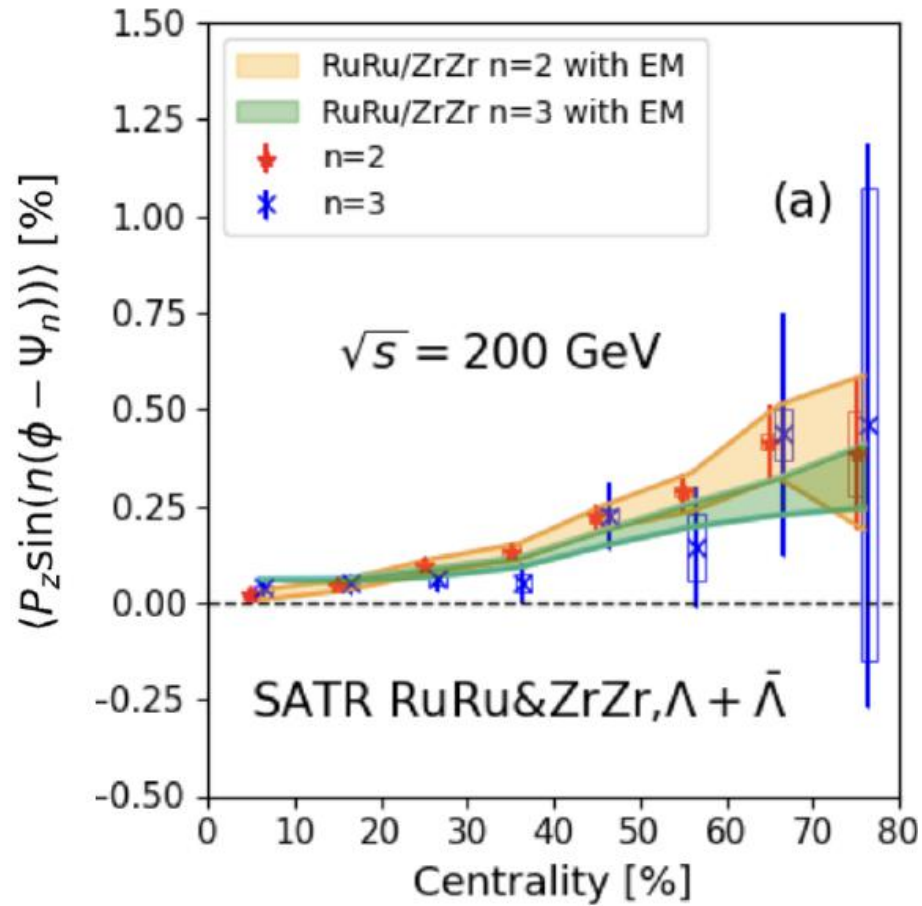
Global polarization merely affected!

Local polarization and weak B field (fixed centrality)



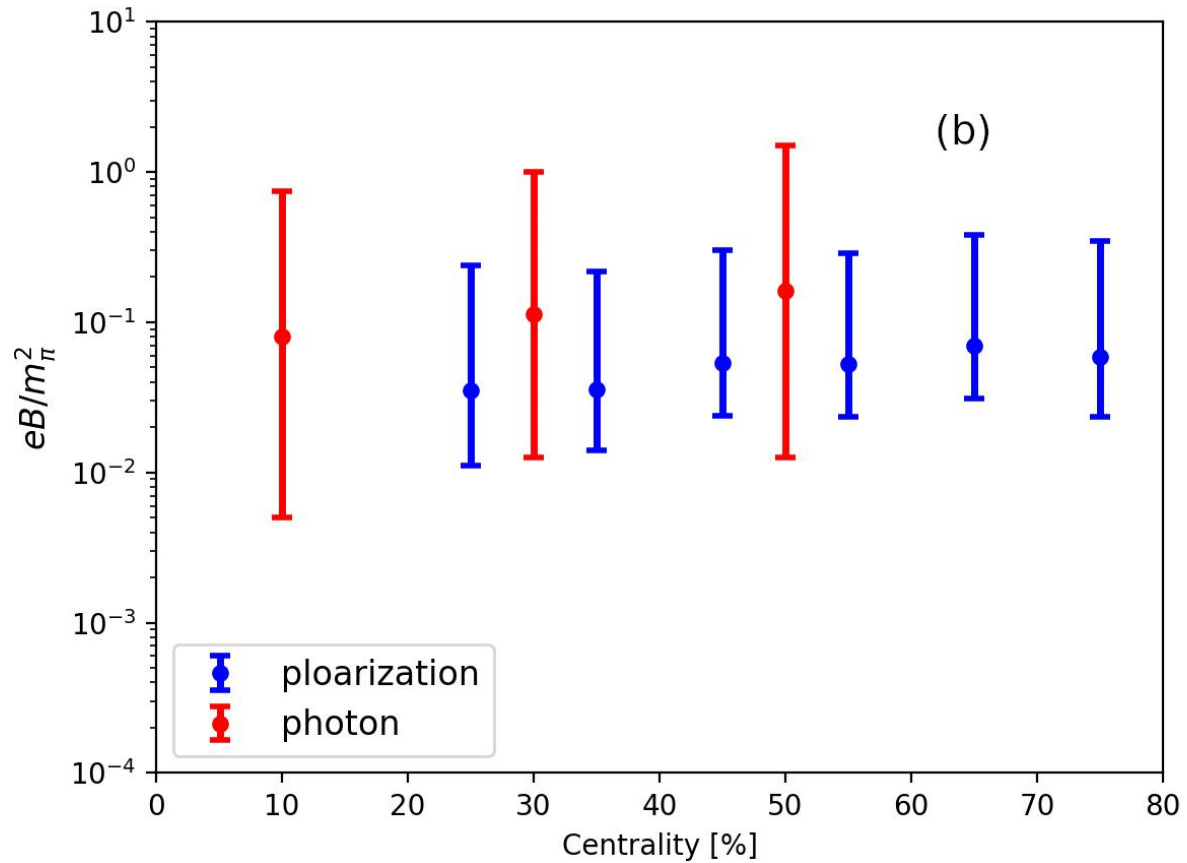
- Sign problem can be solved by a weak B field, in addition to SIP.
- Flip of ordering between $n=2$ and $n=3$ can **only** be found with a finite B field.

Confront experimental data



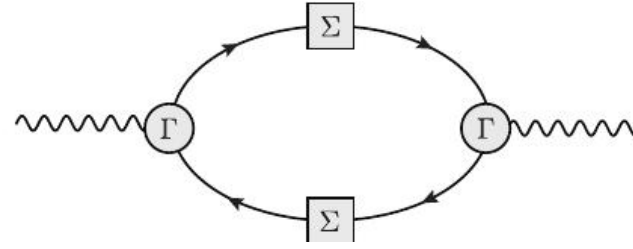
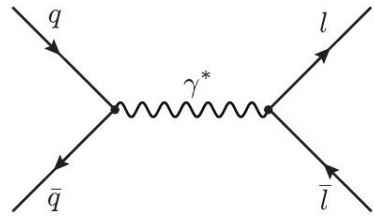
strangeness quark mass = 0.8 GeV

Estimate of eB from Lambda polarization



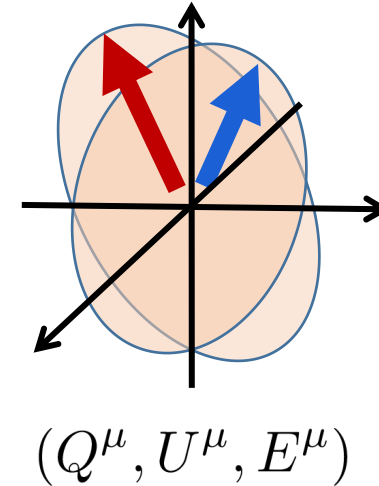
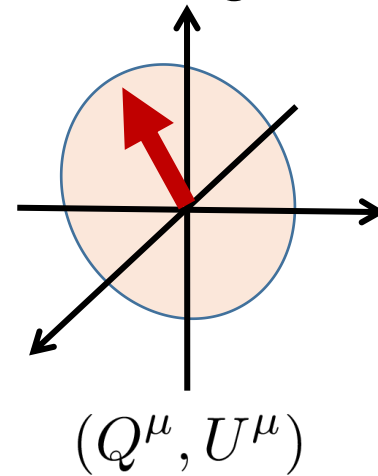
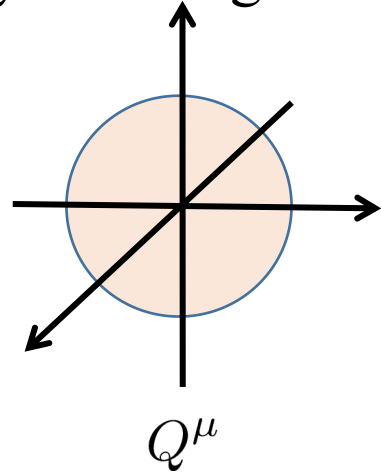
\Rightarrow consistent with the values from direct photon elliptic flow.

In QGP virtual photon polarization

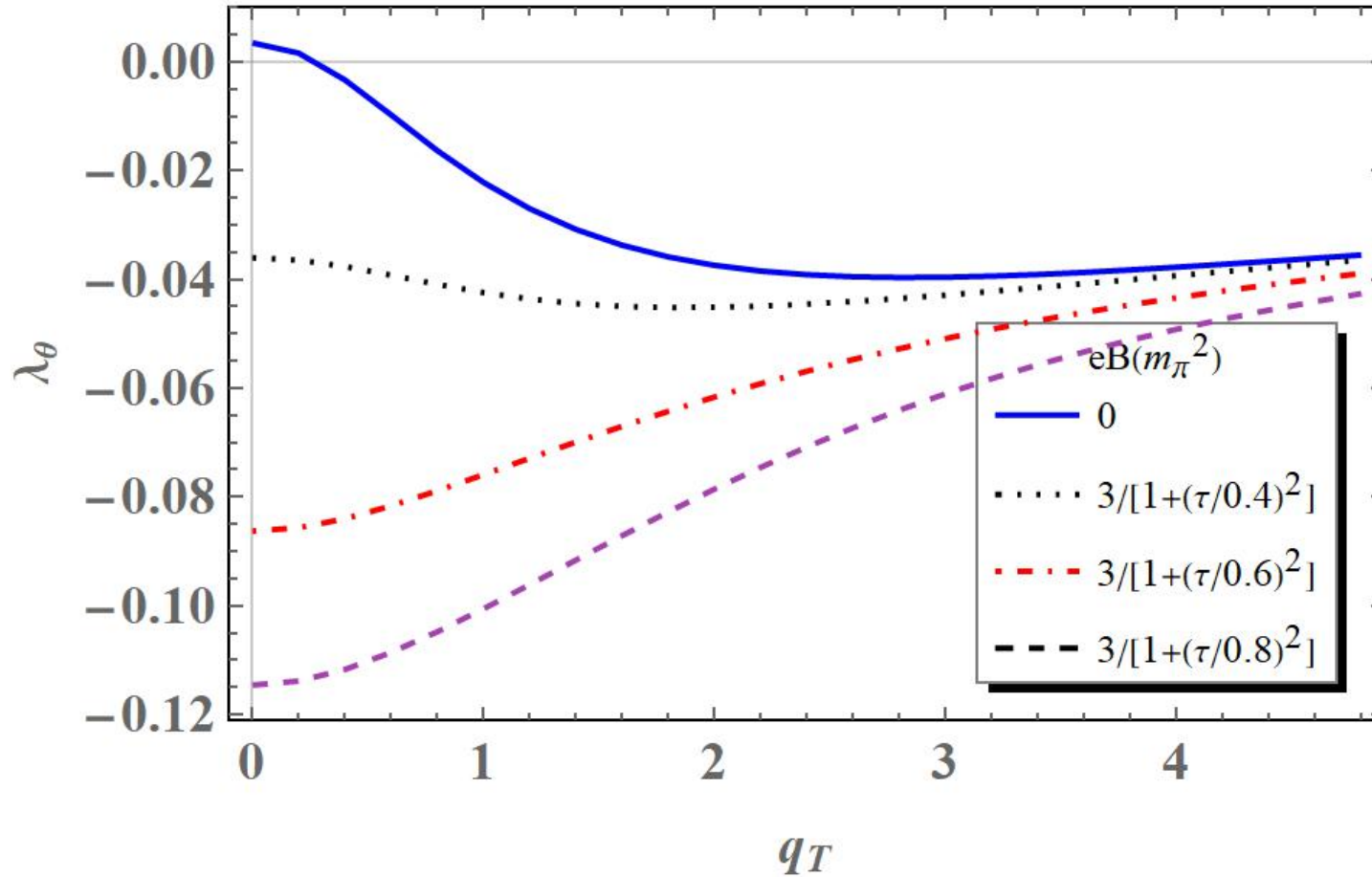


- Virtual photon approximates a intermediate massive spin-1 particle state.
- Polarization states of virtual photon determined by $q\bar{q}$ annihilation in QGP.
- Induced by geometry breaking vectors presenting in QGP.

$$\rho_{ab}^{\gamma^*} \sim \langle \Pi^{\mu\nu}(Q) \rangle$$



Numerical results (Bjorken flow)



$$\lambda_\theta = \frac{\rho_{++} + \rho_{--} - 2\rho_{00}}{\rho_{++} + \rho_{--} + 2\rho_{00}}$$

Helicity frame λ_θ

- $\rho_{00} > 1/3$ due to eB .
- Sensitive to eB life time.

Summary

- There must be (at least) weak EM fields in QGP.
- Weak EM fields result in dissipative corrections in QGP fluid.
- Consequences:
 1. A weak magnetic field can lead to significant direct photon v_2 .
 2. Change sign of local Lambda polarization
 3. Extra polarization of virtual photons.
 4. ...

Back-up slides

Weak magnetic field and QGP dissipation

- dissipation due to shear force

$$T^{\mu\nu} = T_{\text{ideal}}^{\mu\nu} + \pi^{\mu\nu}$$

$$\pi_{\mu\nu} \sim 2\eta \nabla_{\langle\mu} u_{\nu\rangle}$$

$$\delta f_{\pi} \sim \frac{n'_{\text{eq}}}{\chi_e} p^{\mu} p^{\nu} \pi_{\mu\nu}$$

- dissipation due to EM force

$$J^{\mu} = n_c u^{\mu} + \Delta J^{\mu}$$

$$\Delta J^{\mu} = \sigma_{\text{el}} F^{\mu\nu} u_{\nu} = \sigma_{\text{el}} E^{\mu}$$

$$\delta f_{\text{EM}} \sim \frac{n'_{\text{eq}}}{\chi_n} p^{\mu} \Delta J_{\mu}$$

Weak magnetic field and collective dynamics

1. Weak EM field does not change background (neutral) medium evolution.
2. Weak EM field leads to dissipative correction in conserved current,

$$J_{\text{el}}^{\mu} = n_c U^{\mu} + \sigma_{\text{el}} F^{\mu\nu} U_{\nu}$$

3. Charged components evolve slightly differently in weak EM field,

$$\partial_{\mu} \Delta T^{\mu\nu} = J_{\mu} F^{\mu\nu} \rightarrow \Delta u^{\mu} = u_{+}^{\mu} - u_{-}^{\mu} \sim O(Q F^{\mu\nu} / T^2) \neq 0$$

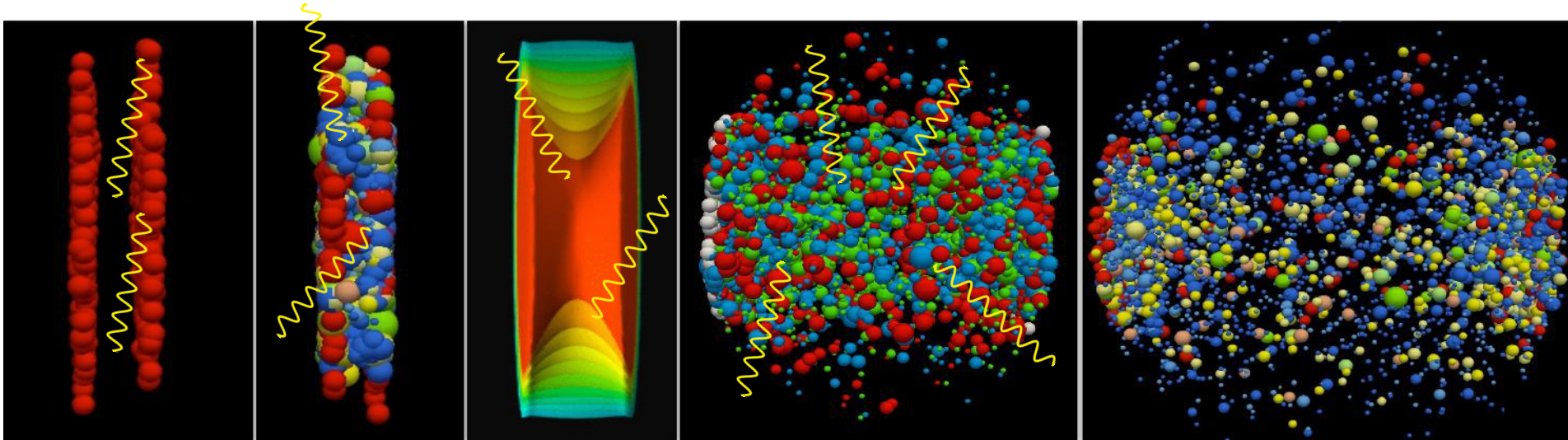
=> QGP bulk evolution is not affected, but charge dependent components feel the presence of the weak EM field, e.g., charged dependent flow

[The STAR collaboration, 2304,03430. U. Gürsoy, D. Kharzeev, and K. Rajagopal, PRC 89, 054905 (2014)]

Direct photons in heavy-ion collisions

[cf., G. David, 1907.08893]

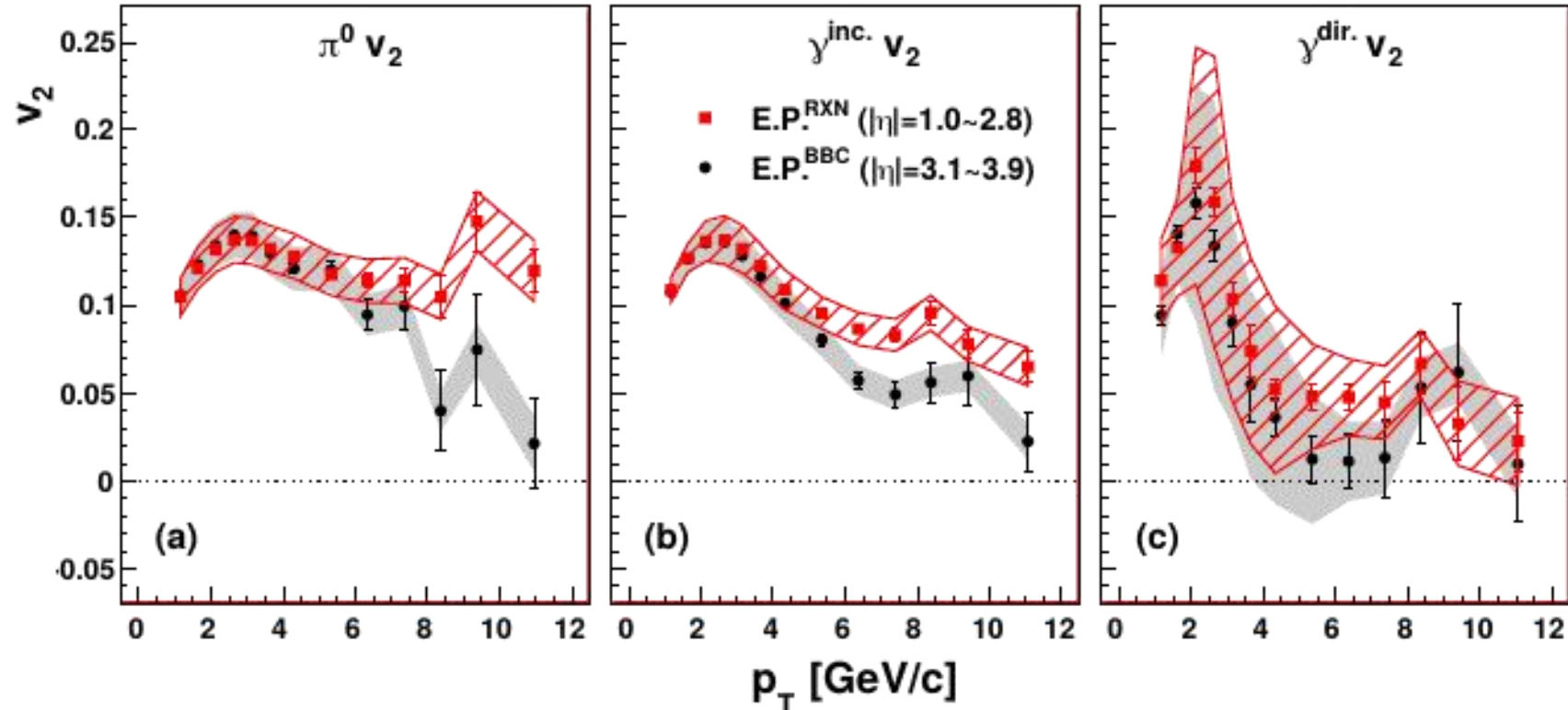
- Produced during the whole system evolution (exclude hadron decay)



- Direct photons = prompt + pre-equ. + **thermal from QGP and HG** + jet etc.

Direct photon puzzle

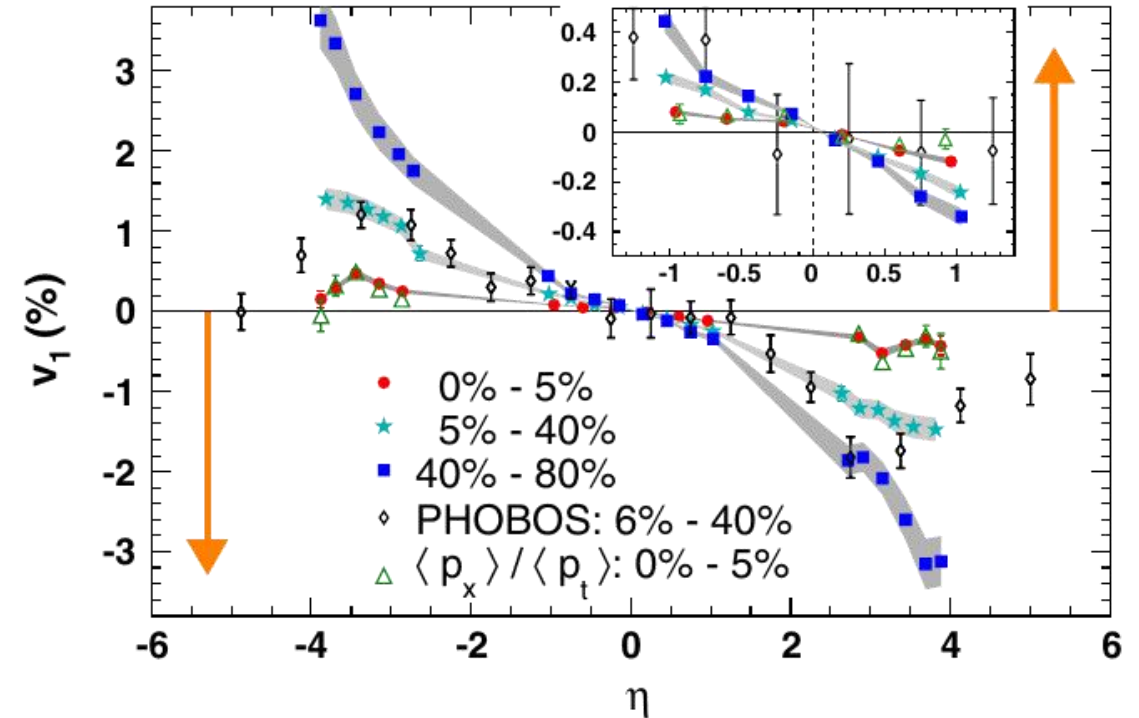
- Experimental observation: $v_2^\gamma \approx v_2^{\text{hadron}}$



[PHENIX collaboration, PRL 109, 122302, (2012)]

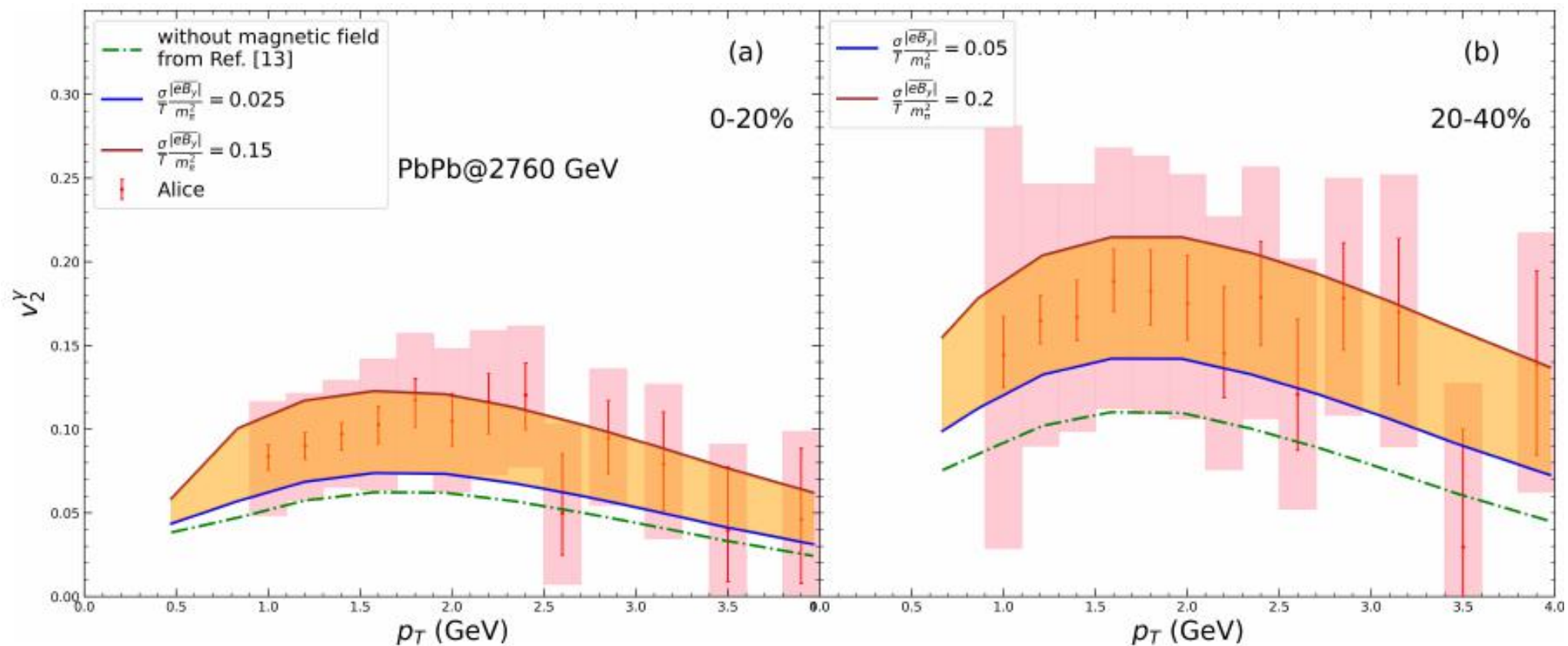
Photon elliptic flow and hadron v_1

- Elliptic flow of photon then requires in the background QGP a v_1 component.
[P. Bozek et al., 1101.3354]
- Exp. measured photons in even rapidity window \Rightarrow one needs a rapidity-odd v_1 in the background.
- Theory: EbE 3+1D hydro + weak B field to solve the direct photon puzzle.

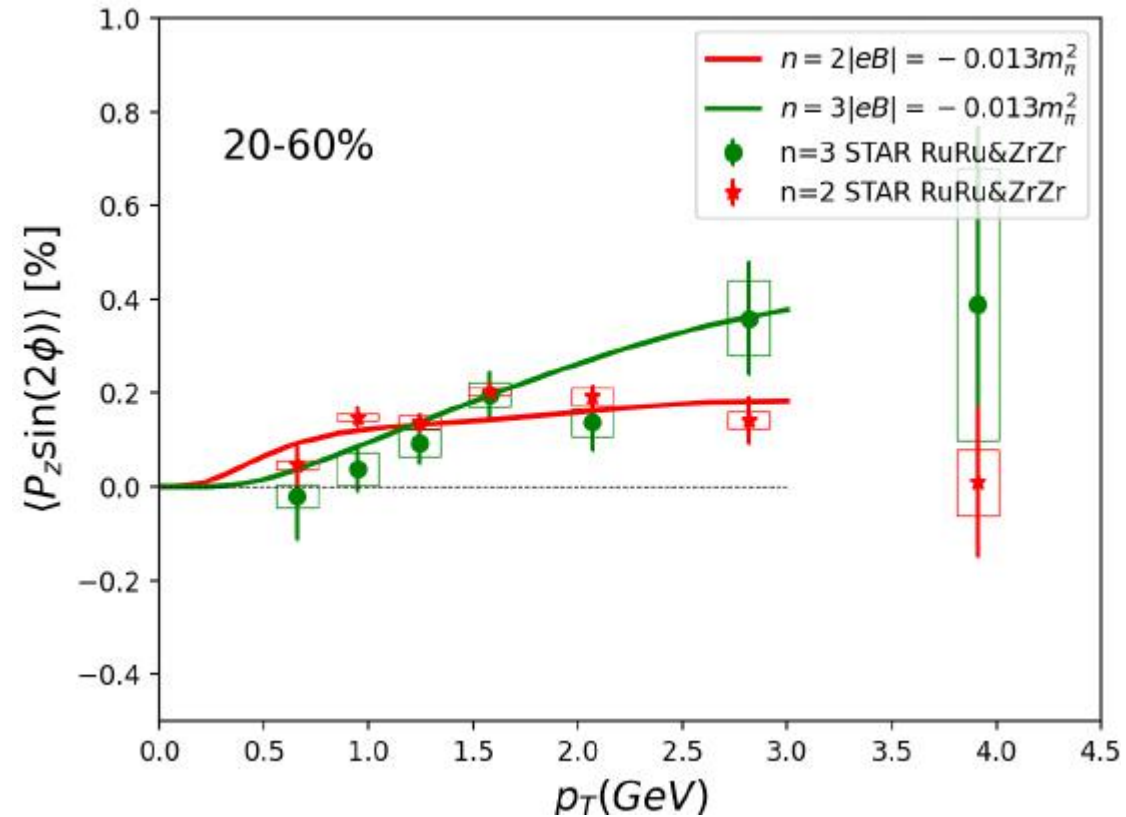


[STAR collaboration, PRL 101, 252301 (2008)]

EbE hydro results: LHC direct photon v2



Hydro results: differential local polarization



strangeness quark mass = 0.5 GeV