



A story of double charm tetraquark

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Outline

- A review of double charm tetraquark
- The molecular picture of the T_{cc}^+
- The production of the T_{cc}^+ in HIC
- Revealing the nature of the T_{cc}^+ in pp collision
- Summary and outlook

A review of double charm tetraquark

● 1964 Quark model

An SU(3) model for strong interaction symmetry and its breaking. Version 2 #1
G. Zweig (CERN) (Feb 21, 1964)
pdf cite claim reference search 657 citations

An SU(3) model for strong interaction symmetry and its breaking. Version 1 #2
G. Zweig (CERN) (Jan 17, 1964)
pdf cite claim reference search 818 citations

A Schematic Model of Baryons and Mesons #6
Murray Gell-Mann (Caltech) (1964)
Published in: *Phys.Lett.* 8 (1964) 214-215
DOI cite claim reference search 4,156 citations

● 1978 tetraquark in baryon-antibaryon system

PHYSICAL REVIEW D VOLUME 17, NUMBER 5 1 MARCH 1978

$Q^2\bar{Q}^2$ resonances in the baryon-antibaryon system

R. L. Jaffe

Center for Theoretical Physics, Laboratory for Nuclear Science and Department of Physics, Massachusetts Institute of Technology,
Cambridge, Massachusetts 02139

(Received 1 September 1977)

Two-quark-two-antiquark mesons which couple strongly to baryon-antibaryon channels are classified. The quantum numbers and masses of prominent states are predicted from the MIT bag model. The couplings of $Q^2\bar{Q}^2$ states to $B\bar{B}$ are estimated using the 3P_0 model and peripherality. Though most $Q^2\bar{Q}^2$ states do not couple strongly to $B\bar{B}$, many prominent resonances remain. Important $Q^2\bar{Q}^2$ resonances in the following processes are enumerated and discussed: elastic $N\bar{N}$ scattering, $N\bar{N} \rightarrow \pi^+\pi^-$, $N\bar{N}$ resonances at or below threshold, and exotic isotensor baryon-antibaryon resonances.

A review of double charm tetraquark

- 2002 the observation of the Ξ_{cc}^+ $ccq \leftrightarrow cc\bar{q}\bar{q}$

First Observation of the Doubly Charmed Baryon Ξ_{cc}^+ #1
SELEX Collaboration • M. Mattson (Carnegie Mellon U.) et al. (Aug, 2002)
Published in: *Phys.Rev.Lett.* 89 (2002) 112001 • e-Print: [hep-ex/0208014](#) [hep-ex]
pdf links DOI cite claim reference search 501 citations

- 2003 the observation of the $X(3872)$ $D\bar{D}^* \leftrightarrow DD^*$

- 2005 double charm tetraquarks by J.-M. Richard and FI. Stancu

Double charm hadrons revisited #5
J.-M. Richard (LPSC, Grenoble), FI. Stancu (Liege U.) (Nov, 2005)
Published in: *Bled Workshops Phys.* 6 (2005) 1, 25-31 • Contribution to: [Mini-Workshop Bled 2005](#), 25-31 •
e-Print: [hep-ph/0511043](#) [hep-ph]
pdf links cite claim reference search 12 citations

A review of double charm tetraquark

● 2003—2020 double charm tetraquark in molecular picture

Coupled-channel analysis of the possible $D^{(*)}D^{(*)}$,
 $\bar{B}^{(*)}\bar{B}^{(*)}$ and $D^{(*)}\bar{B}^{(*)}$ molecular states

$$D\bar{D}^* \leftrightarrow DD^*$$

Ning Li, Zhi-Feng Sun, Xiang Liu, and Shi-Lin Zhu
Phys. Rev. D **88**, 114008 – Published 3 December 2013

Sanchez, Geng, Lu, Hyodo, Valderrama, PRD98(2018)054001,

Xu, Wang, Liu, Liu, PRD99(2019)014027, Wang, Liu, Liu, PRD99(2019)036007,

Yu, Zhou, Chen, Xiao, PRD101(2020)0740270, Ding, Jiang, He, EPJC80(2020)1179.....

double charm tetraquark in compact tetraquark picture

$$ccq \leftrightarrow cc\bar{q}\bar{q}$$

Discovery of the Doubly Charmed Ξ_{cc} Baryon Implies a
Stable $bb\bar{u}\bar{d}$ Tetraquark

Marek Karliner and Jonathan L. Rosner
Phys. Rev. Lett. **119**, 202001 – Published 15 November 2017

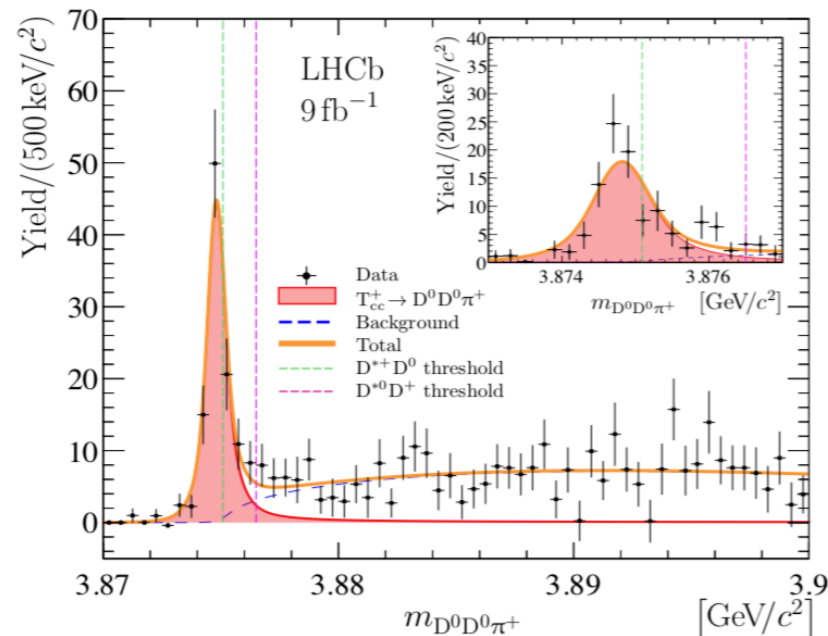
Karliner, Nussinov, JHEP07(2013)153, Cheng, Li, Si, Yao, CPC45(2021)043102,

Gelman, Nussinov, PLB551(2003)296-304, Luo, Chen, Liu, Liu, Zhu, EPJC77(2017)709.....

A review of double charm tetraquark

2021 The observation of $T_{cc}^+(cc\bar{u}\bar{d})$

Breit-Wigner fit LHCb, Nature Phys.18(2022)751-754



$$\delta m \equiv m_{T_{cc}^+} - m_{D^{*+}} - m_{D^0}$$

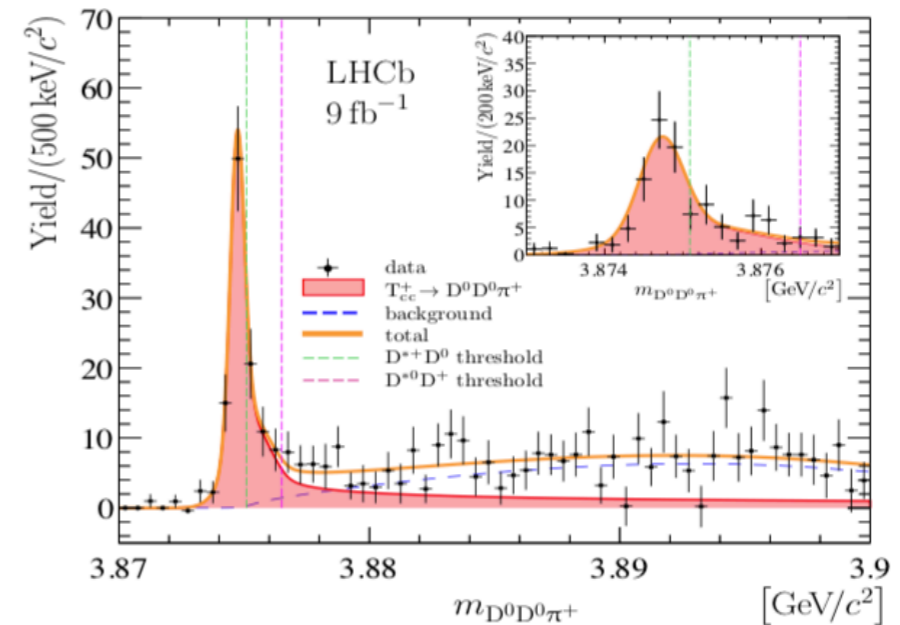
$$\delta m_{\text{BW}} = -273 \pm 61 \text{ keV}$$

$$\Gamma_{\text{BW}} = 410 \pm 165 \text{ keV}$$

No signal in $D^+ D^0 \pi^+$, $D^+ D^+$

→ I=0 isoscalar

Unitarized fit LHCb, Nature Commun. 13(2022)3351



$$\delta m_{\text{pole}} = -360 \pm 40_{-0}^{+4} \text{ keV}$$

$$\Gamma_{\text{pole}} = 48 \pm 2_{-14}^{+0} \text{ keV}$$

$$a = [-(7.16 \pm 0.51) + i(1.85 \pm 0.28)] \text{ fm}$$

$$-r < 11.9(16.9) \text{ fm} \quad 90(95) \% \text{ CL.}$$

$$Z < 0.52(0.58) \quad 90(95) \% \text{ CL.}$$

$$Z = 0 \text{ Composite} \quad Z = 1 \text{ Elementary}$$

The molecular picture of the T_{cc}^+

Why the D^*D molecule?

- Close to the D^*D thresholds
- Approximate 90% of $D^0D^0\pi^+$ events contain a D^{*+}
- $Z < 0.52$

Wave functions for isospin singlet and triplet

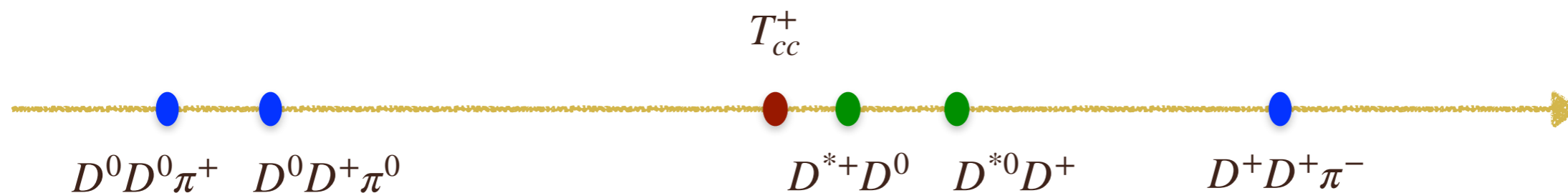
$$|D^*D, I = 0\rangle = -\frac{1}{\sqrt{2}} (D^{*+}D^0 - D^{*0}D^+)$$

$$V_{CT}^{I=0}(D^*D \rightarrow D^*D, J^P = 1^+) = v_0$$

$$|D^*D, I = 1\rangle = -\frac{1}{\sqrt{2}} (D^{*+}D^0 + D^{*0}D^+)$$

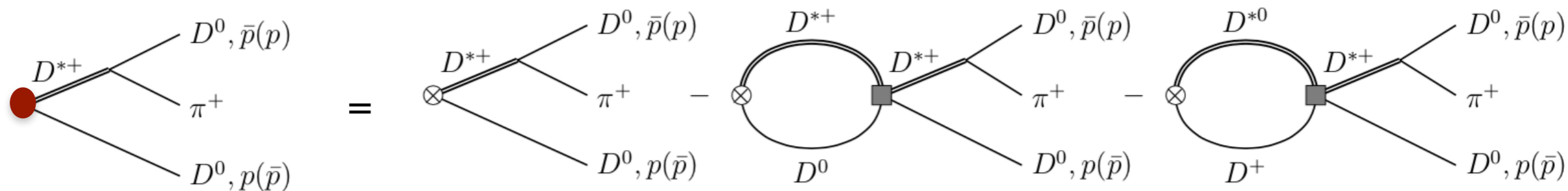
$$V_{CT}^{I=1}(D^*D \rightarrow D^*D, J^P = 1^+) = v_1$$

Three-body cut has to be considered

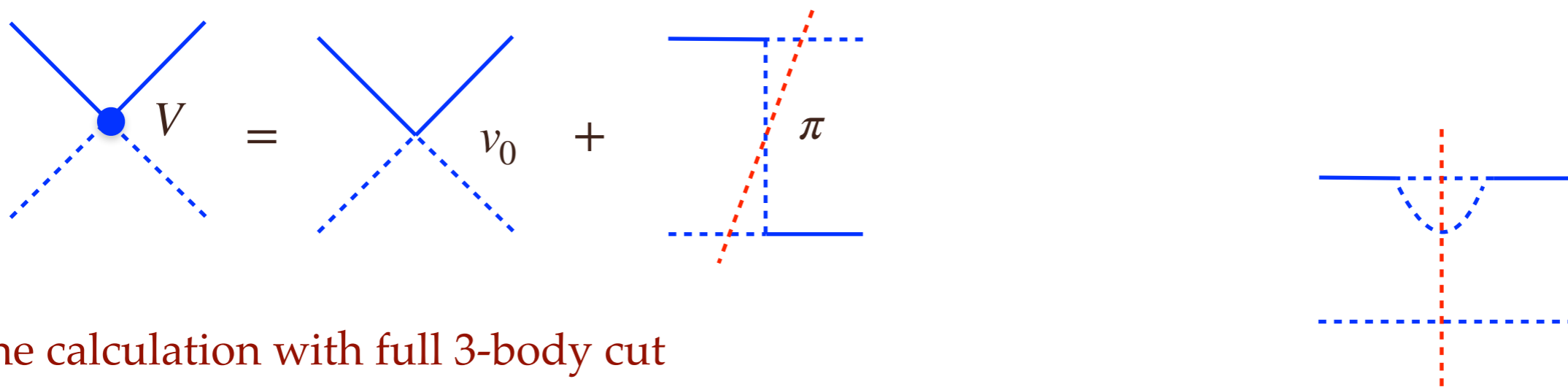
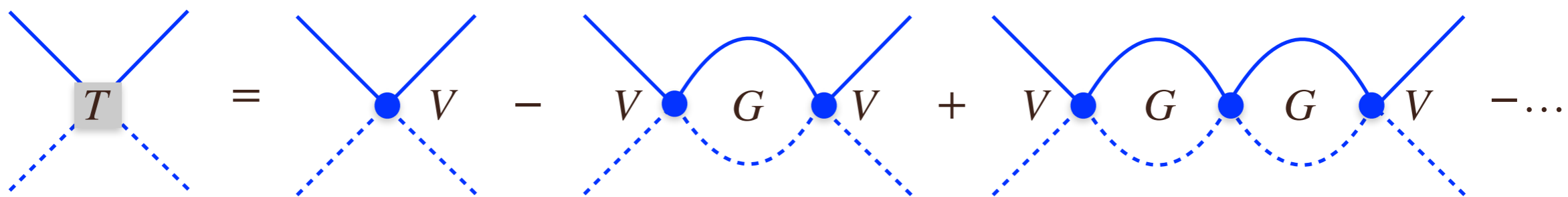


The molecular picture of the T_{cc}^+

$D^0 D^0 \pi^+$ mass distribution



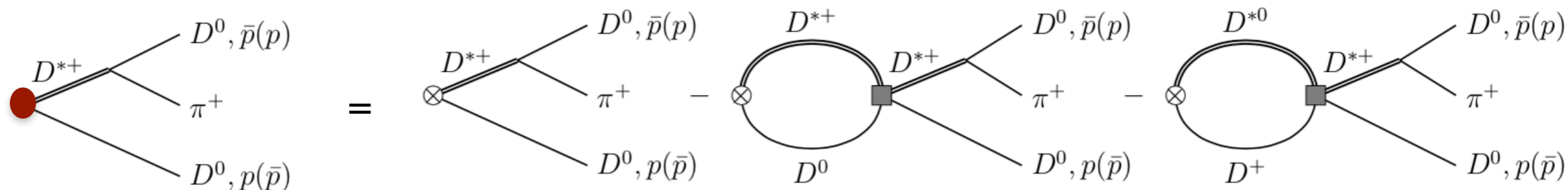
LSE



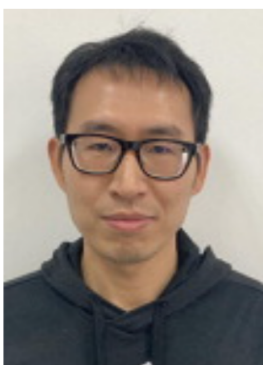
The calculation with full 3-body cut

The molecular picture of the T_{cc}^+

$D^0 D^0 \pi^+$ mass distribution



Schemes	Potential	Pole (keV)	Width (keV)
Scheme I	No OPE $\Gamma_{D^{*+}} = 82.5 \text{ keV}$ $\Gamma_{D^{*0}} = 53.7 \text{ keV}$	$-368_{-42}^{+43} - i(37 \pm 0)$	74
Scheme II	No OPE Dynamical widths of D^*	$-333_{-36}^{+41} - i(18 \pm 1)$	36
Scheme III	OPE Dynamical widths of D^*	$-356_{-38}^{+39} - i(28 \pm 1)$	56



Width is not as large as 400keV.

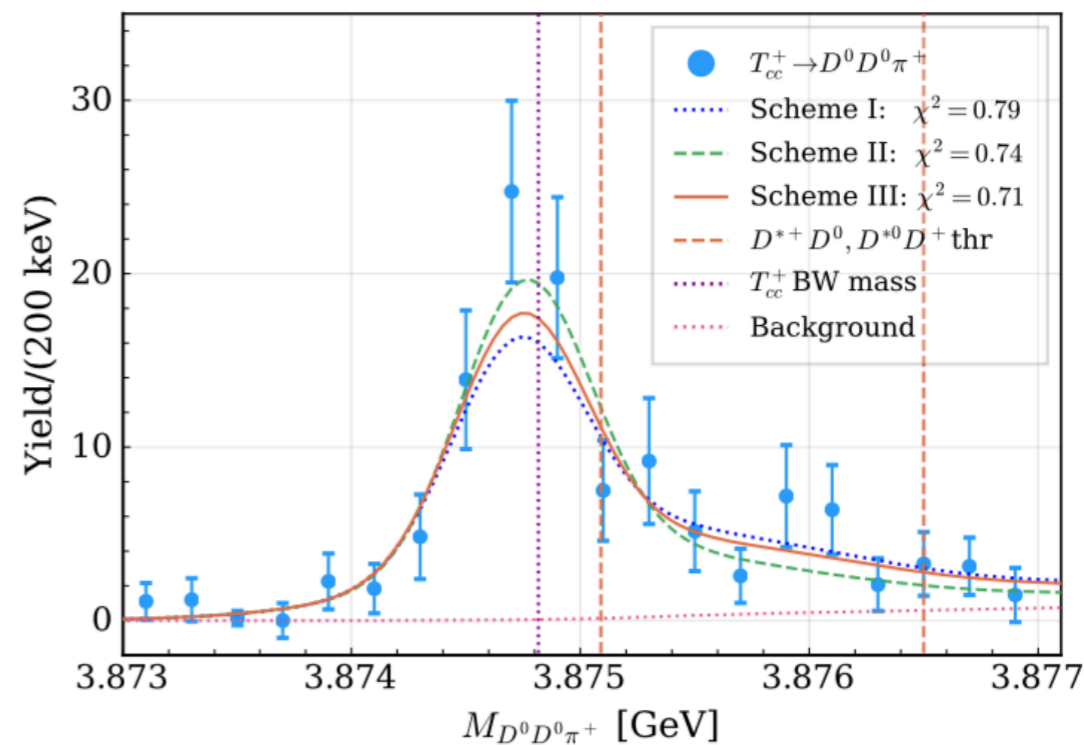
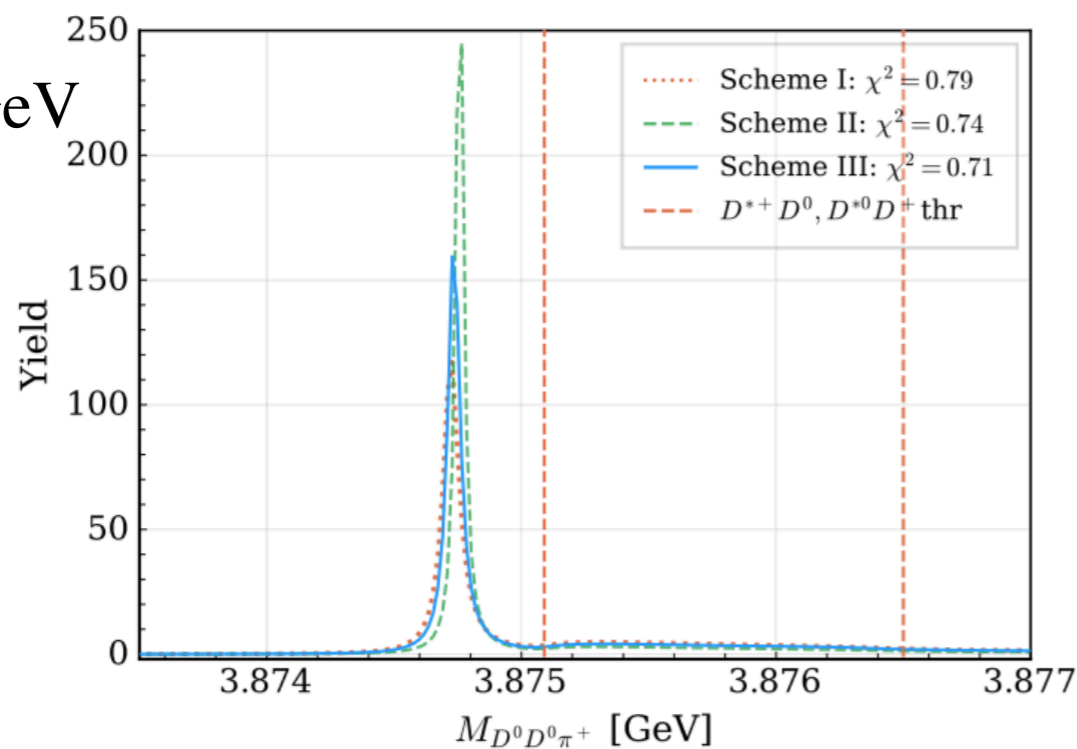
Du, Baru, Dong, Filin, Guo, Hanhart, Nefediev, Nieves, QW, PRD105(2022)014024

The molecular picture of the T_{cc}^+

$D^0 D^0 \pi^+$ mass distribution

Cutoff independent for $\Lambda = [0.3 - 1.2]$ GeV

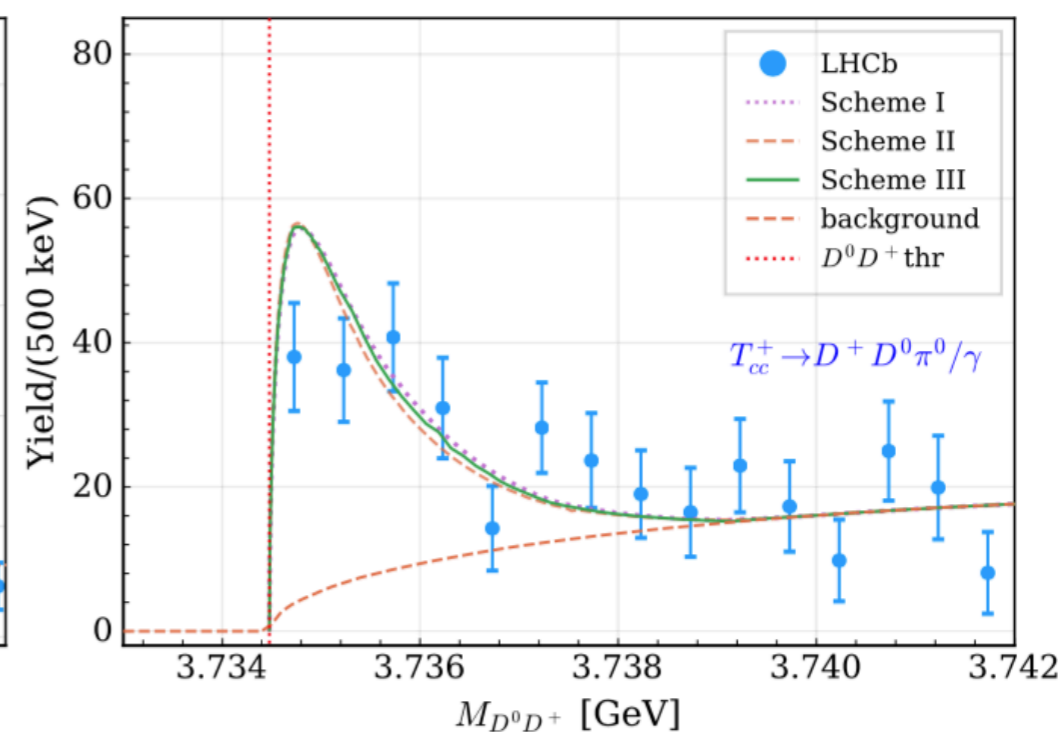
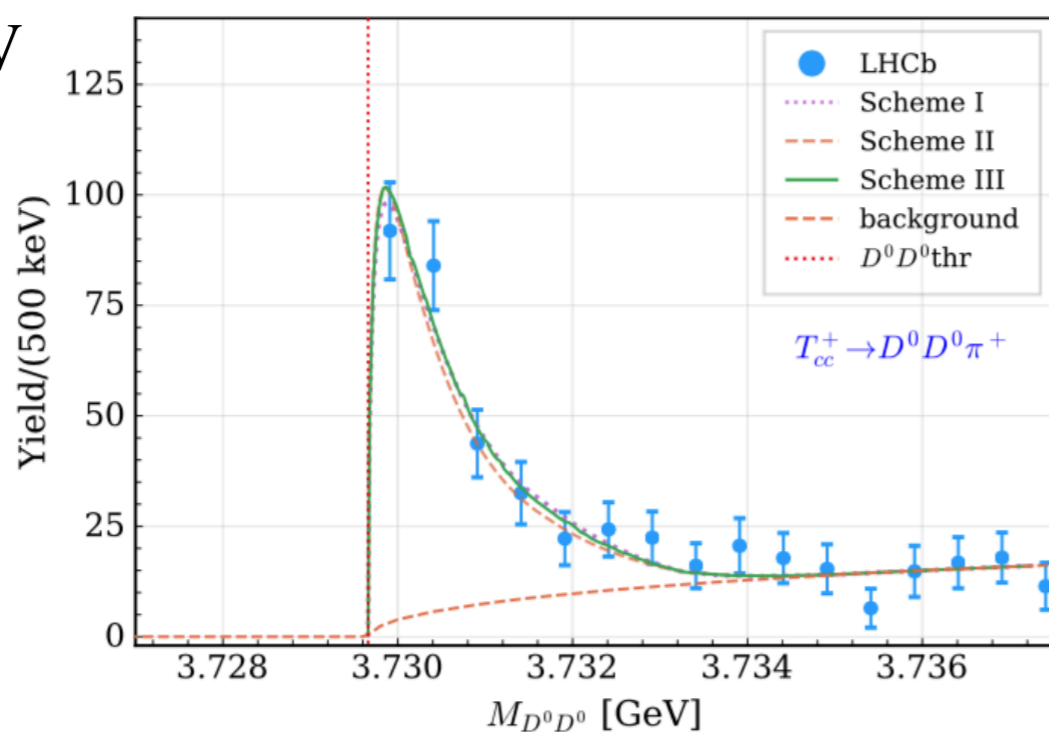
$\Lambda = 0.5$ GeV



$D^0 D^0$ and $D^0 D^+$ mass distribution

Du, et.al., PRD105(2022)014024

$\Lambda = 0.5$ GeV



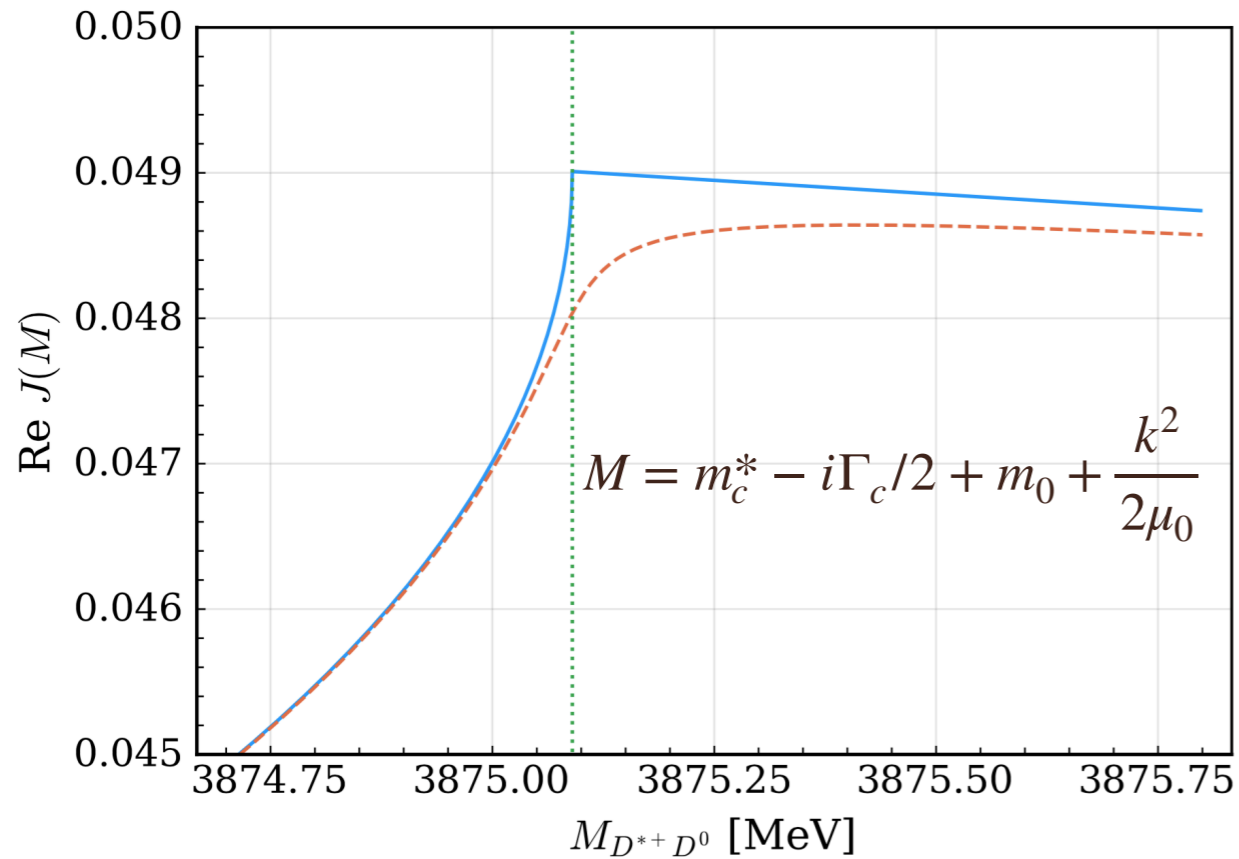
The molecular picture of the T_{cc}^+

Low energy expansion of the scattering amplitude

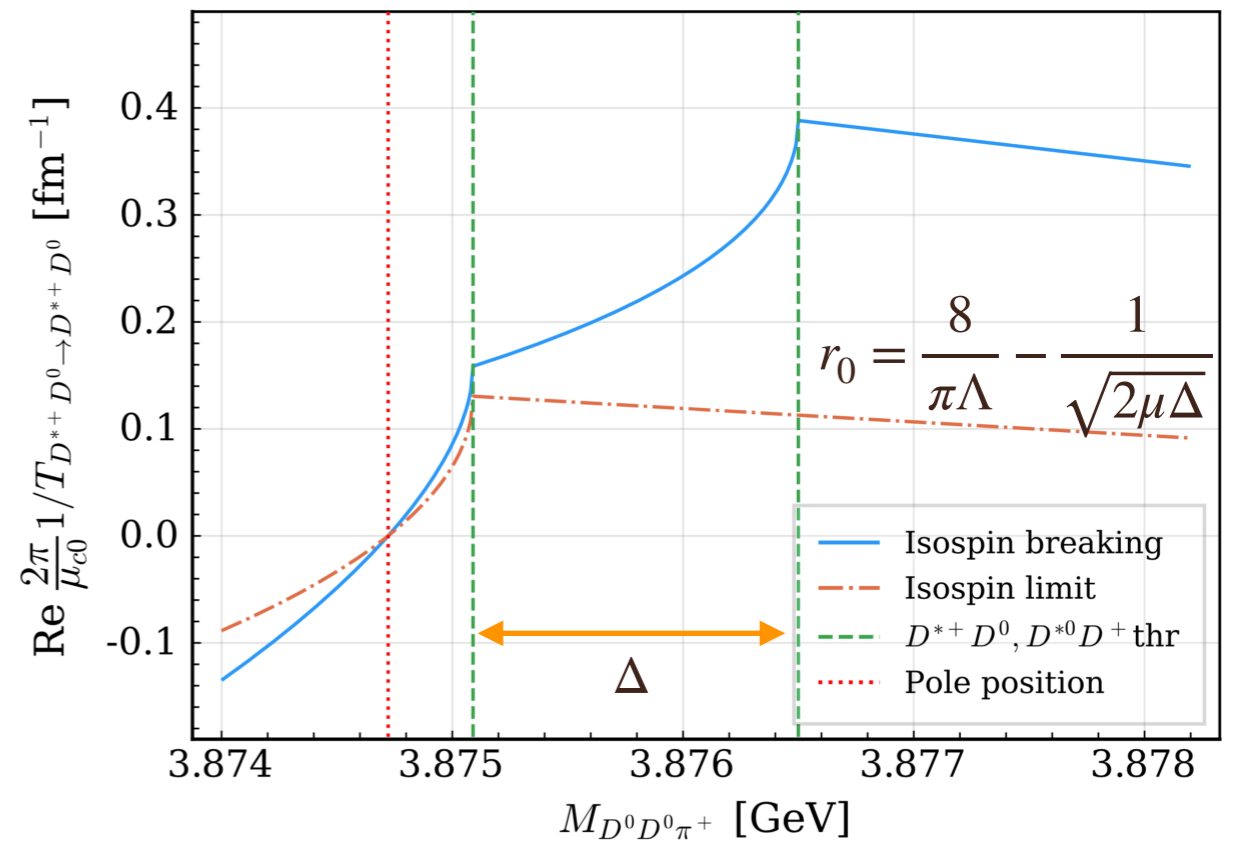
$$T_{D^{*+}D^0 \rightarrow D^{*+}D^0}(k) = -\frac{2\pi}{\mu_{c0}} \left(\frac{1}{a_0} + \frac{1}{2}r_0k^2 - ik + \mathcal{O}(k^4) \right)$$

Effective range $r_0 \propto -\text{Re} \frac{dT^{-1}}{dM} \Big|_{M=M_{\text{thr}}+0^+}$

Width effect



Isospin violation



The molecular picture of the T_{cc}^+

Scheme I: Only contact potentials

$$T_{D^{*+}D^0 \rightarrow D^{*+}D^0}^{-1}(M) = \frac{2}{v_0} + (J_1(M) + J_2(M)) \quad E = M - M_{\text{thr},1}$$

$$J_1(E) = \frac{\Lambda\mu}{\pi^2} - \frac{2\mu^2 E}{\pi^2 \Lambda} + i \frac{\sqrt{2\mu E} \mu}{2\pi} + \mathcal{O}(E^2)$$

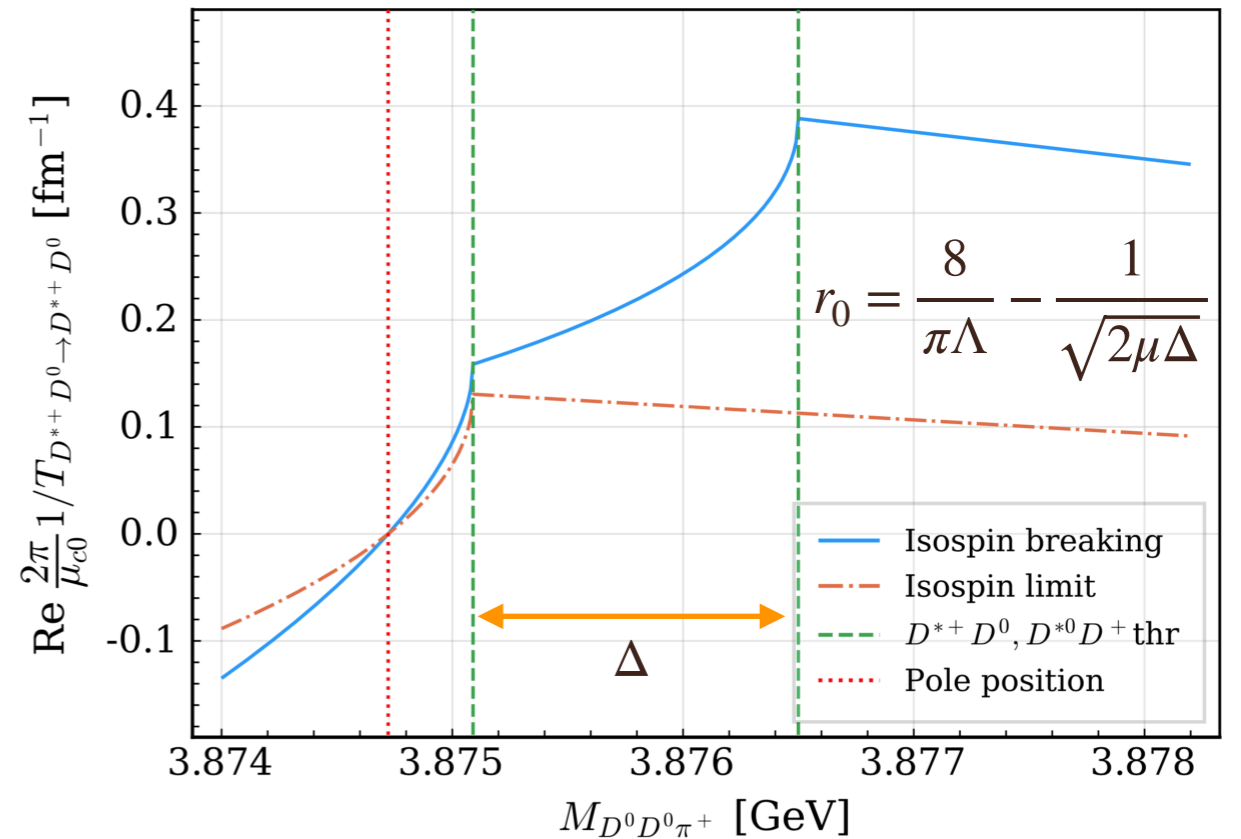
$$J_2(E) = \frac{\Lambda\mu}{\pi^2} - \frac{2\mu^2 E}{\pi^2 \Lambda} + \frac{2\Delta\mu^2}{\pi^2 \Lambda} - \frac{\mu\sqrt{2\mu\Delta}}{2\pi} + \frac{\mu E \sqrt{2\mu\Delta}}{4\pi\Delta} + \mathcal{O}(E^2)$$

$$r_0 = -\frac{2\pi}{\mu^2} \frac{d(J_1(M) + J_2(M))}{dM} \Big|_{M=M_{\text{thr},1}+0^+}$$

$$= -\frac{8}{\pi\Lambda} - \frac{1}{\sqrt{2\mu\Delta}}$$

$$\Delta r_{\text{IV}} \equiv -\sqrt{\frac{1}{2\mu\Delta}} = -3.78 \text{ fm}$$

Isospin violation



The molecular picture of the T_{ee}^+

Compositeness

$$\bar{X}_A = \left(1 + 2 \left| \frac{r'_0}{\text{Re}a_0} \right| \right)^{-1/2}, \quad r'_0 = r_0 - \Delta r_{\text{IV}}$$

$$\Delta r_{\text{IV}} \equiv -\sqrt{\frac{1}{2\mu\Delta}} = -3.78 \text{ fm}$$

Scattering length and effective range

Schemes	a (fm)	r_0 (fm)	r'_0 (fm)	\bar{X}_A
Scheme I	$(-6.31^{+0.36}_{-0.45}) + i(0.05^{+0.01}_{-0.01})$	-2.78 ± 0.01	1.00 ± 0.01	0.87 ± 0.01
Scheme II	$(-6.64^{+0.36}_{0.50}) - i(0.10^{+0.01}_{-0.02})$	-2.80 ± 0.01	0.98 ± 0.01	0.88 ± 0.01
Scheme III	$(-6.72^{+0.36}_{-0.45}) - i(0.10^{+0.03}_{-0.03})$	-2.40 ± 0.01	1.38 ± 0.01	0.84 ± 0.01

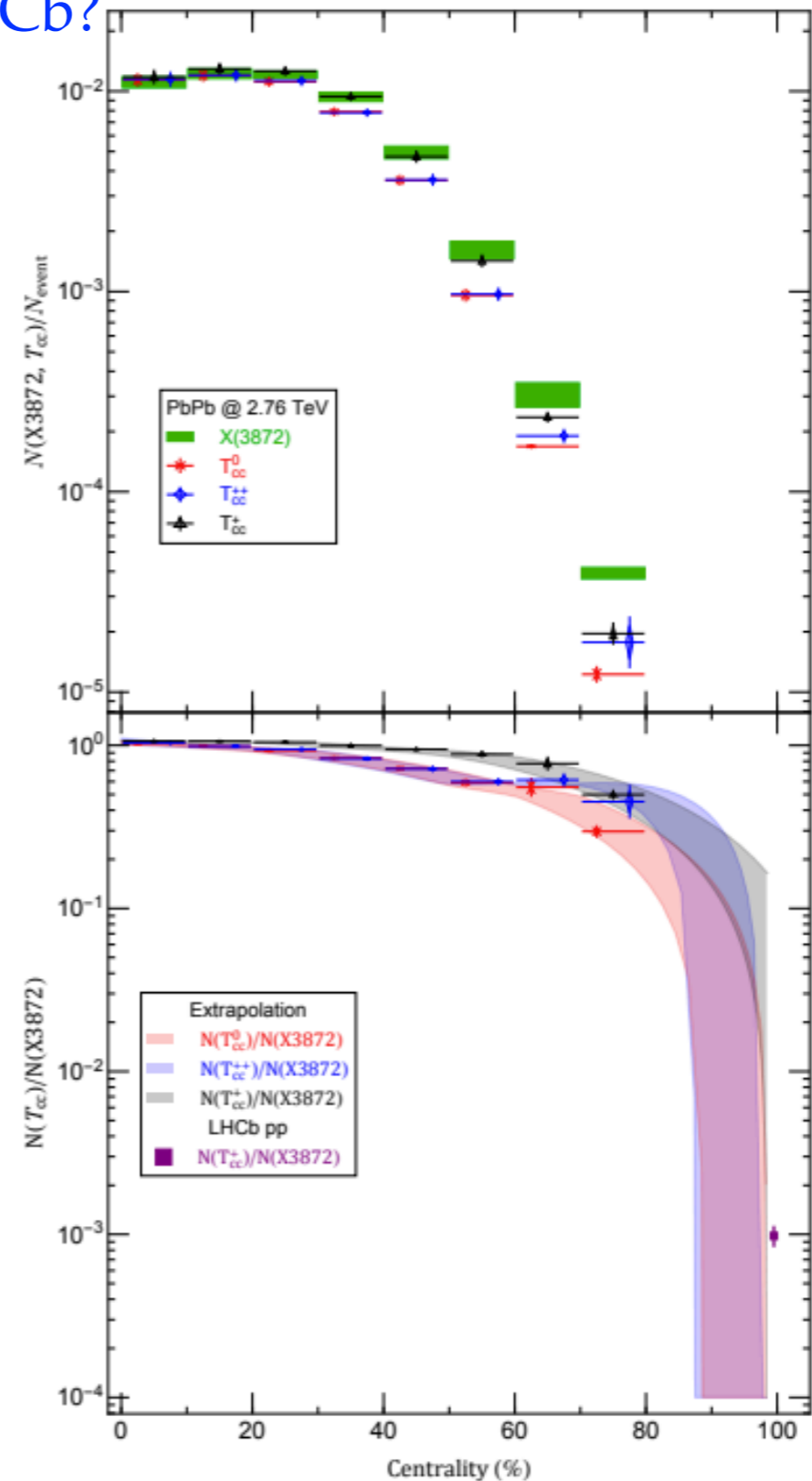
OPE contribute 0.40

The production of the T_{cc}^+ in HIC

P_b – P_b @ 2.76 TeV

Why does the isospin triplet T_{cc} disappear in LHCb?

- Small production in pp collision
 - Isotriplet DD^* state doesn't exist
- ① $I = 1$: $T_{cc}^0, T_{cc}^+, T_{cc}^{++}$
 - ② $I = 0$: T_{cc}^+
 - ③ $R(T_{cc}^{0,+,++}) \equiv N(T_{cc}^{0,+,++})/N(X(3872))$
 - ④ A third order polynomial function for the relative yield ratio
 - ⑤ T_{cc}^+ is about 3 orders smaller at UPC
 - ⑥ Isotriplet partners is even smaller than that of T_{cc}^+

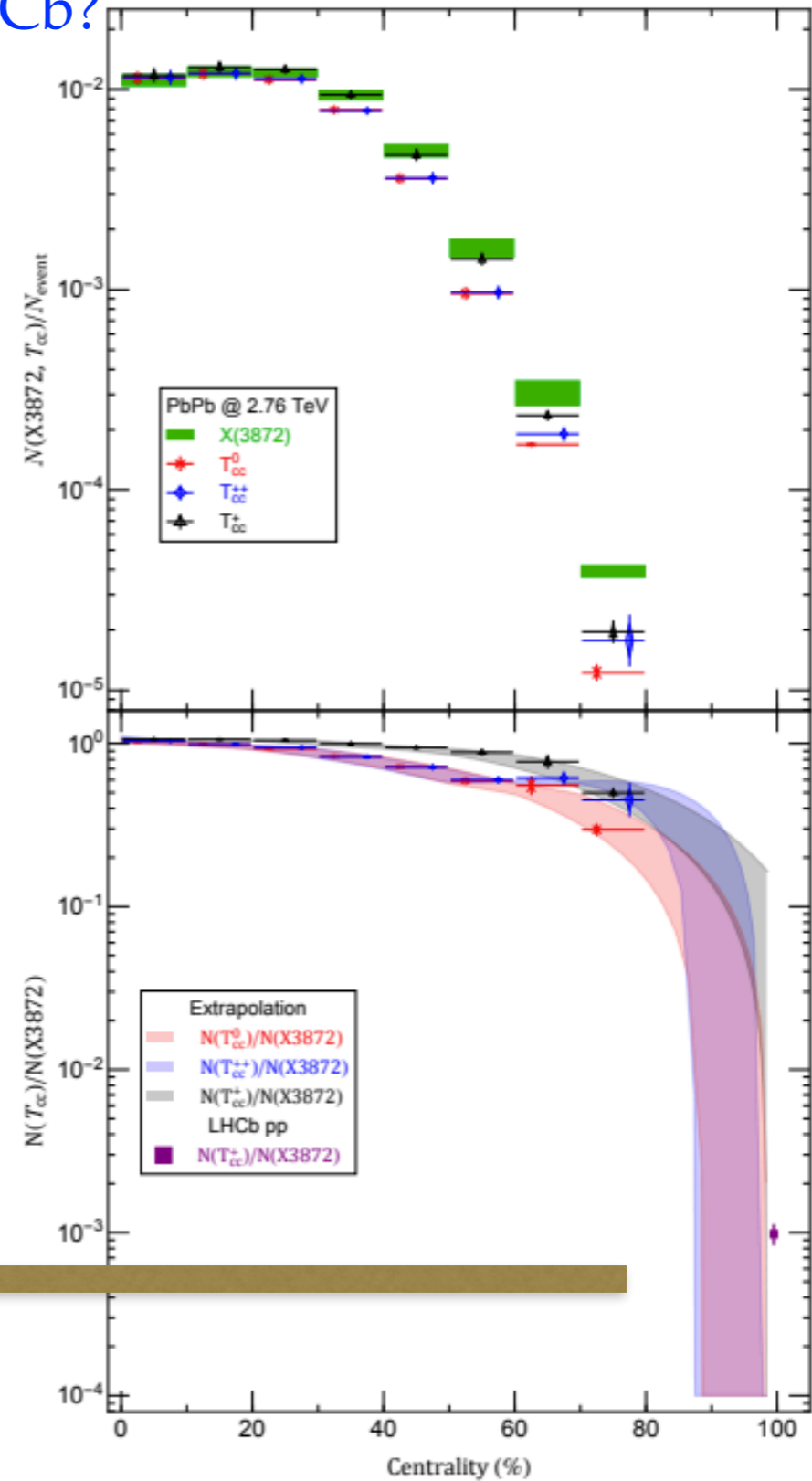
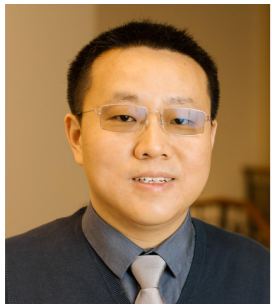


The production of the T_{cc}^+ in HIC

P_b – P_b @ 2.76 TeV

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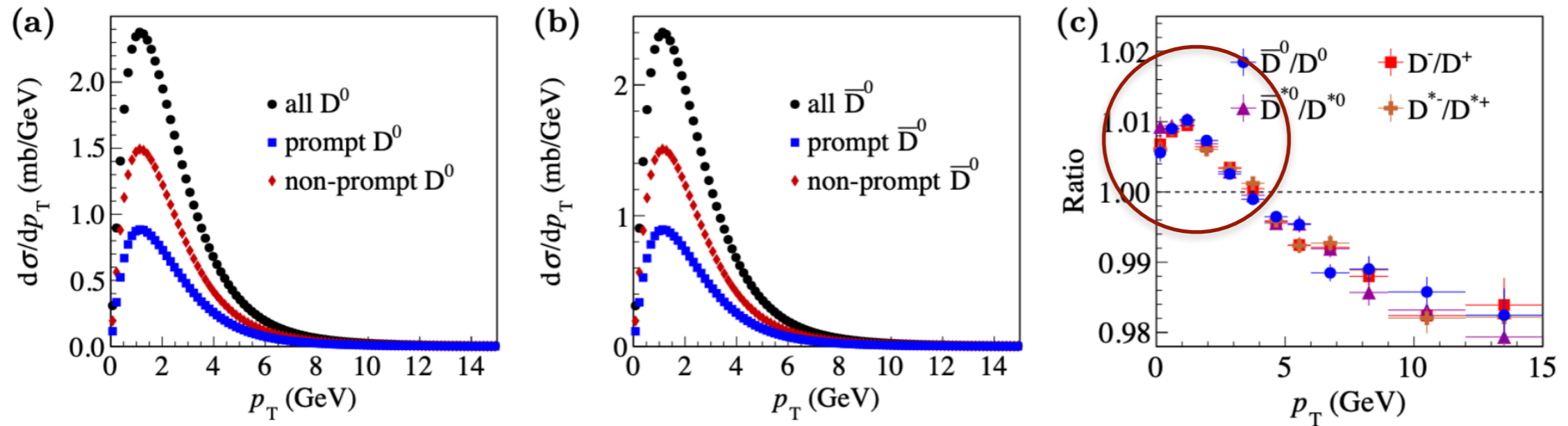


$$\frac{\#T_{cc}^+}{\#X(3872)_{\text{Exp.}}} \simeq 10^{-3}$$

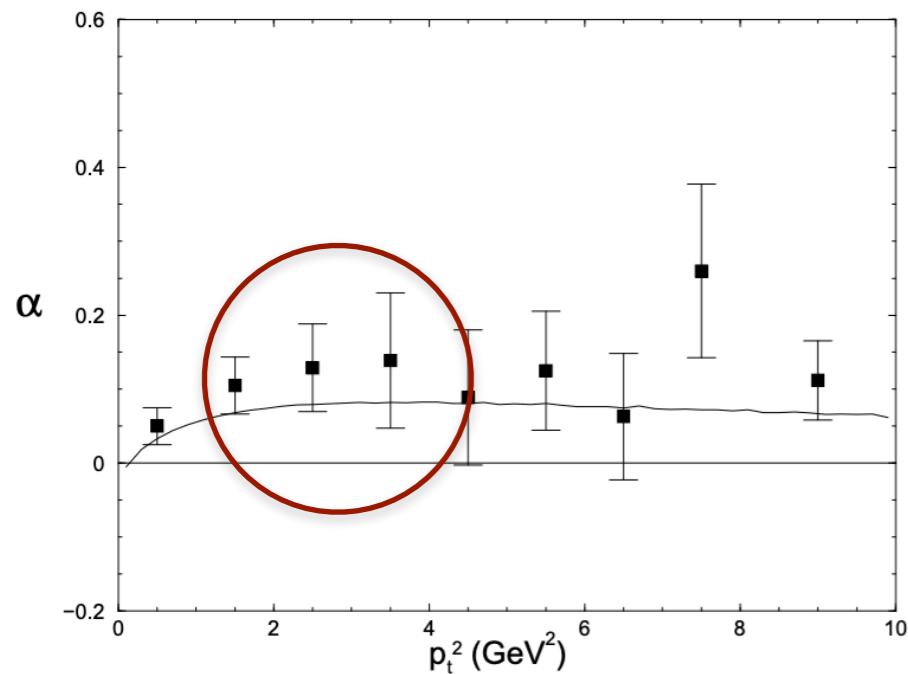
$$\frac{\#T_{cc}^{++}}{\#X(3872)_{\text{Theory}}} \simeq 10^{-4}$$

Revealing the nature of the T_{cc}^+ in pp collision

The production of single charmed meson



Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]



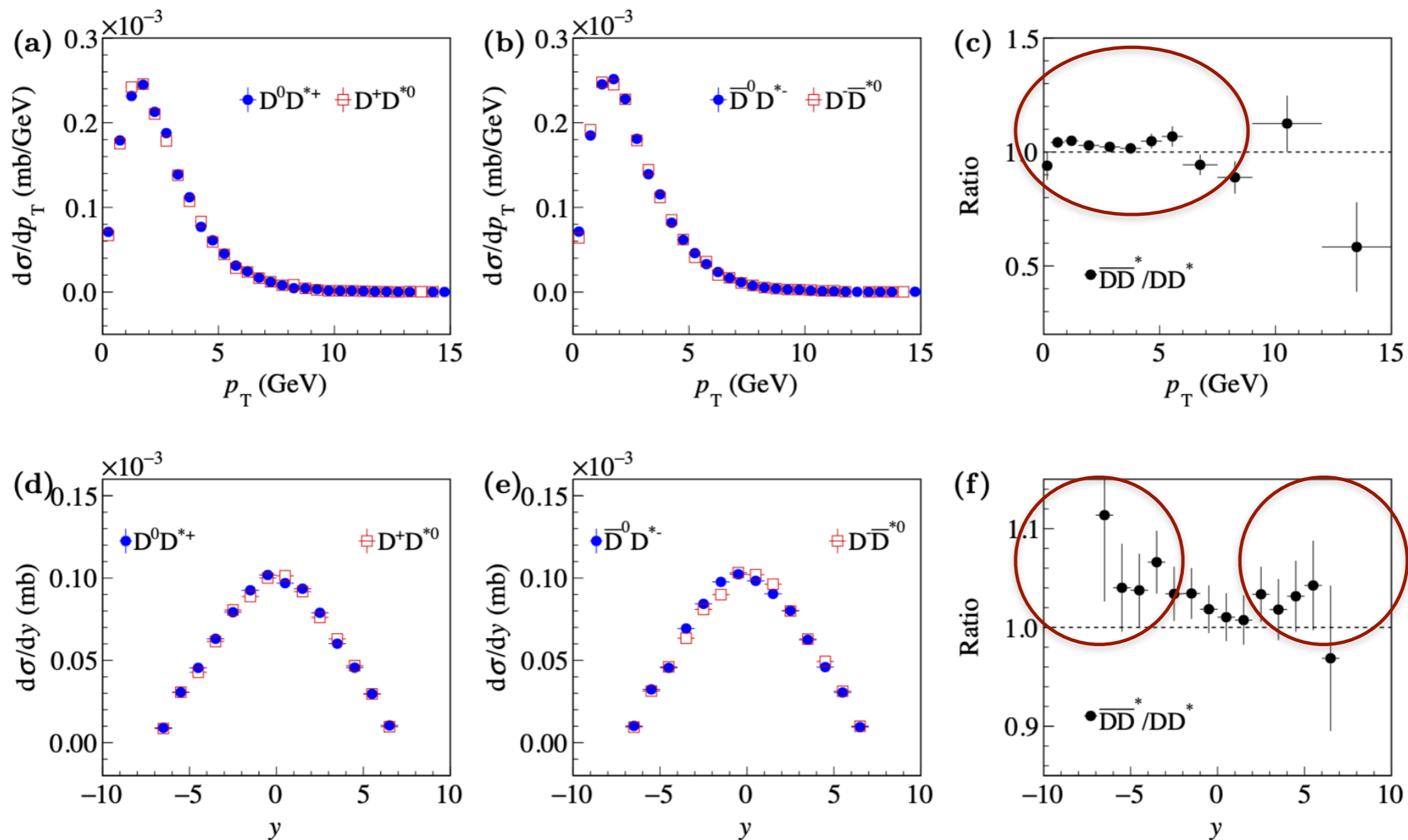
$$\alpha [D^+] = \frac{d\sigma [D^-] - d\sigma [D^+]}{d\sigma [D^-] + d\sigma [D^+]}$$

Heavy quark recombination mechanism

Braaten, Jia, Mehen, PRL89(2002)122002, Chang, Ma, Si, PRD68(2003)014018

Revealing the nature of the T_{cc}^+ in pp collision

The production of prompt (anti)charmed meson pairs

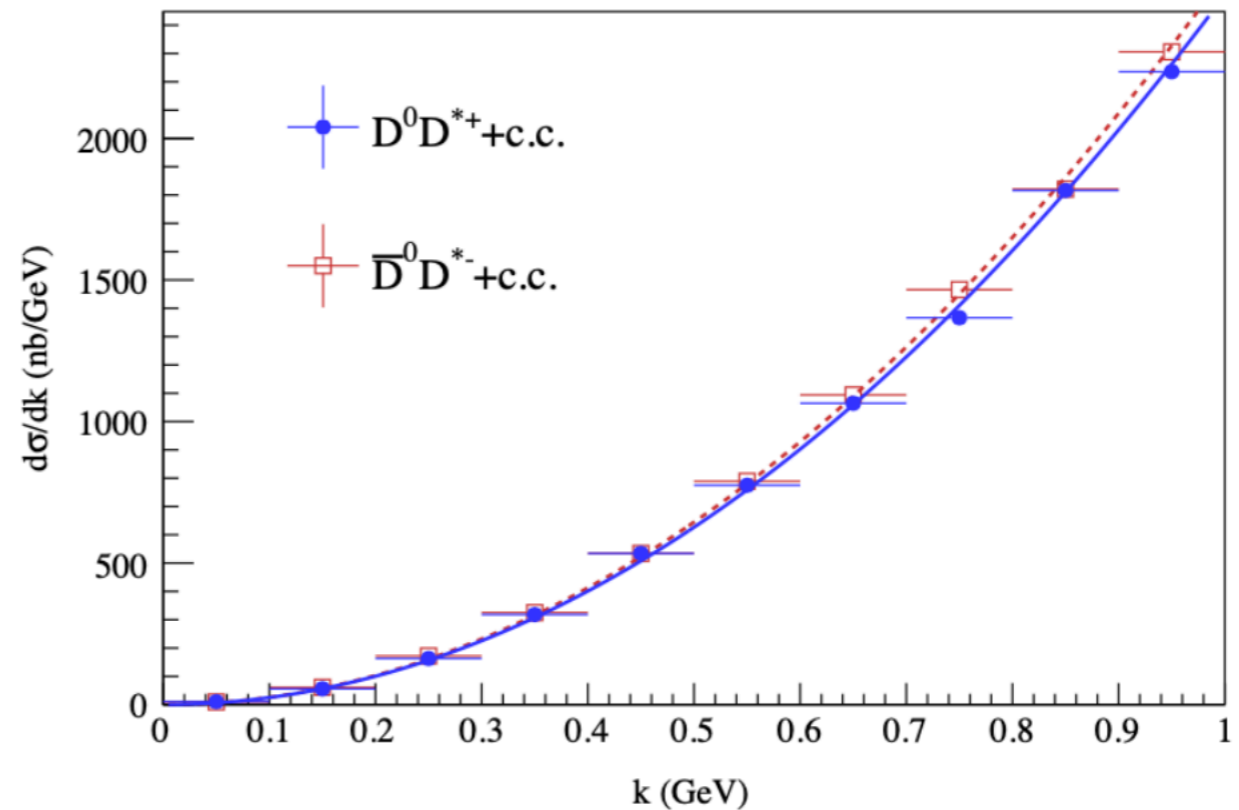
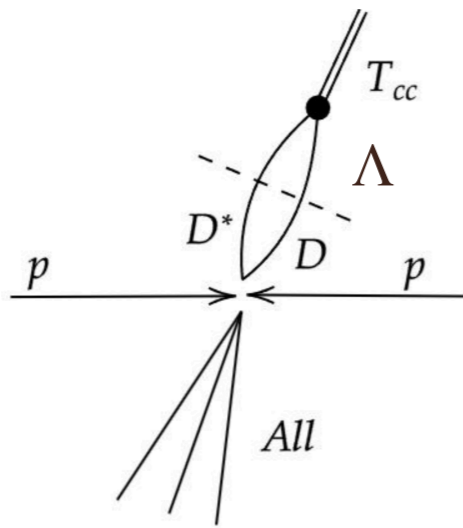


Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The formation of hadronic double charm tetraquark

$$\sigma(T_{cc}) = \frac{1}{4m_D m_{D^*}} g^2 |\mathcal{G}|^2 \left(\frac{d\sigma(DD^*)}{dk} \right) \frac{4\pi^2 \mu}{k^2}$$



Successfully apply to

$X(3872)$: Guo et.al., EPJC74(2014)3063, Albaladejo et.al., CPC41(2017)121001, Yang et.al., CPC45(2021)123101,

Shi et.al., PRD106(2022)114026

P_{cs} : Ling et.al., EPJC81(2021)319

Z_b s: Cao et.al., PRD101(2020)074010

$D_{s0}(2317)$: Guo et.al., JHEP05(2014)138

- Differential cross section behaves as k^2
- Total cross section depends on Λ
- Propose a model independent quantity \mathcal{A}

Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The cross sections and asymmetry in the molecular picture

$$\mathcal{A} \pm \delta_{sys} \pm \delta_{sta}$$

Range(GeV)	$\sigma_{T_{cc}^+} (\sigma_{T_{\bar{c}\bar{c}}^-})$			$\mathcal{A}(\%)$
	$\Lambda = 0.5 \text{ GeV}$	$\Lambda = 1 \text{ GeV}$	$\Lambda = 1.5 \text{ GeV}$	
Full	43.30±0.70 nb (44.13±0.71 nb)	152.42±0.89 nb (156.81±0.91 nb)	313.74±1.03 nb (321.14±1.04 nb)	1.24±0.30±0.20
	LHCb ($2 < y < 4.5$)			
$4 < p_T < 20$ [112]	1.46±0.15 nb (1.45±0.15 nb)	5.27±0.20 nb (5.63±0.20 nb)	11.46±0.23 nb (11.87±0.24 nb)	2.53±2.01±1.79
$p_T > 0$ [11]	8.26±0.44 nb (8.69±0.46 nb)	29.93±0.57 nb (30.82±0.58 nb)	62.28±0.66 nb (64.30±0.67 nb)	1.64±1.03±0.52
	CMS ($ y < 1.2$)			
$10 < p_T < 50(30)$ [113]	0.05±0.02 nb (0.03±0.02 nb)	0.28±0.04 nb (0.20±0.03 nb)	0.55±0.04 nb (0.44±0.04 nb)	-13.42±8.44±2.18
	ATLAS ($ y < 0.75$)			
$10 < p_T < 70$ [114]	0.03±0.02 nb (0.03±0.02 nb)	0.20±0.03 nb (0.13±0.03 nb)	0.38±0.04 nb (0.28±0.03 nb)	-16.87±9.33±10.10

$$\mathcal{A} \equiv \frac{\sigma^- - \sigma^+}{\sigma^- + \sigma^+}$$

$$\mathcal{A} \equiv \omega_1 \mathcal{A}_1 + \omega_2 \mathcal{A}_2 + \omega_3 \mathcal{A}_3$$

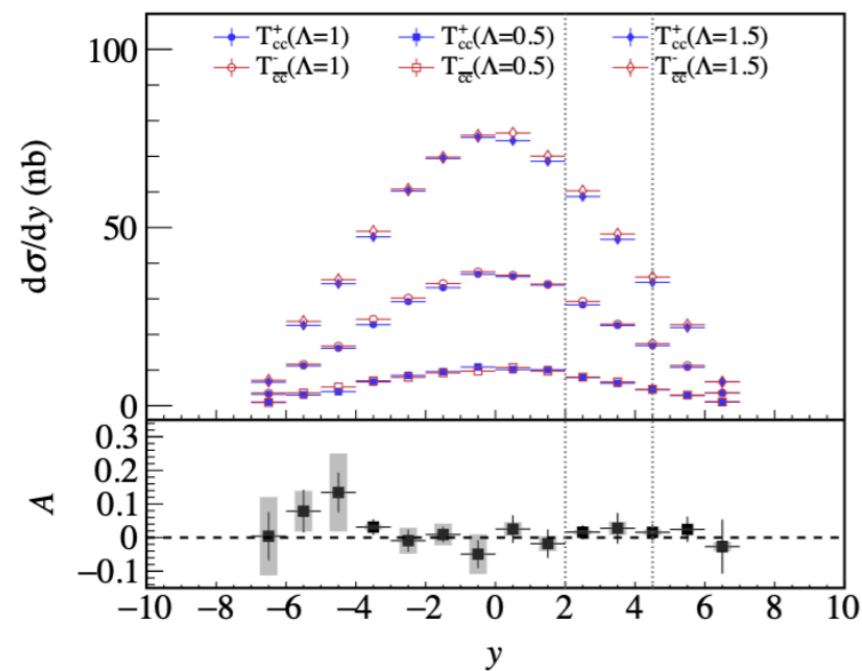
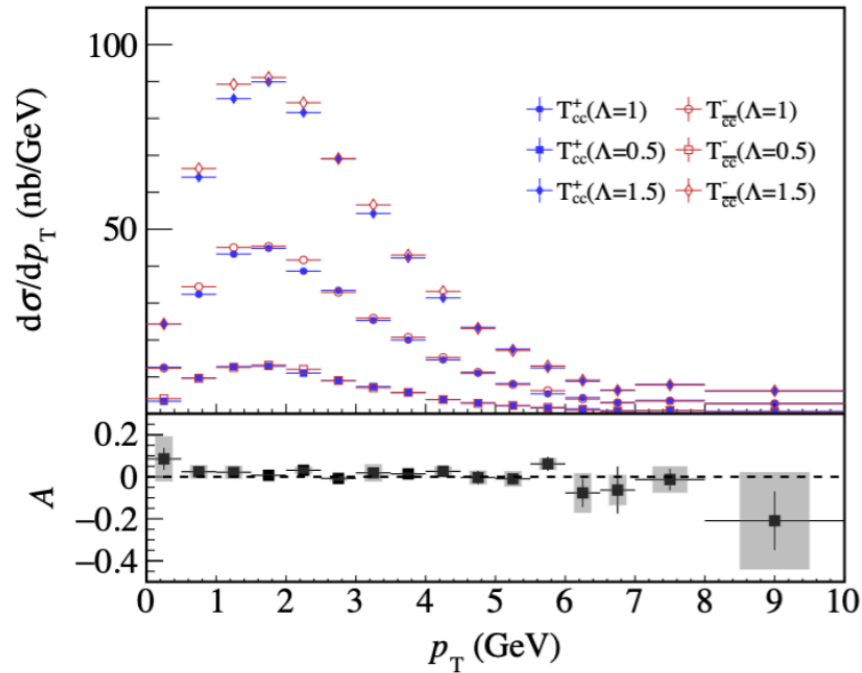
$$\omega_i = \frac{\frac{1}{\delta_i^2}}{\frac{1}{\delta_1^2} + \frac{1}{\delta_2^2} + \frac{1}{\delta_3^2}}$$

$$\delta_{sta} \equiv \omega_1 \delta_1 + \omega_2 \delta_2 + \omega_3 \delta_3$$

$$\delta_{sys} \equiv \sqrt{\frac{\sum_i (\mathcal{A}_i - \mathcal{A})^2}{3}}$$

Revealing the nature of the T_{cc}^+ in pp collision

The p_T and y distributions in the molecular picture

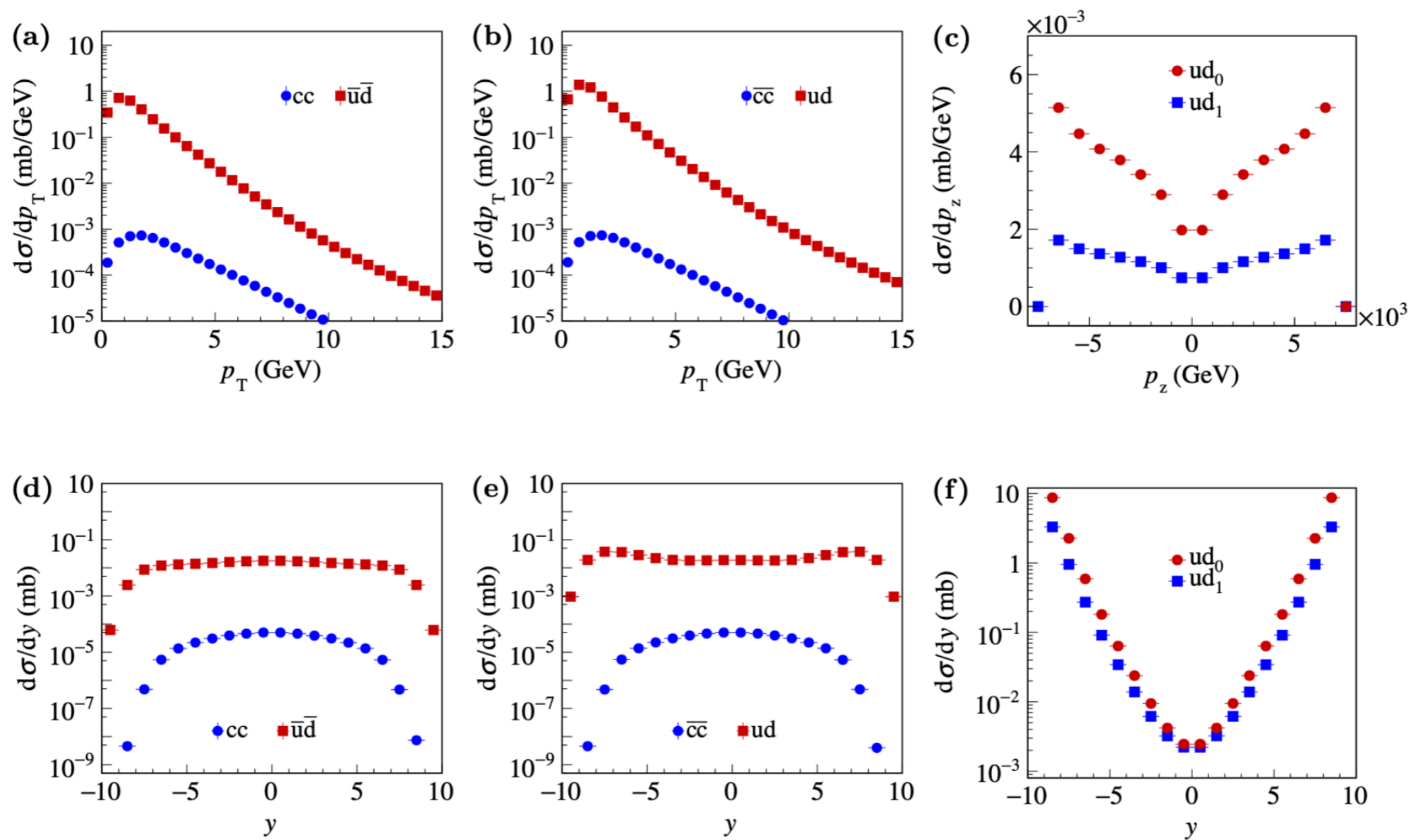


- Cross section increase with the increasing Λ
- Reach maximum value at $p_T = 2$ GeV
- Reach maximum value at $y = 0$
- The asymmetry is positive at low p_T
- The asymmetry is positive at large y
- Cross section depends on the parameter Λ

Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The diquark and antidiquark distributions

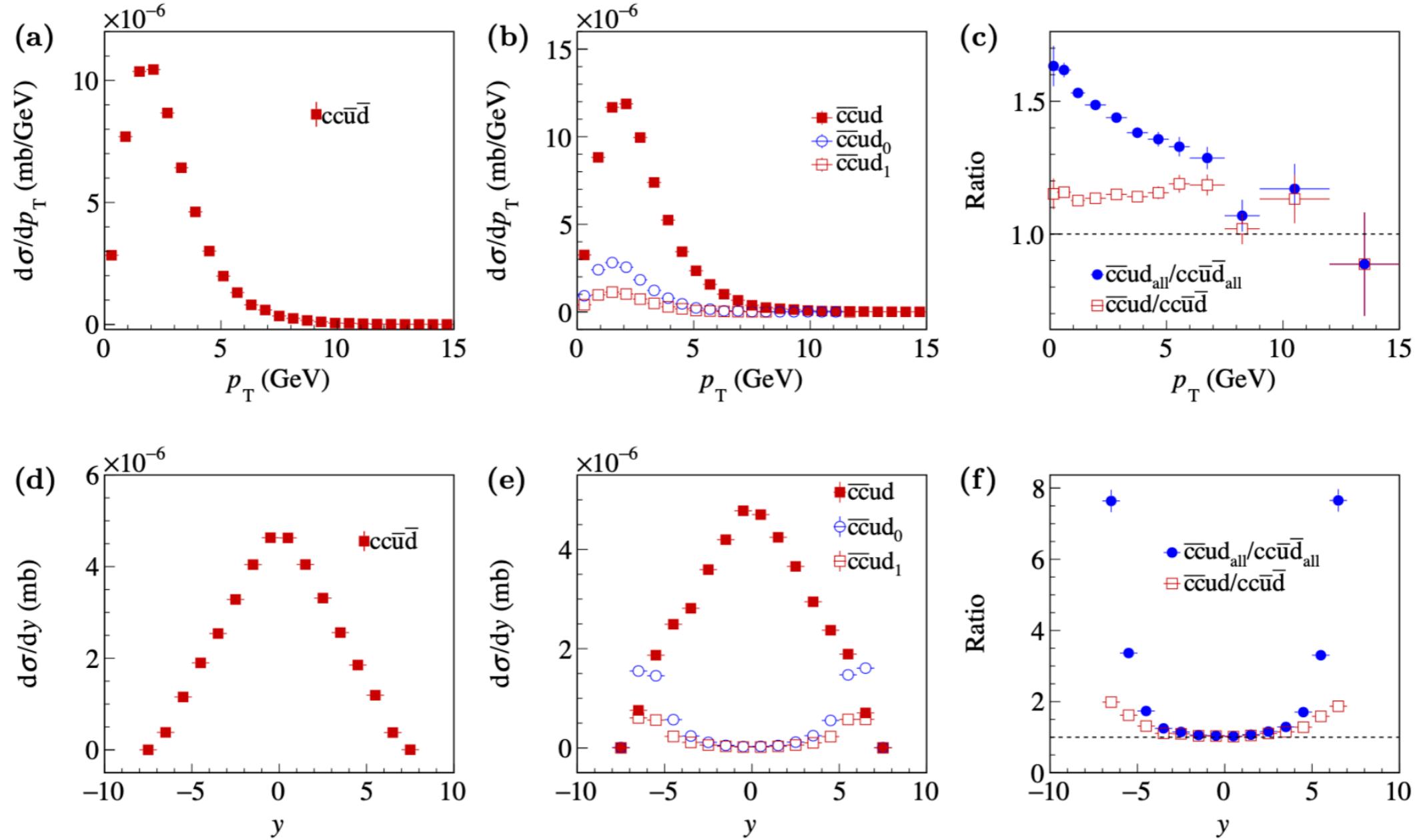


Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The diquark-antidiquark pair distributions

Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

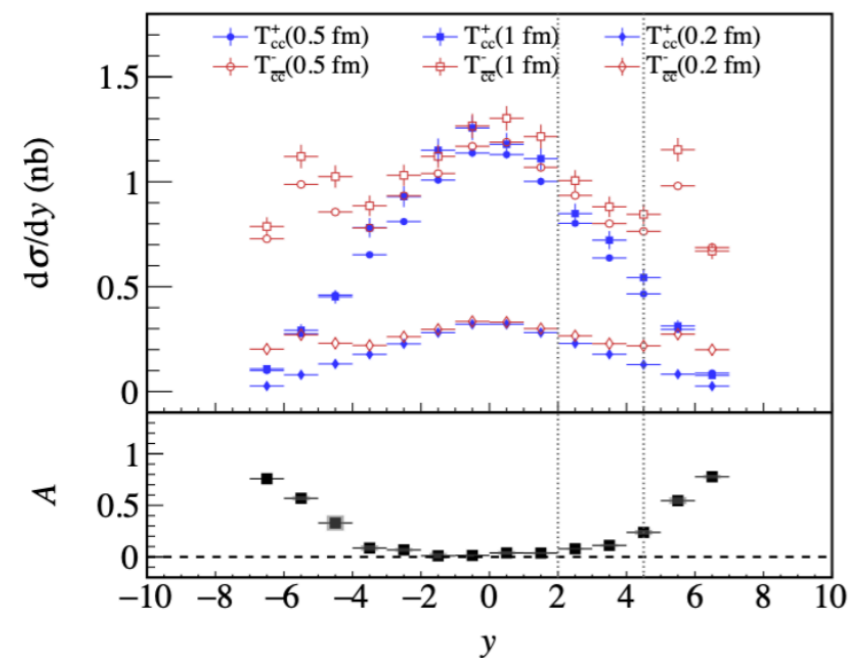
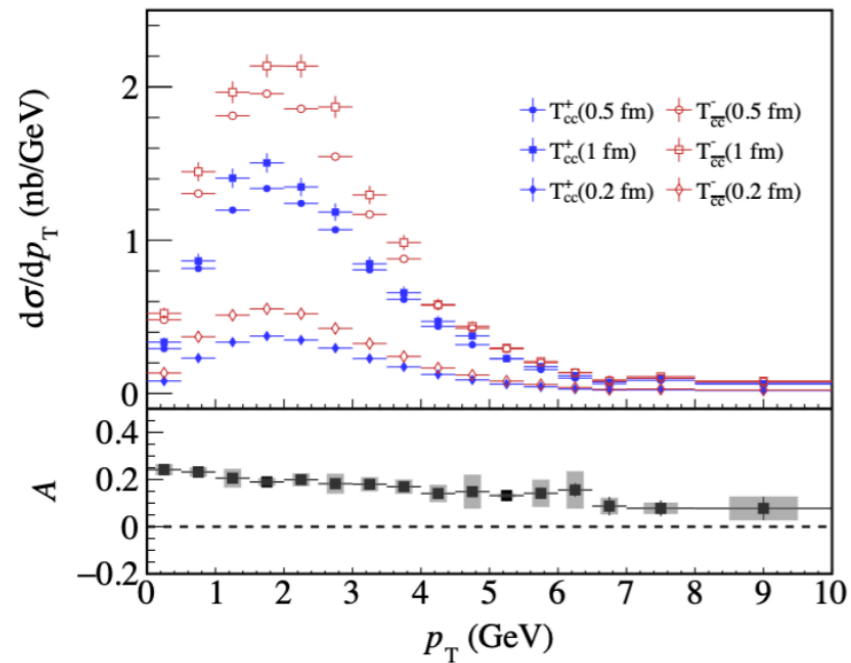


Gaussian Type $\phi(\vec{p}, a) = \left(\frac{1}{\pi}\right)^{\frac{3}{4}} \left(\frac{1}{a}\right)^{\frac{3}{2}} e^{-\frac{p^2}{2a^2}}$

Saturmian Type $\psi(\vec{p}, b) = 2\sqrt{2} \left(\frac{1}{\pi}\right) \left(\frac{1}{b}\right)^{\frac{3}{2}} \frac{1}{\left(\frac{p^2}{b^2} + 1\right)^2}$

Revealing the nature of the T_{cc}^+ in pp collision

The p_T and y distributions in the compact tetraquark picture



- Cross section increases with the size
- Reach maximum value at $p_T = 2$ GeV
- Reach maximum value at $y = 0$ and $y = \pm 6$
- The asymmetry is positive at low p_T and large y
- The asymmetry is more significant
- The large difference region moves to smaller y with smaller c.m. energy

Heavy diquark symmetry

$$cc\bar{u}\bar{d} \leftrightarrow \bar{c}\bar{u}\bar{d}$$

$$\bar{c}\bar{c}ud \leftrightarrow cud$$

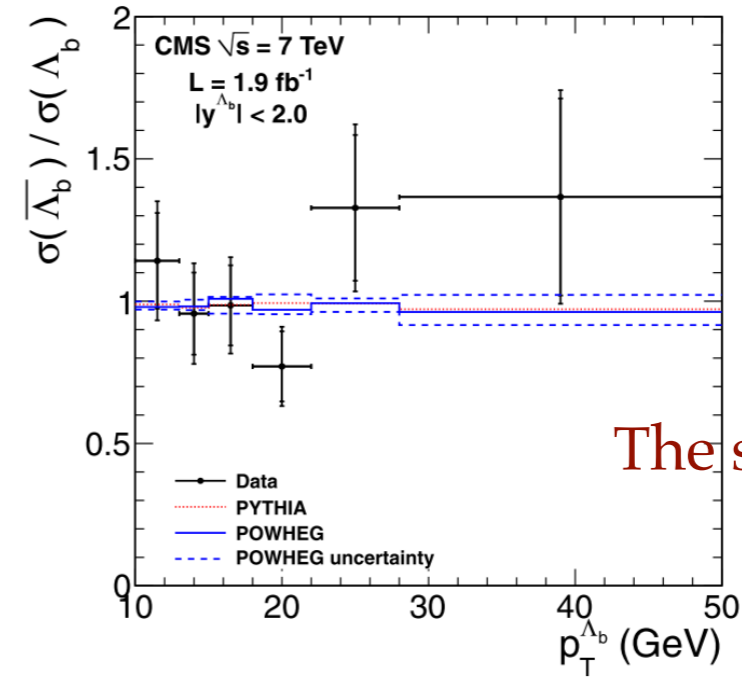
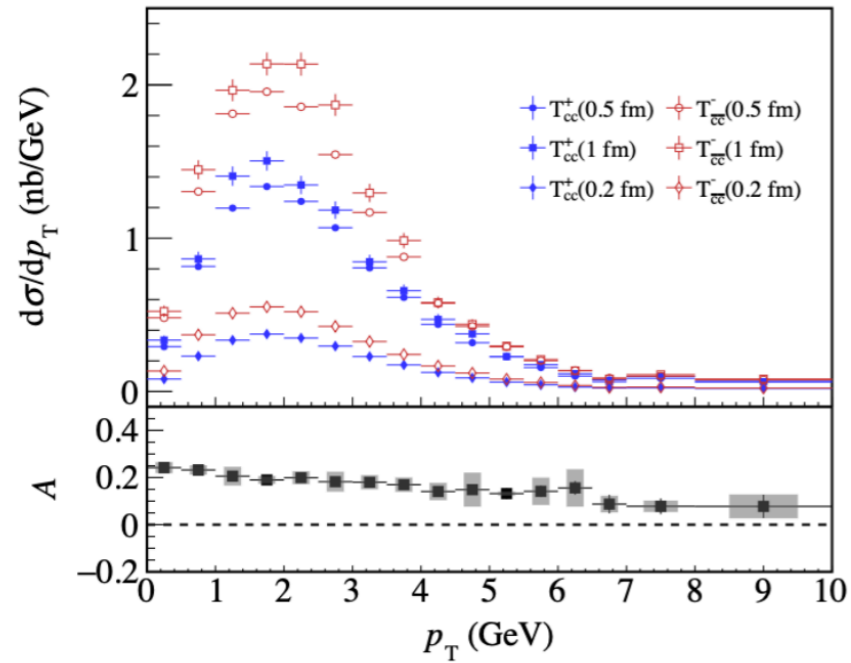


The asymmetry of Λ_c, Λ_b ?

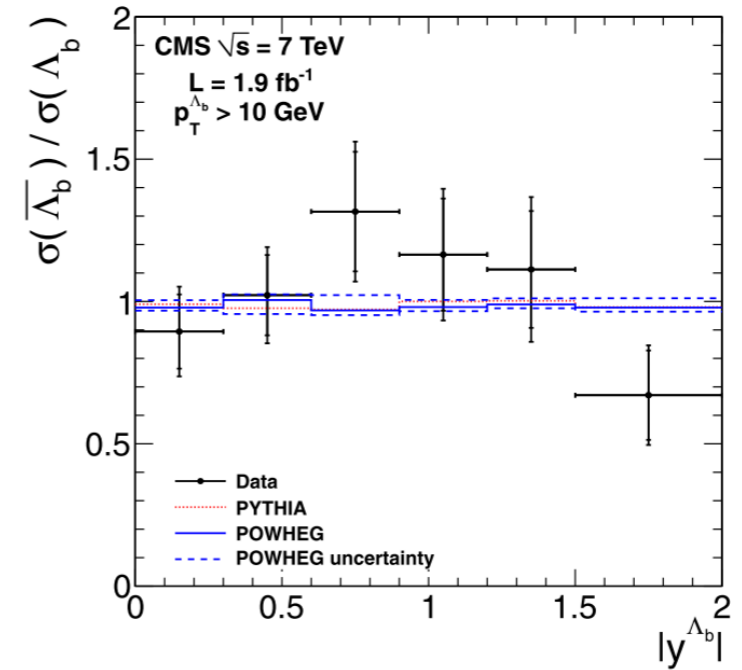
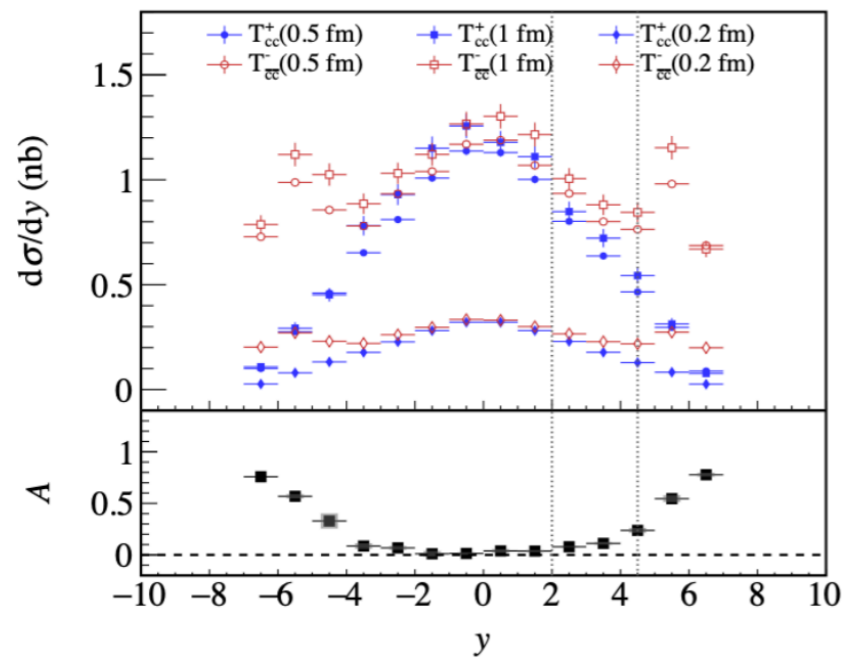
Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The p_T and y distributions in the compact tetraquark picture



The statistic is too low



Revealing the nature of the T_{cc}^+ in pp collision

The cross sections and asymmetry in the compact tetraquark picture

Gaussian Type+1GeV

Range(GeV)	$\sigma_{T_{cc}^+}(\sigma_{T_{\bar{c}\bar{c}}^-})$			$\mathcal{A}(\%)$
	$r = 0.2$ fm	$r = 0.5$ fm	$r = 1$ fm	
Full	1.25 ± 0.005 nb	4.43 ± 0.02 nb	4.88 ± 0.02 nb	$18.73 \pm 0.25 \pm 0.14$
	(1.82 ± 0.01) nb	(6.46 ± 0.02) nb	(7.16 ± 0.02) nb	
LHCb ($2 < y < 4.5$)				
$4 < p_T < 20$ [112]	39.75 ± 0.89 pb	139.88 ± 3.11 pb	163.77 ± 3.65 pb	$7.35 \pm 1.48 \pm 5.24$
	(50.57 ± 1.00) pb	(171.16 ± 3.38) pb	(163.83 ± 3.23) pb	
$p_T > 0$ [11]	0.24 ± 0.002 nb	0.84 ± 0.01 nb	0.91 ± 0.01 nb	$11.42 \pm 0.60 \pm 0.17$
	(0.30 ± 0.002) nb	(1.05 ± 0.01) nb	(1.14 ± 0.01) nb	
CMS ($ y < 1.2$)				
$10 < p_T < 50(30)$ [113]	3.77 ± 0.51 pb	4.73 ± 0.56 pb	4.51 ± 0.53 pb	$-6.62 \pm 8.86 \pm 7.96$
	(1.09 ± 0.15) pb	(3.77 ± 0.51) pb	(4.94 ± 0.67) pb	
ATLAS ($ y < 0.75$)				
$10 < p_T < 70$ [114]	0.92 ± 0.14 pb	3.15 ± 0.46 pb	3.11 ± 0.46 pb	$0.98 \pm 11.04 \pm 15.37$
	(0.69 ± 0.12) pb	(2.83 ± 0.49) pb	(4.88 ± 0.84) pb	

Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Revealing the nature of the T_{cc}^+ in pp collision

The stability of asymmetry

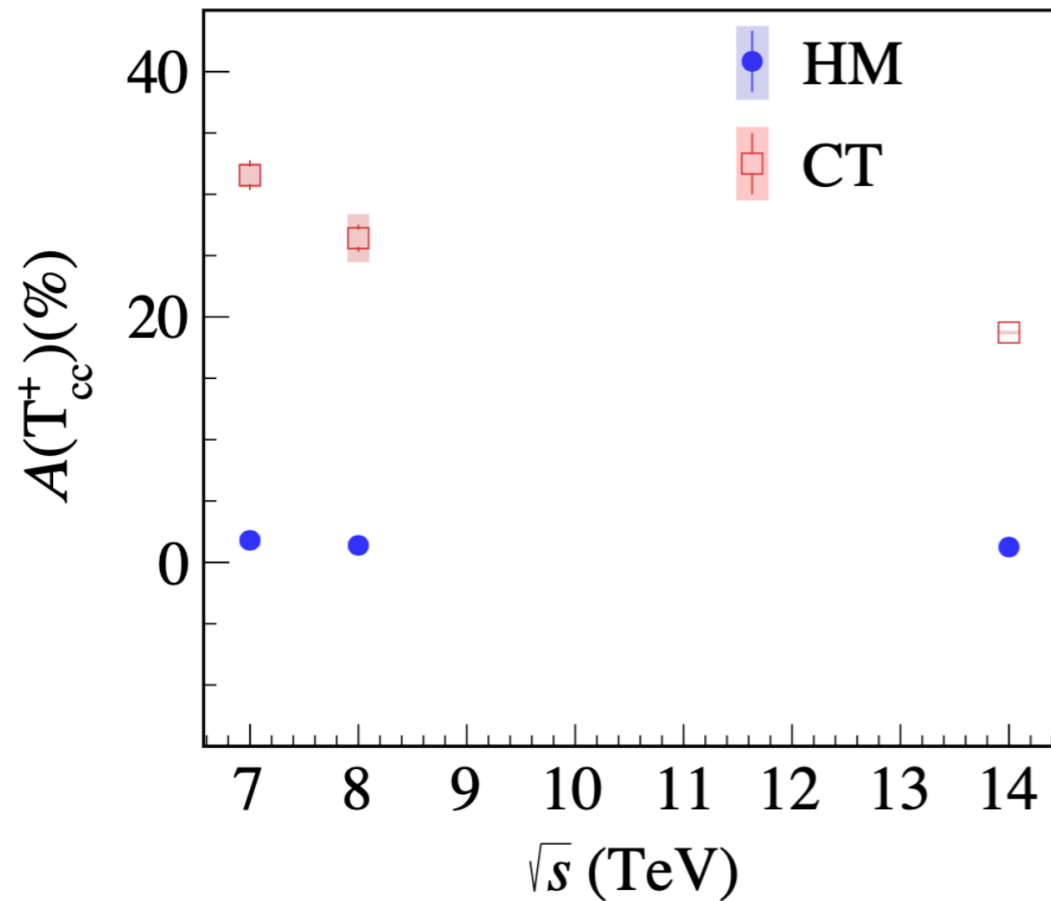
Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

$\mathcal{A}(\%)$	Gaussian 1GeV	Gaussian 0.5GeV	Saturmian 1GeV
Full	$18.73 \pm 0.25 \pm 0.14$	$18.71 \pm 0.67 \pm 0.17$	$18.70 \pm 0.25 \pm 0.08$
$4 < p_T < 20 \text{ GeV}$	$7.35 \pm 1.48 \pm 5.24$	$3.53 \pm 4.15 \pm 2.99$	$6.71 \pm 1.48 \pm 5.17$
$p_T > 0 \text{ GeV}$	$11.42 \pm 0.60 \pm 0.17$	$10.65 \pm 1.64 \pm 0.42$	$11.38 \pm 0.60 \pm 0.16$

- The asymmetry is stable, i.e. the three asymmetries are consistent with each other with uncertainties
- The LHCb acceptance region has the ability to measure the asymmetry

Revealing the nature of the T_{cc}^+ in pp collision

The asymmetries at three different c.m. energies



- Both asymmetries are positive
- The asymmetry of compact tetraquark is larger than that of hadronic molecule
- The asymmetry decrease with the increasing c.m. energy



Hua, Li, QW, Yang, Zhao, Zou, arXiv:2310.04258[hep-ph]

Summary and outlook

- Explore the hadronic molecular nature of the T_{cc}^+ via its line shape
- In the DD^* molecular picture, estimate the yield of T_{cc}^+ and its isospin partners in $Pb - Pb @ 2.76$ TeV.
- Propose asymmetry for distinguishing HM and CT pictures



UPC or HIC?

- Whether it is possible to measure the T_{cc}^+ in HIC or UPC?
- We would naively expect that the asymmetry would decrease with the increasing impact parameter b

Thank you very much for your attention!