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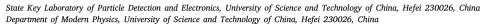
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A thermal bonding method for manufacturing Micromegas detectors

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ABSTRACT

For manufacturing Micromegas detectors, the "bulk" method based on photoetching, was successfully developed and widely used in nuclear and particle physics experiments. However, the complexity of the method requires a considerable number of advanced instruments and processing, limiting the accessibility of this method for production of Micromegas detectors. In view of these limitations with the bulk method, a new method based on thermal bonding technique (TBM) has been developed to manufacture Micromegas detectors in a much simplified and efficient way without etching. This paper describes the TBM in detail and presents performance of the Micromegas detectors built with the TBM. The effectiveness of this method was investigated by testing Micromegas detector prototypes built with the method. Both X-rays and electron beams were used to characterize the prototypes in a gas mixture of argon and CO_2 (7%). A typical energy resolution of ~16% (full width at half maximum, FWHM) and an absolute gain greater than 10^4 were obtained with 5.9 keV X-rays. Detection efficiency greater than 98% and a spatial resolution of ~65 μ m were achieved using a 5 GeV electron beam at the DESY test-beam facility. The gas gain of a Micromegas detector could reach up to 10^5 with a uniformity of better than 10^6 when the size of the avalanche gap was optimized thanks to the flexibility of the TBM in defining the gap. Additionally, the TBM facilitates the exploration of new detector structures based on Micromegas owing to the much-simplified operation with the method.

1. Introduction

The Micromegas detector is a typical micropattern gaseous detector (MPGD) that was invented in 1996 [1]. It has a two-stage structure: a 3–5 mm drift region for primary ionization of charged particles and an avalanche gap of about 0.1 mm. Owing to its excellent energy resolution, rate capability, and possibility for large area, several methods have been developed to manufacture this type of detector. Producing the narrow avalanche gap is the most challenging part of the fabrication process of a Micromegas detector.

Among the different methods of manufacturing Micromegas detectors, the bulk method based on photoetching [2,3] has been well developed and substantiated through its application at various experiments, such as T2 K [4], COMPASS RICH [5], and the ongoing upgrade of the ATLAS NSW [6]. Micromegas detectors built with the bulk method have exhibited excellent performance, such as gas gain of 2×10^4 [2], energy resolution below 20% in argon based mixtures and gain uniformity of 3.6% [7]. However, this method makes pillars that support the avalanche gap by etching photopolymer. So, it heavily depends on the equipment for performing exposure, etching and lamination, and professional operation of advanced devices. To avoid the etching operation, a new method based on a thermal bonding process for manufacturing Micromegas detectors as a result of decade

long effort [8,9]. This method has shown the capability of producing Micromegas detectors and its derivatives in an easier and more efficient way by using simple devices and conventional materials.

This paper is organized as follows. The TBM is described in detail in Section 2, including its basic principle, fabrication process and advantages. In Section 3, performance of Micromegas detectors fabricated with the TBM was intensively studied. The effectiveness of the TBM is demonstrated by the results from the study including energy resolution, avalanche gain, spatial resolution, detection efficiency, and gain uniformity. Moreover, Micromegas detectors built with the TBM have been successfully used in different experiments, for instance, as trackers for muon tomography, neutron beam monitors, and double micro-mesh gaseous detectors (DMM) for single electron detection. These application examples are briefly discussed in Section 4. Finally, optimization of the avalanche gap of a Micromegas detector with the TBM is presented in Section 5. Significant improvements in gas gain and its uniformity are achieved with the optimization.

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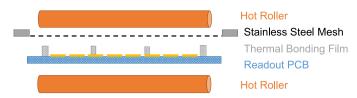


Fig. 1. Schematic of the TBM for Micromegas.

2. Thermal bonding method (TBM)

2.1. Principle of the method

In this method, the avalanche structure of the Micromegas detector is manufactured by bonding the stretched stainless-steel mesh directly onto the readout printed circuit board (PCB) using a hot rolling machine. Fig. 1 shows a schematic of the method. As opposed to etching pillars, the TBM uses two hot rollers to melt and press the spacers made of a thermal bonding film, which have an adhesive–polyester–adhesive triple-layer structure [10], to support the stainless-steel mesh [11] onto the readout anode. Spacer adhesives become solid on cooling and fix the mesh on it, thus forming the micro-gap between them. This whole process involves no chemical agents, hence an etching-free feature of the thermal bonding method. The details of the method will be described in Section 2.2.

The spacers with a diameter ≤ 1 mm are arranged at a distance of ~10 mm on the readout PCB, as depicted in Fig. 2. This special arrangement ensures that the dead area, which is defined as the ratio of the spacer to the active area, is constrained to a small fraction (<1%), similar to the bulk method where pillars with a diameter of 0.2–0.4 mm are set at a distance of 2 mm from each other [2]. Particularly, the large spacer pitch facilitates cleaning of the avalanche region and is prominent to evade sparking in high electric fields. Therefore, a much higher tensile force is required on the mesh to resist the electrostatic force between the mesh and anode.

2.2. Implementation of the manufacturing process

Following the principle mentioned in Section 2.1, the manufacturing process can be implemented step-by-step as shown in Fig. 3 and demonstrated below:

2.2.1. Preparations

The preparation work of the TBM primarily includes stretching of the stainless-steel mesh and cutting the spacers. The mesh is fully stretched to a high tension of more than 25 N/cm (depending on the spacer settings) with a commercial stretching machine, as shown in Fig. 3(a). The small-sized spacers are produced by laser cutting, which is a conventional method for micro- and nanomanufacturing. To avoid melting the adhesive on the film while cutting, the pulse width of the laser machine must be less than 1 ns.

2.2.2. Resistive anode

A germanium layer (grey) coated on the surface of the PCB, as shown in (Fig. 3(b) and (c)), is used as a resistive anode. Its sheet resistivity varies from tens of $M\Omega$ /sq to hundreds of $M\Omega$ /sq corresponding to a film thickness of 500–100 nm.

2.2.3. Pre-setting for the spacers and borders

The spacers and borders, made of a thermal bonding film, are set on the readout PCB in advance by pre-heating. As shown in Fig. 3(b), the spacer is picked and pressed on the PCB anode by using a needle (in right hand). Further, a hot air pencil, controlled by the left hand, is used to flush the spacer with hot air for a few seconds, so that the adhesive on the spacer will melt slightly and attach onto the PCB. Moreover, the borders (white) are set following a similar operation, and the completed subsystem is shown in Fig. 3(c).

2.2.4. Thermal bonding process

As shown in Fig. 3(d), the tensile mesh is placed directly on the PCB where the spacer and border have been set. They are inserted together between two hot rollers (orange) and driven by them to move back and forth for several rounds at an empirical temperature $\sim 150\,$ °C. Therefore, the thermal bonding adhesive on the spacer and border is melted and the mesh is fixed onto the readout PCB after cooling to room temperature (20–30 °C).

2.2.5. Assembly

After the thermal bonding process, the Micromegas detector is finally assembled by trimming the mesh extending out of the sensitive area, installing the drift electrode and gas chamber, soldering the read-out connectors, and other associated processes. Fig. 3(f) shows an image of assembled detector, which has a sensitive area of $150 \times 150 \text{ mm}^2$.

2.3. Advantages of the TBM

It is obvious from the above implementation of the method, the TBM avoids using any environmentally harmful chemical agents, which benefits the long-term development and application of the method. The simple process and minimal requirement of equipment establishes it as an economical and easy to operate process in the laboratory, while simultaneously maintaining good performance of the fabricated detectors

Additionally, the thickness of the avalanche gap is determined by the thermal bonding film, and is therefore, convenient to adjust the gap for optimization and specific applications. The commercially available thermal bonding films have different thicknesses, ranging from several tens of microns to hundreds of microns, allowing the fabrication of avalanche gap as minimum as ${\sim}30~\mu m$.

However, it is worth noting that, the avalanche gap is related to the conditions such as temperature and pressure in the thermal-bonding process as well. Our studies show that the adhesive layers of the spacer should be fully squeezed while heating to press the spacer to a thickness that is identical to the unmelting middle layer. This helps to achieve maximum uniformity of the TBM. The impact of this effect will be discussed in Section 5.

3. Performance of TBM micromegas detectors

To ascertain the feasibility of the TBM, several Micromegas prototypes with a sensitive area of $150\times150~\text{mm}^2$ were fabricated and tested (MM01–MM05). The mesh has 400 lines per inch (LPI) with 19 μm wire diameter and 49% opening rate. A 250 μm thick thermal bonding film, configured as a stack of 75–100–75 μm , was used to develop the spacers. An avalanche gap of $\sim\!110~\mu m$ was measured with a microscope in the cross-section of a sawed detector after solidification of the glue that had been filled into the gap. The prototypes were designed as two-dimensional (2D) readout PCB with strips each having a 400 μm pitch. As shown in Fig. 4, two orthogonal layers of strips were set in the inner layer of the readout PCB. The widths of the strips were 80 μm and 320 μm for the upper (X direction) and lower (Y direction) layers, respectively. The drift gaps between the drift cathode and the mesh were 5 mm and a gas mixture of argon and CO2 (7%) was used as the detector active medium.

3.1. Testing with an X-ray source

The basic performance of the prototypes was assessed in the laboratory, using a gas mixture of argon and CO_2 (7%), with the 55 Fe radioactive source that emits 5.9 keV X-rays. An energy resolution of $\sim 16\%$ (FWHM) (shown in Fig. 5) and a high gain ($> 10^4$) (shown in Fig. 6), were the typical performance of these prototypes. The energy resolution was extracted by fitting the energy spectrum with a sum of

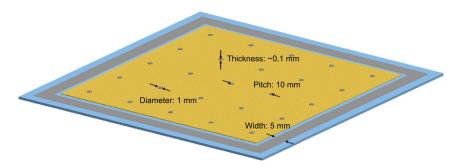


Fig. 2. Schematic design of the spacer arrangement.

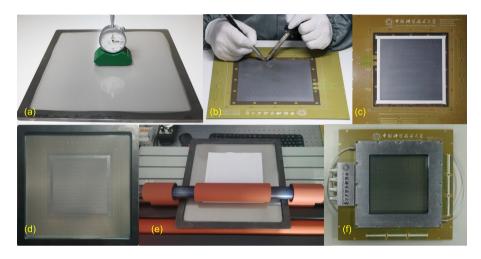


Fig. 3. Fabrication process for Micromegas using the TBM. (a) The stainless-steel woven mesh is stretched over a tension of 25 N/cm and glued on a metal frame. (b) A manual method for pre-setting the spacers on the readout PCB anode. (c) The view after setting borders (the thermal bonding film in white colour, 3–6 mm in width), which fix the mesh and define the active area for the Micromegas. (d) Stacking the stretched mesh on the PCBs for thermal bonding. (e) Manufacturing the avalanche structure of the Micromegas by thermal bonding using a hot roller. (f) Completed view of the thermal bonding Micromegas after assembling the drift electrode, gas chambers, connectors, and other associated components.

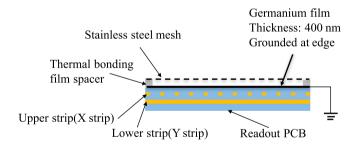


Fig. 4. Schematic representation of the Micromegas prototype structure with a 2D readout.

three Gaussian functions, named K_{alpha} , K_{beta} of $^{55}{\rm Mn}$ and an escape peak of argon.

The uniformity of gain is defined as the ratio of the variance to the mean of the peak values of the energy spectra and was tested for every $20 \times 20 \text{ mm}^2$ block on the detector sensitive area. As shown in Fig. 7, the dots represent the positions of X-ray irradiation for scanning the uniformity of the detector gain. The scanned uniformity was typically ~16%, as shown in Fig. 8.

3.2. Testing with electron beam

The detection efficiency and spatial resolution of the Micromegas prototypes were investigated with the 5 GeV electron beam at Deutsches Elektronen–Synchrotron (DESY), Germany [12]. Three prototypes

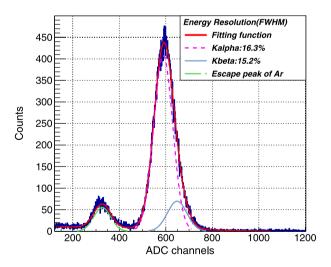


Fig. 5. X-ray spectrum of the 55 Fe source at a gain of ~ 5000 .

(MM01, MM02, and MM03) were installed in the beam line, as shown in Fig. 9. A dedicated front-end electronics (FEE) system for the MPGDs, based on the AGET ASIC chip [13,14], was developed to read out the strip signals. Each channel can record 512 sampling points with adjustable sampling frequency from 3 MHz to 100 MHz and optional peaking time from 50 ns to 1 μs . However, 40 MHz sampling frequency and 1 μs peaking time were selected to record the data in this test.

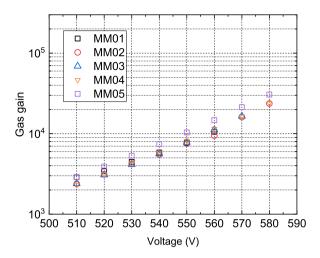


Fig. 6. Dependence of gain on mesh voltage.

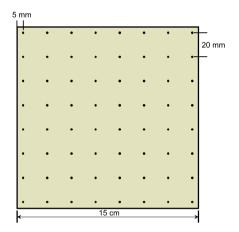


Fig. 7. Schematic diagram showing the scanning positions (dots) in the sensitive area of the prototypes.

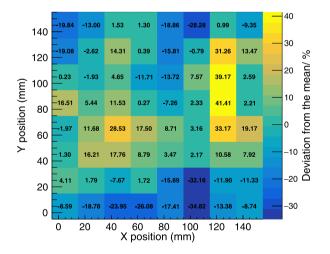


Fig. 8. Map the uniformity of the gain of one of the prototypes ($\sim 16\%$ at mean gas gain of 19082.2 and R.M.S. of 3083.5).

The electrons behave as the minimum ionization particles (MIPs) in the 5 mm drift region of the detectors. The distribution of the energy deposition measured as the sum of ADCs of the fired strips, was fitted with a Landau function convoluted with a Gaussian function accounting for the detector energy resolution. One example of the measured energy

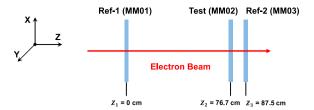


Fig. 9. Electron beam test setup.

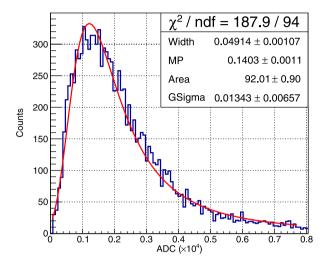


Fig. 10. Detector response to the electron beam.

distribution along with the fitting is illustrated in Fig. 10 for MM02. The mesh voltage was 540 V and the gas gain was $\sim 2 \times 10^4$ for this prototype.

Further, to estimate the spatial resolution, MM01 and MM03 were used as reference tracking detectors and MM02 was the test detector. The deviations of the tested MM02 (\triangle X, here take the deviation of X strip as an example) were defined as the differences between the hit positions (X_t) recorded by MM02 and the expected positions (X_{exp}) derived from the reference trackers. The X_{exp} can be simply calculated by a \cdot Z_t + b, where Z_t is the installing position of MM02 in the Z direction, and a and b are the slope and intercept of the reference tracker, respectively.

The spatial resolution can be estimated by the following equation

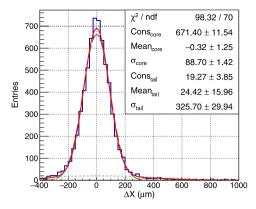
$$\delta^{2}\left(\Delta \; X\right) = (1 + \left(\frac{Z_{2} - Z_{t}}{Z_{2} - Z_{1}}\right)^{2} + \left(\frac{Z_{1} - Z_{t}}{Z_{2} - Z_{1}}\right)^{2})\delta^{2}\left(X_{t}\right) \tag{1}$$

where, Δ $X=X_t-X_{exp}$ is the residual, $\delta\left(X_t\right)$ is the square root of the variance of the residual represents the spatial resolution of the (MM02) detector under test. For the electron beam test, the detectors were placed at $Z_1=0\, cm,\, Z_t=76.7\, cm,$ and $Z_2=87.5\, cm,$ respectively, as shown in Fig. 9. Thus,

$$\delta \left(\Delta \mathbf{X} \right) = 1.36\delta \left(\mathbf{X}_{t} \right) \tag{2}$$

The δ (Δ X) is obtained by fitting the residual distribution with a double Gaussian function, as shown in Fig. 11. Angular divergence of the electron beam and multiple scattering of the electrons both contribute to the tail events in the residual distribution, which are accounted for by the wide component of the double Gaussian function (dotted line). The narrow component or the core Gaussian distribution (dashed line) has a standard deviation of 88.7 \pm 1.42 μ m and hence a spatial resolution δ (X_t) of 65.22 \pm 1.04 μ m.

The detection efficiency is defined as the ratio of effective events recorded by the test prototype (MM02) to coincidence events of reference trackers (MM01 and MM03), where the effective events of MM02



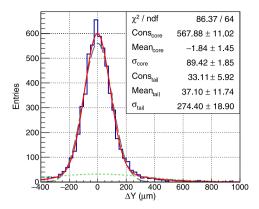


Fig. 11. Deviation distributions of the Micromegas prototypes in the X (left) and Y (right) directions.

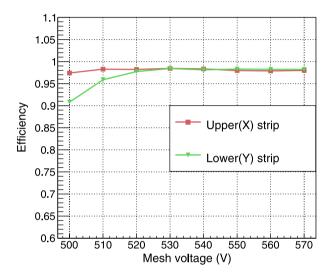


Fig. 12. Detection efficiency of thermal bonding Micromegas.

were selected with the condition of Δ X < 10 · δ (Δ X). Fig. 12 shows the detection efficiency of the individual X and Y strips. The squares represent the efficiency of the X strips, and the inverted triangles represent the efficiency of the Y strips. Efficiencies of both strips reach higher than 98% on the plateau area when the mesh voltage exceed 530 V, which corresponds to a gas gain of ~5000.

4. Applications of TBM micromegas detectors

The best way to ascertain the stability and performance of the detectors is to use them in experiments. Several Micromegas detectors have been designed and manufactured using the TBM for different applications. A six-layer tracker system was built as a cosmic ray telescope using Micromegas detectors. In this telescope, each detector has a sensitive area of $90 \times 90 \text{ mm}^2$ and has been operated periodically for approximately two years, which indicates good stability of these detectors [15].

A 2D readout Micromegas-based neutron detector has been employed as beam monitor in the China Spallation Neutron Source (CSNS) facility, which is an advanced accelerator-based neutron facility proposed to serve both scientific research and industrial applications [16]. Using this novel detector, a 2D neutron beam spot distribution at the Back-n beam line was obtained for the first time [17].

Additionally, owing to the flexibility of the TBM, a new type of detector, having high-gain and low ion-backflow (IBF), with a double micro-mesh gaseous structure (DMM) has been developed using TBM for single photoelectron detection [18,19]. A schematic diagram of

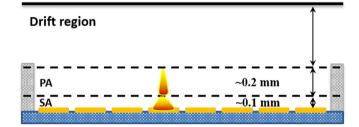


Fig. 13. Schematic diagram of DMM.

the DMM structure is shown in Fig. 13. In this, two stacked meshes are bonded to produce a two-stage cascade avalanche structure. The observed high gain ($>10^6$) and low IBF ratio (\sim 0.05%) of the DMM indicate that it can be potentially used in gaseous photomultiplier tubes (Gas-PMTs) and other applications, for instance, readout of the time projection chamber detectors for the future collider experiments [20, 21]

5. Micromegas detector optimization with TBM

Following the idea mentioned in Section 2.3, the temperature and pressure for thermal bonding were increased to completely squeeze the spacers and five optimized prototypes (OptimizedMM01-05) were fabricated and subsequently tested. The avalanche gaps were measured to be $\sim\!100~\mu m$, which corresponded to the thickness of the polyester layer of the 75–100–75 μm spacers. Consequently, a higher gas gain was obtained compared with the previously obtained gain for the avalanche gap of $\sim\!110~\mu m$. As shown in Fig. 14, the gas gains of all the optimized Micromegas prototypes can easily exceed 5×10^4 at a working voltage above 560 V, and some of them can even reach a gain as high as 10^5 without breakdown.

To further evaluate the impact of changing the avalanche gap length from ${\sim}110~\mu m$ to ${\sim}100~\mu m$, a simulation was performed to study the correlation between the gas gain and avalanche gap. The gas gain can be calculated using the following formula:

$$M = e^{\int [\alpha(x) - \eta(x)] dx}$$
(3)

where $\alpha(x)$ and $\eta(x)$ represent the Townsend and attachment coefficients along the moving path dx. The Townsend and attachment coefficients for the gas mixture of Ar/CO2 (93%/7%) 1 atm were computed using the Magboltz program [22] for various electric field strengths. The gain variation with the avalanche gap sizes at different mesh voltages is displayed in Fig. 15. This result shows that the detector with ~100 μ m avalanche gap has higher gain than that with ~110 μ m gap at the same mesh voltage. The optimized prototypes (with ~100 μ m gap) attained a high mesh voltage, similar to that of the previous

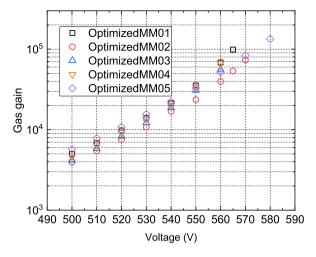


Fig. 14. Gas gains of five optimized Micromegas prototypes. The vertical scale is drawn identical to Fig. 6 for comparison.

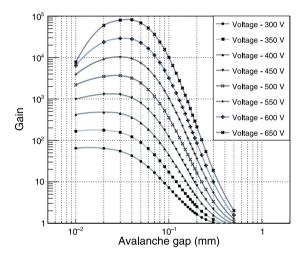


Fig. 15. Gain distribution variation along with the avalanche gap size.

prototypes (with $\sim \! 110~\mu m)$ before undergoing breakdown, and thus, achieved a higher gain.

Concurrently, squeezing the spacers to their unmelting layer helps to improve the uniformity of the avalanche gap, which leads to an expected significant improvement in the uniformity of the gain. As shown in Fig. 16, the value of 8.1% was obtained as compared to the previous result of 16% under the same gain conditions. Additionally, the non-uniformities at different gas gain of ~ 20000 (Fig. 16), ~ 5000 , and ~ 1000 (Fig. 17) were investigated for the OptimizedMM01, and the values of 8.1%, 6.3%, and 5.5% were obtained, respectively. The uniformity within one detector inherently results from the variation of the avalanche gap. Further, the relative differences at low gains become small, which leads to the refinement of the detector uniformity.

6. Conclusions

The TBM, an innovative method for Micromegas fabrication, has been proposed and extensively developed in the past decade. In this TBM, a large distance between the spacers makes it easier to keep the avalanche region clean and reduce the sparking probability. Good performance characteristic has been achieved in terms of gas gain, energy resolution, spatial resolution, and detection efficiency with the Micromegas prototypes. A high gain of $\sim 10^5$, an energy resolution of 16%, and a uniformity of $\sim 6.3\%$ (at gain of ~ 10000) were obtained

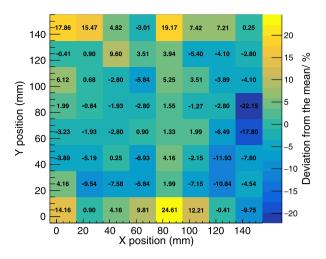


Fig. 16. Map the uniformity of gain of the prototypes (OptimizedMM01, 8.1% at mean gas gain of 19950.1 and R.M.S. of 1610.5).

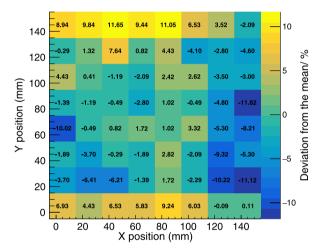


Fig. 17. Map the uniformity of prototypes (OptimizedMM01, 5.5% at mean gas gain of 996.9 and R.M.S. of 54.3).

with 5.9 keV X-rays as the ionizing radiation. Moreover, better than 98% detection efficiency and a spatial resolution of $\sim\!65~\mu m$ were observed with a 5 GeV electron beam. The excellent stability of the Micromegas detectors built with the TBM and the capability of the method to make new detector structures have been demonstrated in various applications, such as tracker systems, neutron monitors, etc.

However, in the TBM, manual setting of the spacers is the most timeconsuming step. An alternative method of automatically setting spacers is being developed, which is critical for the mass production of large area detectors.

CRediT authorship contribution statement

Jianxin Feng: Validation, Methodology, Investigation, Formal analysis, Data curation, Writing - original draft, Writing - review & editing. Zhiyong Zhang: Conceptualization, Methodology, Investigation, Writing - review & editing, Supervision, Project administration. Jianbei Liu: Conceptualization, Methodology, Investigation, Writing - review & editing, Supervision, Project administration. Binbin Qi: Writing - reviewing, Data curation, Editing. Anqi Wang: Writing - reviewing, Data curation, Editing. Ming Shao: Writing - review & editing. Yi Zhou: Writing - review & editing.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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