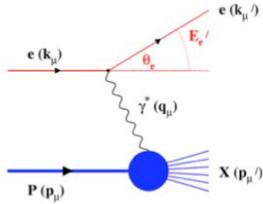


Precision studies of the DVCS process at JLab

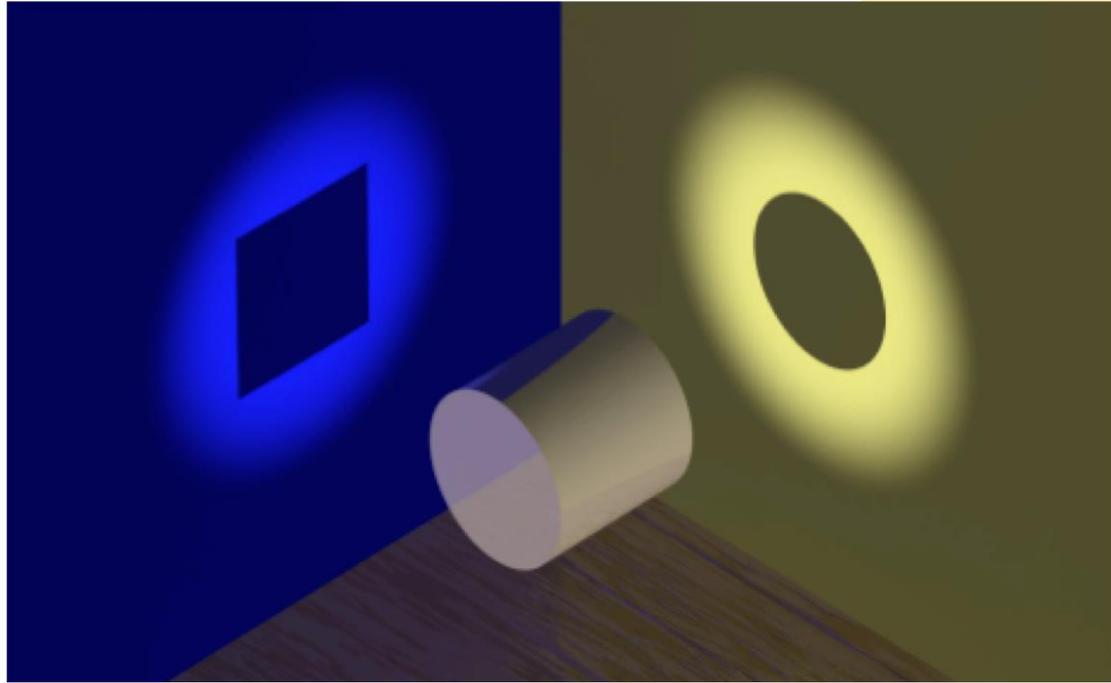
Julie Roche

- Generalized Parton Distributions (GPD) are accessible through many exclusive processes. So far, the most studied channels are Deeply Virtual Compton Scattering (DVCS) and Deep Virtual Meson production (DVMP).
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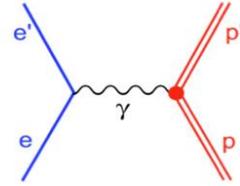
DIS Parton Distribution Functions



No information on the spatial location of the constituents



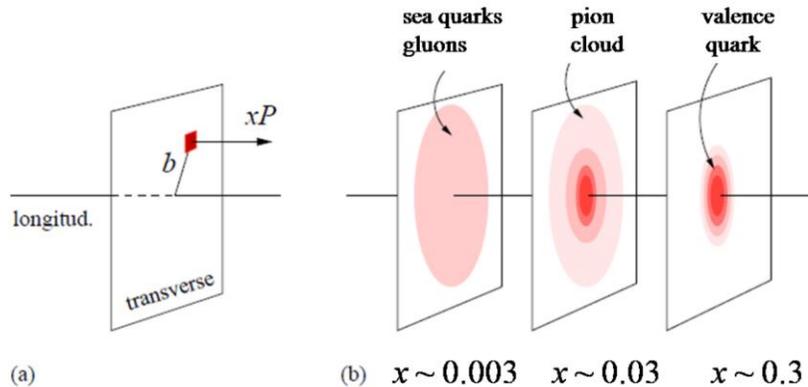
Elastic Form Factors



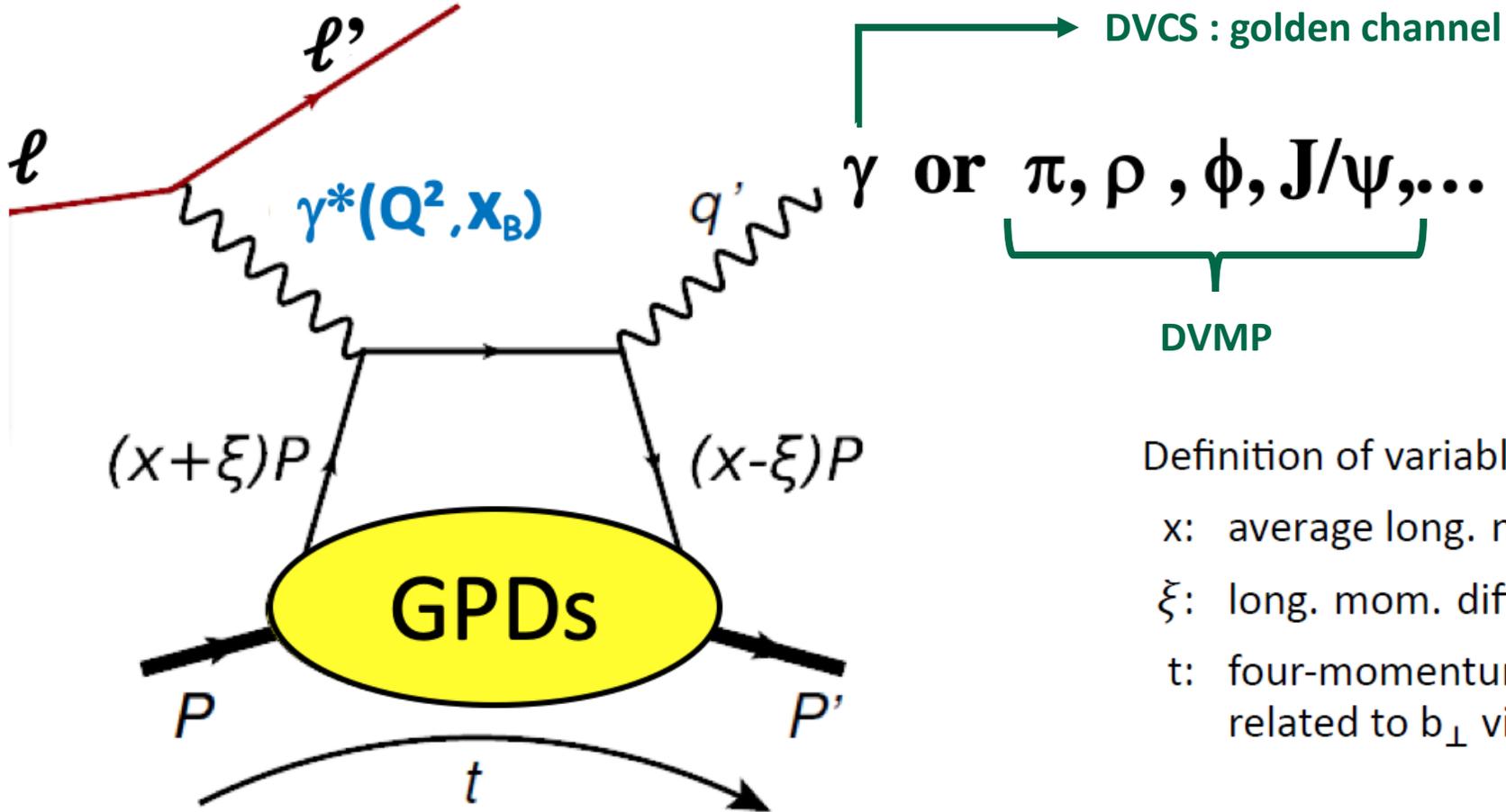
No information about the underlying dynamics of the system

Generalized Parton Distribution Function :

3-D imaging of the nucleon with access to **correlations** between **transverse spatial distribution and longitudinal momentum distributions.**



Exclusive reactions: handbag diagram



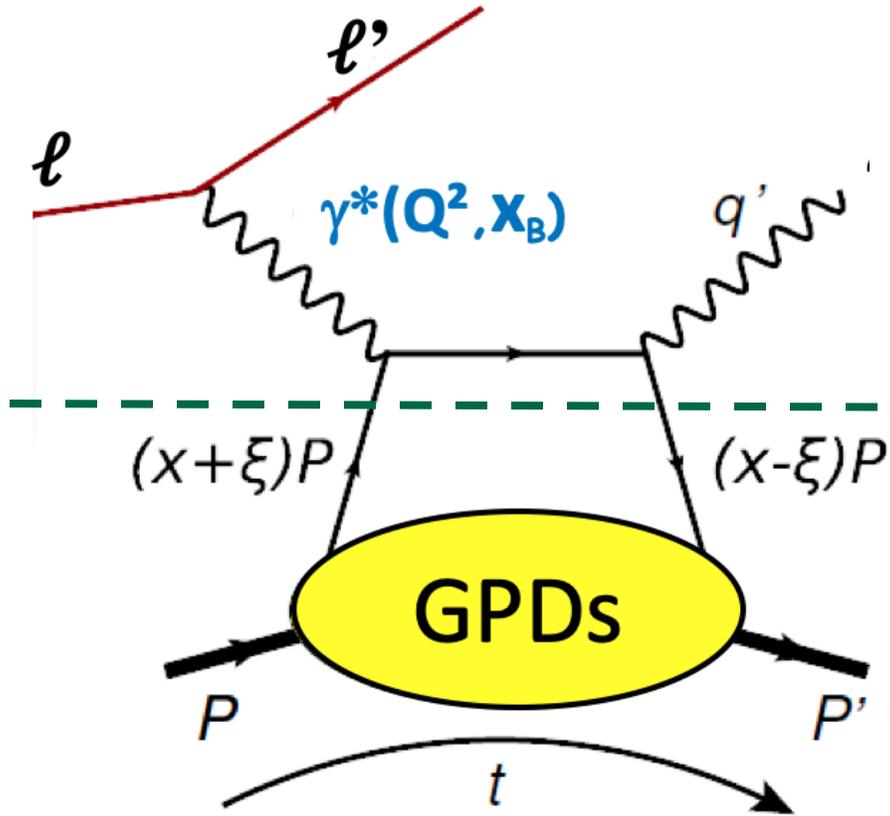
Definition of variables:

x : average long. momentum

ξ : long. mom. difference $\simeq x_B/(2 - x_B)$

t : four-momentum transfer
related to b_{\perp} via Fourier transform

In the Bjorken limit: $Q^2 = \frac{-q^2}{\nu} \rightarrow \infty \left. \vphantom{Q^2} \right\} x_B = \frac{Q^2}{2M\nu} \text{ fixed}$



Hard process
LO: QED
NLO: QCD perturbative

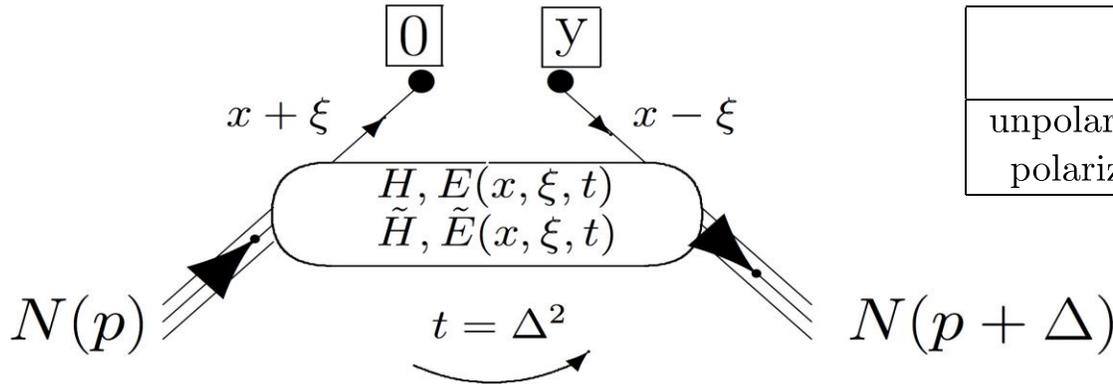
D. Mueller *et al*, Fortsch. Phys. 42 (1994)
X.D. Ji, PRL 78 (1997), PRD 55 (1997)
A. V. Radyushkin, PLB 385 (1996), PRD 56 (1997)

Soft process
Non perturbative QCD
described by GPD(x, ξ, t)

x is not accessible, one measures $CFF(\xi, t) = CFF(x_B, t)$

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i \pi H(x = \xi, \xi, t)$$

Generalized Parton Distributions



	Nucleon Helicity	
	conserving	non-conserving
unpolarized GPD	H	E
polarized GPD	\tilde{H}	\tilde{E}

$\lim_{t \rightarrow 0} (GPD) \rightarrow PDF$

DIS

$$H^q(x, 0, 0) = q(x), -\bar{q}(-x)$$

$$\tilde{H}^q(x, 0, 0) = \Delta q(x), \Delta \bar{q}(-x)$$

GPD first moments \rightarrow Form Factors

Elastic

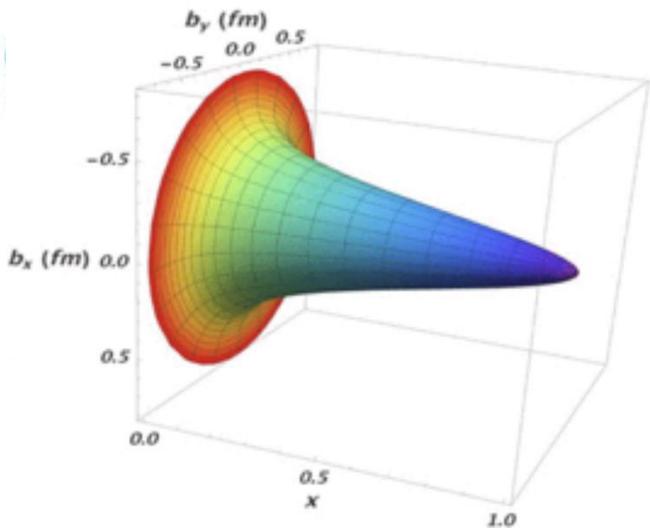
$$\int_{-1}^{+1} dx H^q(x, \xi, t) = F_1^q(t) \quad \int_{-1}^{+1} dx \tilde{H}^q(x, \xi, t) = g_A^q(t)$$

$$\int_{-1}^{+1} dx E^q(x, \xi, t) = F_2^q(t) \quad \int_{-1}^{+1} dx \tilde{E}^q(x, \xi, t) = h_A^q(t)$$

No relation for the GPD E and \tilde{E}



GPDs and hadronic physics issues

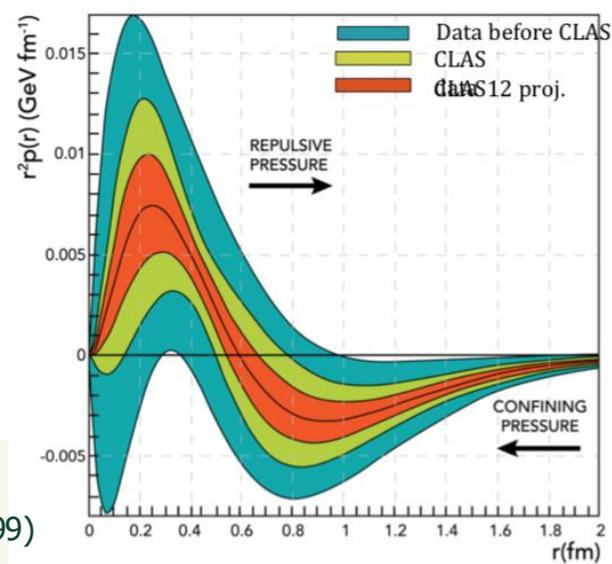
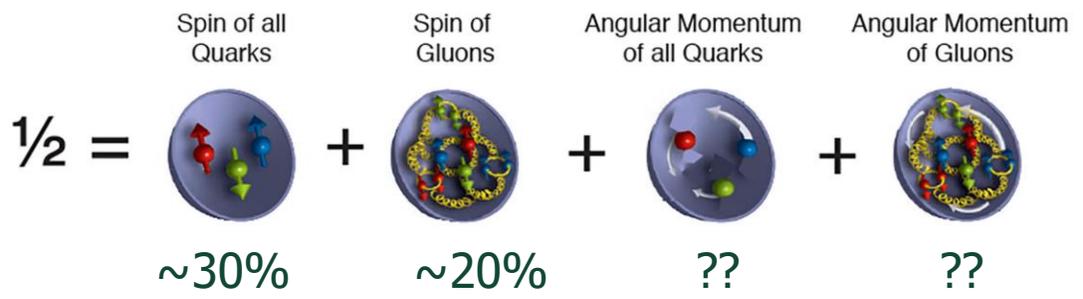


Nucleon tomography

R. Dupré et al. PRD95 (2017)
 Local fits from Hermes, CLAS & Hall A data
 Assume H contribution only
 Model dependent assumptions for x_B dependence

Contribution of the **angular momentum of quark** to the proton spin (Ji's sum rule: PRL 78(4):610-613, 1997)

$$\frac{1}{2} = \underbrace{\frac{1}{2} \Delta \Sigma + L_q}_{J_q} + J_g \Rightarrow J_q = \frac{1}{2} \int_{-1}^1 dx x [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$



V. Burkert,
 L. Elouadrhiri,
 F.X. Girod
 Nature 557
 (2018, 7705, 396-399)

Access to the **mechanical properties/gravitational form factors** of the proton

$$\mathcal{H} = \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi + i\epsilon} = \mathcal{P} \int_{-1}^{+1} dx \frac{H(x, \xi, t)}{x - \xi} - i \pi H(x = \xi, \xi, t)$$

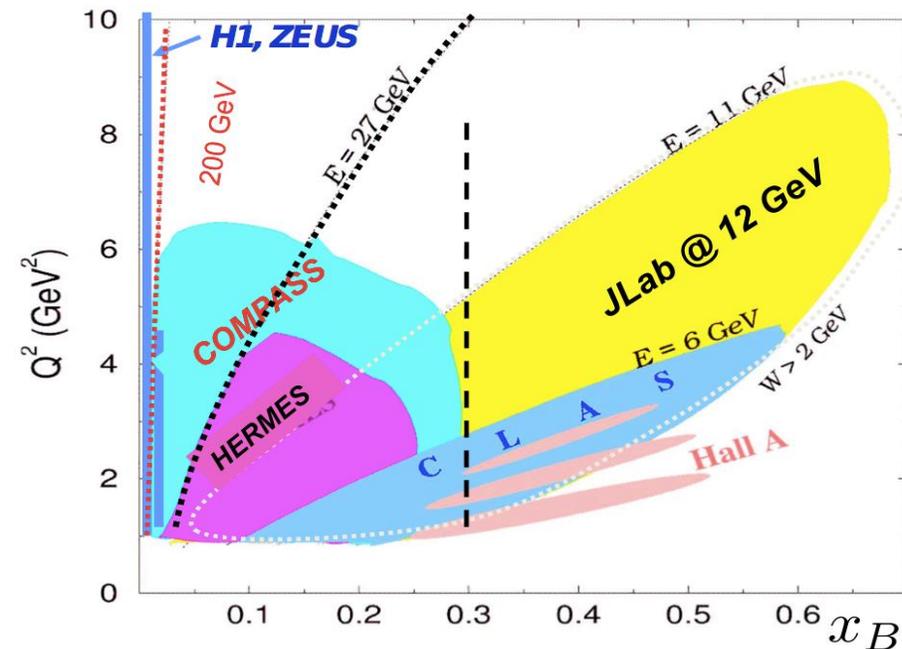
$$D(t) = -\text{Re}(\mathcal{H}) + \int dx \frac{\text{Im} \mathcal{H}}{x + \xi}$$



JLab Exclusive reactions for GPDs program

Tables by C. Munoz-Camacho

Measurement	Hall
DVCS Polarized beam and/or target	A,B,C
nDVCS Deuterium/He3	B,C, A(Solid)
DVCS w/ e+	B, C
TCS	A (Solid), B, C
Excl. π^0	A,B,C
Excl. π^-	A (Solid), (B)
Excl. ϕ, η	B
L/T separation (K, π^+)	C
WACS (γ, π^0)	A, C
Backwards π^0	C



In the valence region (JLab 6 and JLab 12)

Partially complimentary, overlapping

- Hall A/C: Test the validity of the formalism
 - high accuracy (~5%)
 - limited kinematic
- Hall B: Map the GPDS
 - limited accuracy (15+%)
 - wide kinematic range

JLab Exclusive reactions for GPDs program

This talk focuses on
recent
DVCS and DVMP- π^0 using the Hall A/C scheme
and the future of this niche technique.

	Experiment	PAC	Goal	Results
6 GeV	E00-110	PAC18	1 st dedicated DVCS experiment at JLab	PRL97 (2006) , PRC83 (2011) , PRC92 (2015)
	E03-106	PAC24	1 st neutron DVCS experiment	PRL99 (2007)
	E07-007	PAC31	DVCS Rosenbluth-like separation (proton)	PRL117 (2016) , Nature Commun. 8 (2017)
	E08-025	PAC33	DVCS Rosenbluth-like separation (neutron)	PRL118 (2017) , Nature Physics 16 (2020)
12 GeV	E12-06-114	PAC30+38+41+47	1 st 12 GeV experiment	PRL127 (2021) , PRL128 (2022)
	E12-13-010	PAC40	DVCS Rosenbluth-like separation (proton)	<i>Scheduled 2023-2024</i>

Precision studies of the DVCS process at JLab

Julie Roche

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F. Georges,¹ M. N. H. Rashad,² A. Stefanko,³ M. Dlamini,⁴ B. Karki,⁴ S. F. Ali,⁵ P-J. Lin,¹ H-S Ko,^{1,6} N. Israel,⁴ D. Adikaram,⁷ Z. Ahmed,⁸ H. Albataineh,⁹ B. Aljawrneh,¹⁰ K. Allada,¹¹ S. Allison,² S. Alsalmi,¹² D. Androic,¹³ K. Aniol,¹⁴ J. Annand,¹⁵ H. Atac,¹⁶ T. Averett,¹⁷ C. Ayerbe Gayoso,¹⁷ X. Bai,¹⁸ J. Bane,¹⁹ S. Barcus,¹⁷ K. Bartlett,¹⁷ V. Bellini,²⁰ R. Beminiwatha,²¹ J. Bericic,⁷ D. Biswas,²² E. Brash,²³ D. Bulumulla,² J. Campbell,²⁴ A. Camsonne,⁷ M. Carmignotto,⁵ J. Castellano,²⁵ C. Chen,²² J-P. Chen,⁷ T. Chetry,⁴ M. E. Christy,²² E. Cisbani,²⁶ B. Clary,²⁷ E. Cohen,²⁸ N. Compton,⁴ J. C. Cornejo,^{17,3} S. Covrig Dusa,⁷ B. Crowe,²⁹ S. Danagoulian,¹⁰ T. Danley,⁴ F. De Persio,²⁶ W. Deconinck,¹⁷ M. Defurne,³⁰ C. Desnault,¹ D. Di,¹⁸ M. Duer,²⁸ B. Duran,¹⁶ R. Ent,⁷ C. Fanelli,¹¹ G. Franklin,³ E. Fuchey,²⁷ C. Gal,¹⁸ D. Gaskell,⁷ T. Gautam,²² O. Glamazdin,³¹ K. Gnanvo,¹⁸ V. M. Gray,¹⁷ C. Gu,¹⁸ T. Hague,¹² G. Hamad,⁴ D. Hamilton,¹⁵ K. Hamilton,¹⁵ O. Hansen,⁷ F. Hauenstein,² W. Henry,¹⁶ D. W. Higinbotham,⁷ T. Holmstrom,³² T. Horn,^{5,7} Y. Huang,¹⁸ G. M. Huber , C. E. Hyde ,² H. Ibrahim,³³ C-M. Jen,³⁴ K. Jin,¹⁸ M. Jones,⁷ A. Kabir,¹² C. Keppel,⁷ V. Khachatryan,^{7,35,36} P. M. King,⁴ S. Li,³⁷ W. B. Li,⁸ J. Liu,¹⁸ H. Liu,³⁸ A. Liyanage,²² J. Magee,¹⁷ S. Malace,⁷ J. Mammei,³⁹ P. Markowitz,²⁵ E. McClellan,⁷ M. Mazouz ,⁴⁰ F. Meddi,²⁶ D. Meekins,⁷ K. Mesik,⁴¹ R. Michaels,⁷ A. Mkrtchyan,⁵ R. Montgomery,¹⁵ C. Muñoz Camacho,^{1,*} L. S. Myers,⁷ P. Nadel-Turonski,⁷ S. J. Nazeer,²² V. Nelyubin,¹⁸ D. Nguyen,¹⁸ N. Nuruzzaman,²² M. Nycz,¹² O. F. Obretch,²⁷ L. Ou,¹¹ C. Palatchi,¹⁸ B. Pandey,²² S. Park,³⁵ K. Park,² C. Peng,⁴² R. Pomatsalyuk,³¹ E. Pooser,⁷ A. J. R. Puckett,²⁷ V. Punjabi,⁴³ B. Quinn,³ S. Rahman,³⁹ P. E. Reimer,⁴⁴ J. Roche ,⁴ I. Sapkota,⁵ A. Sarty,⁴⁵ B. Sawatzky,⁷ N. H. Saylor,⁴⁶ B. Schmookler,¹¹ M. H. Shabestari,⁴⁷ A. Shahinyan,⁴⁸ S. Sirca,⁴⁹ G. R. Smith,⁷ S. Sooriyaarachchilage,²² N. Sparveris,¹⁶ R. Spies,³⁹ T. Su,¹² A. Subedi,⁴⁷ V. Sulkosky,⁵⁰ A. Sun,³ L. Thorne,³ Y. Tian,⁵¹ N. Ton,¹⁸ F. Tortorici,²⁰ R. Trotta,⁵² G. M. Urciuoli,²⁶ E. Voutier,¹ B. Waidyawansa,⁷ Y. Wang,¹⁷ B. Wojtsekhowski,⁷ S. Wood,⁷ X. Yan,⁵³ L. Ye,⁴⁷ Z. Ye,¹⁸ C. Yero,²⁵ J. Zhang,¹⁸ Y. Zhao,³⁵ and P. Zhu⁵⁴

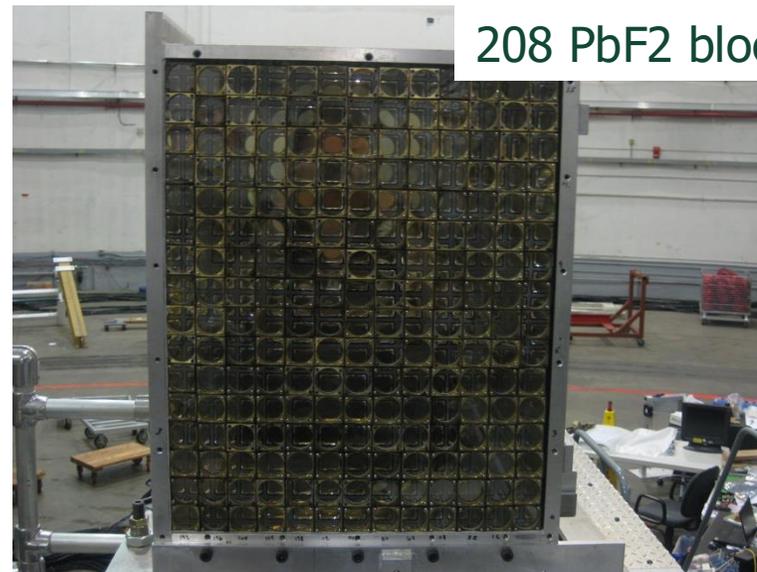
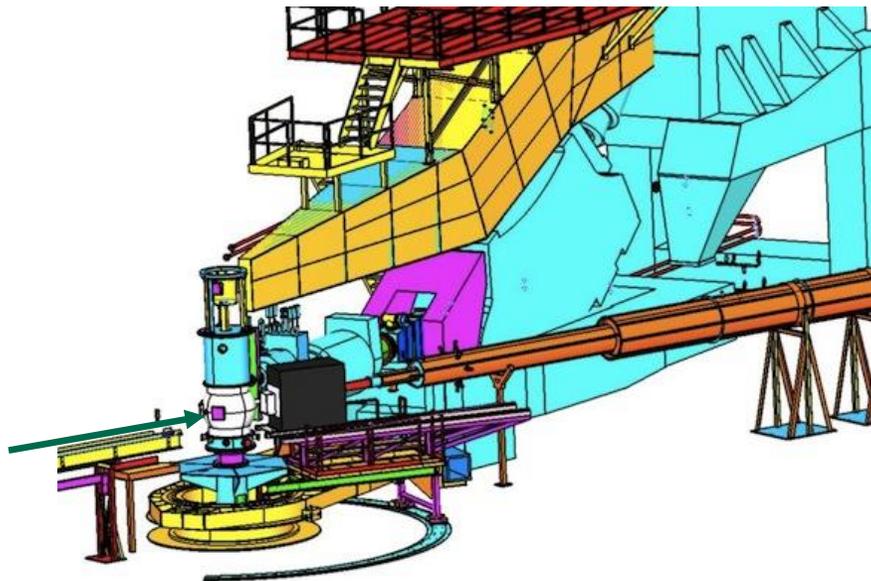


Students/Postdocs Spoke-persons

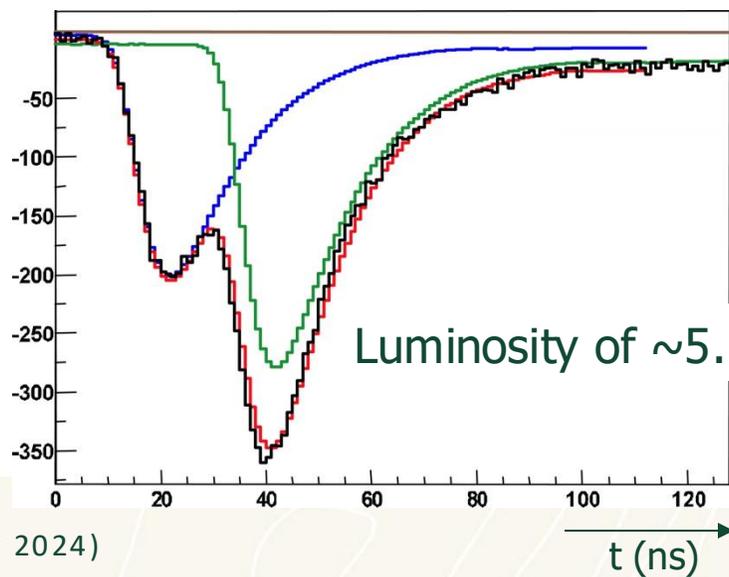
The Jefferson Lab Hall A/C technical staff
The Jefferson Lab Accelerator staff



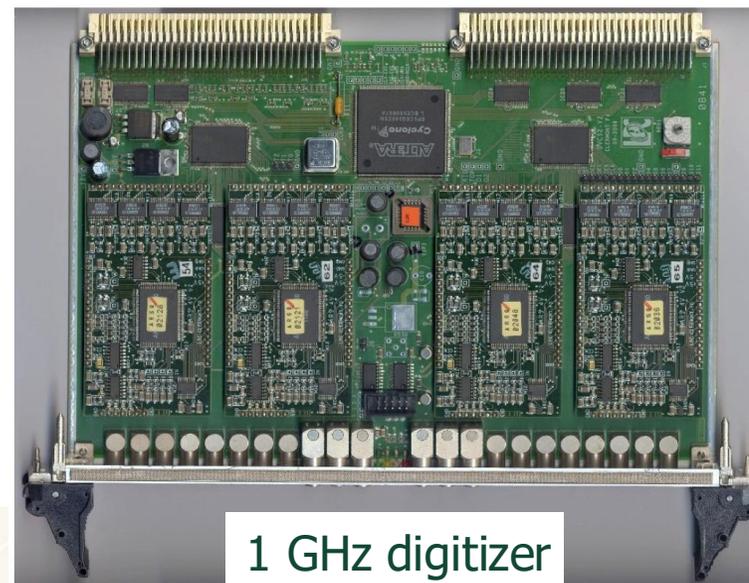
Dedicated apparatus e.g. the Hall A scheme



208 PbF2 blocks

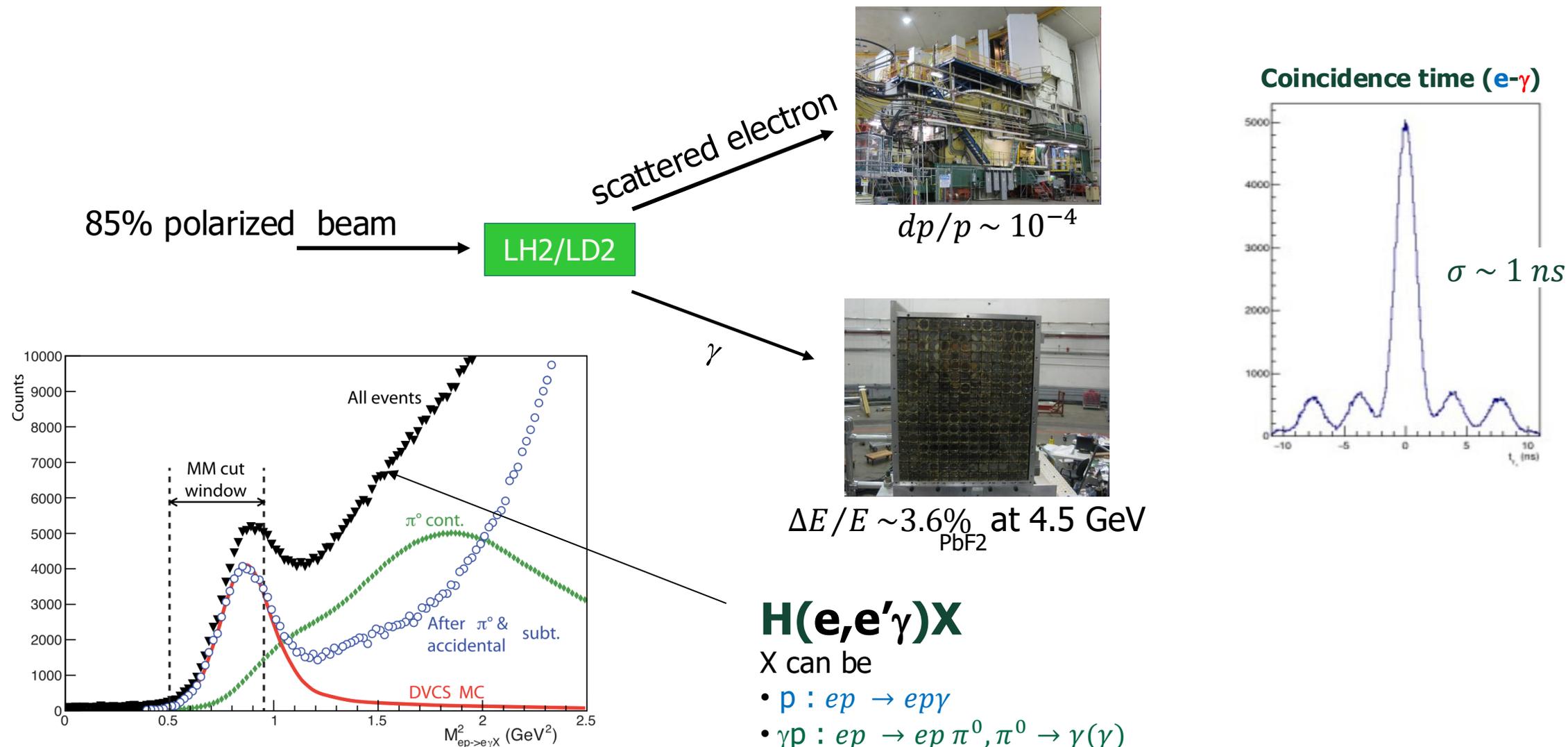


Luminosity of $\sim 5 \cdot 10^{37}$ Hz/cm²



1 GHz digitizer

Exclusivity : the DVCS@Hall A scheme



$H(e, e'\gamma)X$

X can be

- $p : ep \rightarrow ep\gamma$
- $\gamma p : ep \rightarrow ep\pi^0, \pi^0 \rightarrow \gamma(\gamma)$
- $\gamma\gamma p : ep \rightarrow ep\pi^0, \pi^0 \rightarrow \gamma\gamma$
- $N\pi : ep \rightarrow eN\gamma\pi$

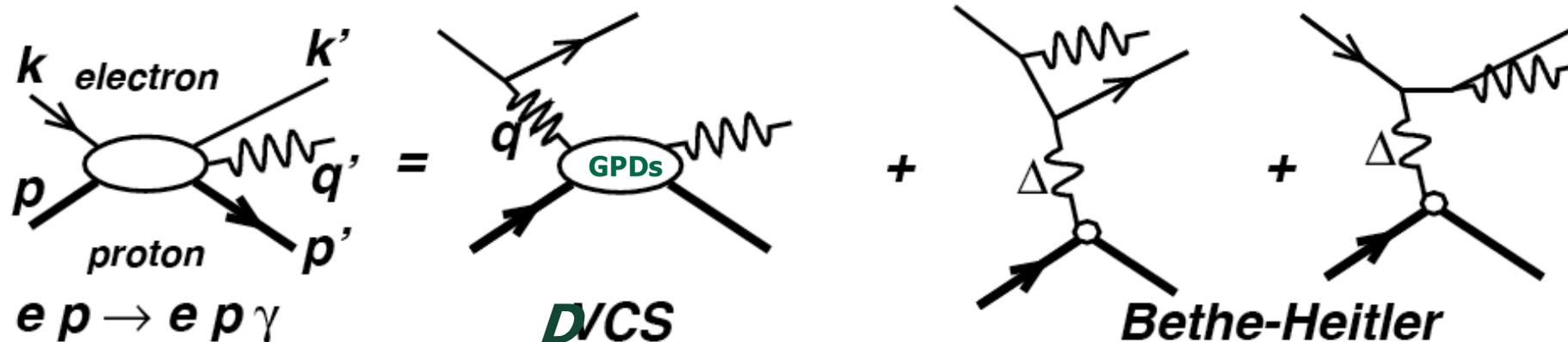
...

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Measuring DVCS to access GPDs information



$$\frac{d^4\sigma(lp \rightarrow lp\gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{\text{BH}} + \underbrace{d\sigma_{\text{unpol}}^{\text{DVCS}} + \mathbf{P}_1 d\sigma_{\text{pol}}^{\text{DVCS}}}_{\text{Bilinear combinations of GPD/CFFs}} + \underbrace{e_1 (\text{Re}(\mathbf{I}) + \mathbf{P}_1 \text{Im}(\mathbf{I}))}_{\text{Linear combinations of GPD/CFFs and FFs}}$$

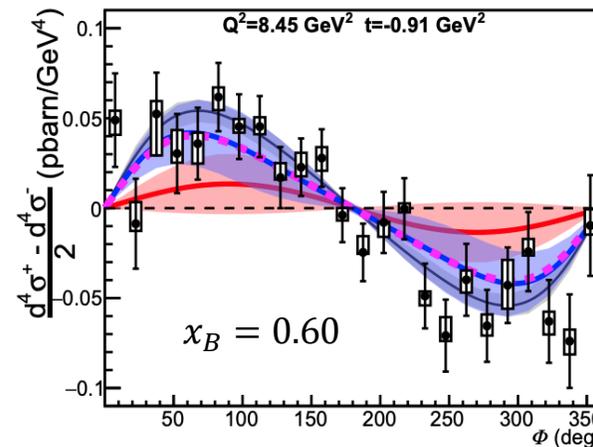
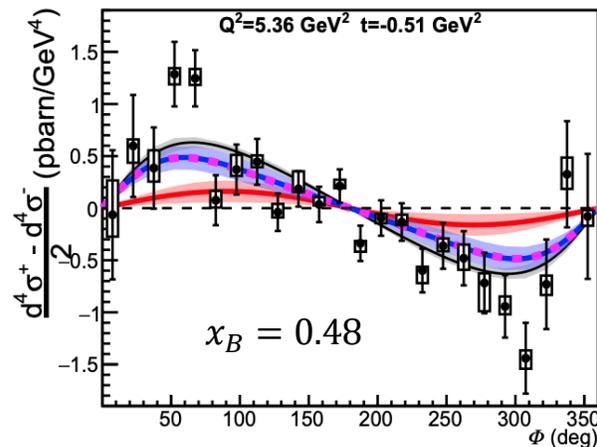
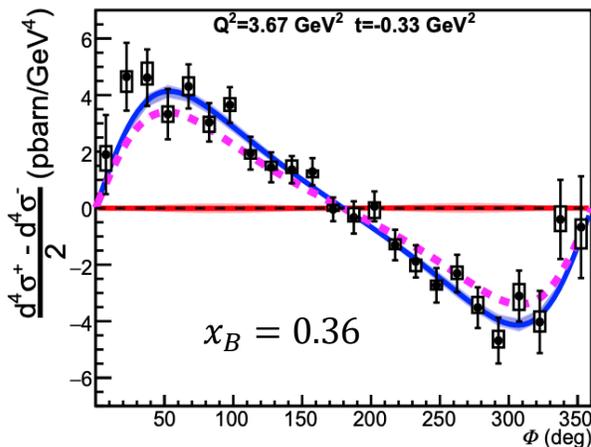
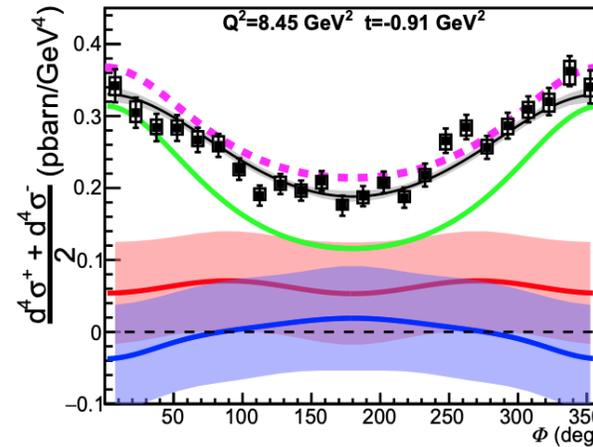
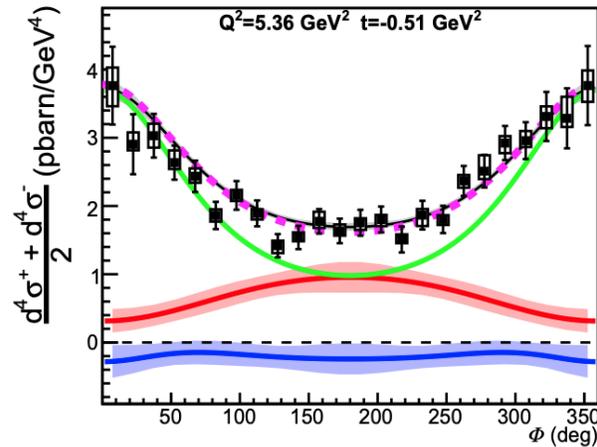
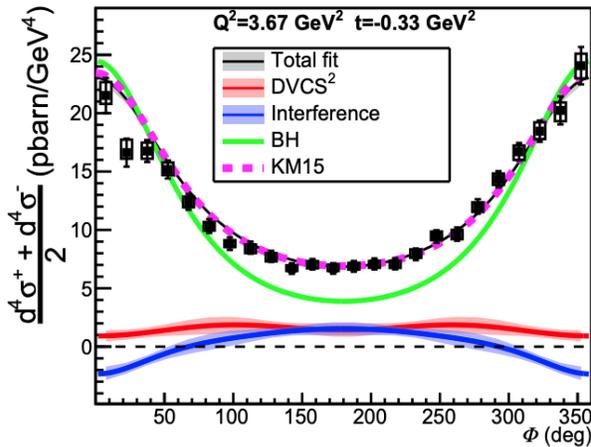
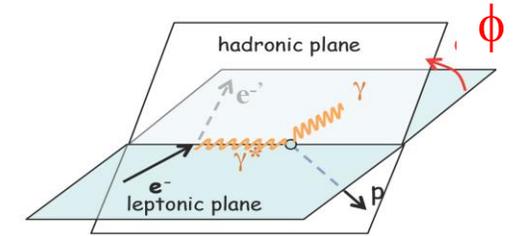
Known if
Nucleon FFs are
known

Bilinear combinations
of GPD/CFFs

Linear combinations
of GPD/CFFs and FFs

\mathbf{P}_1 : polarization target or beam
 e_1 : charge of the lepton beam

$$\frac{d^4\sigma(lp \rightarrow lp\gamma)}{dx_B dQ^2 d|t| d\phi} = \underbrace{d\sigma^{\text{BH}}}_{\text{BH}} + \underbrace{d\sigma_{\text{unpol}}^{\text{DVCS}} + \mathbf{P}_1 d\sigma_{\text{pol}}^{\text{DVCS}}}_{\text{DVCS}^2} + \underbrace{e_1 (\text{Re}(\mathbf{I}) + \mathbf{P}_1 \text{Im}(\mathbf{I}))}_{\text{Interference}}$$



[KM15] K. Kumericki and D. Mueller, *EPJ Web Conf.* 112 (2016) 01012

Braun-Manashov-Müller-Pimay formalism, *Phys. Rev. D* 89, 074022 (2014).



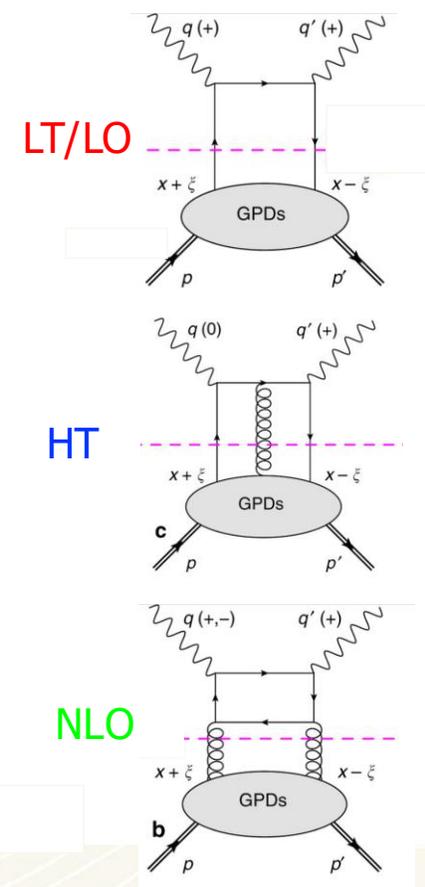
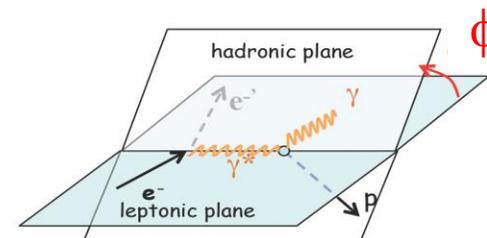
How to parametrize the DVCS cross-sections?

$$\frac{d^4\sigma(\text{lp} \rightarrow \text{lp}\gamma)}{dx_B dQ^2 d|t| d\phi} = d\sigma^{\text{BH}} + d\sigma_{\text{unpol}}^{\text{DVCS}} + \mathbf{P}_1 d\sigma_{\text{pol}}^{\text{DVCS}} + \mathbf{e}_1 (\text{Re}(I) + \mathbf{P}_1 \text{Im}(I))$$

$$\begin{aligned} d\sigma^{\text{BH}} &\propto c_0^{\text{BH}} + c_1^{\text{BH}} \cos \phi + c_2^{\text{BH}} \cos 2\phi \\ d\sigma_{\text{unpol}}^{\text{DVCS}} &\propto c_0^{\text{DVCS}} + c_1^{\text{DVCS}} \cos \phi + c_2^{\text{DVCS}} \cos 2\phi \\ d\sigma_{\text{pol}}^{\text{DVCS}} &\propto s_1^{\text{DVCS}} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

$$s_1^I = F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} + kF_2 \mathcal{E}$$

World-wide GPDs analysis include more or less terms:
both in terms of harmonics (c_i 's and s_i 's) and
in term of GPD/CFFs.



DVCS in Hall A@Jlab program

1st Generation (2004)

Q^2 dependence study (of red terms)

2nd Generation (2010)

Beam energy dependence study at fixed x_B and Q^2

- Separate C_0^{DVCS} from C_0^I
- Separate HT and NLO from LT/LO coefficients

3rd Generation (2014-2016)

Multiple x_B and Q^2 measurements

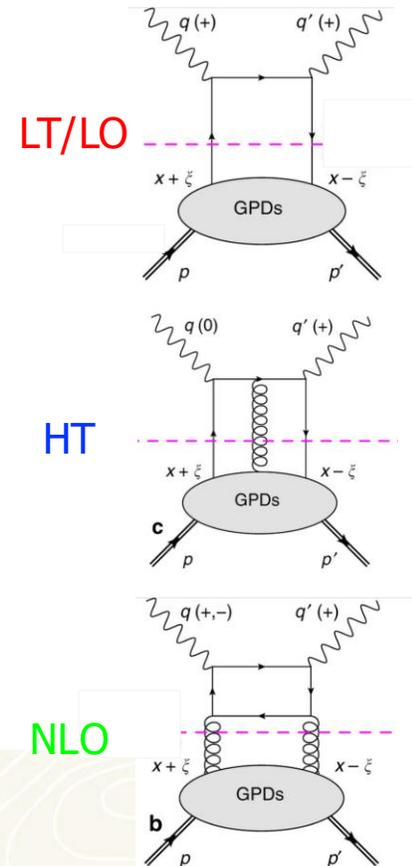
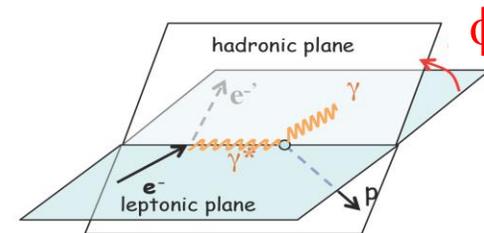
- Experimental extraction of the CFFs as a function of x_B
- Importance of considering all CFFs when extracting CFFs

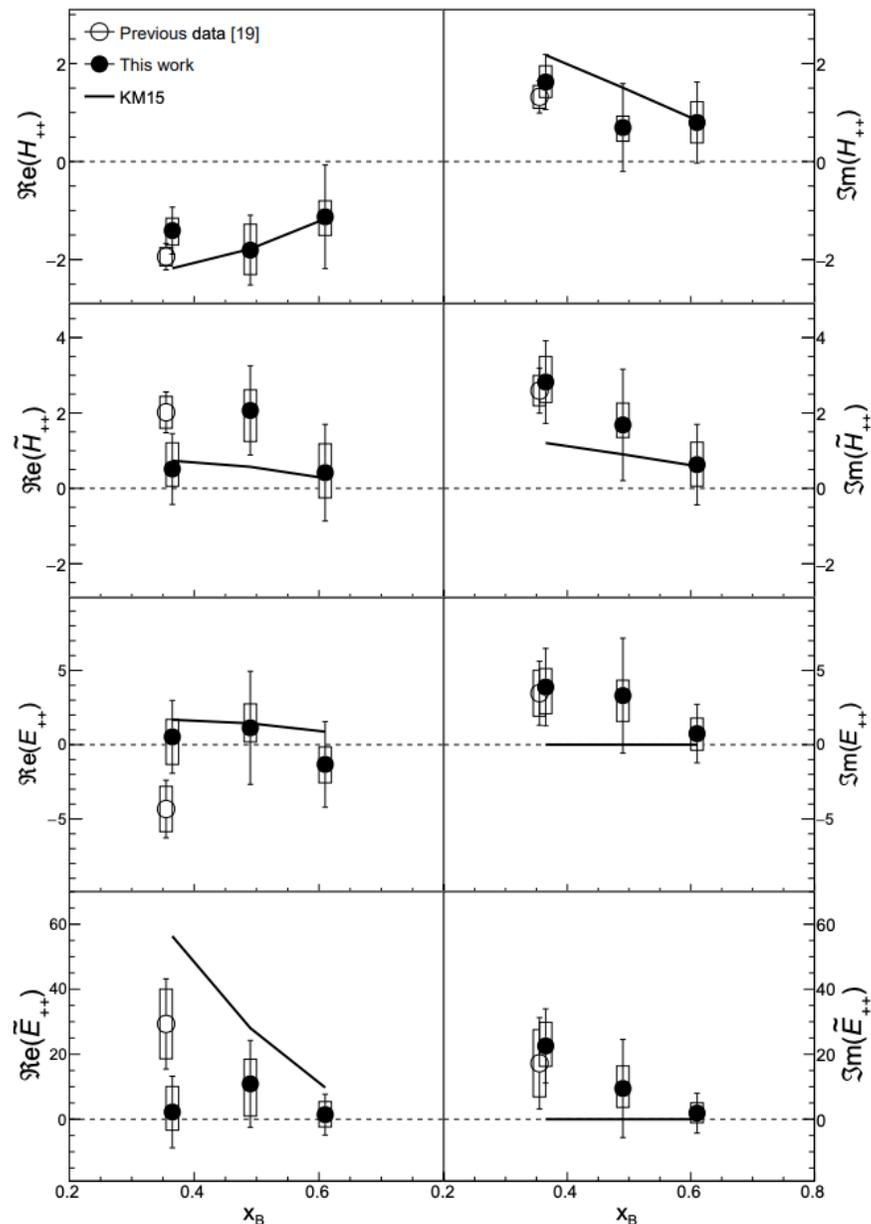
Results:

**off LH2 and LD2 (neutron)
on DVCS and DVMP- π^0**

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$

$$s_1^I = F_1 \mathcal{H} + \xi(F_1 + F_2) \tilde{\mathcal{H}} + k F_2 \mathcal{E}$$





The average t values are -0.281 GeV^2 for [19] and $-0.345, -0.702, -1.050 \text{ GeV}^2$ at $x_B=0.36, 0.48, 0.60$, respectively for this work..

DVCS in Hall A: 12 GeV results

[This work] **F. Georges, *Phys.Rev.Lett.* 128 (2022) 25, 252002**

Error bars: statistical

Error boxes: systematic

[19] M. Defurne et al., *Phys. Rev. C* 92, 055202 (2015)

[KM15] K. Kumericki and D. Mueller, *EPJ Web Conf.* 112 (2016) 01012

The precise measurement of cross-sections at the same x_B - Q^2 bin but multiple beam energies is essential for this extraction.

Also demonstrated in

M. Defurne et al., *Nat. Commun.* 8, 1408 (2017).

B. Kriesten et al., *Phys. Rev. D* 101, 054021 (2020).

M. Čuić et al., *Phys. Rev. Lett.* 125, 232005 (2020).

$\text{CFF}_{\lambda\lambda'}$ λ : polarization state of virtual photon (0,+,-)

λ' : polarization state of outgoing real photon (+,-)

Fit has 24 CFF

$$(\tilde{H}, H, \tilde{E}, E) \otimes (\Re, \Im) \otimes (+, +, 0, +, +, -)$$

but **only the results from the LO ones (++) are shown.**

Fits performed at constant x_B and t over Q^2 and ϕ bins.

No Q^2 evolution of the CFFs.



A combined neutron and proton targets data allows for flavor separation of the GPDs.

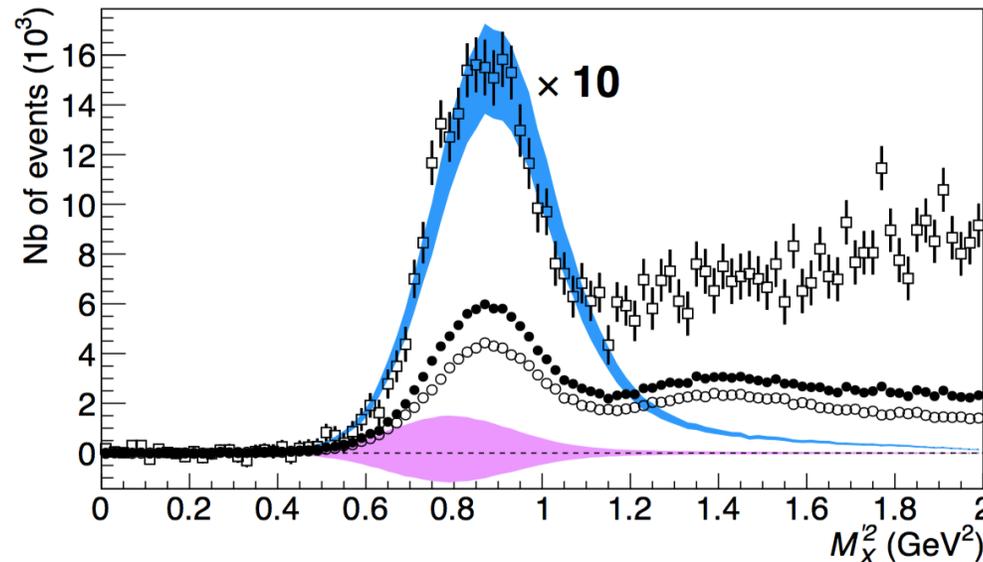
Neutron data are uniquely sensitive to the elusive GPD E (no connection to PDFs).

Below the two pions threshold: $D(e, e' \pi^0)X = d(e, e' \pi^0)d + n(e, e' \pi^0)n + p(e, e' \pi^0)p$.

- LD2
- LH2
- LD2-LH2

■ $d(e, e' \pi^0)d$
■ $n(e, e' \pi^0)n$
 separated by

$$\Delta M_X^2 = t(1 - M/M_d) \approx t/2$$



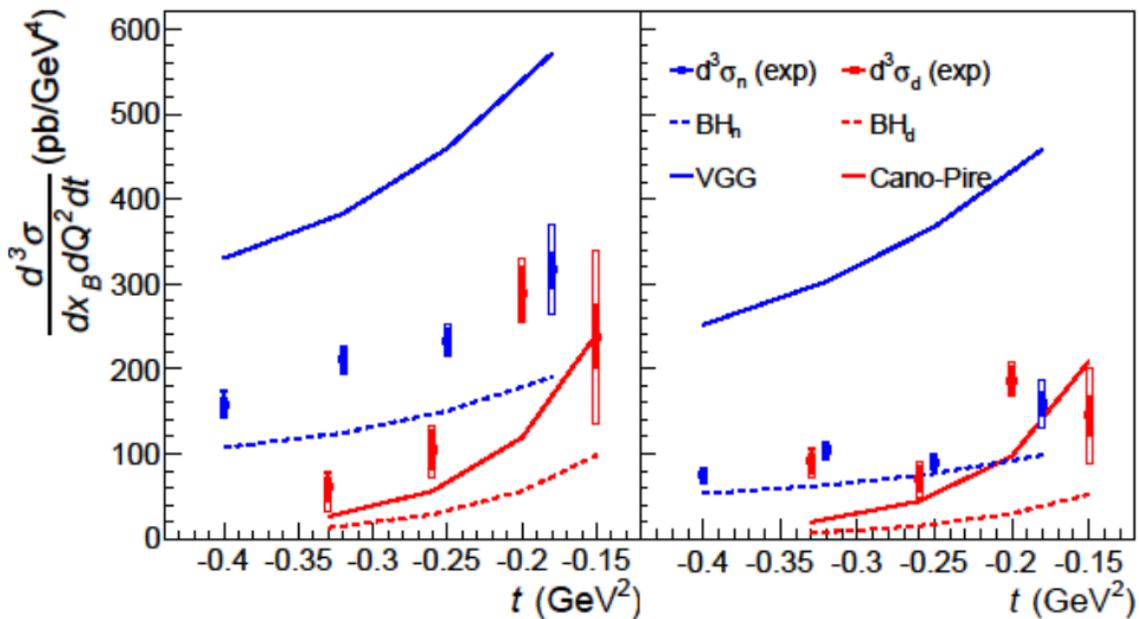
Computed for $n(e, e' \pi^0)X$

Figure from M. Mazouz PRL 118 (2017) 22, 222002



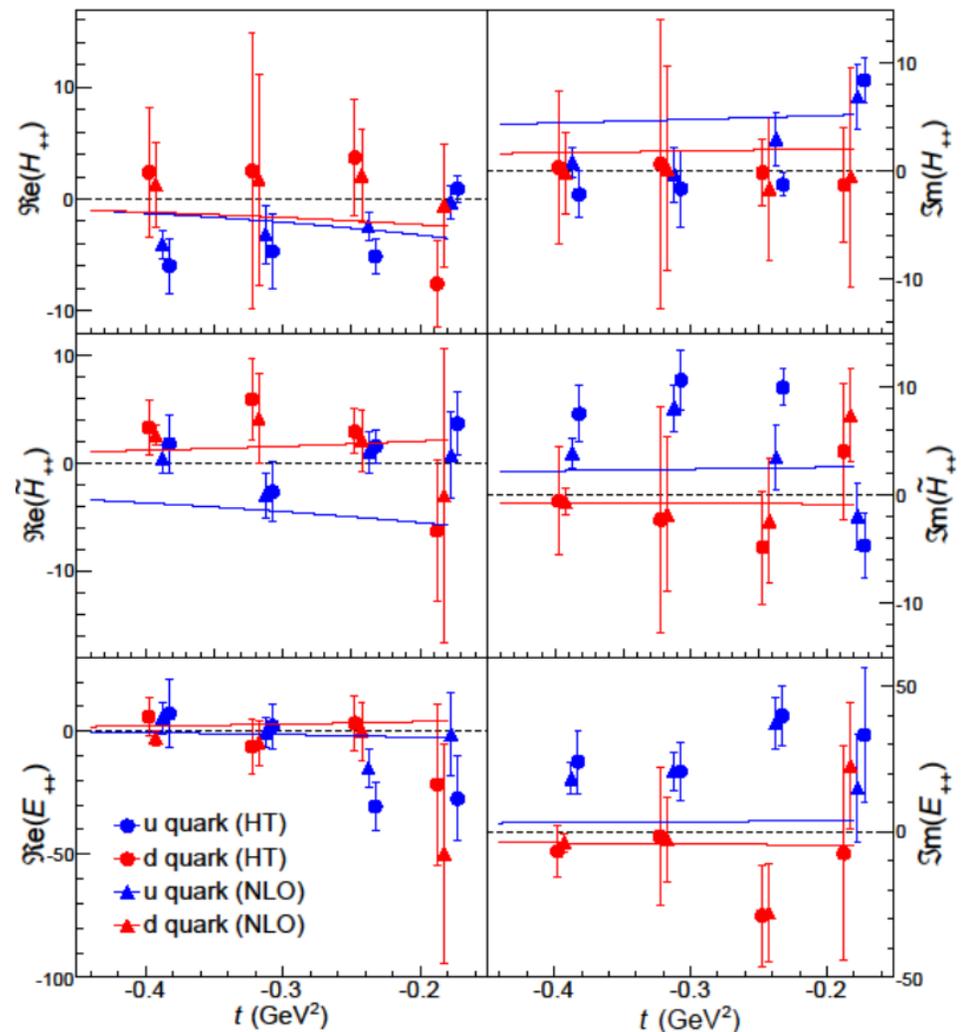
Flavor separation of Compton Form Factors

Cross section measurements from E08-205

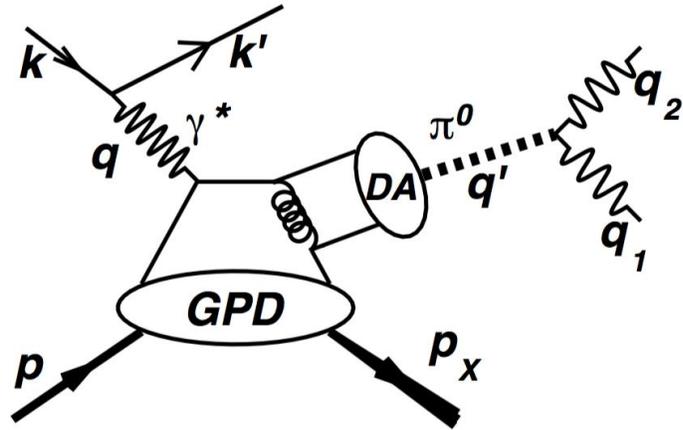


$Q^2 = 1.9 \text{ GeV}^2, x_B = 0.36$

Benali et al, Nature Phys. 16, 191 (2020)



M. Defurne et al. PRL 117, 26 (2015)



4 chiral-even GPDs
4 chiral-odd GPDs (not seen in DVCS)

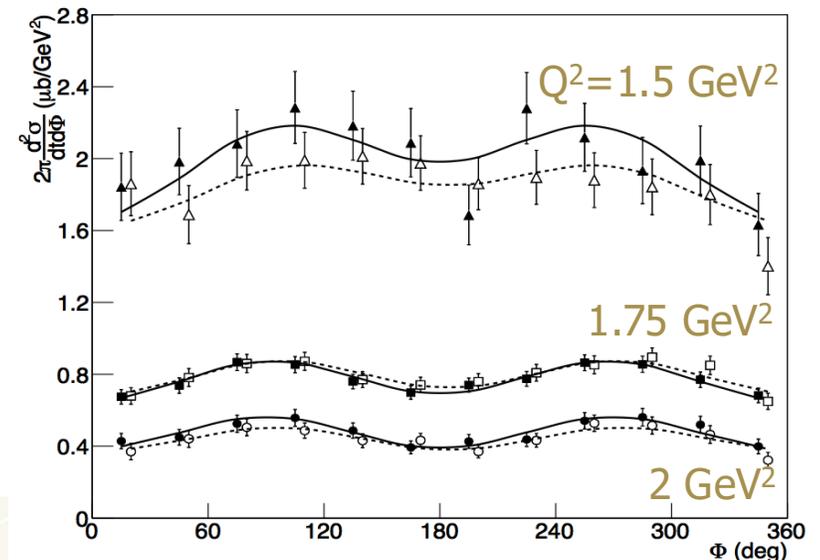
Leading twist, leading order factorization is only proven for $d\sigma_{\perp}/dt$ (Collins et al. 1997)

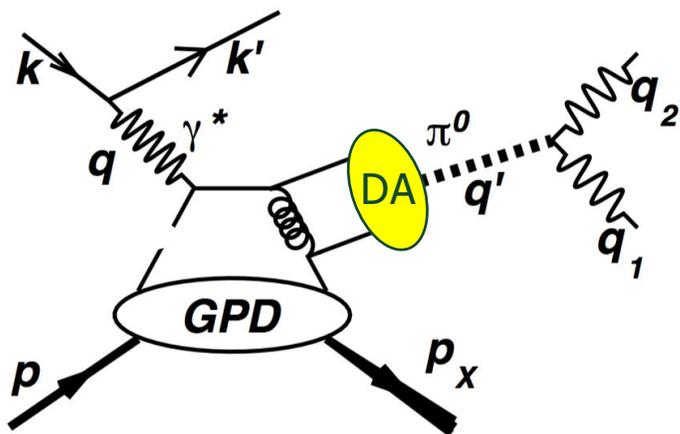
$$\frac{d^4\sigma}{dt d\phi dQ^2 dx_B} = \frac{1}{2\pi} \Gamma_{\gamma^*}(Q^2, x_B, E_e) \left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{TL}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) \right]$$

Setting	E (GeV)	Q^2 (GeV ²)	x_B	ϵ
2010-Kin1	(3.355 ; 5.55)	1.5	0.36	(0.52 ; 0.84)
2010-Kin2	(4.455 ; 5.55)	1.75	0.36	(0.65 ; 0.79)
2010-Kin3	(4.455 ; 5.55)	2	0.36	(0.53 ; 0.72)

Dominance of $d\sigma_T/dt$ observed (directly or not) also at

- Hermes & Hall C π^+
- Hall B, Hall A π^0



GPDs through DVMP- π^0 

Factorization is only exact for longitudinal photons (Collins et al., 1997).

At large Q^2 , QCD predicts that $\sigma_L \rightarrow Q^2 \sigma_T$

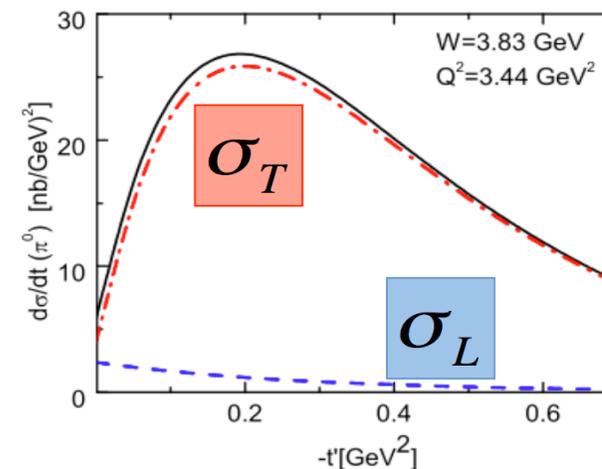
Effective transverse factorization schemes exploit and explain the observed $\sigma_T > \sigma_L$ for existing high x data.

S. Goloskokov and P. Kroll (Eur.Phys.J A47, 112(2011))

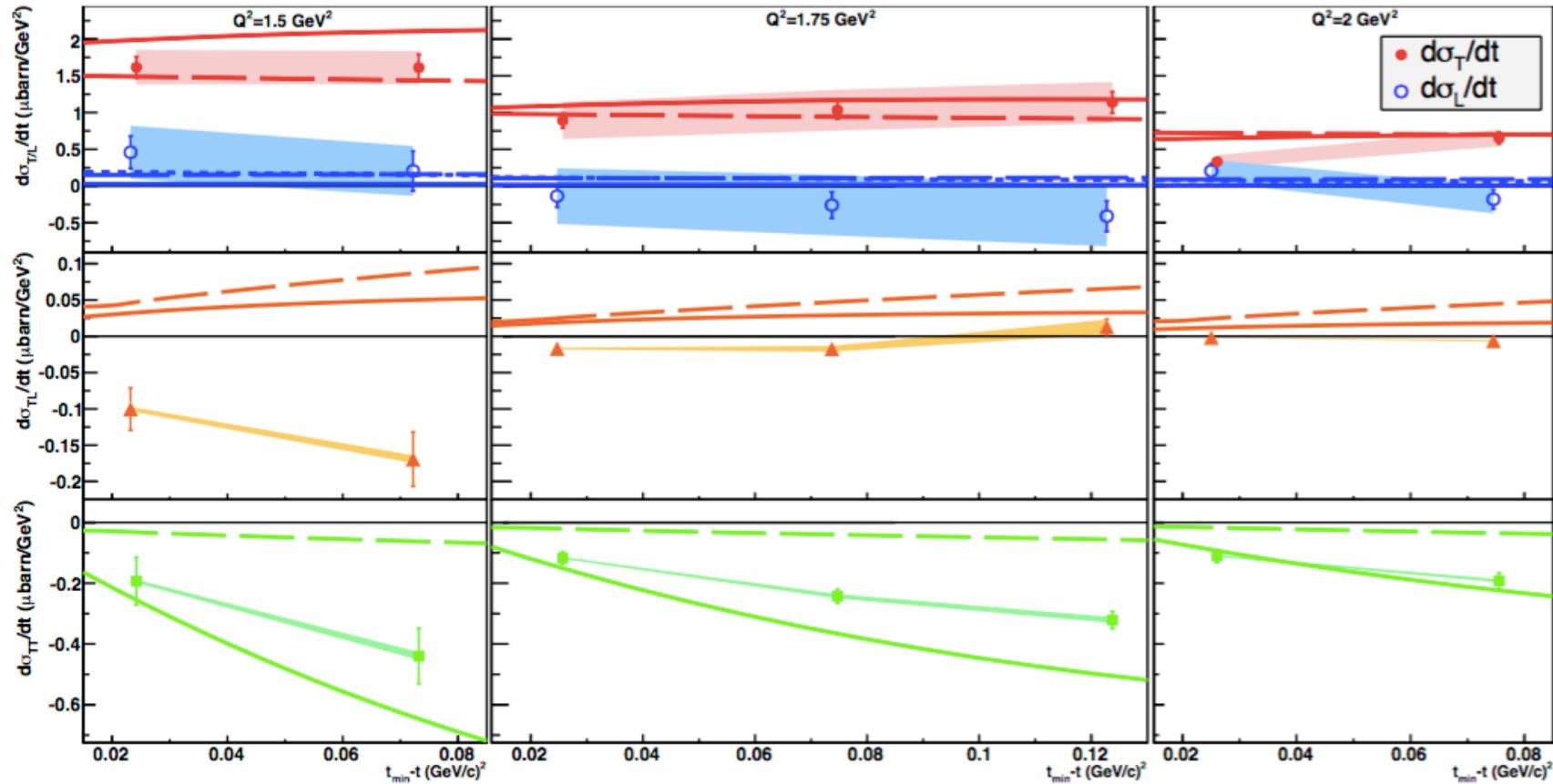
S. Liuti and G. Golstein (Phys.Rev.D79, 054014 (2009))

Twist 3 Distribution Amplitudes (DA) couple with transversity GPDs.

The dominant contribution is not the leading twist contribution



M. Defurne et al. PRL 117, 26 (2015)



Models with $d\sigma_T$ factorization scheme

- — — G-H-L (JPG:NPP39 '12)
- — — G-K (EPJA47 '11)

We observed:

- Small $d\sigma_L$, large $d\sigma_T$: models ok with these
- Wrong sign and t dependence for $d\sigma_{TL}$ and $d\sigma_{TT}$
- sizeable $d\sigma_{TL} \Rightarrow d\sigma_L$ is small but not null

Precision studies of the DVCS process at JLab

Julie Roche

- GPDs are accessible through many exclusive processes. So far, the most studied channels are Deeply Virtual Compton Scattering (DVCS) and Deep Virtual Meson production (DVMP).
- Our JLab Hall A/C collaboration measures DVCS and DVMP- π^0 absolute cross-section using a “simple” dedicated apparatus since 2006.
- Our precise (5%) measurements are essential to interpret the data.
- We just completed a new set of measurements in Hall C/JLab using a new PbWO_4 calorimeter and expanding our kinematic reach.

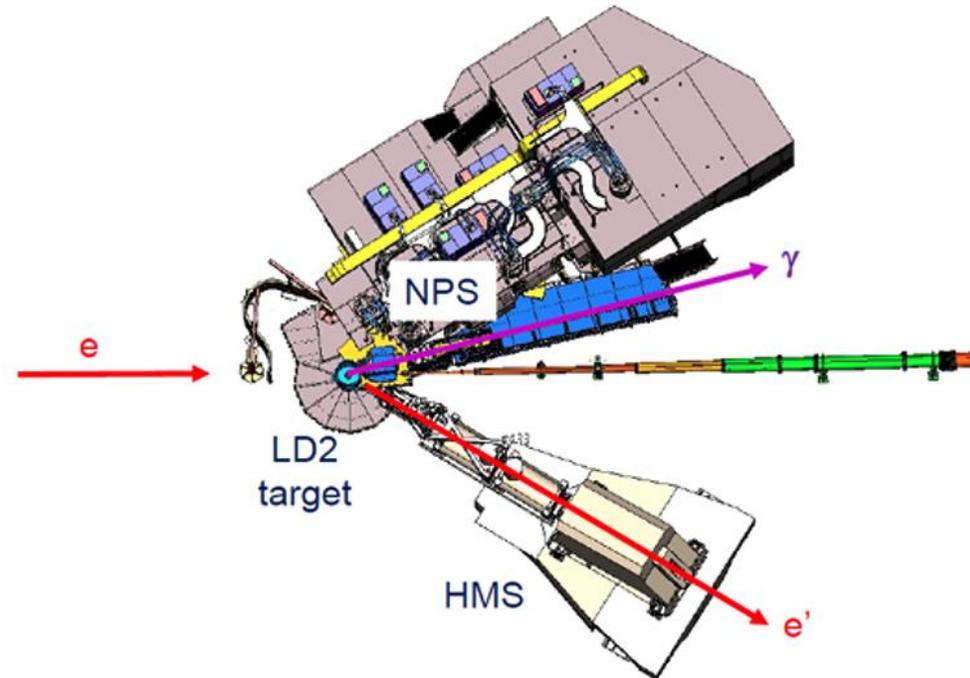
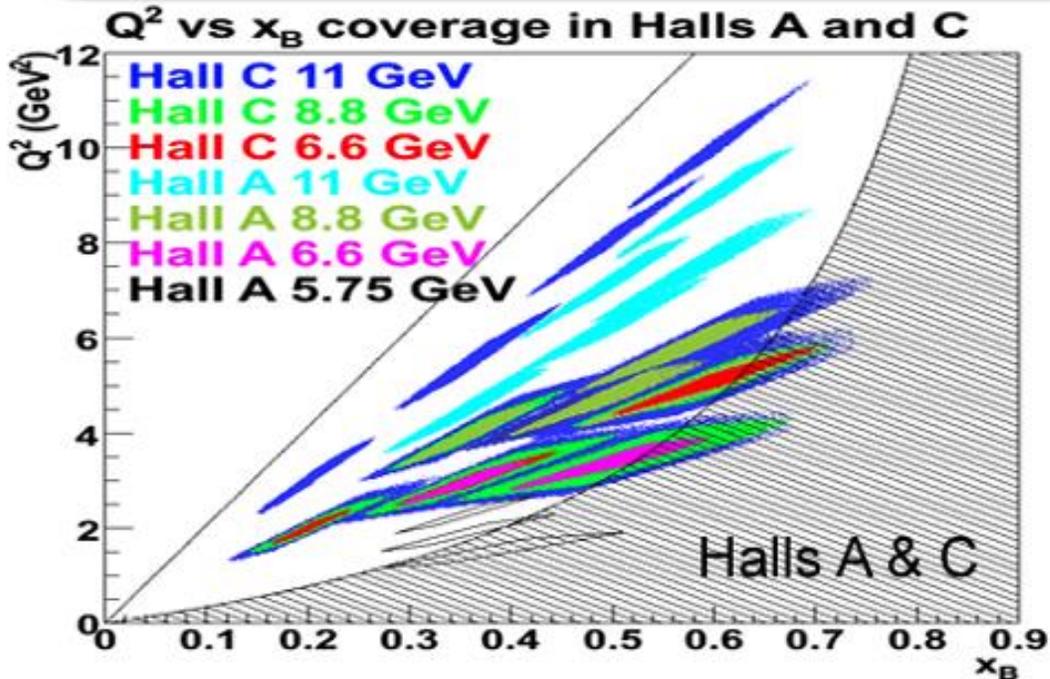
Neutral Particle Spectrometer (Hall C/JLab)

E12-13-010 DVCS measurements follow up on measurements in Hall A:

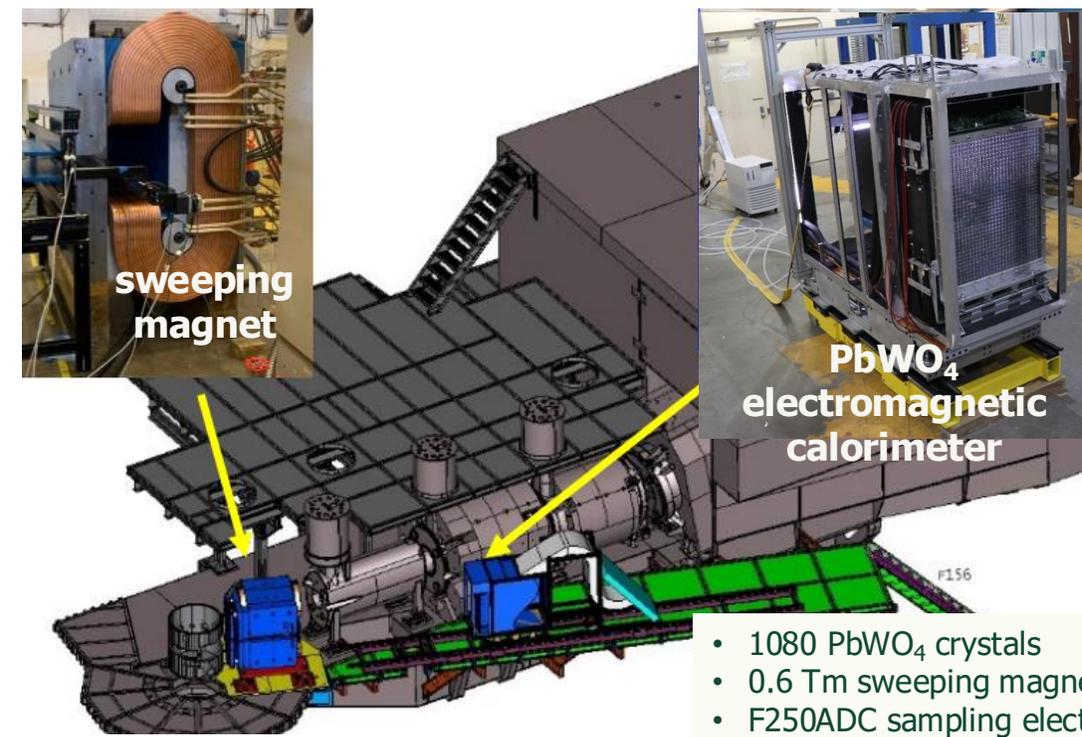
- Scaling of the Compton Form Factor
- Rosenbluth-like separation of DVCS:

$$\sigma = |BH|^2 + \text{Re}\left[\underset{\sim E_{beam}^2}{DVCS^\perp BH} \right] + \underset{\sim E_{beam}^3}{|DVCS|^2}$$
- L/T separation of π^0 production

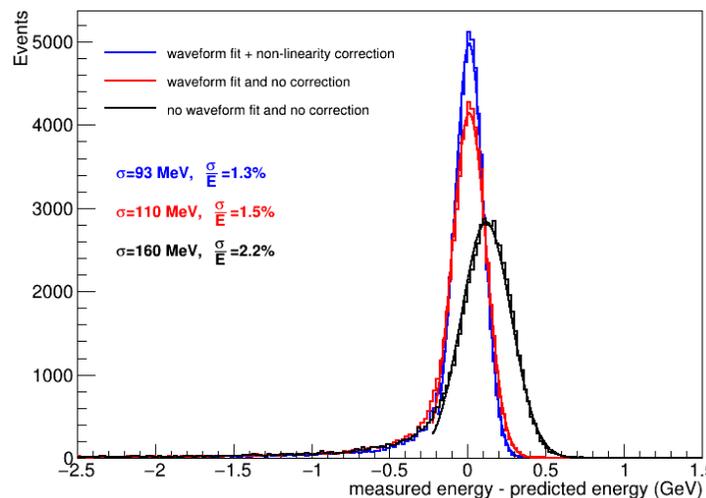
- The NPS calorimeter consists of 1080 PbWO_4 crystals, the preferred material for high-resolution calorimetry, also at EIC – NPS has the largest set of PbWO_4 crystals in an operating calorimeter in the US
- The NPS Science Program consists of ten approved experiments. **4** experiments have been running in parallel from Sept '23 to May '24.



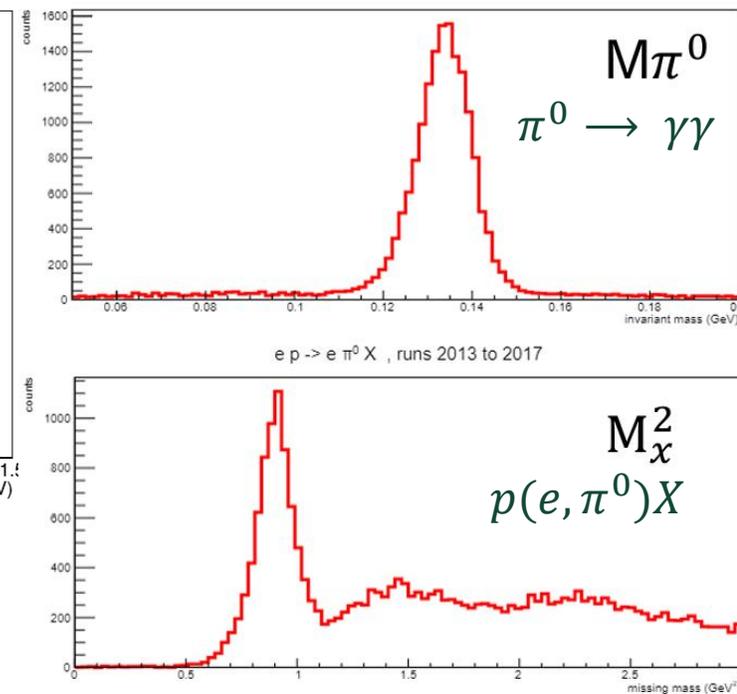
Neutral Particle Spectrometer (Hall C/JLab)



Elastic NPS Energy Resolution



π^0 Resolution



All channels perform well at very high luminosity on LH2 and LD2 ($\approx 8 \times 10^{37} \text{ cm}^2/\text{s}$). The expected resolution energy resolution was achieved (1.3% at 7.3 GeV).

Precision studies of the DVCS process at JLab

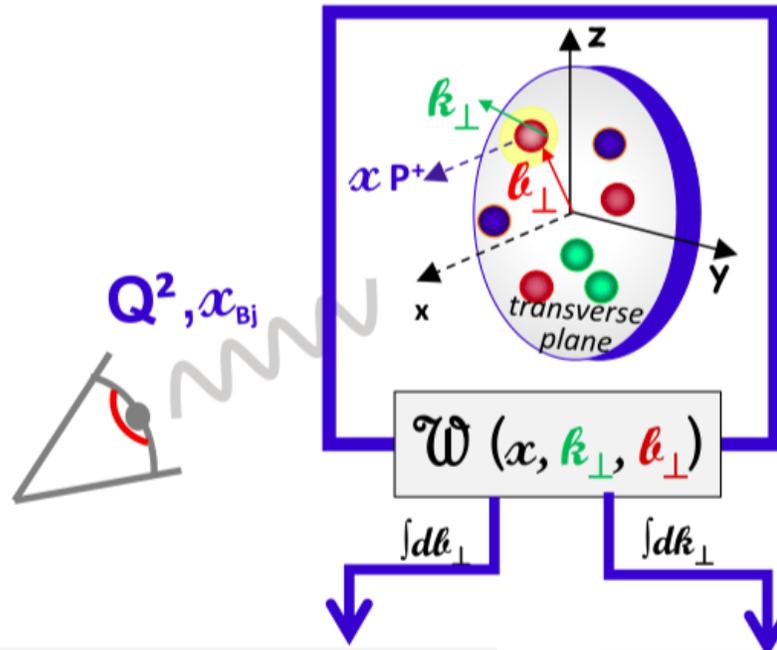
- Our JLab Hall A/C collaboration measures DVCS and DVMP- π^0 absolute cross-section using a “simple” dedicated apparatus since 2006.
- Our precise (5%) measurements are essential to interpret the data.
 - proton, neutron and exclusive pi-zero absolute cross section measurement
 - total and beam helicity correlated cross section measurement
 - cover the valence region ($x_B=0.2$ to 0.6) with Q^2 in the 1-10 GeV^2 range
 - repeated (x_B - Q^2) measurement at different beam energies
- We just completed a new set of measurements in Hall C/JLab using a new PbWO_4 calorimeter and expanding our kinematic reach.

Thank you for your attention!

Thank you for your attention!

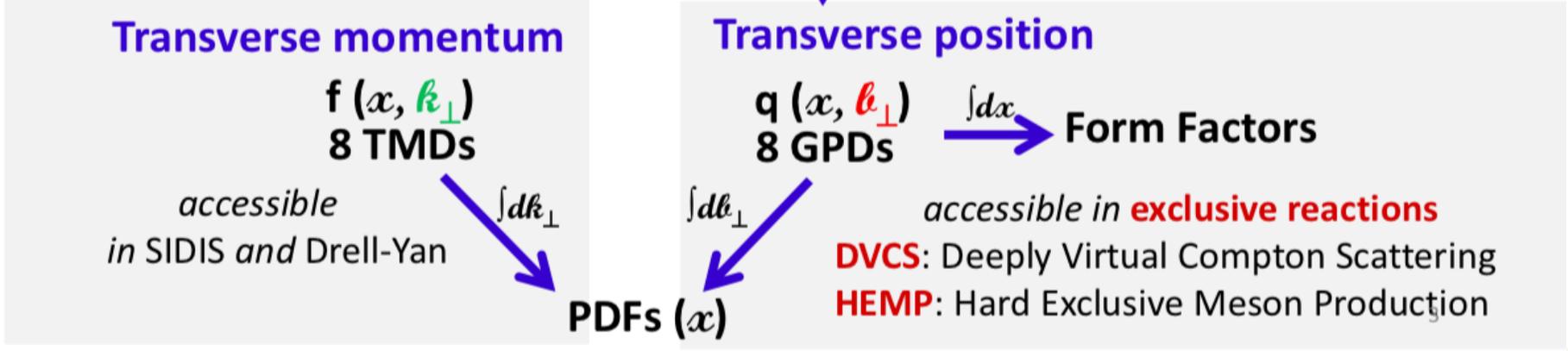
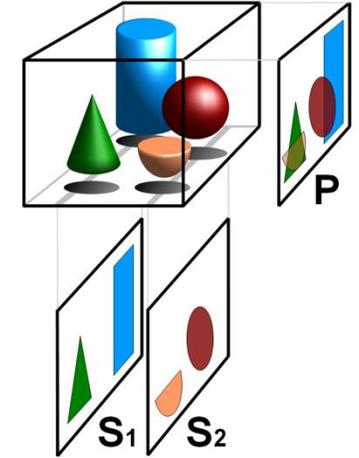
Toward a more complete description of the nucleon

Slide from N. d'Hose, APS-April meeting 2019.

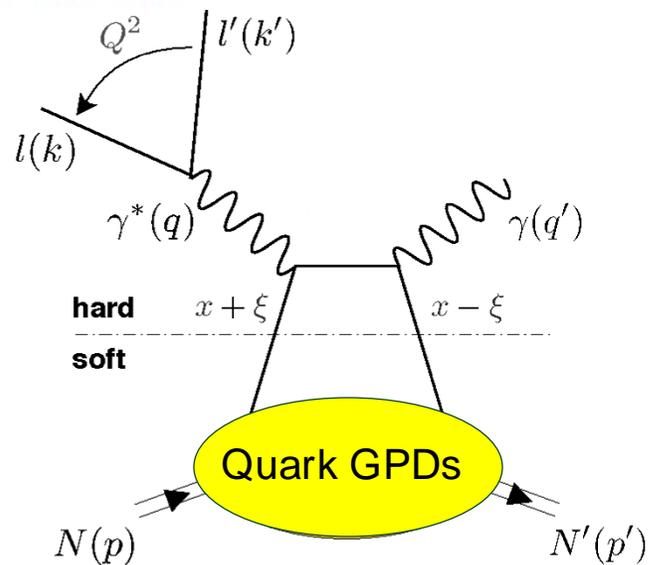


Quantum tomography of the nucleon

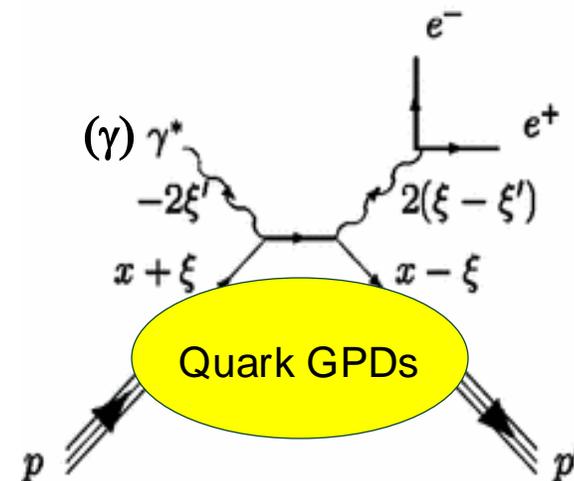
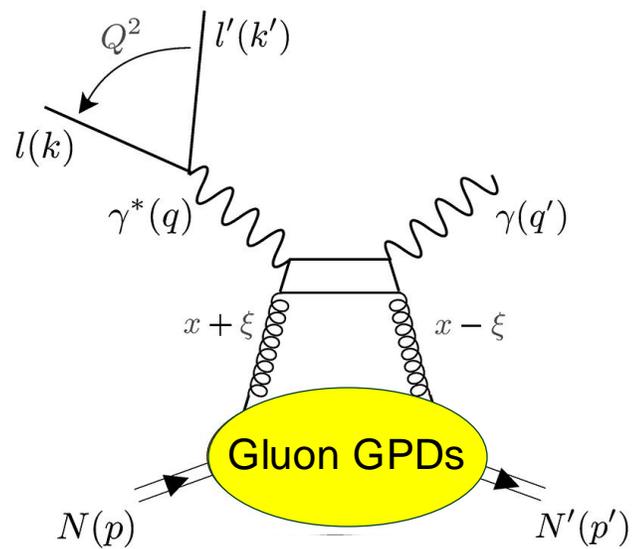
Ji, PRL91 (2003)
 Belitsky, Ji, Yuan, PRD69 (2004)
 Lorcé et al, JHEP1105 (2011)



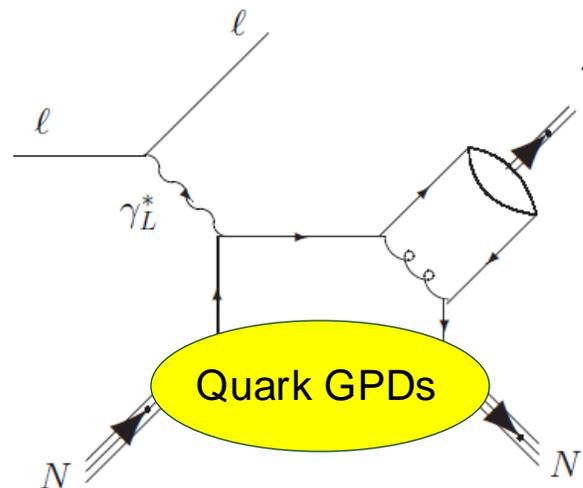
Universality of GPDs



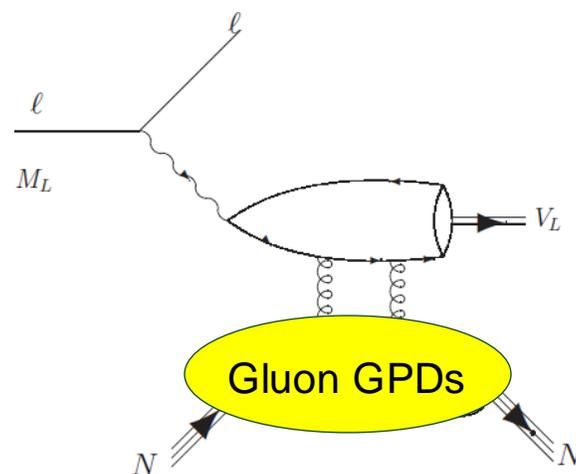
DVCS

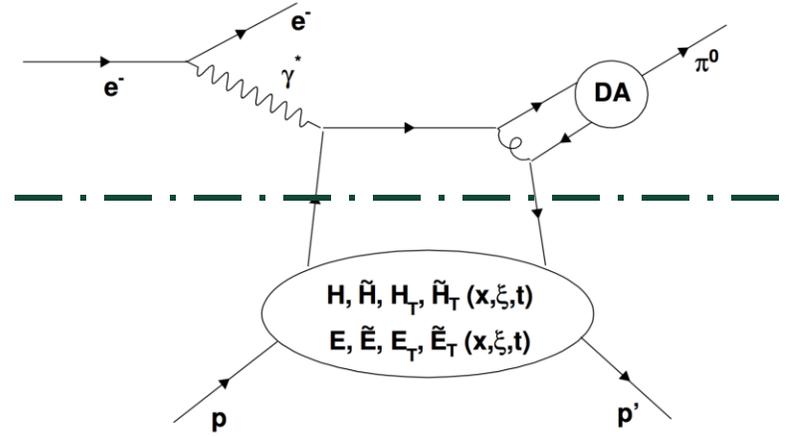
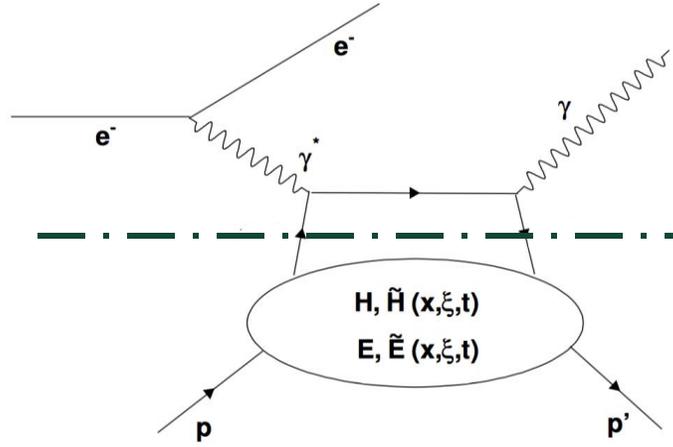


(TCS), DDVCS

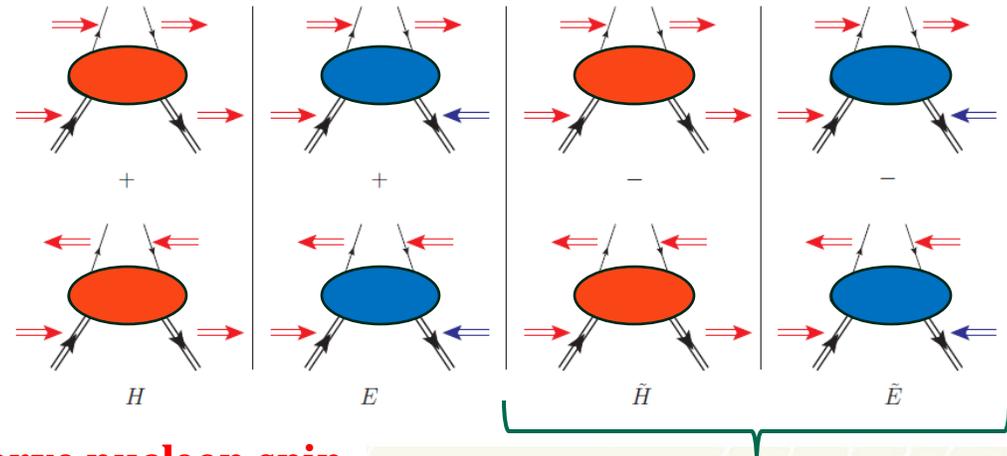


DVMP





Chiral even GPDs:
(helicity of the parton is conserved)



Chiral even GPDs

+

Chiral-odd GPDs:

(helicity of the parton can flip)

conserve nucleon spin
flip nucleon spin (24)

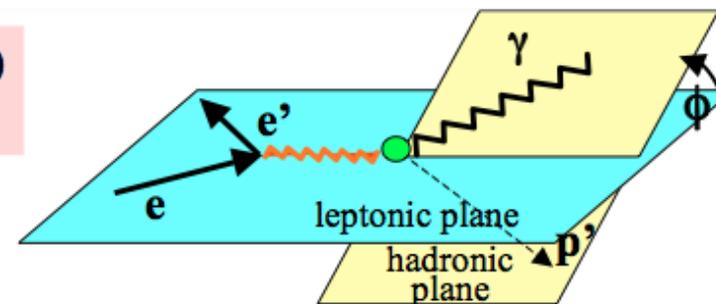
parton helicity sensitive

Sensitivity to different GPDs (at leading order and twist)

$$\Delta\sigma = d^5\vec{\sigma} - d^5\overleftarrow{\sigma}$$

$$\xi = x_B/(2-x_B)$$

$$k = -t/4M^2$$



Polarized **beam**, unpolarized **proton** target:

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im} \{ F_1 H + \xi(F_1 + F_2) \tilde{H} + kF_2 E \} d\phi$$

Kinematically suppressed

$$\rightarrow H_p, \tilde{H}_p, E_p$$

Unpolarized beam, **longitudinal proton** target:

$$\Delta\sigma_{UL} \sim \sin\phi \operatorname{Im} \{ F_1 \tilde{H} + \xi(F_1 + F_2)(H + \dots) \} d\phi$$

$$\rightarrow H_p, \tilde{H}_p$$

Unpolarized beam, **transverse proton** target:

$$\Delta\sigma_{UT} \sim \sin\phi \operatorname{Im} \{ k(F_2 H - F_1 E) + \dots \} d\phi$$

$$\rightarrow H_p, E_p$$

Polarized **beam**, unpolarized **neutron** target:

$$\Delta\sigma_{LU} \sim \sin\phi \operatorname{Im} \{ F_1 H + \xi(F_1 + F_2) \tilde{H} - kF_2 E \} d\phi$$

$$\rightarrow H_n, \tilde{H}_n, E_n$$

Suppressed because $F_1(t)$ is small

Suppressed because of **cancellation** between PPD's of **u** and **d** quarks

$$H_p(x, \xi, t) = 4/9 H_u(x, \xi, t) + 1/9 H_d(x, \xi, t)$$

$$H_n(x, \xi, t) = 1/9 H_u(x, \xi, t) + 4/9 H_d(x, \xi, t)$$

The DVCS program worldwide

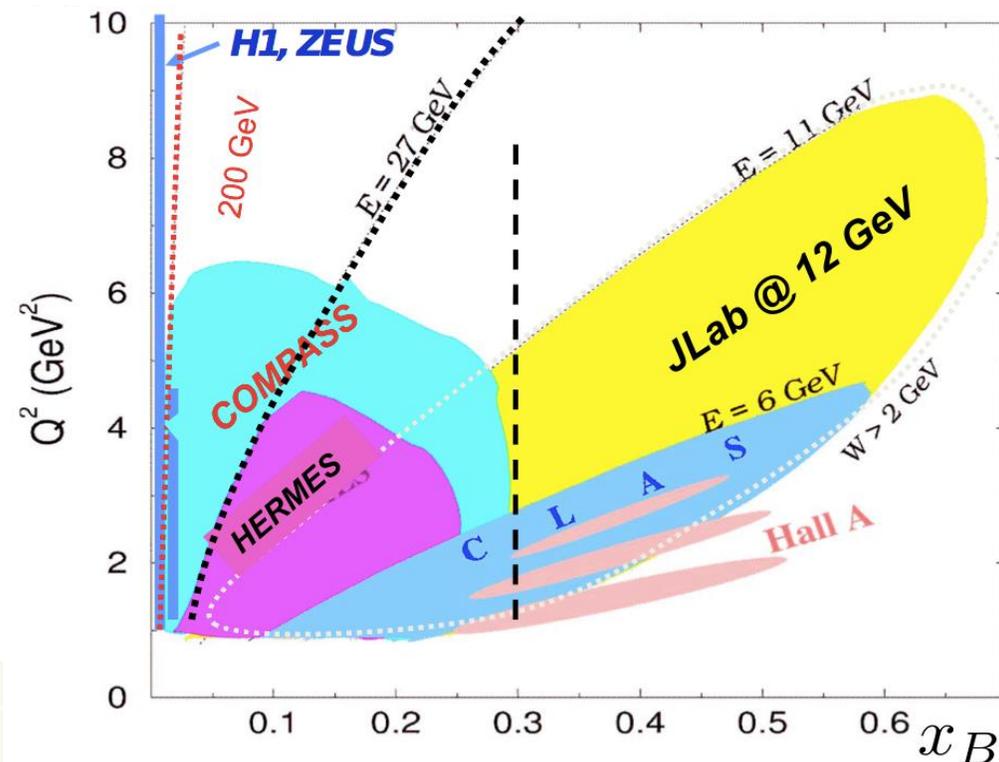
Experimental timeline

- Pioneering results from non-dedicated experiments (Hall B and Hermes): ~ 2001
- First round of dedicated experiments (Hall A/B, Hermes, H1&ZEUS): ~ 2005
- Second round of dedicated experiments (Halls A/B): ~ 2010
- Compelling DVCS program at JLab-12 GeV and Compass: 2015 and later
- EIC program...

In the valence region (JLab 6 and JLab 12)

Partially complimentary, overlapping

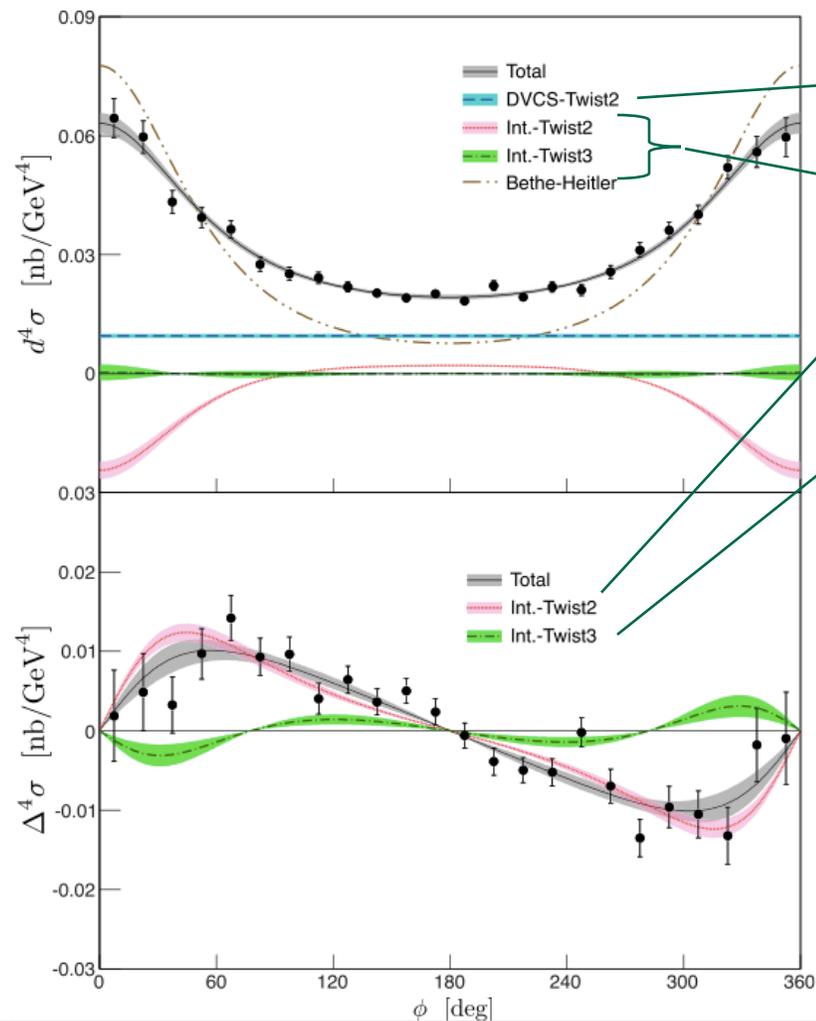
- Hall A/C: Test the validity of the formalism
 - high accuracy ($\sim 5\%$)
 - limited kinematic
- Hall B: Map the GPDS
 - limited accuracy (15+%)
 - wide kinematic range



DVCS Hall A@Jlab 1rst generation

PRC C92, Nov '15

$$x_B = 0.37, \quad Q^2 = 2.36 \text{ GeV}^2, \quad -t = 0.32 \text{ GeV}^2$$



$$d\sigma^{BH} \propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi$$

$$d\sigma_{unpol}^{DVCS} \propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi$$

$$d\sigma_{pol}^{DVCS} \propto s_1^{DVCS} \sin \phi$$

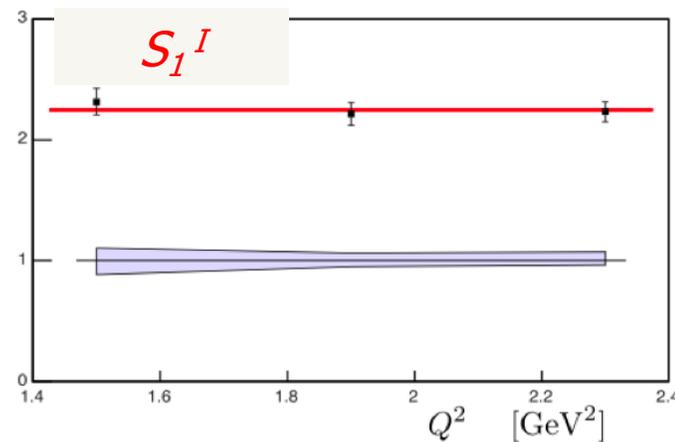
$$\text{Re } I \propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi$$

$$\text{Im } I \propto s_1^I \sin \phi + s_2^I \sin 2\phi$$

LT/LO

HT

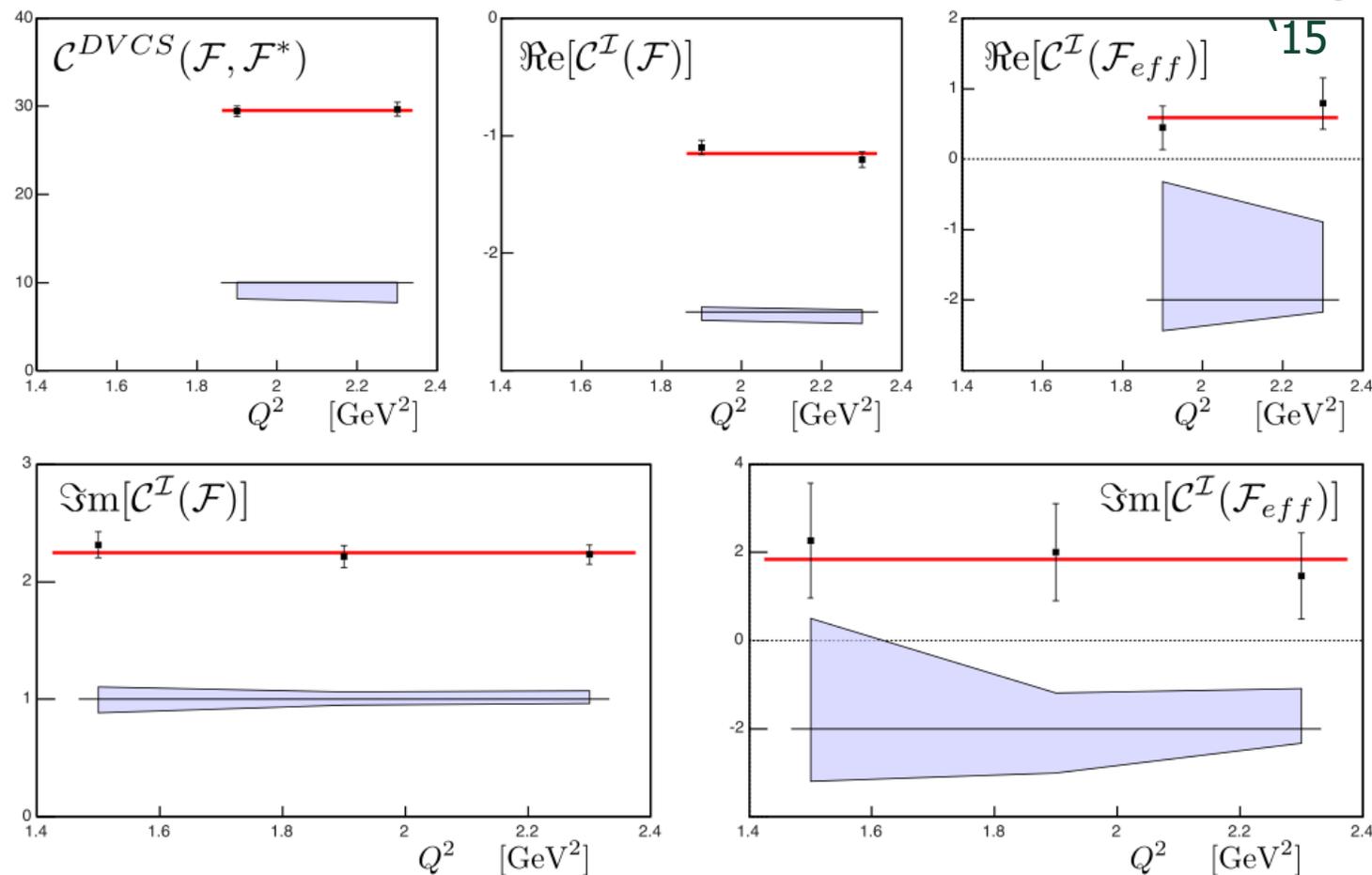
NLO



No Q^2 dependence within this limited range => leading twist dominance
Needs to be checked over a larger Q^2 bite

Hall A E00-110: cross section Q^2 dependence

PRC C92, Nov

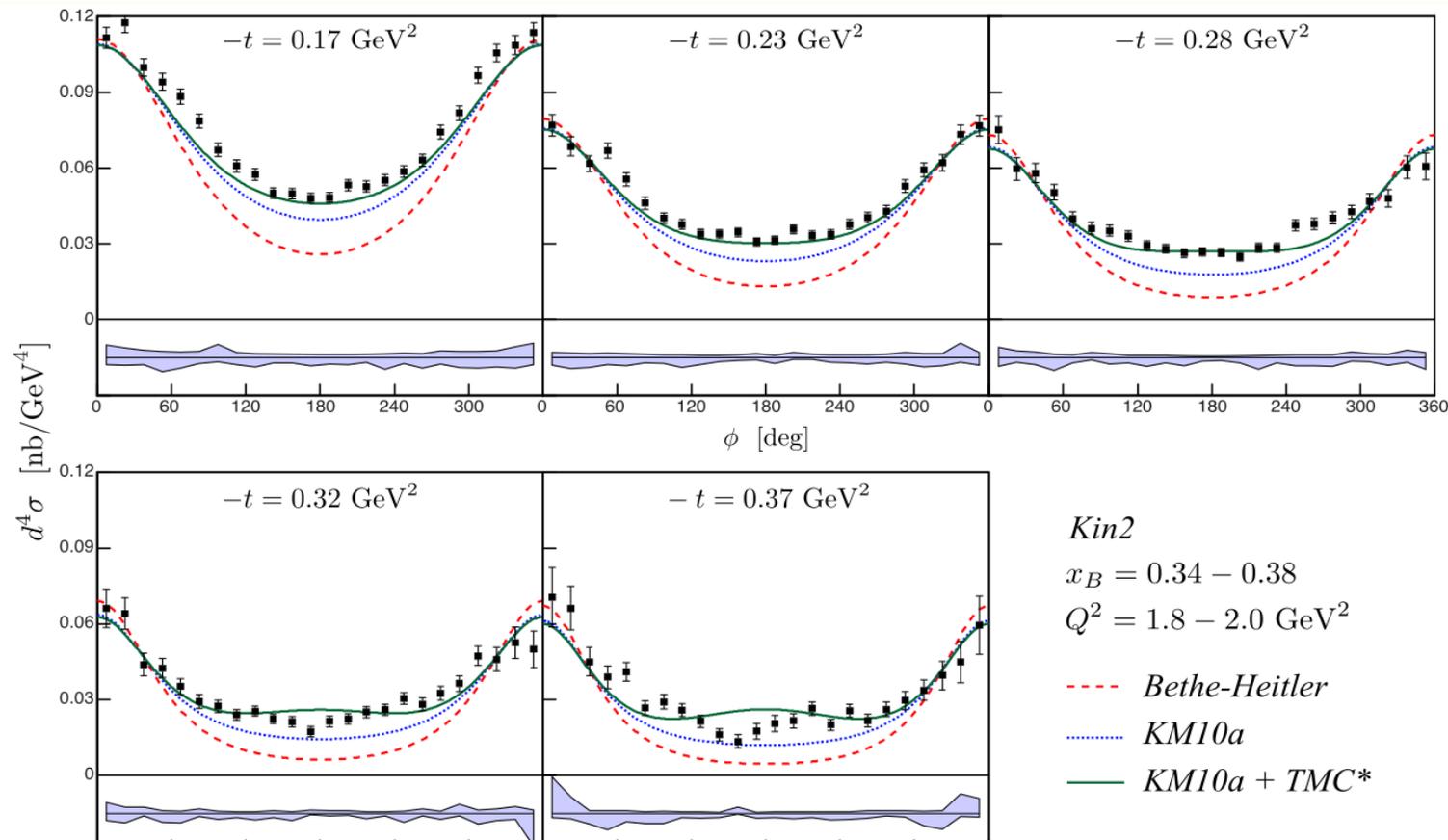


No Q^2 dependence within this limited range => leading twist dominance

Need to be checked over a larger Q^2 bite

Hall A E00-110 cross sections: higher twist corrections

PRC C92, Nov '15



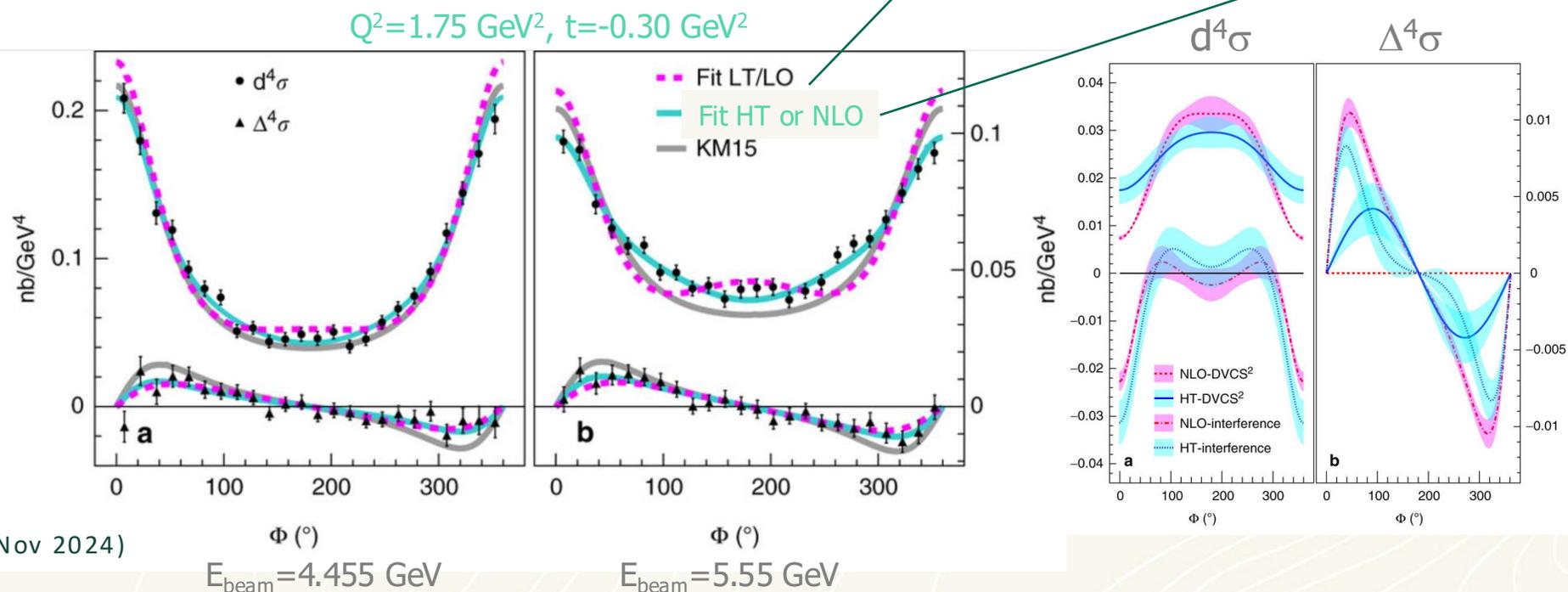
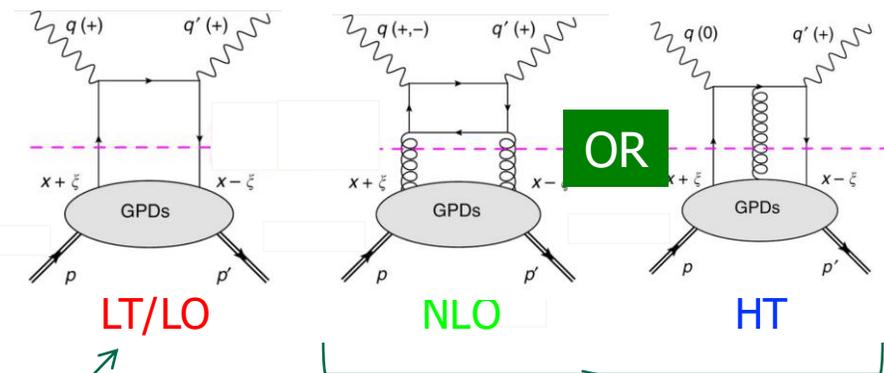
Higher twist corrections might be necessary to fully explain experimental data
 Confirmation of the significant deviation from BH => Need to measure T^2_{DVCS}

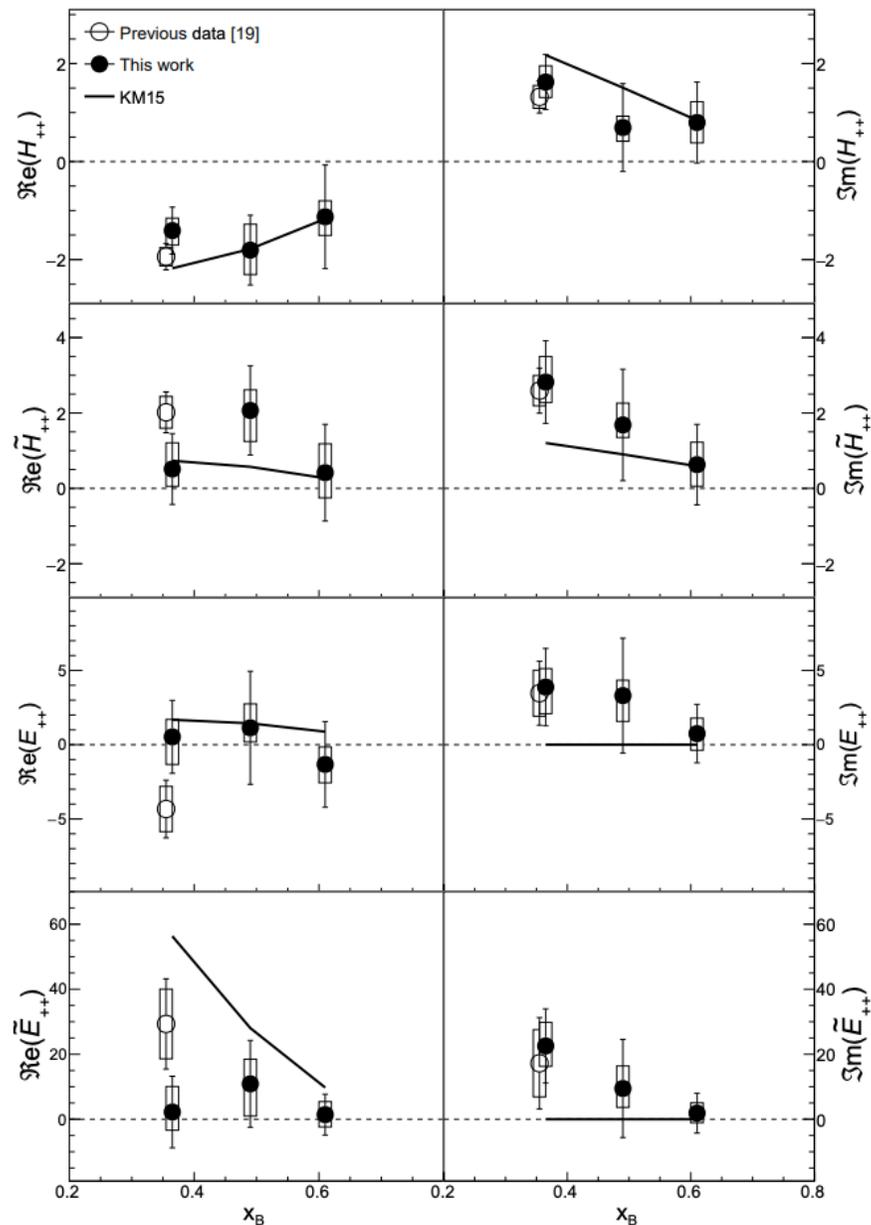
Hall A E07-007 (JLab 6 GeV)

Goal: To separate the BH.DVCS interference contribution from the DVCS² contribution.

Nature Commun. 8 (2017) no.1, 1408

$$\begin{aligned} d\sigma^{BH} &\propto c_0^{BH} + c_1^{BH} \cos \phi + c_2^{BH} \cos 2\phi \\ d\sigma_{unpol}^{DVCS} &\propto c_0^{DVCS} + c_1^{DVCS} \cos \phi + c_2^{DVCS} \cos 2\phi \\ d\sigma_{pol}^{DVCS} &\propto s_1^{DVCS} \sin \phi \\ \text{Re } I &\propto c_0^I + c_1^I \cos \phi + c_2^I \cos 2\phi + c_3^I \cos 3\phi \\ \text{Im } I &\propto s_1^I \sin \phi + s_2^I \sin 2\phi \end{aligned}$$





DVCS in Hall A: 12 GeV results

[This work] **F. Georges,**
***Phys.Rev.Lett.* 128 (2022) 25, 252002**

Error bars: statistical

Error boxes: systematic

$\text{CFF}_{\lambda\lambda'}$ λ : polarization state of virtual photon (0,+,-)

λ' : polarization state of outgoing real photon (+,-)

Fit has 24 CFF

$$(\tilde{H}, H, \tilde{E}, E) \otimes (\Re, \Im) \otimes (+, +, 0, +, -)$$

but **only the results from the LO ones (++)** are shown.

Fits performed at constant x_B and t over Q^2 and ϕ bins.

No Q^2 evolution of the CFFs.

[19] M. Defurne et al., *Phys. Rev. C* 92, 055202 (2015)

[KM15] K. Kumericki and D. Mueller, *EPJ Web Conf.* 112 (2016) 01012

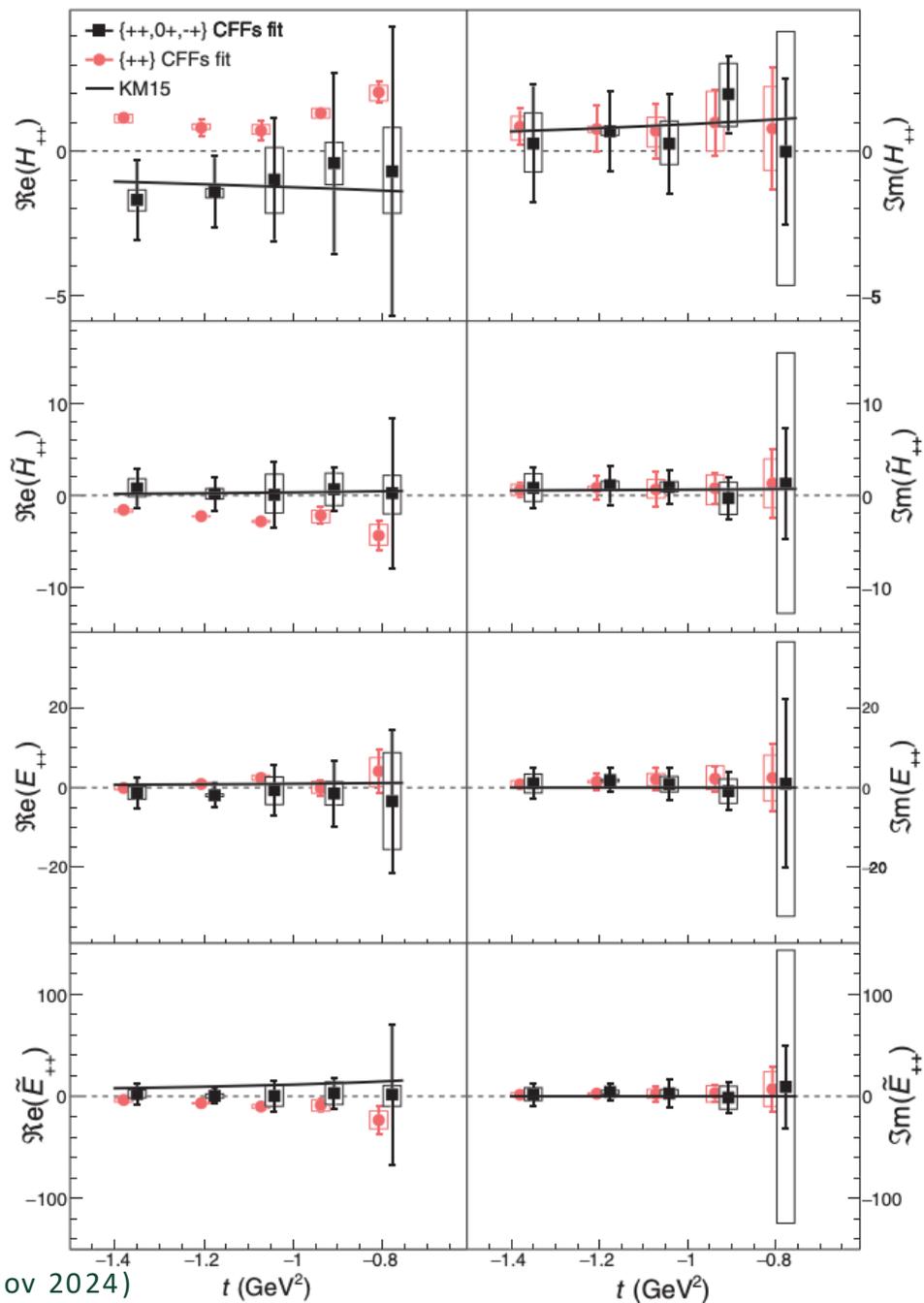
The precise measurement of cross-sections at multiple values of center-of-mass energy (\sqrt{s}) but same Q^2 and it is essential for this extraction.

Also demonstrated in

M. Defurne et al., *Nat. Commun.* 8, 1408 (2017).

B. Kriesten et al., *Phys. Rev. D* 101, 054021 (2020).

M. Čuić et al., *Phys. Rev. Lett.* 125, 232005 (2020).



DVCS in Hall A: 12 GeV results

F. Georges, *Phys.Rev.Lett.* 128 (2022) 25, 252002

Results shown at $x_B=0.60$

Fits performed at constant x_B and t over Q^2 and ϕ bins.

No Q^2 evolution of the CFFs.

■ Fit has 24 CFF

$$(\tilde{H}, H, \tilde{E}, E) \otimes (\Re, \Im) \otimes (+, +, 0+, +, -)$$

but **only the results from the LO ones (++) are shown.**

● Fit has only 8 CFF

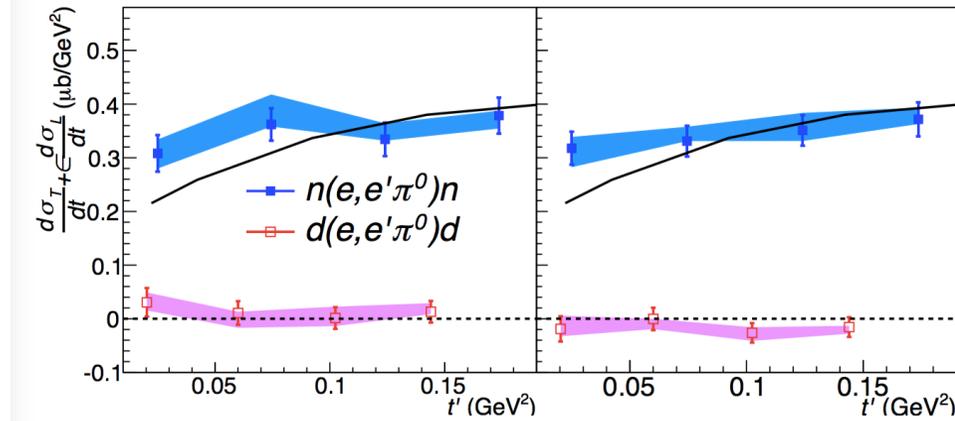
$$(\tilde{H}, H, \tilde{E}, E) \otimes (\Re, \Im) \otimes (+, +)$$

**Fitting all CFFs is essential
to get realistic estimates
of their uncertainties.**

π^0 DVCS2n results: fully separated contributions

M. Mazouz PRL 118 (2017) 22, 222002

$Q^2=1.75 \text{ GeV}^2$ and $x_B=0.36$

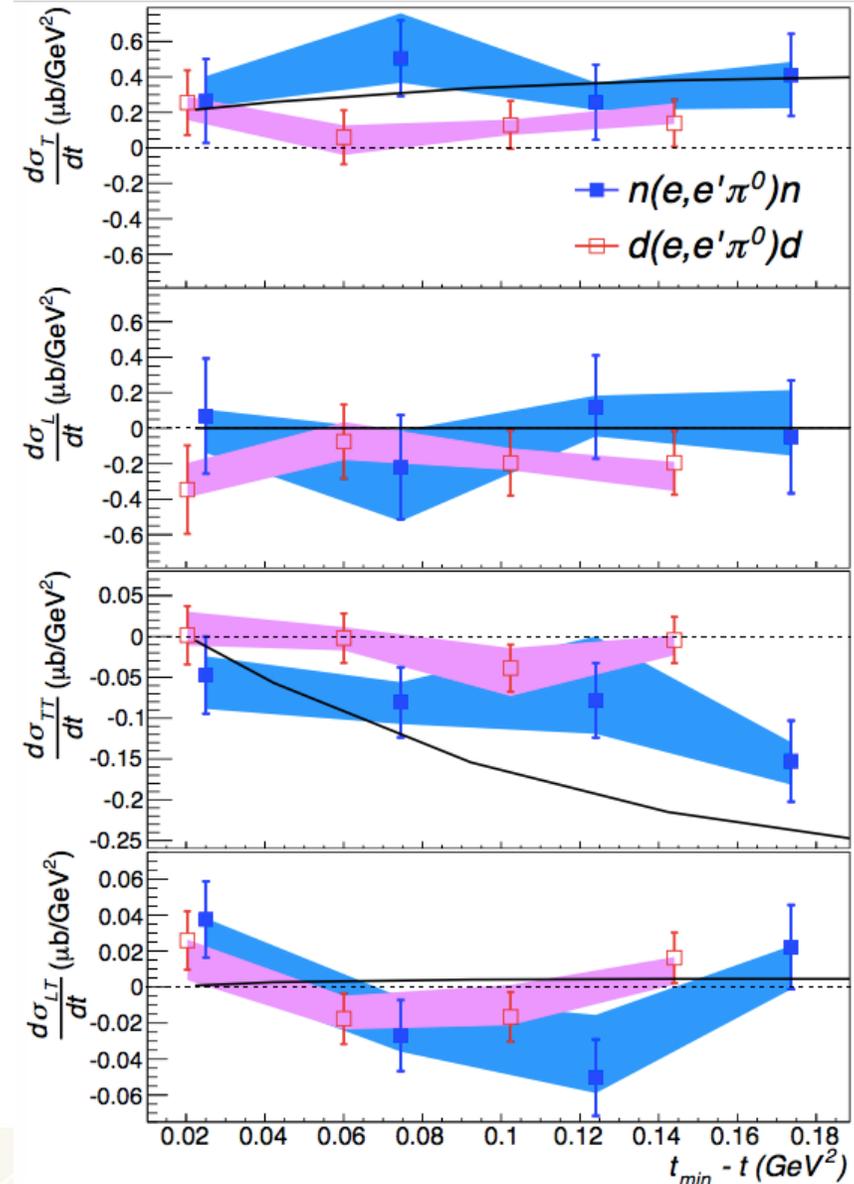


— Goloskokov and Kroll
Eur Phys J A47 (2012)

$$\frac{d\sigma_T}{dt} = \Lambda \left[(1 - \xi^2) |\langle H_T \rangle|^2 - \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 \right]$$

$$\frac{d\sigma_{TT}}{dt} = \Lambda \frac{t'}{8M^2} |\langle \bar{E}_T \rangle|^2 .$$

$$\bar{E}_T = 2\tilde{H}_T + E_T$$



π^0 DVCS2n results: flavor separation

M. Mazouz PRL 118 (2017) 22, 222002

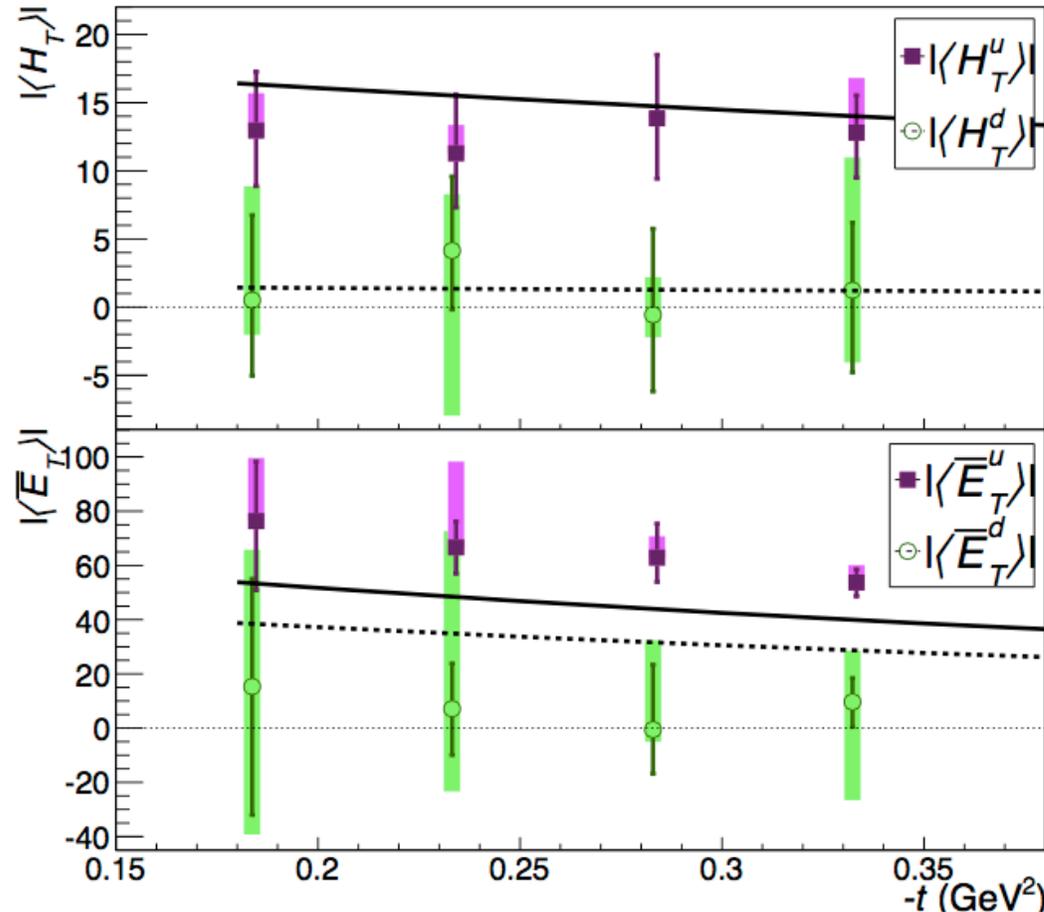
$$|\langle H_T^{p,n} \rangle|^2 = \frac{1}{2} \left| \frac{2}{3} \langle H_T^{u,d} \rangle + \frac{1}{3} \langle H_T^{d,u} \rangle \right|^2$$

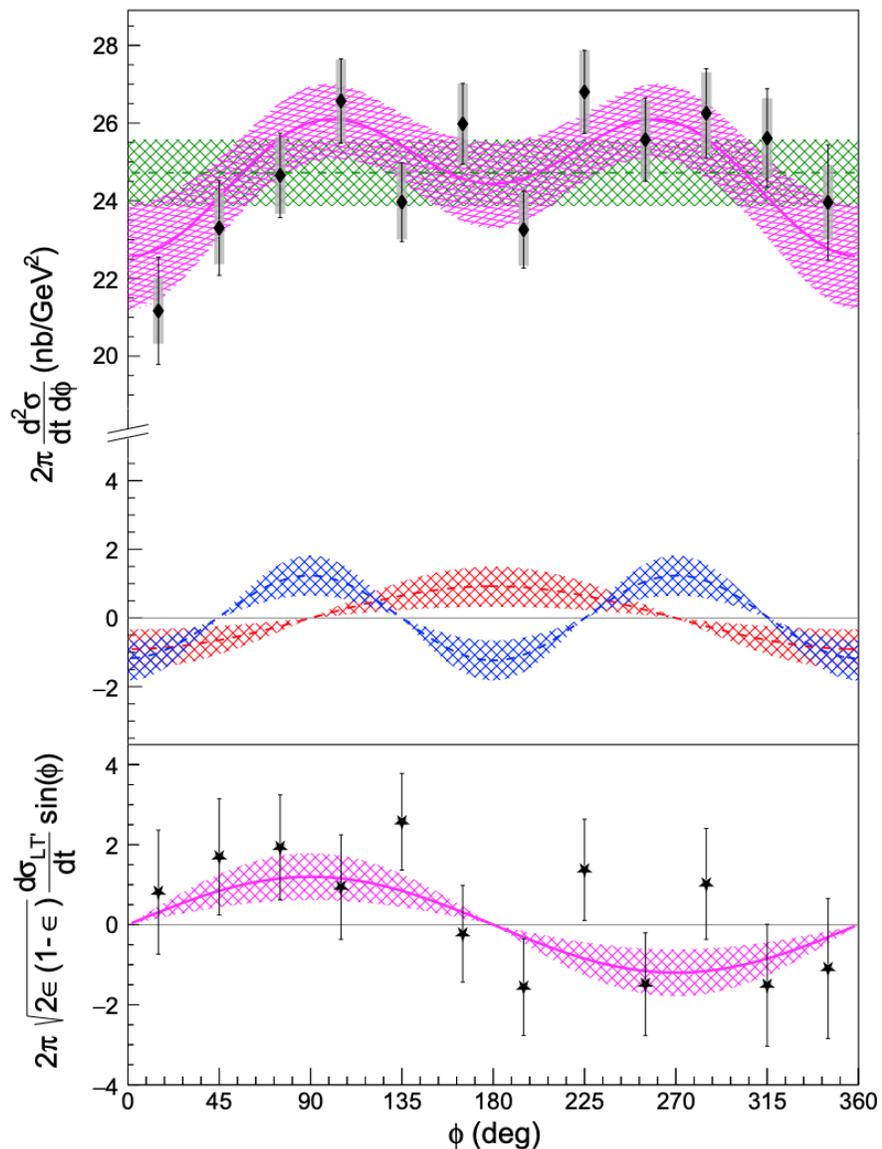
$Q^2 = 1.75 \text{ GeV}^2, x_B = 0.36$

account for the unknown phase variation between u and the d amplitude $\gamma^*q \rightarrow q'\pi^0$ convoluted with $(H,E)_T$

Goloskokov and Kroll
Eur Phys J A47 (2012)

— u quark
- - - d quark



$Q^2=8.31 \text{ GeV}^2, t'=0.15 \text{ GeV}^2, x_B=0.60$


$ep \rightarrow ep\pi^0$ in Hall A : 12 GeV result

M. Dlamini *et al*, Phys. Rev. Lett **127**, 152301

At sufficient high Q^2 , meson production should be understandable in terms of the “handbag” diagram.

But the factorization is only exact for longitudinal virtual photons (Collins, Frankfurt, Strikman, 1997). Effective factorization of the transverse part exists.

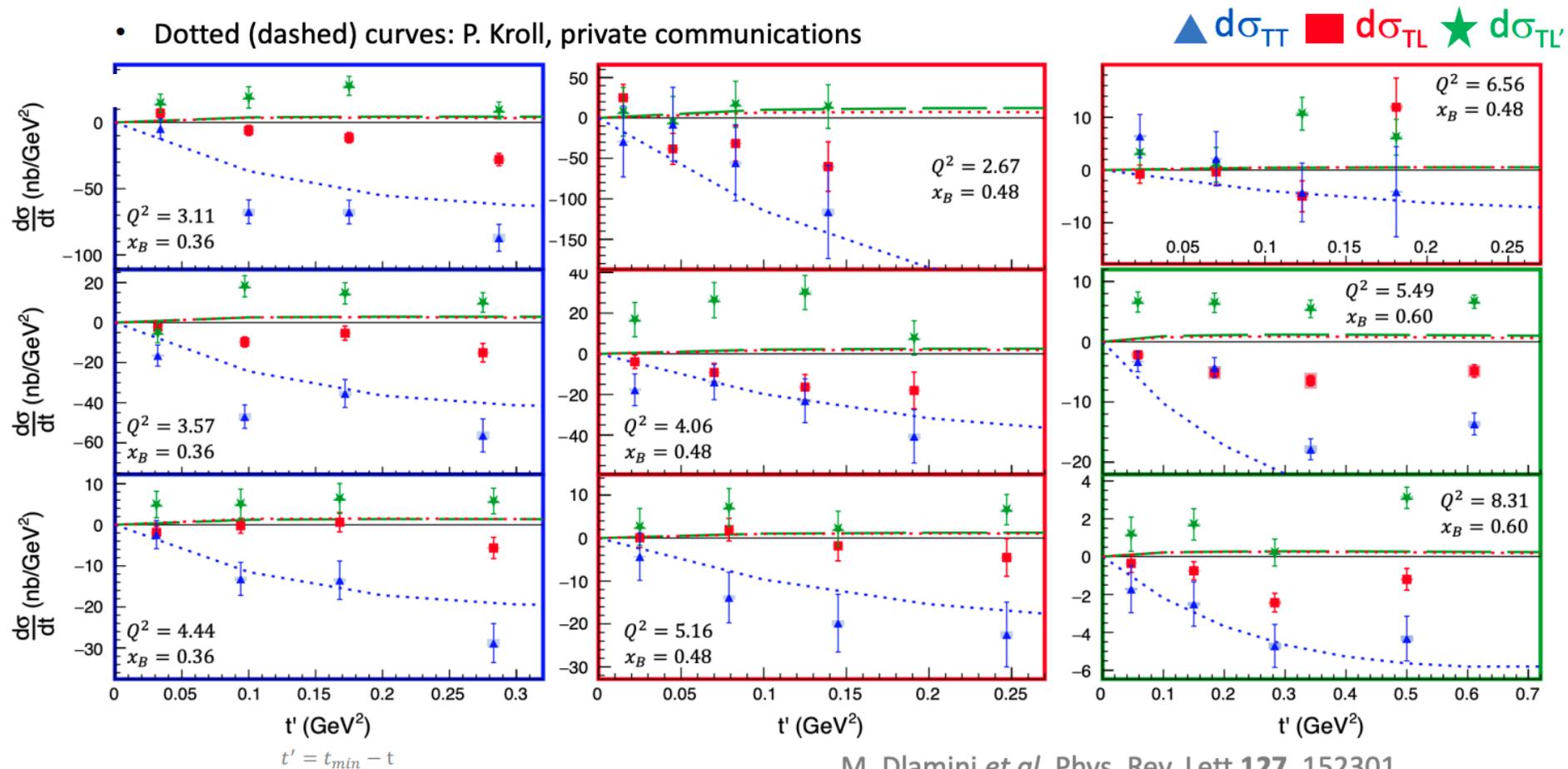
Asymptotic QCD predicts that

$$\sigma_L \rightarrow Q^{-6}, \quad \sigma_T \rightarrow Q^{-8} \quad \text{and} \quad \frac{\sigma_T}{\sigma_L} \rightarrow \frac{1}{Q^2}$$

$$\frac{d^4\sigma}{dQ^2 dx_B dt d\phi} = \frac{1}{2\pi} \frac{d^2 \Gamma_\gamma}{dQ^2 dx_B}(Q^2, x_B, E)$$

$$\left[\frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt} + \sqrt{2\epsilon(1+\epsilon)} \frac{d\sigma_{LT}}{dt} \cos(\phi) + \epsilon \frac{d\sigma_{TT}}{dt} \cos(2\phi) + h \sqrt{2\epsilon(1-\epsilon)} \frac{d\sigma_{LT'}}{dt} \sin(\phi) \right]$$

$d\sigma_U = d\sigma_T + \epsilon d\sigma_L$

M. Dlamini *et al*, Phys. Rev. Lett **127**, 152301

17

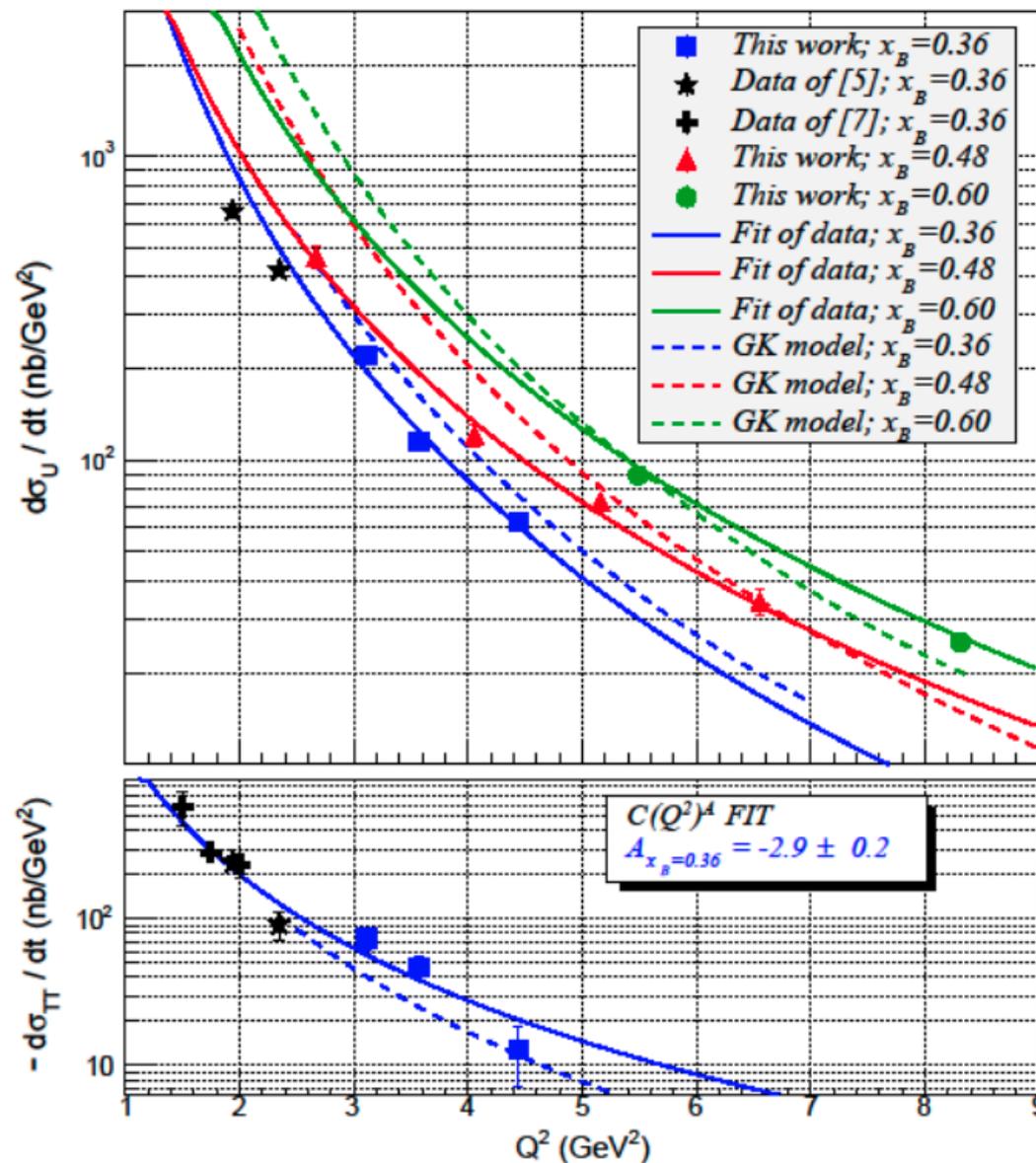
- σ_{TT} larger than σ_{TL} and $\sigma_{TL'}$: hint of **dominance of the transverse amplitude** as suggested by the GK model
- σ_{TL} and $\sigma_{TL'}$ underestimated by the GK model: larger contribution of the longitudinal amplitude than the one expected by GK
- **sign difference between GK and σ_{TL}** (Hall B & COMPASS results agree with GK)

$$\langle t' \rangle = 0.1 \text{ GeV}^2$$

$ep \rightarrow ep\pi^0$ in Hall A : 12 GeV result

M. Dlamini *et al*, Phys. Rev. Lett **127**, 152301

$$\frac{d\sigma_u}{dt} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt}$$



dashed lines: P. Kroll private communication

■ This work, $x_B = 0.36$

▲ This work, $x_B = 0.48$

● This work, $x_B = 0.60$

★ E. Fuchey *et al*, Phys. Rev. C 83, 025201 (2011)

⊕ M. Defurne *et al*, Phys. Rev. Lett. 117, 262001 (2016)

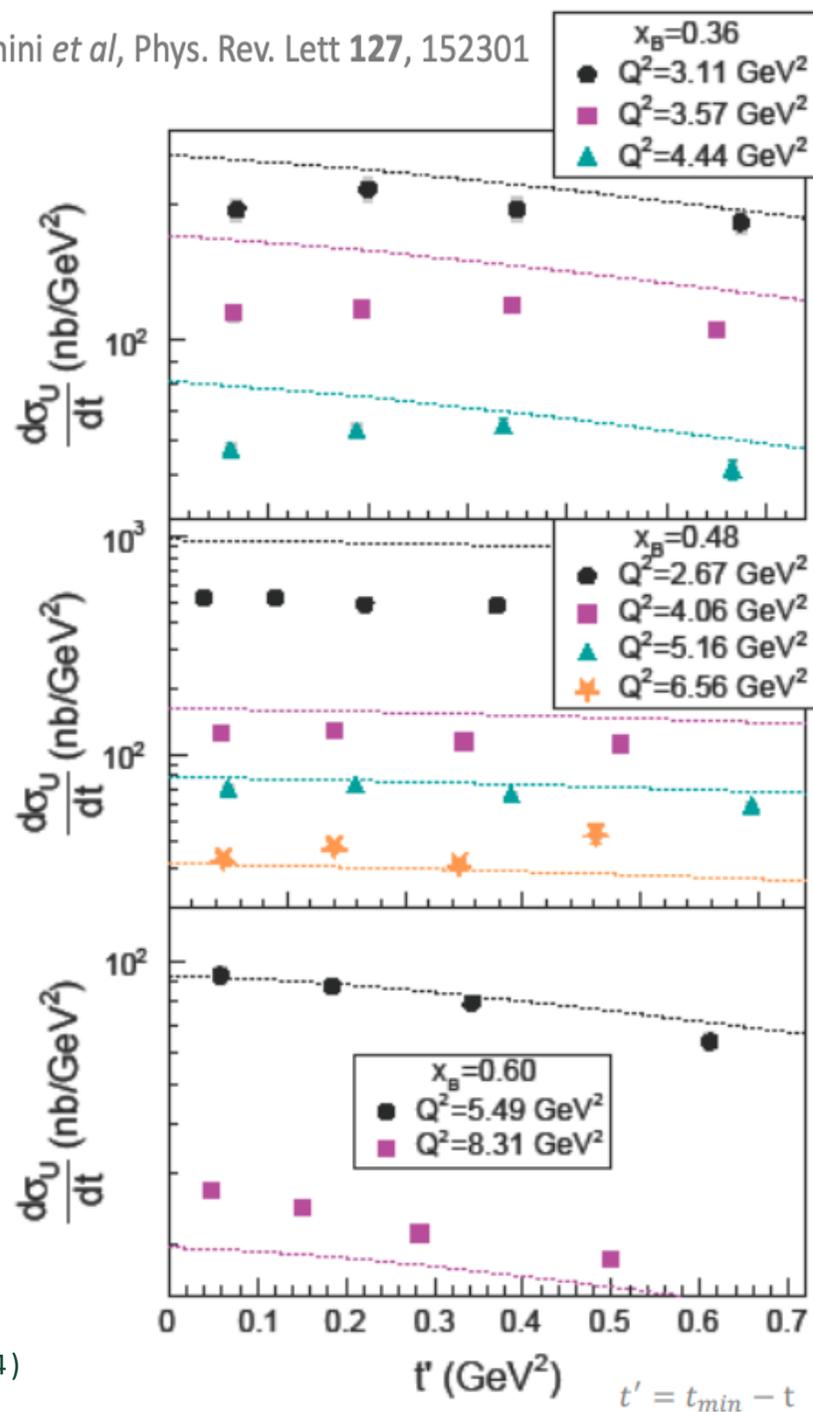
➤ $C(Q^2)^A \exp(-Bt')$ fit to experimental results of $d\sigma_U$ in different $x_B \rightarrow$ solid curves

$$x_B = 0.36 \rightarrow A = -3.3 \pm 0.1$$

$$x_B = 0.48 \rightarrow A = -2.9 \pm 0.1$$

$$x_B = 0.60 \rightarrow A = -3.1 \pm 0.1$$

➤ Q^2 dependence closer to Q^{-6} , rather than Q^{-8} as expected for σ_T at high Q^2



$$\frac{d\sigma_u}{dt} = \frac{d\sigma_T}{dt} + \epsilon \frac{d\sigma_L}{dt}$$

dashed lines: P. Kroll private communication

E12-13-010: precision DVCS/ π^0 cross sections



Simplest process: $e + p \rightarrow e' + p + \gamma$ (DVCS)

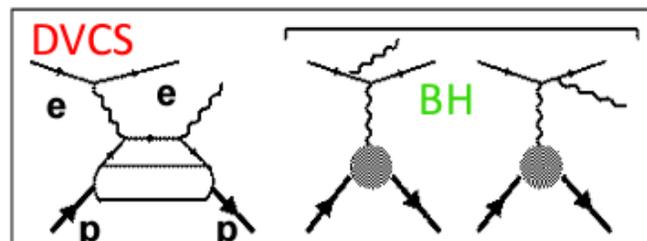
E12-13-010 DVCS measurements follow up on measurements in Hall A:

- Scaling of the Compton Form Factor
- Rosenbluth-like separation of DVCS:

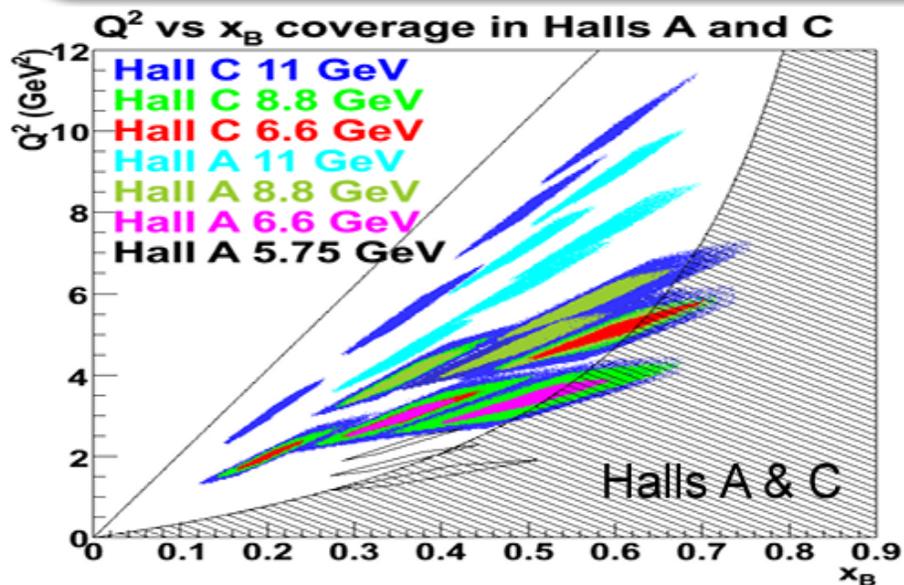
$$\sigma = |BH|^2 + \text{Re}[DVCS^\perp BH] + |DVCS|^2$$

$\sim E_{beam}^2$ $\sim E_{beam}^3$

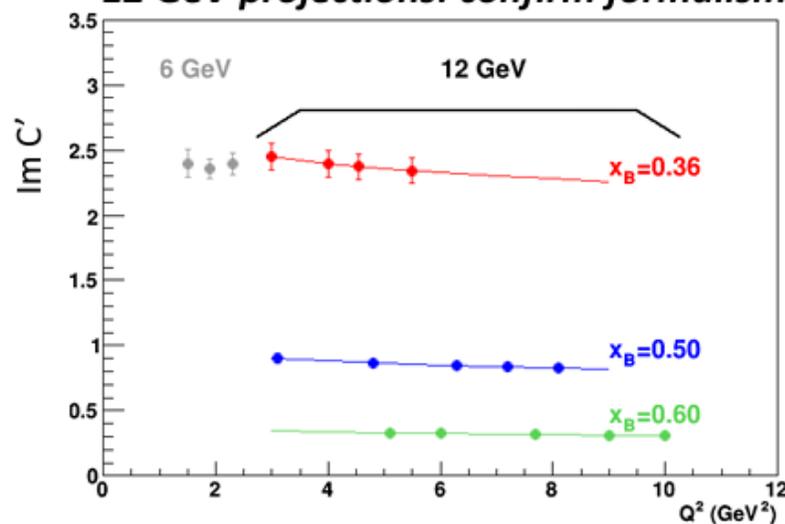
- L/T separation of π^0 production



Hall A data for Compton form factor (over *limited* Q^2 range) agree with hard-scattering



12 GeV projections: confirm formalism

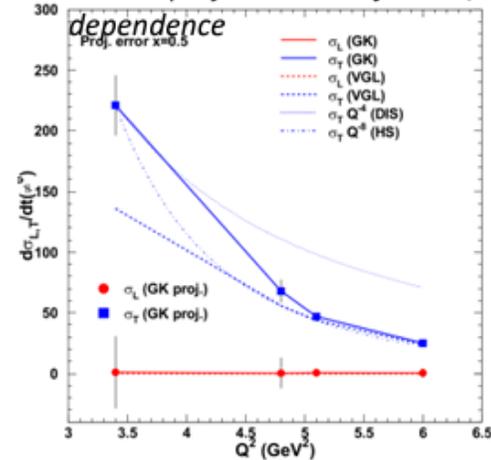


Extracting the real part of CFFs from DVCS requires measuring the cross section at multiple beam energies (DVCS²-Interference separation)

π^0 Exclusive Cross Sections

- Relative L/T contribution to π^0 cross section important in probing transversity
- Results from Hall A at 6 GeV Jlab suggest that the longitudinal cross section in π^0 production is non-zero up to $Q^2 = 2 \text{ GeV}^2$

12 GeV projections: confirm Q^2/t



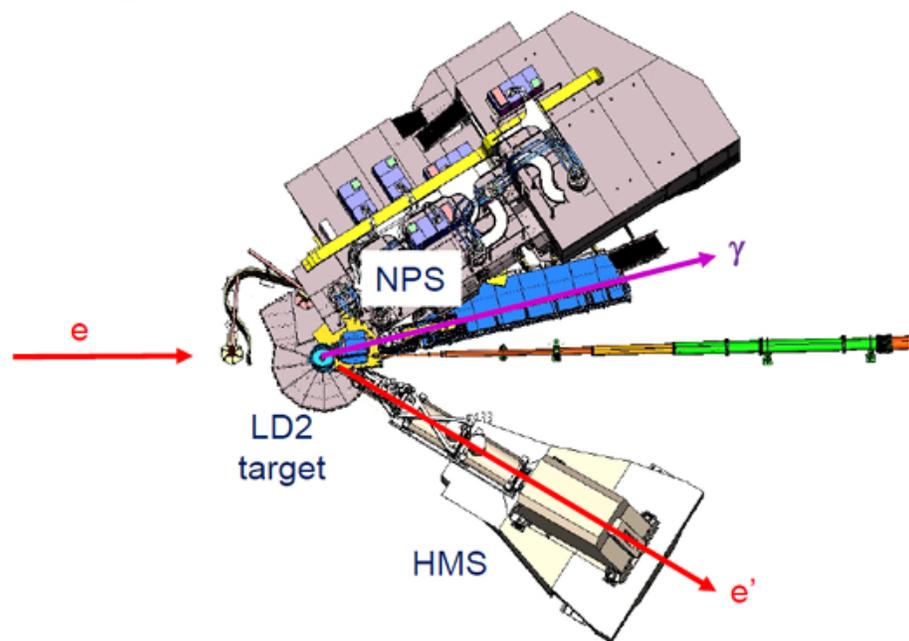
E12-13-010 provides also data on σ_T and σ_L at higher Q^2 for reliable interpretation of 12 GeV GPD data

E12-22-006: DVCS off the Neutron



Probe **flavor dependence of GPDs** with precision nDVCS cross sections

Measurement of the $N \rightarrow e' \gamma X$ reaction ($N=p, n, d$) using an LD₂ target in Hall C



With NPS and HMS in Hall C reach $\sim x2$ -
12 better nDVCS & dDVCS separation
than previous 6 GeV experiment

Projected Impact on flavor dependence of CFFs

- Simultaneous fit of E12-13-010 (p) and E12-22-006 (n)
- Real and imaginary parts of CFFs H and \tilde{H} and E (u & d) as free parameters (nDVCS not sensitive to \tilde{E})

