

QCD relations on gravitational form factors and hadron mass decompositions

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KT, JHEP03, 013 ('23)

KT, in preparation

Contents

1. Gravitational form factors
energy-momentum tensor, GPD
2. $\bar{C}_{q,g}$ at NNLO QCD
trace anomaly constraints
3. Mass decompositions at NNLO QCD
N, π

Symmetric energy-momentum tensor

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2$$

$$T^{\mu\nu} = T^{\nu\mu}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

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$$\left(= \frac{1}{2} \bar{\psi} \gamma^{\mu} i \partial^{\nu} \psi - F^{\mu\rho} \partial^{\nu} A_{\rho} + \frac{\eta^{\mu\nu}}{4} F^2 + \partial_{\lambda} X^{[\lambda\mu]\nu} \right)$$

$$\sum_n \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \phi_n)} \partial^{\nu} \phi_n - \eta^{\mu\nu} \mathcal{L}$$

$$T^{\mu\nu} = T^{\nu\mu}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2$$

translation

$$T^{\mu\nu}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

Lorentz tr.

$$M^{\mu\nu\lambda} = x^{\nu} T^{\mu\lambda} - x^{\lambda} T^{\mu\nu}$$

$$\partial_{\mu} M^{\mu\nu\lambda} = 0$$

scale tr.

$$D^{\mu} = x_{\nu} T^{\mu\nu}$$

$$\partial_{\mu} D^{\mu} = T_{\mu}^{\mu}$$

conformal tr.

$$C^{\mu\nu} = \left(2x^{\rho} x^{\nu} - \eta^{\rho\nu} x^2 \right) T_{\rho}{}^{\mu}$$

$$\partial_{\mu} C^{\mu\nu} = 2x^{\nu} T_{\mu}^{\mu}$$

Symmetric energy-momentum tensor

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2$$

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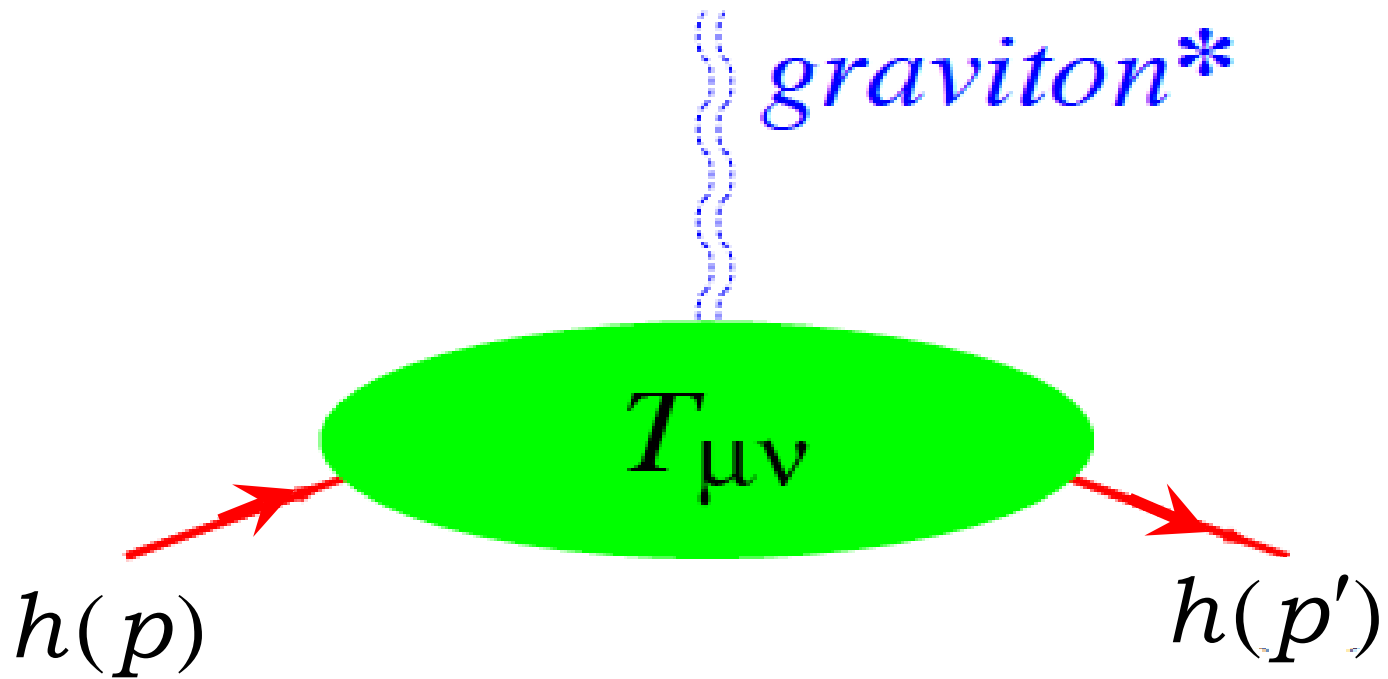
$$\underbrace{\hspace{10em}}_{\sum_n \frac{\partial \mathcal{L}}{\partial (\partial_{\mu} \phi_n)} \partial^{\nu} \phi_n - \eta^{\mu\nu} \mathcal{L}}$$

$$T_{\mu\nu}(x) = \frac{2}{\sqrt{-g}} \frac{\delta S}{\delta g^{\mu\nu}(x)} \Big|_{g^{\mu\nu} \rightarrow \eta^{\mu\nu}}$$

$$T^{\mu\nu} = T^{\nu\mu}$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$\begin{aligned}
T^{\mu\nu} &= \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \\
&\equiv T_q^{\mu\nu} + T_g^{\mu\nu}
\end{aligned}$$



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$$\langle N(p') | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p') \left[A_{q,g}(t) \gamma^{(\mu} P^{\nu)} + B_{q,g}(t) \frac{P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}}{2M} + D_{q,g}(t) \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{q,g}(t) M \eta^{\mu\nu} \right] u(p)$$

$$P = \frac{p + p'}{2}$$

$$\Delta = p' - p$$

$$t = \Delta^2$$

Ji, PRL78, 610 ('97)

Polyakov, Schweitzer, IJMPA33, 1830025('18)

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

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$P = \frac{p + p'}{2}$
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$$A_q(0) + A_g(0) = 1$$

$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^{\mu} p^{\nu}$$

$$\frac{1}{2} (A_q(0) + B_q(0) + A_g(0) + B_g(0)) = \frac{1}{2}$$

$$\frac{\langle N(p) S | J^i | N(p) S \rangle}{\langle N(p) S | N(p) S \rangle} = \frac{1}{2} S^i$$

$$B_q(0) + B_g(0) = 0$$

$$J^i = \frac{1}{2} \epsilon^{ijk} \int d^3x M^{+jk}$$

$$M^{\mu\rho\sigma} = x^{\rho} T^{\mu\sigma} - x^{\sigma} T^{\mu\rho}$$

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$$M^{\mu\rho\sigma} = x^{\rho} T^{\mu\sigma} - x^{\sigma} T^{\mu\rho}$$

$$\bar{C}_q(t) + \bar{C}_g(t) = 0$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$D \equiv D_q(0) + D_g(0)$$

“D term”: the last unknown global property

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

mass & energy distribution

angular momentum distribution

$$\langle N(p') | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p') \left[A_{q,g}(t) \gamma^{(\mu} P^{\nu)} + B_{q,g}(t) \frac{P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}}{2M} + D_{q,g}(t) \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{q,g}(t) M \eta^{\mu\nu} \right] u(p)$$

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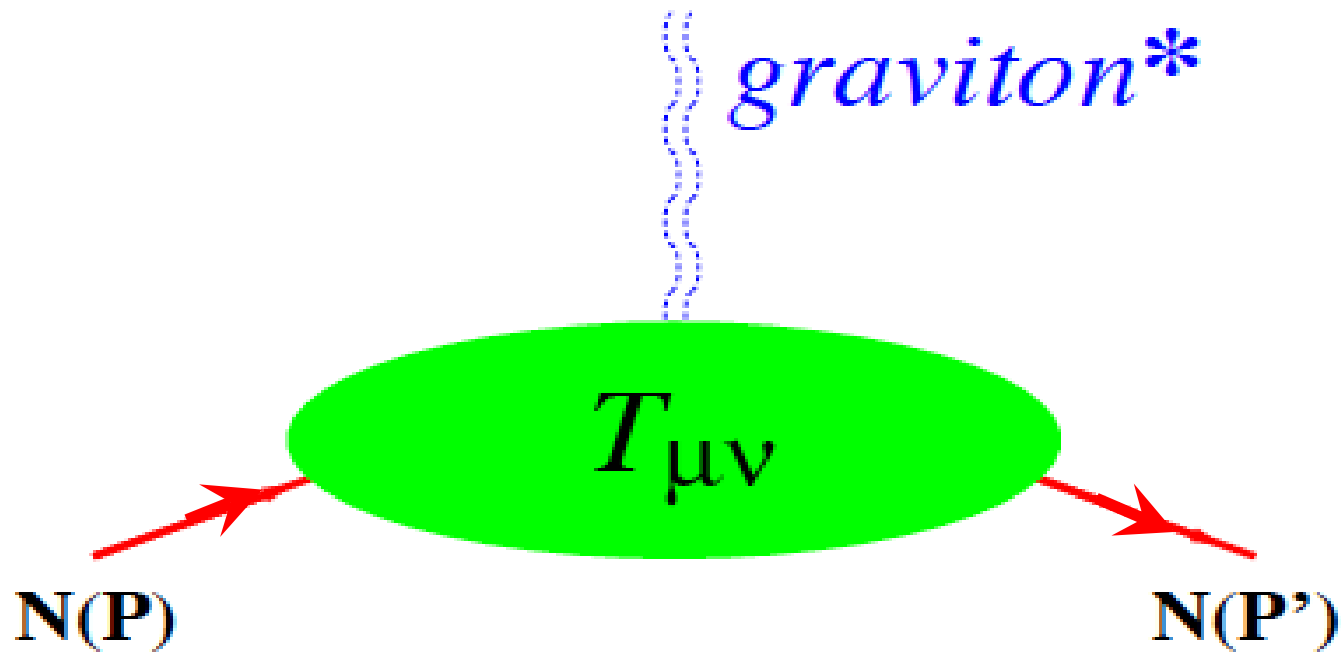
force & pressure distribution

mass & pressure distribution

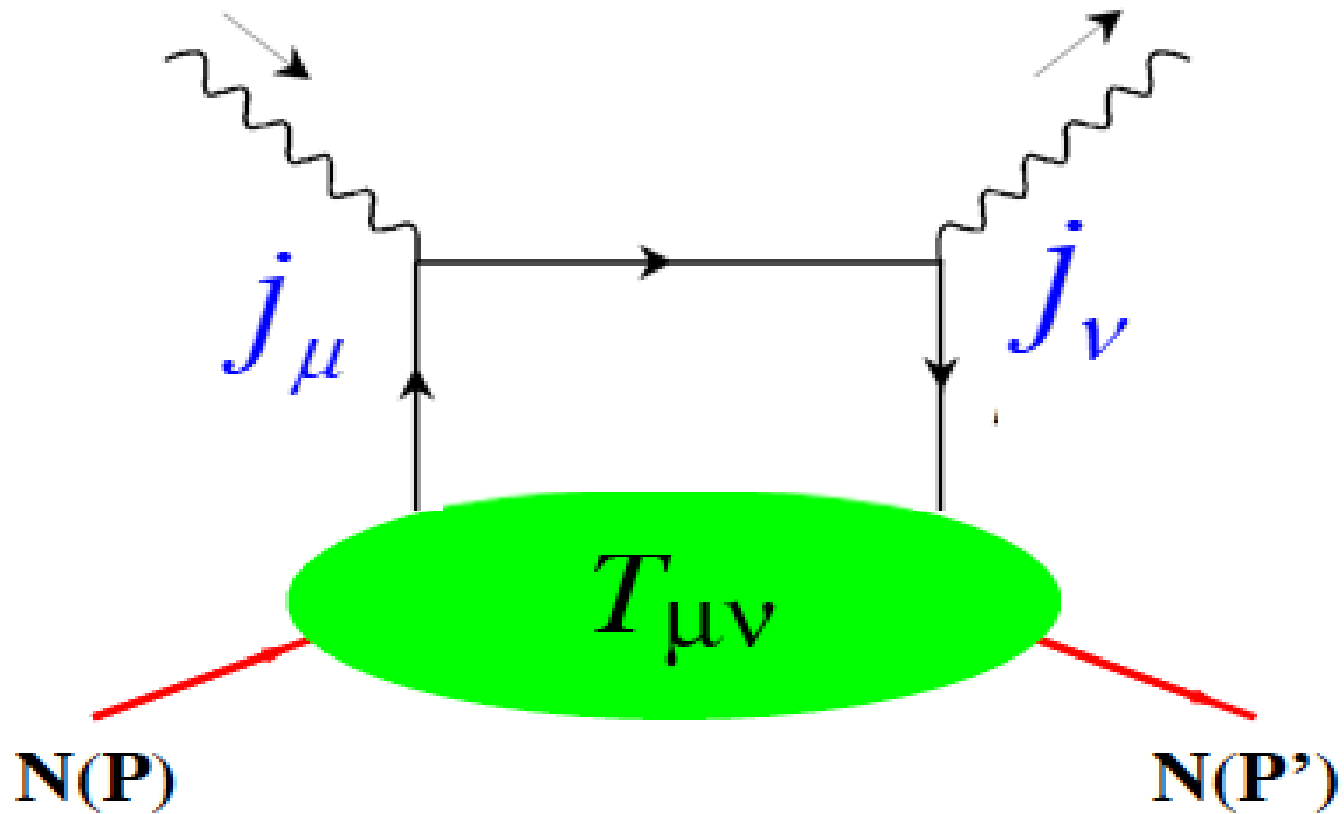
$$T^{\mu\nu} = \begin{bmatrix} \text{energy density} & \text{momentum density} & & \\ T^{00} & T^{01} & T^{02} & T^{03} \\ \text{momentum density} & \text{momentum flux} & & \\ T^{10} & T^{11} & T^{12} & T^{13} \\ T^{20} & T^{21} & T^{22} & T^{23} \\ T^{30} & T^{31} & T^{32} & T^{33} \end{bmatrix}$$

shear stress

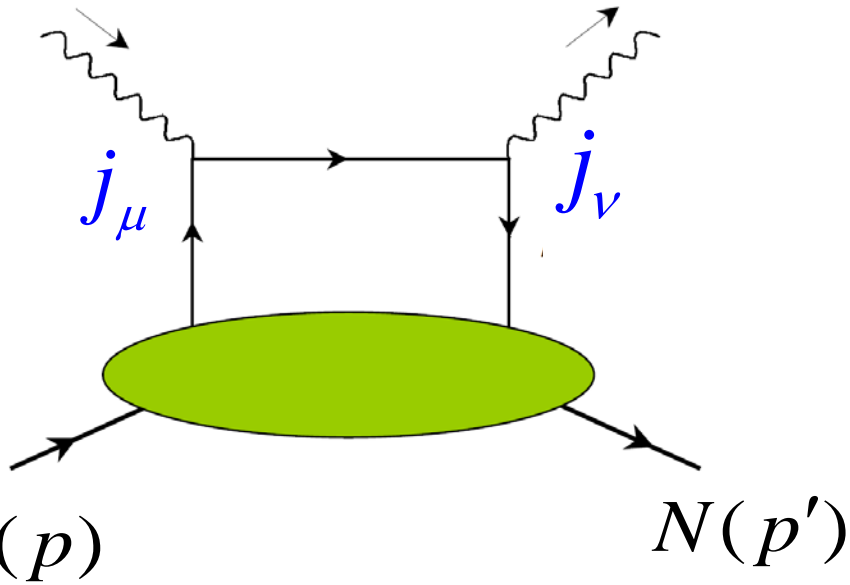
pressure



$$\begin{aligned}
 T^{\mu\nu} &= \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \\
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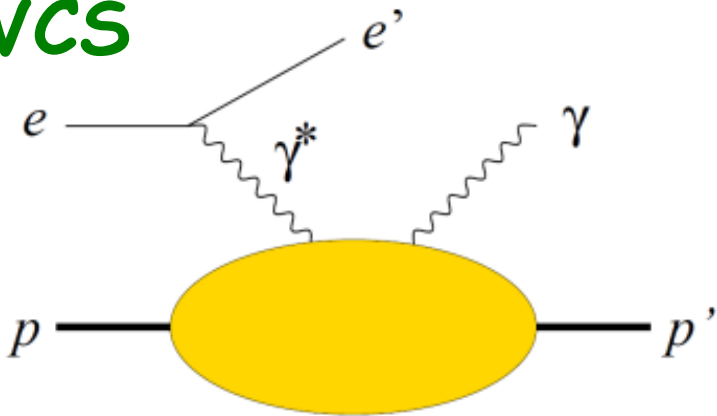


$$j_\mu(x) j_\nu(0) \sim \sum_i C_i(x) O_i(0)$$

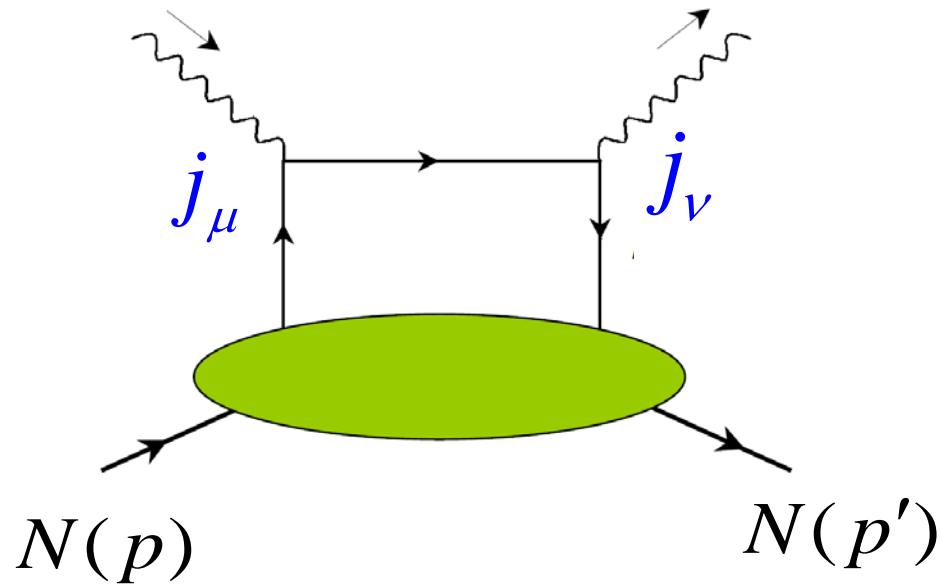
$$C_{\mu\nu;\alpha\beta}^q(x) T_q^{\alpha\beta}$$

$$C_{\mu\nu;\alpha\beta}^g(x) T_g^{\alpha\beta}$$

DVCS



JLab, EIC...

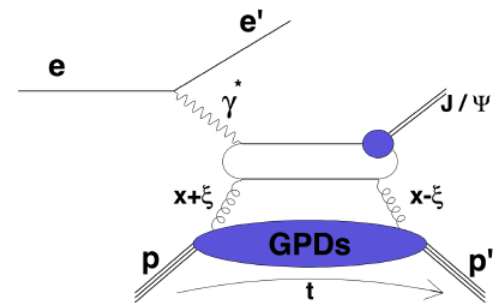
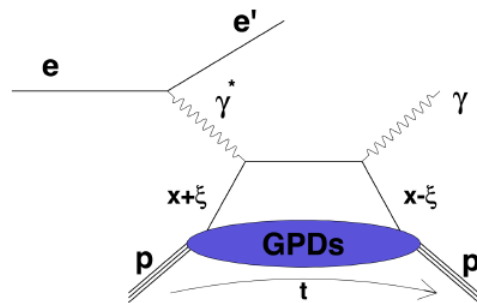
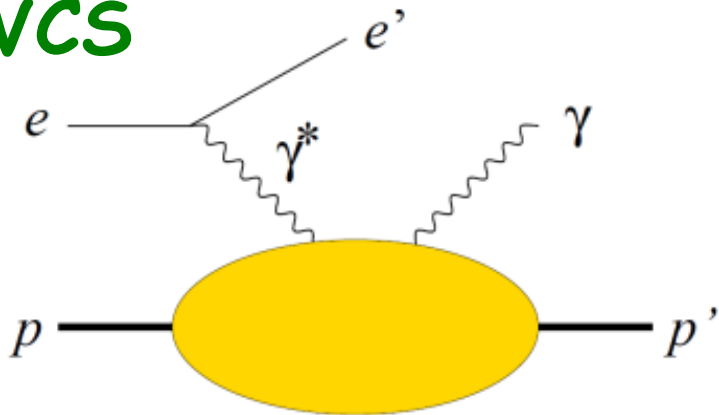


$$j_\mu(x) j_\nu(0) \sim \sum_i C_i(x) O_i(0)$$

$$C_{\mu\nu;\alpha\beta}^q(x) T_q^{\alpha\beta}$$

$$C_{\mu\nu;\alpha\beta}^g(x) T_g^{\alpha\beta}$$

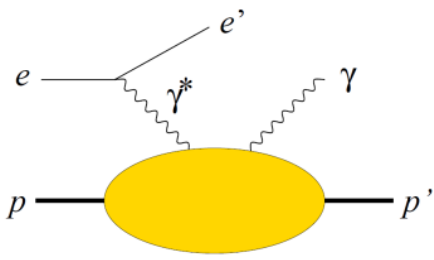
DVCS



JLab, EIC...

DVCS

$$P = \frac{p + p'}{2}$$



JLab, HERMES, COMPASS, EIC

$$\int \frac{dz^-}{2\pi} e^{ixPz^-} \langle N(p') | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \psi(\frac{z^-}{2}) | N(p) \rangle = \frac{1}{P^+} \left[H^q(x, \eta, t) \bar{u}(p') \gamma^+ u(p) + E^q(x, \eta, t) \bar{u}(p') \frac{i\sigma^{+\alpha} (p' - p)_\alpha}{2M} u(p) \right]$$

$$\int \frac{dz^-}{2\pi} e^{ixPz^-} \langle N(p') | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | N(p) \rangle = \frac{1}{P^+} \left[\tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p' - p)^+}{2M} u(p) \right]$$

$$\Delta^\mu = p'^\mu - p^\mu \rightarrow 0$$

$$\left(t = \Delta^2 \rightarrow 0, \quad \eta = \frac{-\Delta \cdot n}{2P \cdot n} \rightarrow 0 \right)$$

$$H^q(x, 0, 0) = q(x)$$

$$\int_{-1}^1 dx H^q(x, \eta, t) = F_1^q(t), \quad \int_{-1}^1 dx E^q(x, \eta, t) = F_2^q(t)$$

$$\int_{-1}^1 dx x H^q(x, \eta, t) = A_q(t) + 4\eta^2 D_q(t), \quad \int_{-1}^1 dx x E^q(x, \eta, t) = B_q(t) - 4\eta^2 D_q(t)$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

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force & pressure distribution

mass & pressure distribution

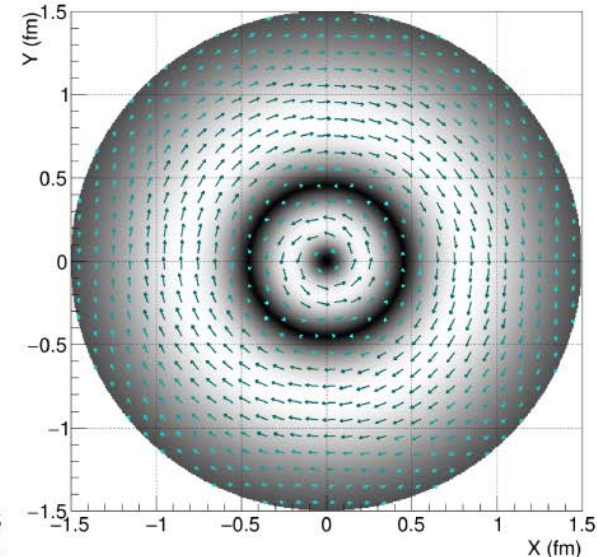
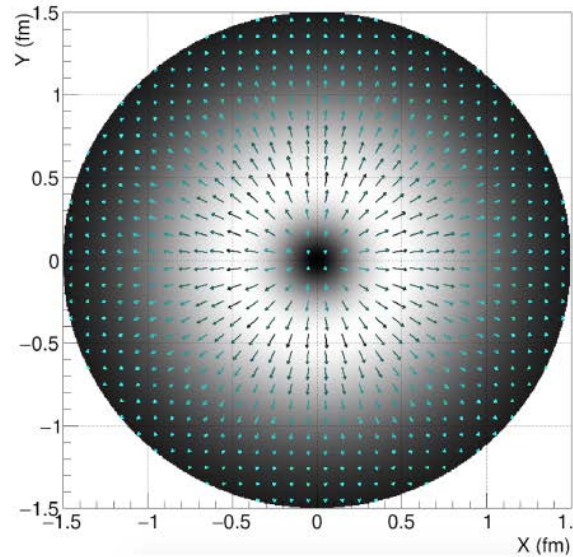
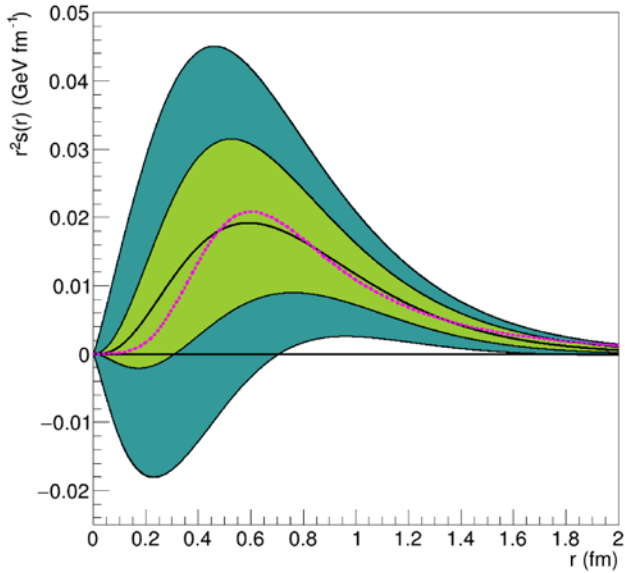
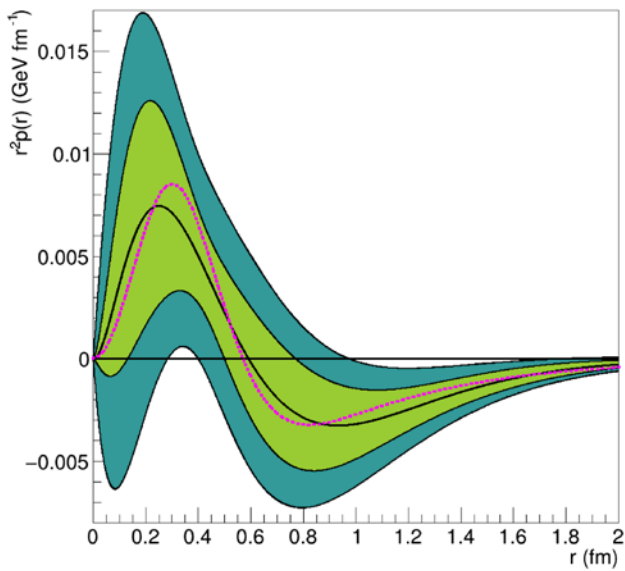
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shear stress

pressure

V. D. Burkert et al, Nature 557 ('18) 396

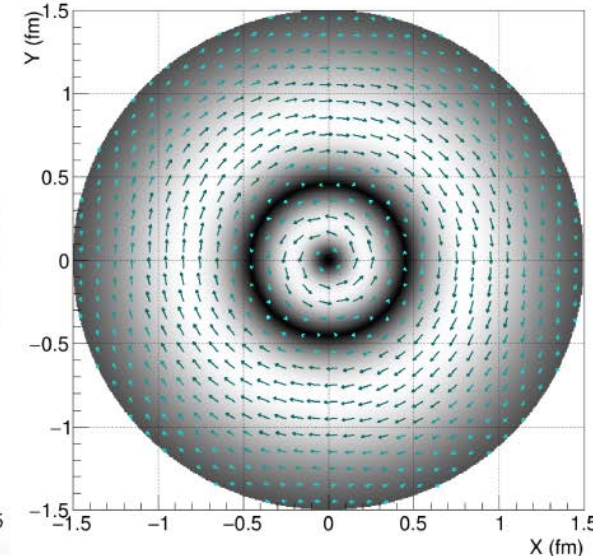
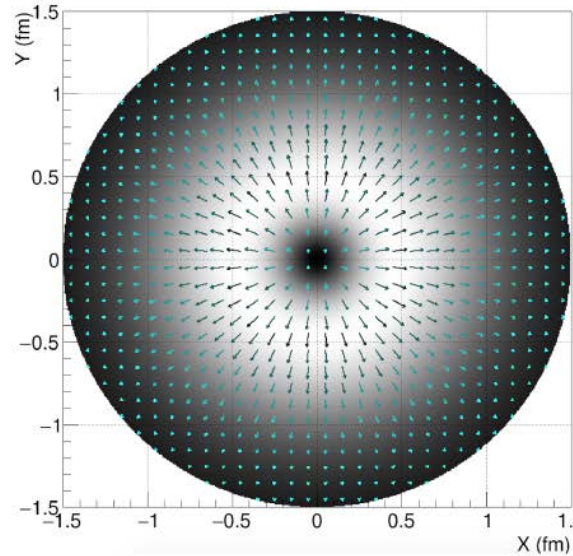
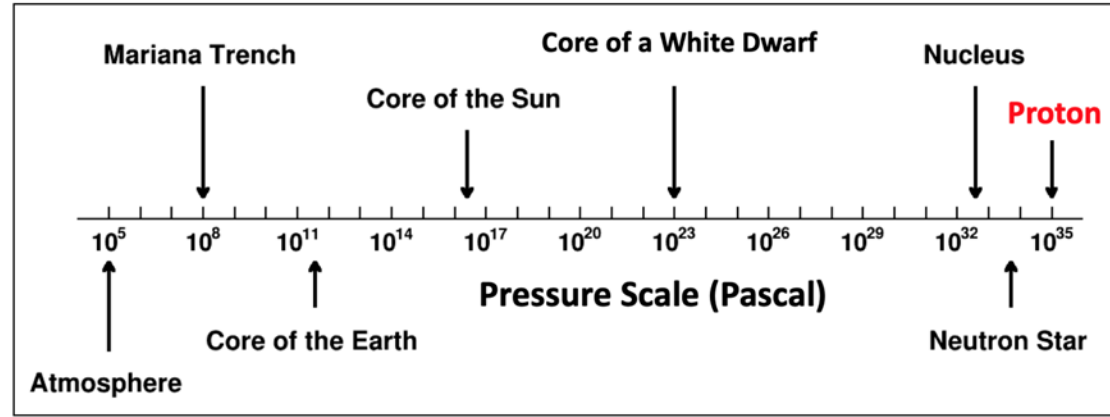
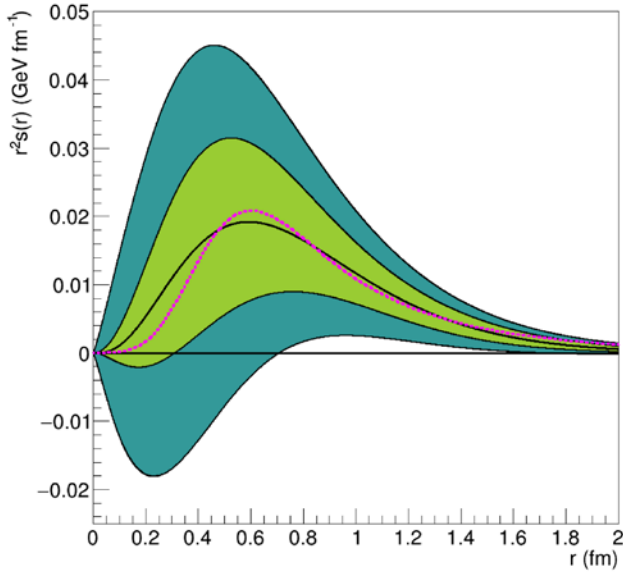
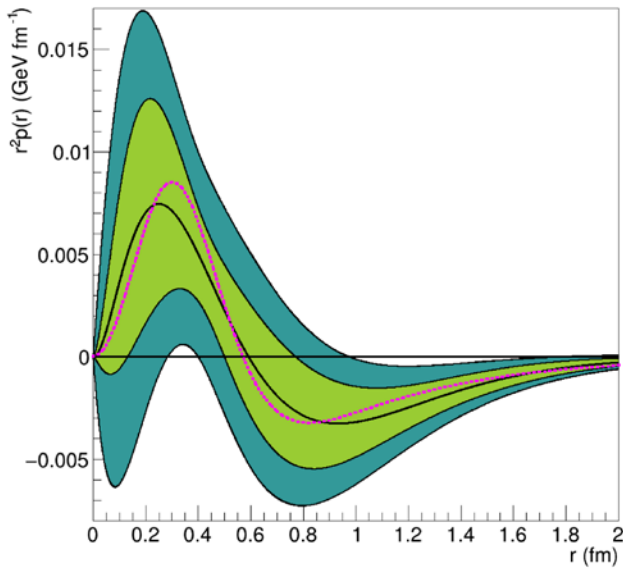
V. D. Burkert et al, 2303.08347



$$\langle N(p') | T^{ik} | N(p) \rangle \sim (\Delta^i \Delta^k - \delta^{ik} \Delta^2) D(t), \quad \langle T^{ij} \rangle(r) = \left(\frac{r^i r^j}{r^2} - \frac{1}{3} \delta^{ij} \right) s(r) + \delta^{ij} p(r)$$

V. D. Burkert et al, Nature 557 ('18) 396

V. D. Burkert et al, 2303.08347



$$\langle N(p') | T^{ik} | N(p) \rangle \sim (\Delta^i \Delta^k - \delta^{ik} \Delta^2) D(t), \quad \langle T^{ij} \rangle(r) = \left(\frac{r^i r^j}{r^2} - \frac{1}{3} \delta^{ij} \right) s(r) + \delta^{ij} p(r)$$

GPD: $H_q(x, \xi, t) = \int \frac{dy^-}{4\pi} e^{ixP^+y^-} \langle \pi^0(p') | \bar{\psi}(-y/2) \gamma^+ \psi(y/2) | \pi^0(p) \rangle \Big|_{y^+=0, \vec{y}_\perp=0}, \quad P^+ = \frac{(p+p')^+}{2}$

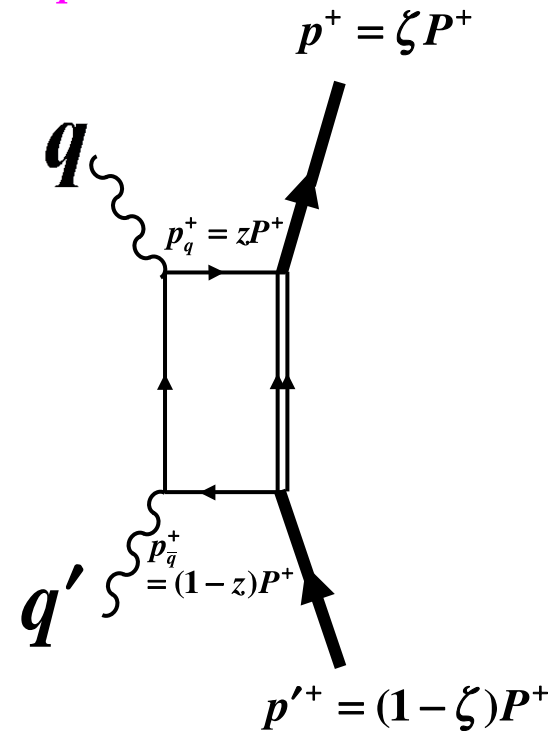
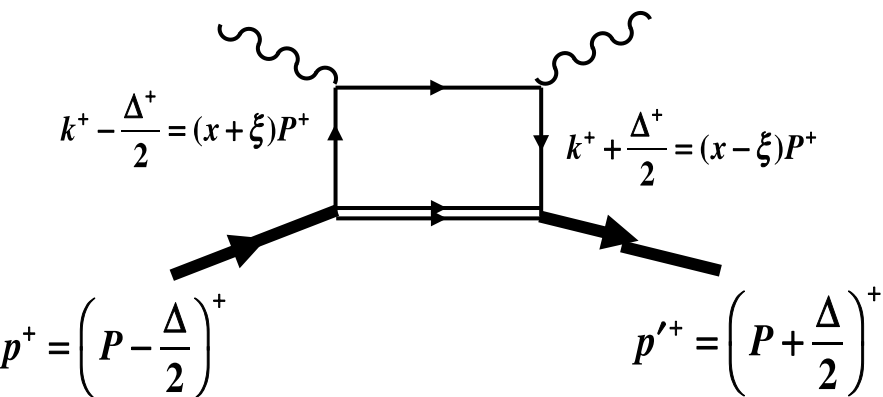
GDA: $\Phi_q(z, \zeta, s) = \int \frac{dy^-}{2\pi} e^{izP^+y^-} \langle \pi^0(p) \pi^0(p') | \bar{\psi}(-y/2) \gamma^+ \psi(y/2) | 0 \rangle \Big|_{y^+=0, \vec{y}_\perp=0}$

$H_q^h(x, \xi, t)$



$\Phi_q^{hh}(z, \zeta, W^2)$

s-t crossing



$\gamma\gamma^* \rightarrow \pi^0 \pi^0$

Spacelike gravitational form factors and radii for pion

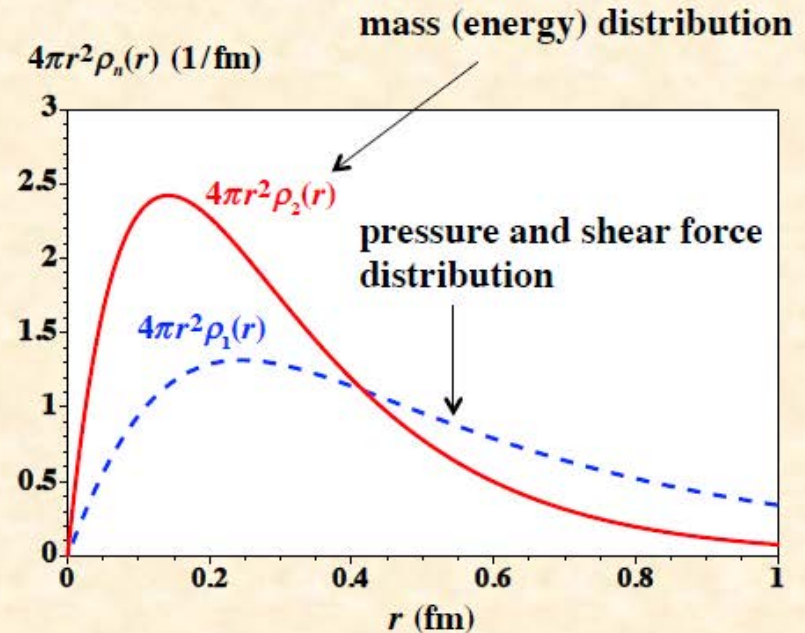
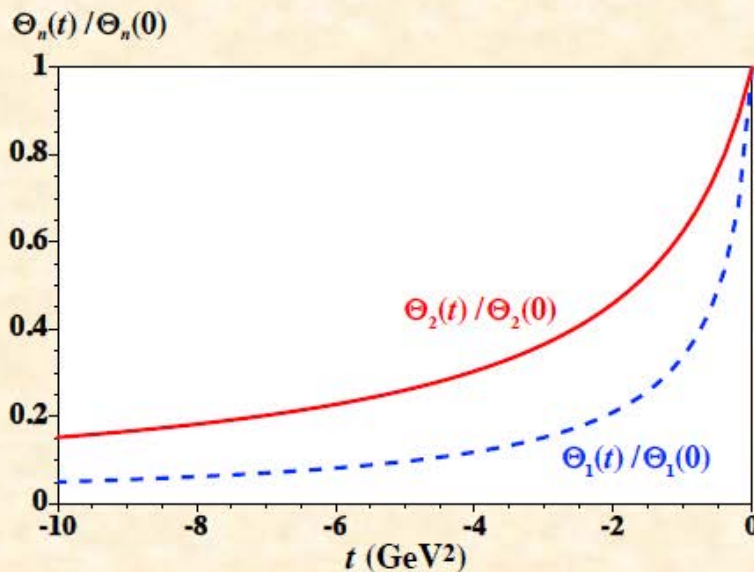
$$F(s) = \Theta_1(s), \Theta_1(s), \quad F(t) = \int_{4m_\pi^2}^{\infty} ds \frac{\text{Im}F(s)}{\pi(s-t-i\epsilon)}, \quad \rho(r) = \frac{1}{(2\pi)^3} \int d^3q e^{-i\vec{q}\cdot\vec{r}} F(q) = \frac{1}{4\pi^2} \frac{1}{r} \int_{4m_\pi^2}^{\infty} ds e^{-\sqrt{s}r} \text{Im}F(s)$$

This is the first report on gravitational radii of hadrons from actual experimental measurements.

$$\sqrt{\langle r^2 \rangle_{\text{mass}}} = 0.32 \sim 0.39 \text{ fm}, \quad \sqrt{\langle r^2 \rangle_{\text{mech}}} = 0.82 \sim 0.88 \text{ fm}$$

First finding on gravitational radius from actual experimental measurements

$$\Leftrightarrow \sqrt{\langle r^2 \rangle_{\text{charge}}} = 0.672 \pm 0.008 \text{ fm}$$



$$\Theta_2(t) = 4A^\pi(t), \quad \Theta_1(t) = -D^\pi(t)$$

Spacelike gravitational form factors and radii for pion

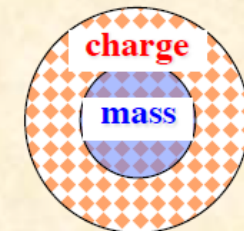
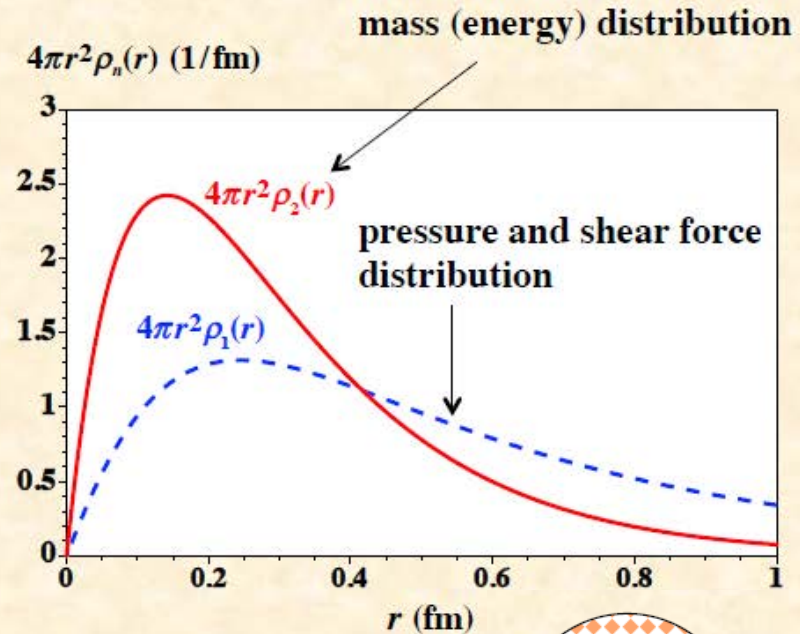
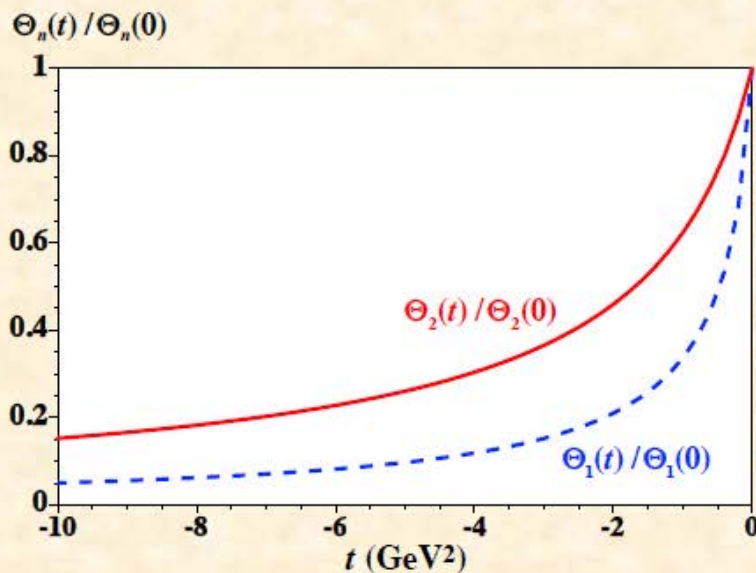
$$F(s) = \Theta_1(s), \Theta_1(s), \quad F(t) = \int_{4m_\pi^2}^{\infty} ds \frac{\text{Im}F(s)}{\pi(s-t-i\epsilon)}, \quad \rho(r) = \frac{1}{(2\pi)^3} \int d^3q e^{-i\vec{q}\cdot\vec{r}} F(q) = \frac{1}{4\pi^2} \frac{1}{r} \int_{4m_\pi^2}^{\infty} ds e^{-\sqrt{s}r} \text{Im}F(s)$$

This is the first report on gravitational radii of hadrons from actual experimental measurements.

$$\sqrt{\langle r^2 \rangle_{\text{mass}}} = 0.32 \sim 0.39 \text{ fm}, \quad \sqrt{\langle r^2 \rangle_{\text{mech}}} = 0.82 \sim 0.88 \text{ fm}$$

First finding on gravitational radius from actual experimental measurements

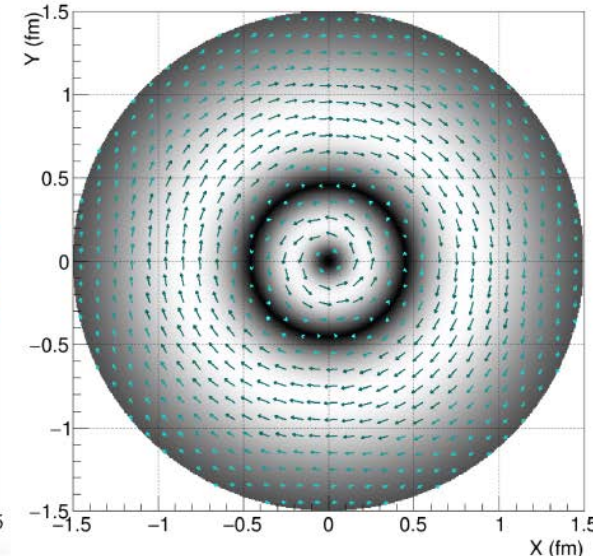
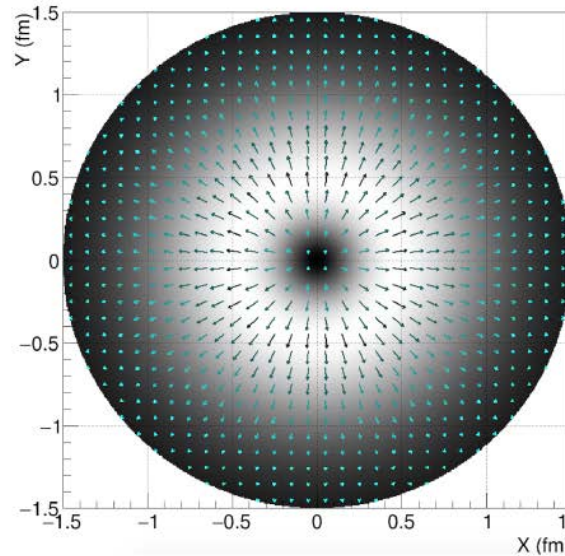
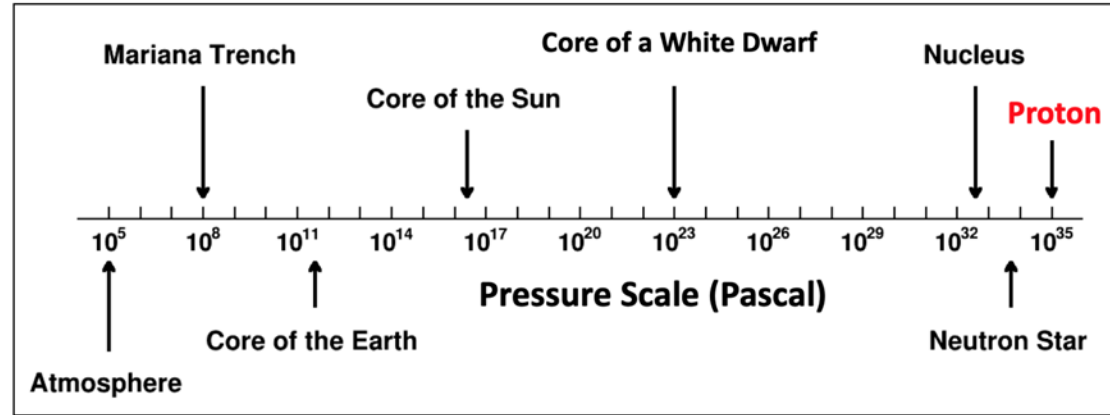
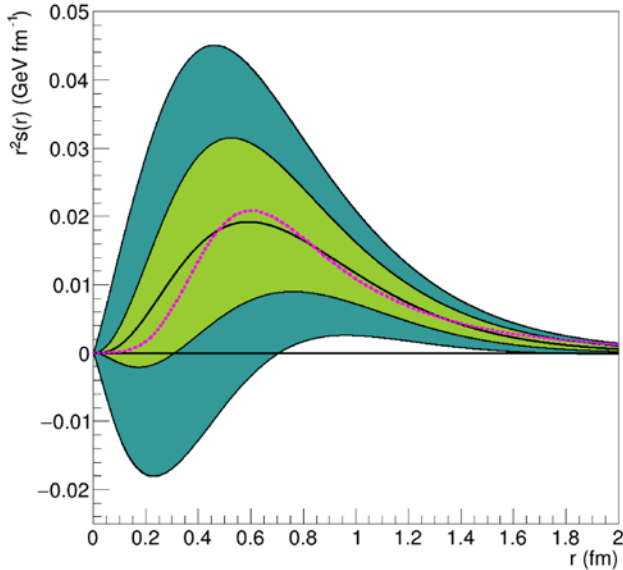
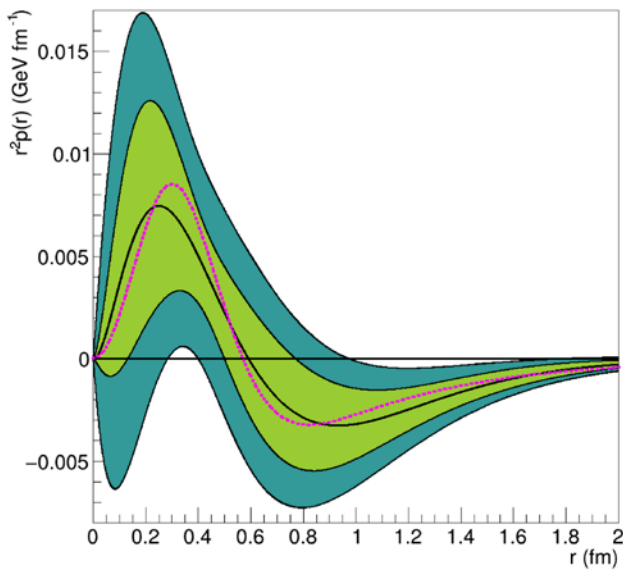
$$\Leftrightarrow \sqrt{\langle r^2 \rangle_{\text{charge}}} = 0.672 \pm 0.008 \text{ fm}$$



$$\Theta_2(t) = 4A^\pi(t), \quad \Theta_1(t) = -D^\pi(t)$$

V. D. Burkert et al, Nature 557 ('18) 396

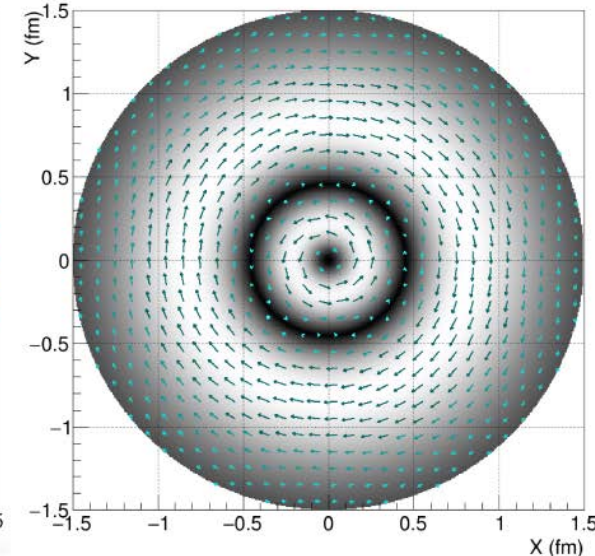
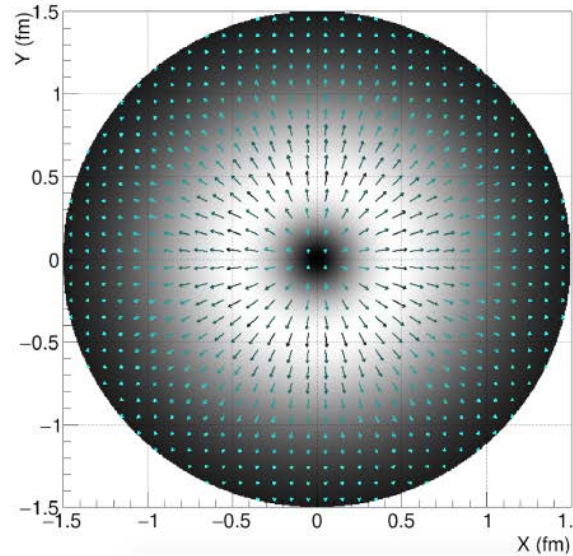
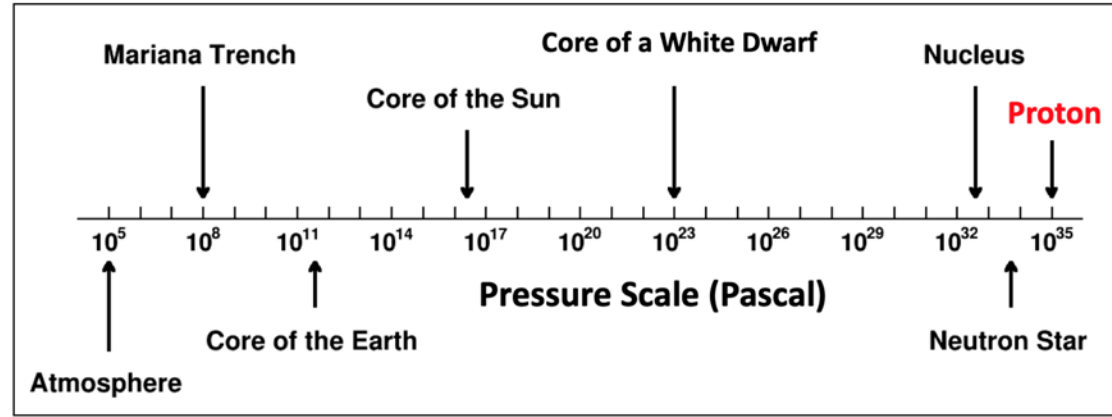
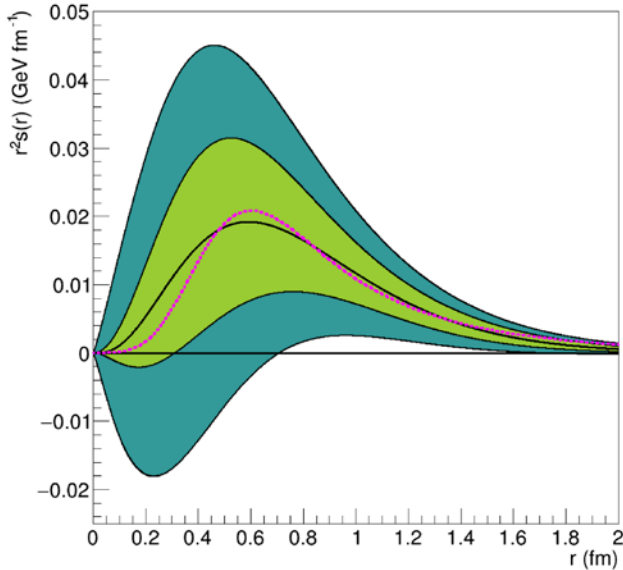
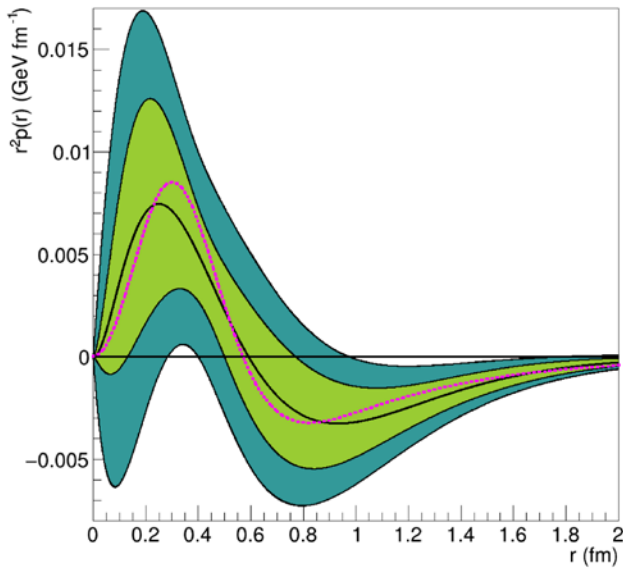
V. D. Burkert et al, 2303.08347



$$\langle N(p') | T^{ik} | N(p) \rangle \sim (\Delta^i \Delta^k - \delta^{ik} \Delta^2) D(t), \quad T^{ij}(r) = \left(\frac{r^i r^j}{r^2} - \frac{1}{3} \delta^{ij} \right) s(r) + \delta^{ij} p(r)$$

V. D. Burkert et al, Nature 557 ('18) 396

V. D. Burkert et al, 2303.08347



$$\langle N(p') | T_q^{ik} | N(p) \rangle \sim (\Delta^i \Delta^k - \delta^{ik} \Delta^2) D_q(t) - 4M^2 \delta^{ik} \bar{C}_q(t)$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\langle N(p') | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p') \left[A_{q,g}(t) \gamma^{(\mu} P^{\nu)} + B_{q,g}(t) \frac{P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}}{2M} + D_{q,g}(t) \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{q,g}(t) M \eta^{\mu\nu} \right] u(p)$$

$P = \frac{p + p'}{2}$
 $\Delta = p' - p$
 $t = \Delta^2$

$$A_q(0) + A_g(0) = 1$$

$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^{\mu} p^{\nu}$$

$$\frac{1}{2} (A_q(0) + B_q(0) + A_g(0) + B_g(0)) = \frac{1}{2}$$

$$\frac{\langle N(p) S | J^i | N(p) S \rangle}{\langle N(p) S | N(p) S \rangle} = \frac{1}{2} S^i$$

$$B_q(0) + B_g(0) = 0$$

$$J^i = \frac{1}{2} \epsilon^{ijk} \int d^3x M^{+jk}$$

$$M^{\mu\rho\sigma} = x^{\rho} T^{\mu\sigma} - x^{\sigma} T^{\mu\rho}$$

$$\bar{C}_q(t) + \bar{C}_g(t) = 0$$

$$\partial_{\mu} T^{\mu\nu} = 0$$

$$D \equiv D_q(0) + D_g(0)$$

“D term”: the last unknown global property

① mass decomposition

Ji, PRD52 271 ('95)

Lorce, Moutarde, Trawinski, EPJC79, 89 ('19)

Metz, Pasquini, Rodini, PRD102, 114042 ('20)

Ji, Liu, Schafer, NPB971, 115537 ('21)

② pressure

$$-\bar{C}_{q,g} \frac{M}{V}$$

Lorce, EPJC78, 120 ('18)

Liu, PRD104, 076010 ('21)

③ nucleon's transverse spin sum rule

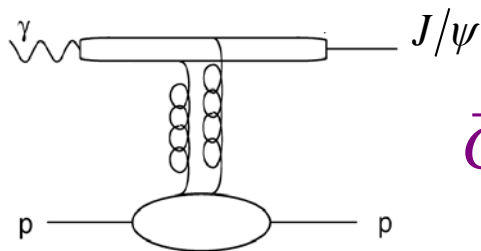
Hatta, KT, Yoshida,

JHEP 02 ('13) 003

$$J_{q,g} = \frac{1}{2}(A_{q,g} + B_{q,g}) + \frac{p^3}{2(p^0 + M)} \bar{C}_{q,g}$$

④ $\gamma p \rightarrow J/\psi p$ near threshold

JLab, EIC



$$\bar{C}_g \quad (= -\bar{C}_q)$$

Y. Hatta, D. Yang, PRD98, 074003

Y. Hatta, A. Rajan, D. Yang,

PRD100, 014032

Studies for $\bar{C}_{q,g}$ themselves

QCD EOMs $(i\not{D} - m)\psi = 0$, $D_\nu F^{\mu\nu} = g\bar{\psi}\gamma^\mu\psi$

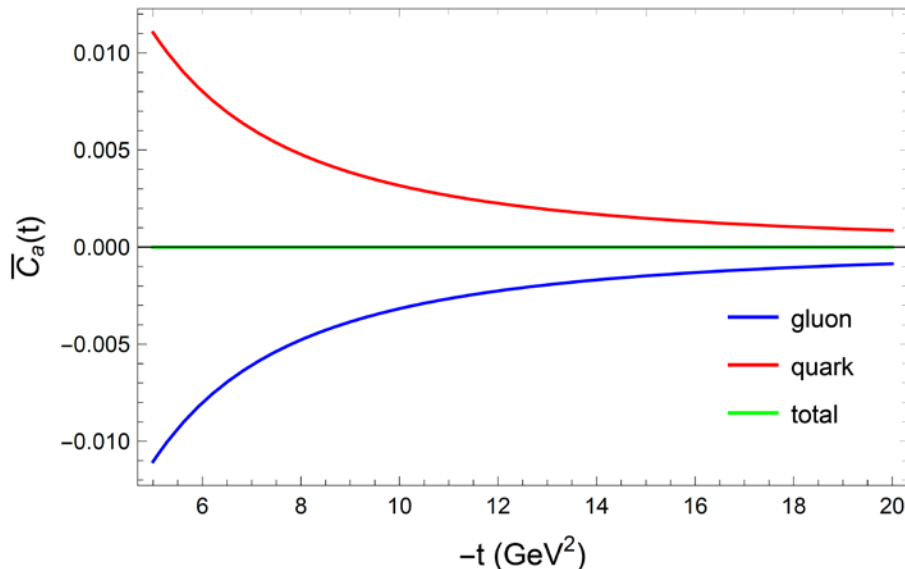
$$\partial_\nu T_q^{\mu\nu} = -\bar{\psi} g F^{\mu\nu} \gamma_\nu \psi, \quad \partial_\nu T_g^{\mu\nu} = -F_a^{\mu\nu} D_{ab}^\rho F_{\rho\nu}^b$$

KT, PRD98,
034009 ('18)

$$\Delta^\mu \bar{u}(p') u(p) M \bar{C}_q(t) = \langle N(p') | \bar{\psi} i g F^{\mu\nu} \gamma_\nu \psi | N(p) \rangle$$

$$\Delta^\mu \bar{u}(p', S') u(p, S) M \bar{C}_g(t) = \langle N(p') | F_a^{\mu\nu} i D_{ab}^\rho F_{\rho\nu}^b | N(p) \rangle$$

pQCD for large t



Tong, Ma, Yuan,
PLB823, 136751 ('21)

Tong, Ma, Yuan,
JHEP10, 046 ('22)

① mass decomposition

Ji, PRD52 271 ('95)

Lorce, Moutarde, Trawinski, EPJC79, 89 ('19)

Metz, Pasquini, Rodini, PRD102, 114042 ('20)

Ji, Liu, Schafer, NPB971, 115537 ('21)

② pressure

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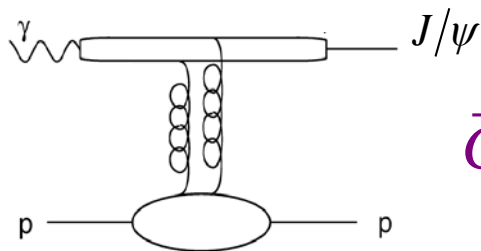
③ nucleon's transverse spin sum rule

Hatta, KT, Yoshida, JHEP 02 ('13) 003

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④ $\gamma p \rightarrow J/\psi p$ near threshold

JLab, EIC



$$\bar{C}_g (= -\bar{C}_q)$$

Y. Hatta, D. Yang, PRD98, 074003
Y. Hatta, A. Rajan, D. Yang, PRD100, 014032

at $t = 0$:

$$\bar{C}_q(0, \mu \sim 0.4 \text{ GeV}) = 0.25$$

Bag model [Ji, Melnitchouk, Song, PRD56, 5511 ('97)]

$$\bar{C}_q(0, \mu = 2 \text{ GeV}) \approx -0.11$$

Phenomenological [Lorce, EPJC78, 120 ('18)]

$$\bar{C}_q(0, \mu \sim 0.63 \text{ GeV}) = 0.014$$

Instanton [Polyakov, Son, JHEP09, 156 ('18)]

$$\bar{C}_q(0, \mu = 1 \text{ GeV}) = -0.021 \pm 0.008$$

LCSR [Azizi, Ozdem, EPJC80, 104 ('20)]

$$\bar{C}_q(0, \mu \rightarrow \infty) \simeq -0.15$$

Trace anomaly [Hatta, Rajan, KT, JHEP12, 008 ('18)]

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\langle N(p') | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p') \left[A_{q,g}(t) \gamma^{(\mu} P^{\nu)} + B_{q,g}(t) \frac{P^{(\mu} i \sigma^{\nu)\alpha} \Delta_{\alpha}}{2M} + D_{q,g}(t) \frac{\Delta^{\mu} \Delta^{\nu} - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{q,g}(t) M \eta^{\mu\nu} \right] u(p)$$

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$$D \equiv D_q(0) + D_g(0)$$

“D term”: the last unknown global property

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$\bar{C}_q(0) \left(= -\bar{C}_g(0) \right) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

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$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \quad \left(\beta(g) \equiv \mu \frac{dg}{d\mu}, \gamma_m(g) = -\frac{\mu}{m} \frac{dm}{d\mu} \right)$$

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

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$$\eta_{\mu\nu} T_{q,g}^{\mu\nu}$$

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

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$\eta_{\mu\nu} T_{q,g}^{\mu\nu}$

1&2-loop

Hatta, Rajan, KT, JHEP 12 ('18) 008

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

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$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

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$\eta_{\mu\nu} T_{q,g}^{\mu\nu}$

1&2-loop

Hatta, Rajan, KT, JHEP 12 ('18) 008

3-loop (& all orders)

KT, JHEP 01 ('19) 120

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$\bar{C}_q(0) \left(= -\bar{C}_g(0) \right) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

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$\eta_{\mu\nu} T_{q,g}^{\mu\nu}$

1&2-loop

Hatta, Rajan, KT, JHEP 12 ('18) 008

3-loop (& all orders)

KT, JHEP 01 ('19) 120

4-loop

Ahmed, Chen, Czakon, JHEP 01 ('23) 077

trace anomaly separately for q, g

$$\eta_{\mu\nu} T_q^{\mu\nu} = m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right)$$

$$\eta_{\mu\nu} T_g^{\mu\nu} = \frac{\alpha_s}{4\pi} \left(\frac{14}{3} C_F m\bar{\psi}\psi - \frac{11}{6} C_A F^2 \right)$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi \quad C_A = N_c, \quad C_F = \frac{N_c^2 - 1}{2N_c}$$

trace anomaly separately for q, g $\overline{\text{MS}}$ scheme Hatta, Rajan, KT, JHEP

$$\eta_{\mu\nu} T_q^{\mu\nu} = m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right) + \left(\frac{\alpha_s}{4\pi} \right)^2 \left[\left(C_F \left(\frac{61C_A}{27} - \frac{68n_f}{27} \right) - \frac{4C_F^2}{27} \right) m\bar{\psi}\psi + \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2 \right]$$

$$\eta_{\mu\nu} T_g^{\mu\nu} = \frac{\alpha_s}{4\pi} \left(\frac{14}{3} C_F m\bar{\psi}\psi - \frac{11}{6} C_A F^2 \right) + \left(\frac{\alpha_s}{4\pi} \right)^2 \left[\left(C_F \left(\frac{812C_A}{27} - \frac{22n_f}{27} \right) + \frac{85C_F^2}{27} \right) m\bar{\psi}\psi + \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2 \right]$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi \quad C_A = N_c, \quad C_F = \frac{N_c^2 - 1}{2N_c}$$

$$\begin{aligned}
 \eta_{\mu\nu} T_q^{\mu\nu} &= m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right) + \left(\frac{\alpha_s}{4\pi} \right)^2 \left[\left(C_F \left(\frac{61C_A}{27} - \frac{68n_f}{27} \right) - \frac{4C_F^2}{27} \right) m\bar{\psi}\psi + \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2 \right] \\
 &+ \left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\left(\frac{64\zeta(3)}{9} - \frac{8305}{729} \right) C_F^2 - \frac{2}{243} (864\zeta(3) + 1079) C_A C_F \right) - \frac{8}{729} (972\zeta(3) + 143) C_A C_F^2 \right. \right. \\
 &+ \left. \left. \left(\frac{32\zeta(3)}{9} + \frac{6611}{729} \right) C_A^2 C_F - \frac{76}{243} C_F n_f^2 + \frac{8}{729} (648\zeta(3) - 125) C_F^3 \right\} m\bar{\psi}\psi \right. \\
 &+ \left. \left\{ n_f \left(\left(\frac{52\zeta(3)}{9} - \frac{401}{324} \right) C_A C_F + \left(\frac{134}{27} - 4\zeta(3) \right) C_A^2 + \left(\frac{2407}{1458} - \frac{16\zeta(3)}{9} \right) C_F^2 \right) + n_f^2 \left(-\frac{697C_A}{729} - \frac{169C_F}{1458} \right) \right\} F^2 \right]
 \end{aligned}$$

$$\begin{aligned}
 \eta_{\mu\nu} T_g^{\mu\nu} &= \frac{\alpha_s}{4\pi} \left(\frac{14}{3} C_F m\bar{\psi}\psi - \frac{11}{6} C_A F^2 \right) + \left(\frac{\alpha_s}{4\pi} \right)^2 \left[\left(C_F \left(\frac{812C_A}{27} - \frac{22n_f}{27} \right) + \frac{85C_F^2}{27} \right) m\bar{\psi}\psi + \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2 \right] \\
 &+ \left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\left(\frac{368\zeta(3)}{9} - \frac{25229}{729} \right) C_F^2 - \frac{2}{243} (4968\zeta(3) + 1423) C_A C_F \right) + \left(\frac{32\zeta(3)}{3} - \frac{91753}{1458} \right) C_A C_F^2 \right. \right. \\
 &+ \left. \left. \left(\frac{294929}{1458} - \frac{32\zeta(3)}{9} \right) C_A^2 C_F - \frac{554}{243} C_F n_f^2 + \left(\frac{95041}{729} - \frac{64\zeta(3)}{9} \right) C_F^3 \right\} m\bar{\psi}\psi \right. \\
 &+ \left. \left\{ n_f \left(\left(\frac{1123}{162} - \frac{52\zeta(3)}{9} \right) C_A C_F + \left(4\zeta(3) + \frac{293}{36} \right) C_A^2 + \frac{16}{729} (81\zeta(3) - 98) C_F^2 \right) + n_f^2 \left(\frac{655C_A}{2916} - \frac{361C_F}{729} \right) - \frac{2857C_A^3}{108} \right\} F^2 \right]
 \end{aligned}$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi \quad C_A = N_c, \quad C_F = \frac{N_c^2 - 1}{2N_c}$$

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

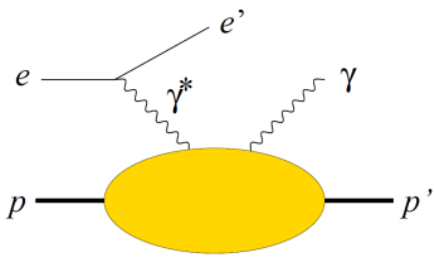
$$\bar{C}_q(0) \left(= -\bar{C}_g(0) \right) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \quad \left(\beta(g) \equiv \mu \frac{dg}{d\mu}, \gamma_m(g) = -\frac{\mu}{m} \frac{dm}{d\mu} \right)$$

DVCS

$$P = \frac{p + p'}{2}$$



JLab, HERMES, COMPASS, EIC

$$\int \frac{dz^-}{2\pi} e^{ixPz^-} \langle N(p') | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \psi(\frac{z^-}{2}) | N(p) \rangle = \frac{1}{P^+} \left[H^q(x, \eta, t) \bar{u}(p') \gamma^+ u(p) + E^q(x, \eta, t) \bar{u}(p') \frac{i\sigma^{+\alpha} (p' - p)_\alpha}{2M} u(p) \right]$$

$$\int \frac{dz^-}{2\pi} e^{ixPz^-} \langle N(p') | \bar{\psi}(-\frac{z^-}{2}) \gamma^+ \gamma_5 \psi(\frac{z^-}{2}) | N(p) \rangle = \frac{1}{P^+} \left[\tilde{H}^q(x, \eta, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}^q(x, \eta, t) \bar{u}(p') \frac{\gamma_5 (p' - p)^+}{2M} u(p) \right]$$

$$\Delta^\mu = p'^\mu - p^\mu \rightarrow 0$$

$$\left(t = \Delta^2 \rightarrow 0, \quad \eta = \frac{-\Delta \cdot n}{2\bar{P} \cdot n} \rightarrow 0 \right)$$

$$H^q(x, 0, 0) = q(x)$$

$$\int_{-1}^1 dx H^q(x, \eta, t) = F_1^q(t), \quad \int_{-1}^1 dx E^q(x, \eta, t) = F_2^q(t)$$

$$\int_{-1}^1 dx x H^q(x, \eta, t) = A_q(t) + 4\eta^2 D_q(t), \quad \int_{-1}^1 dx x E^q(x, \eta, t) = B_q(t) - 4\eta^2 D_q(t)$$

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$\bar{C}_q(0) \left(= -\bar{C}_g(0) \right) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \quad \left(\beta(g) \equiv \mu \frac{dg}{d\mu}, \gamma_m(g) = -\frac{\mu}{m} \frac{dm}{d\mu} \right)$$

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$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

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$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \quad \left(\beta(g) \equiv \mu \frac{dg}{d\mu}, \gamma_m(g) = -\frac{\mu}{m} \frac{dm}{d\mu} \right)$$

$$\int_0^1 dx x [q(x, \mu) + \bar{q}(x, \mu)] = A_q(t=0, \mu)$$

$$A_q(0, \mu) = \frac{n_f}{4C_F + n_f} + \frac{4C_F A_q(0, \mu_0) + n_f (A_q(0, \mu_0) - 1)}{4C_F + n_f} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} + \dots$$

$$\bar{C}_q(0) \quad (= -\bar{C}_g(0)) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$\begin{aligned} A_q(0, \mu) &= \int_0^1 dx x [q(x, \mu) + \bar{q}(x, \mu)] \\ &= \frac{n_f}{4C_F + n_f} + \frac{4C_F A_q(0, \mu_0) + n_f (A_q(0, \mu_0) - 1)}{4C_F + n_f} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} + \dots \end{aligned}$$

$$\eta_{\mu\nu} T_q^{\mu\nu} = m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right) + \dots$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi$$

$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^\mu p^\nu$$

$$\bar{C}_q(0) \quad (= -\bar{C}_g(0)) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

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$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi$$

$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^\mu p^\nu$$

$$2M^2 = \frac{\beta(g)}{2g} \langle N(p) | F^2 | N(p) \rangle + (1 + \gamma_m(g)) \langle N(p) | m\bar{\psi}\psi | N(p) \rangle$$

$$\begin{aligned}
\bar{C}_q(0, \mu) = & -\frac{1}{4} \left(\frac{n_f}{4C_F + n_f} + \frac{2n_f}{3\beta_0} \right) + \frac{1}{4} \left(\frac{2n_f}{3\beta_0} + 1 \right) \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} \\
& - \frac{4C_F A_q(\mu_0) + n_f (A_q(\mu_0) - 1)}{4(4C_F + n_f)} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} \\
& + \frac{\alpha_s(\mu)}{4\pi} \left[\frac{n_f}{4\beta_0} \left(-\frac{34C_A}{27} - \frac{49C_F}{27} \right) + \frac{\beta_1 n_f}{6\beta_0^2} \right] \\
& + \left[\frac{n_f (34C_A + 157C_F)}{108\beta_0} + \frac{C_F}{3} - \frac{\beta_1 n_f}{6\beta_0^2} \right] \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} - \frac{1}{4} A_q^{\text{NLO}}(\mu) \\
& + \left(\frac{\alpha_s(\mu)}{4\pi} \right)^2 \left(\frac{n_f^2}{\beta_0} \left[\frac{697C_A}{1458} + \frac{169C_F}{2916} \right] + n_f \left[\frac{17\beta_1 C_A}{54\beta_0^2} + \frac{\beta_2}{6\beta_0^2} + \frac{49\beta_1 C_F}{108\beta_0^2} \right] \right. \\
& + \frac{1}{\beta_0} \left\{ \left(\frac{401}{648} - \frac{26\zeta(3)}{9} \right) C_A C_F + \left(2\zeta(3) - \frac{67}{27} \right) C_A^2 + \left(\frac{8\zeta(3)}{9} - \frac{2407}{2916} \right) C_F^2 \right\} - \frac{\beta_1^2}{6\beta_0^3} \\
& + \left[-\frac{n_f^2}{\beta_0} \left(\frac{697C_A}{1458} + \frac{1789C_F}{2916} \right) + n_f \left(-\frac{17\beta_1 C_A}{54\beta_0^2} - \frac{\beta_2}{6\beta_0^2} - \frac{157\beta_1 C_F}{108\beta_0^2} + \frac{\beta_1^2}{6\beta_0^3} - \frac{17C_F}{27} \right) \right. \\
& + \left. \frac{n_f}{\beta_0} \left\{ \left(\frac{26\zeta(3)}{9} + \frac{4315}{648} \right) C_A C_F + \left(\frac{67}{27} - 2\zeta(3) \right) C_A^2 + \left(\frac{11803}{2916} - \frac{8\zeta(3)}{9} \right) C_F^2 \right\} \right. \\
& + \left. \frac{61C_A C_F}{108} - \frac{C_F^2}{27} \right] \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} - \frac{1}{4} A_q^{\text{NNLO}}(\mu) ,
\end{aligned}$$

$$\begin{aligned}
\bar{C}_q(0, \mu) = & -\frac{1}{4} \left(\frac{n_f}{4C_F + n_f} + \frac{2n_f}{3\beta_0} \right) + \frac{1}{4} \left(\frac{2n_f}{3\beta_0} + 1 \right) \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} \\
& - \frac{4C_F A_q(\mu_0) + n_f (A_q(\mu_0) - 1)}{4(4C_F + n_f)} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} A_q(\mu_0) \\
& + \frac{\alpha_s(\mu)}{4\pi} \left[\frac{n_f}{4\beta_0} \left(-\frac{34C_A}{27} - \frac{49C_F}{27} \right) + \frac{\beta_1 n_f}{6\beta_0^2} \right] \langle N(p) | m\bar{\psi}\psi | N(p) \rangle \\
& + \left[\frac{n_f (34C_A + 157C_F)}{108\beta_0} + \frac{C_F}{3} - \frac{\beta_1 n_f}{6\beta_0^2} \right] \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} - \frac{1}{4} A_q^{\text{NLO}}(\mu) \\
& + \left(\frac{\alpha_s(\mu)}{4\pi} \right)^2 \left(\frac{n_f^2}{\beta_0} \left[\frac{697C_A}{1458} + \frac{169C_F}{2916} \right] + n_f \left[\frac{17\beta_1 C_A}{54\beta_0^2} + \frac{\beta_2}{6\beta_0^2} + \frac{49\beta_1 C_F}{108\beta_0^2} \right] \right. \\
& + \frac{1}{\beta_0} \left\{ \left(\frac{401}{648} - \frac{26\zeta(3)}{9} \right) C_A C_F + \left(2\zeta(3) - \frac{67}{27} \right) C_A^2 + \left(\frac{8\zeta(3)}{9} - \frac{2407}{2916} \right) C_F^2 \right\} - \frac{\beta_1^2}{6\beta_0^3} \\
& + \left[-\frac{n_f^2}{\beta_0} \left(\frac{697C_A}{1458} + \frac{1789C_F}{2916} \right) + n_f \left(-\frac{17\beta_1 C_A}{54\beta_0^2} - \frac{\beta_2}{6\beta_0^2} - \frac{157\beta_1 C_F}{108\beta_0^2} + \frac{\beta_1^2}{6\beta_0^3} - \frac{17C_F}{27} \right) \right. \\
& + \left. \frac{n_f}{\beta_0} \left\{ \left(\frac{26\zeta(3)}{9} + \frac{4315}{648} \right) C_A C_F + \left(\frac{67}{27} - 2\zeta(3) \right) C_A^2 + \left(\frac{11803}{2916} - \frac{8\zeta(3)}{9} \right) C_F^2 \right\} \right. \\
& + \left. \frac{61C_A C_F}{108} - \frac{C_F^2}{27} \right] \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} - \frac{1}{4} A_q^{\text{NNLO}}(\mu) ,
\end{aligned}$$

$A_q(\mu_0) = \int_0^1 dx x [q(x, \mu_0) + \bar{q}(x, \mu_0)]$ global QCD analysis at NNLO

$$A_q(\mu_0 = 1.3 \text{ GeV}) = 0.613$$

CT18
(MMHT2014, NNPDF)

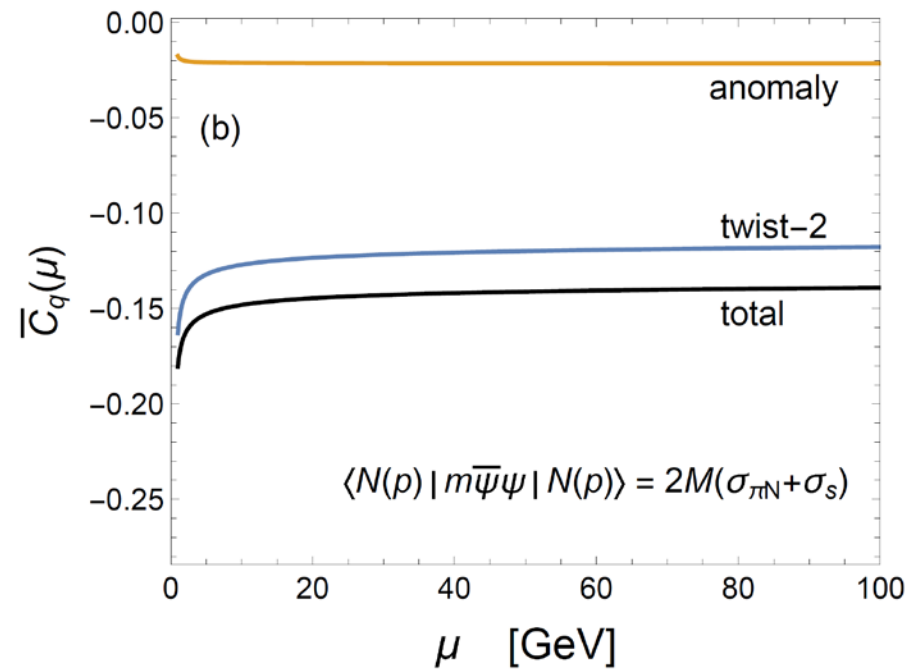
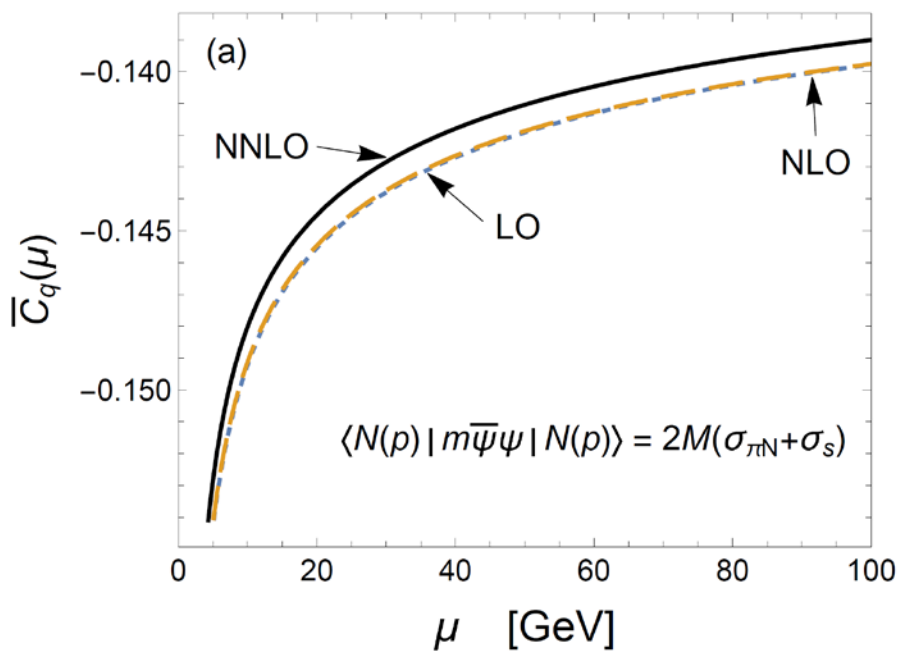
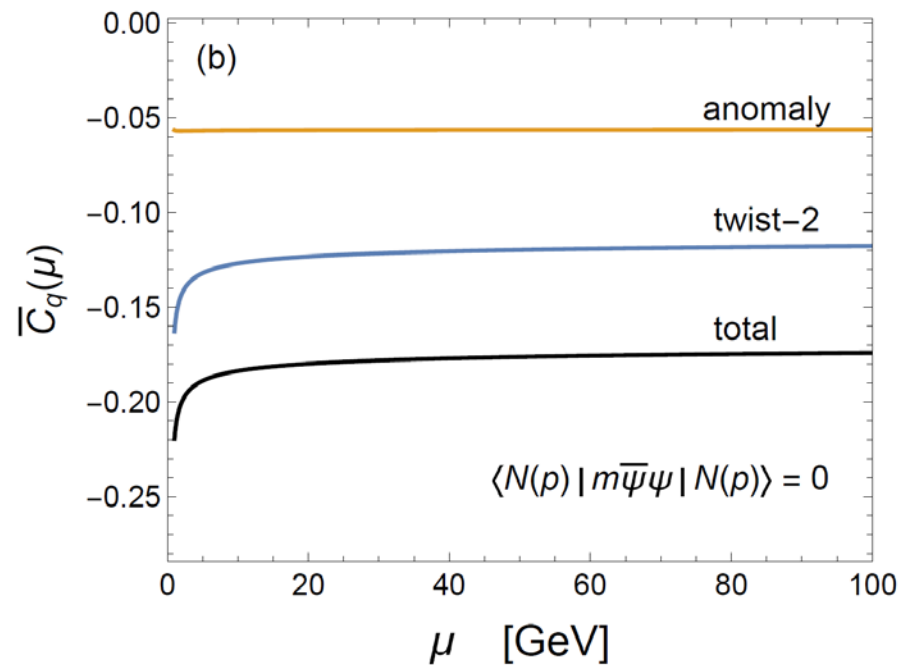
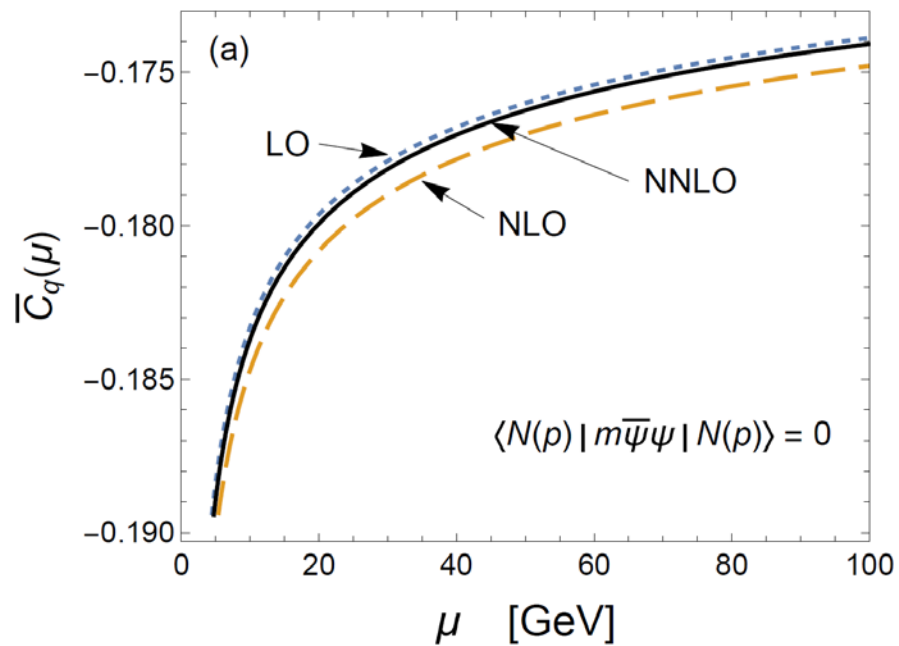
$$\langle N(p) | m \bar{\psi} \psi | N(p) \rangle = \langle N(p) | m_u \bar{u} u + m_d \bar{d} d + m_s \bar{s} s | N(p) \rangle = 2M (\sigma_{\pi N} + \sigma_s)$$

$$\sigma_{\pi N} = \frac{1}{2M} \langle N(p) | \frac{m_u + m_d}{2} (\bar{u} u + \bar{d} d) | N(p) \rangle = 59.1 \pm 3.5 \text{ MeV}$$

Hoferichter, Elvira, Kubis, Meißner, PRL115, 092301

$$\sigma_s = \frac{1}{2M} \langle N(p) | m_s \bar{s} s | N(p) \rangle = 45.6 \pm 6.2 \text{ MeV}$$

Alexandrou, et al., PRD102, 054517



$$\bar{C}_q(0) \quad (= -\bar{C}_g(0)) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

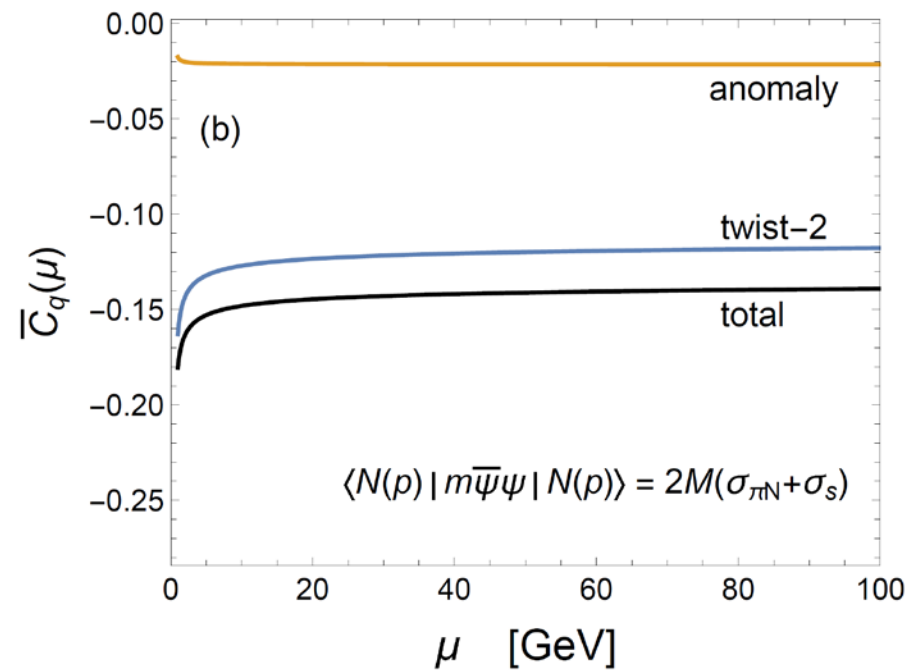
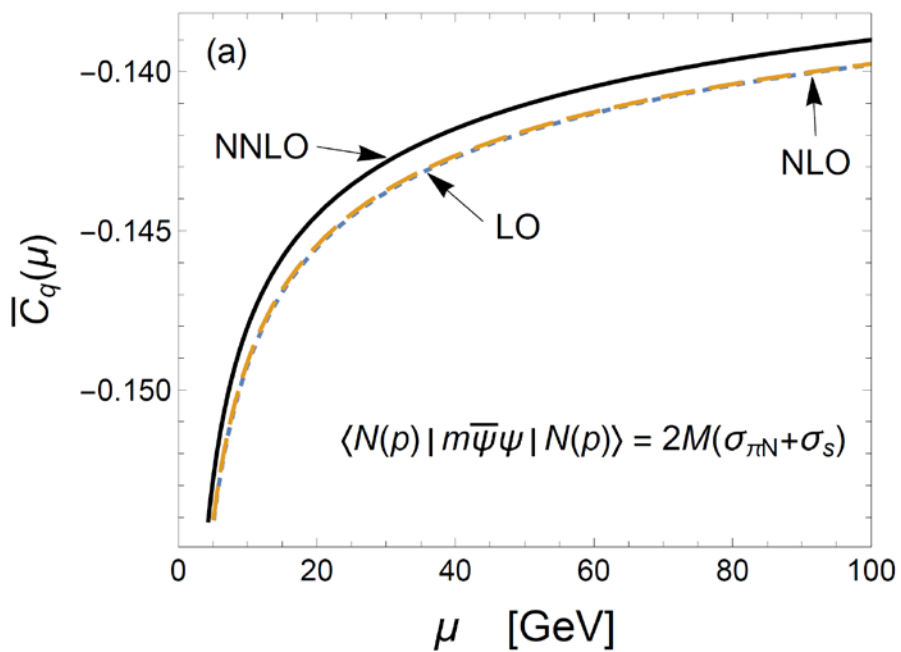
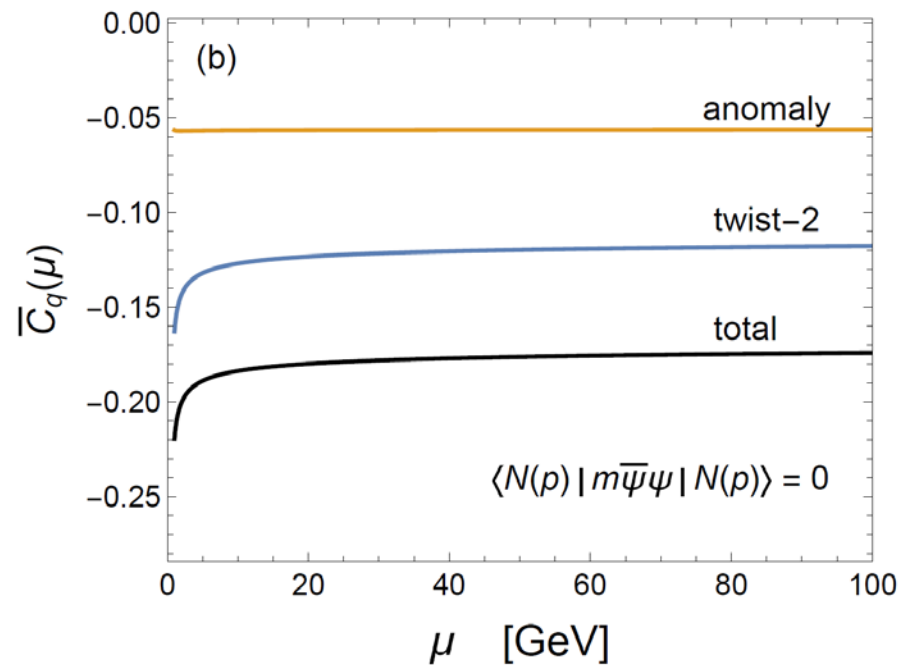
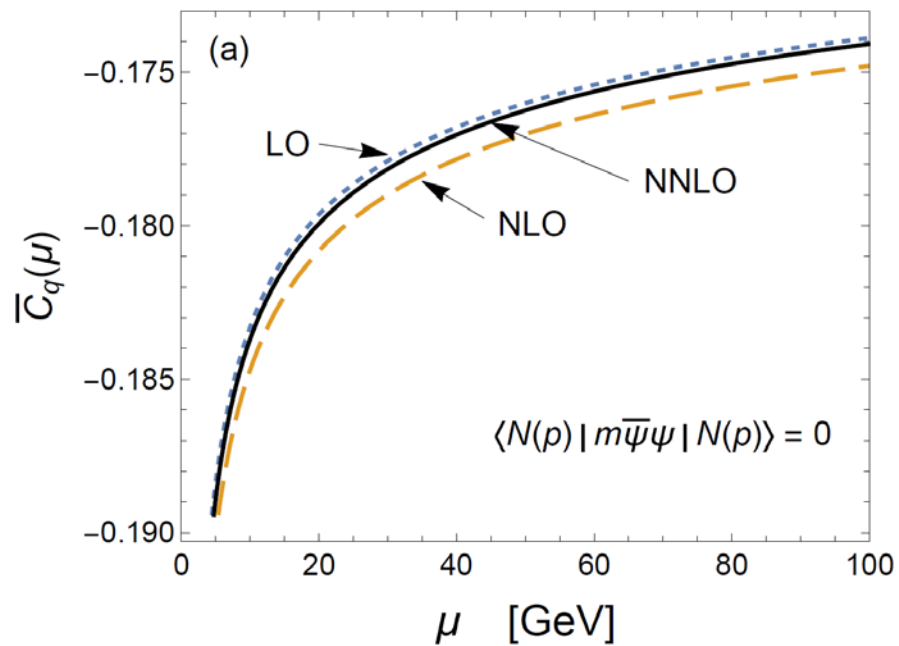
$$\begin{aligned} A_q(0, \mu) &= \int_0^1 dx x [q(x, \mu) + \bar{q}(x, \mu)] \\ &= \frac{n_f}{4C_F + n_f} + \frac{4C_F A_q(0, \mu_0) + n_f (A_q(0, \mu_0) - 1)}{4C_F + n_f} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} + \dots \end{aligned}$$

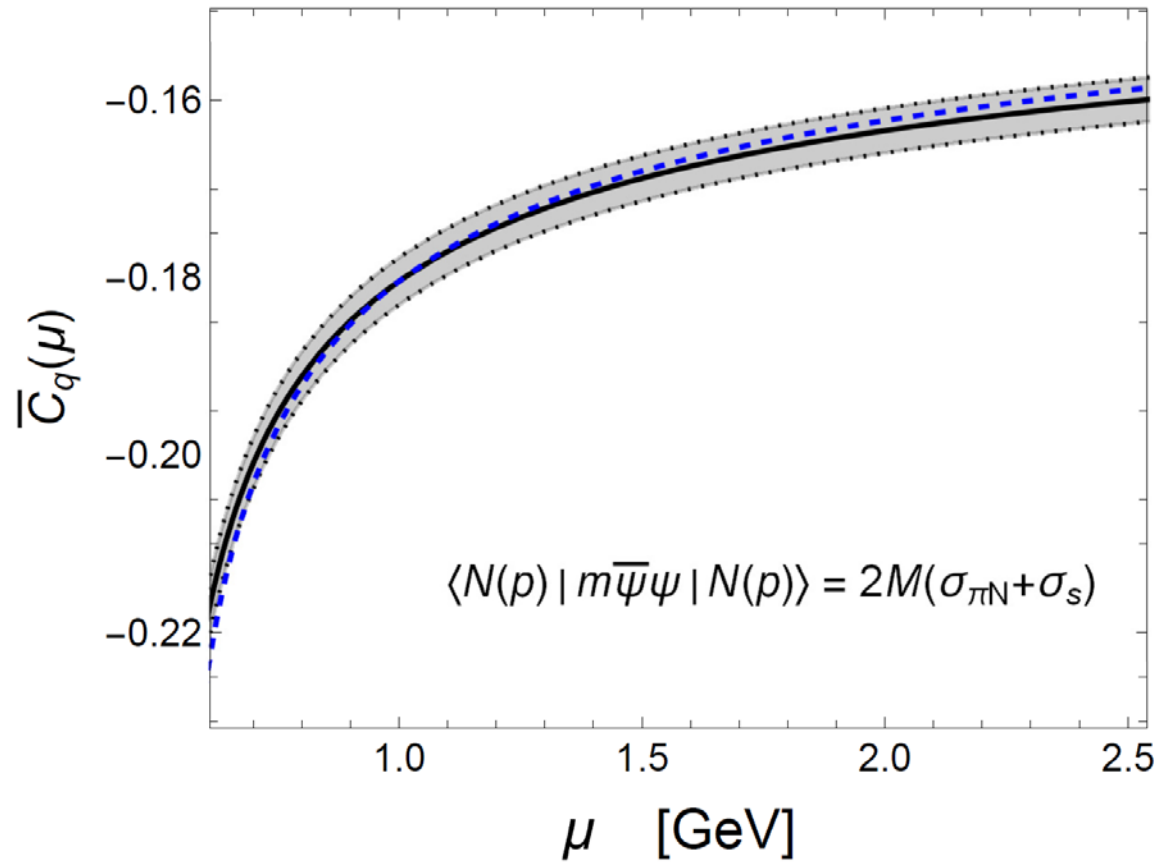
$$\eta_{\mu\nu} T_q^{\mu\nu} = m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right) + \dots$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m\bar{\psi}\psi$$

$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^\mu p^\nu$$

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$$\bar{C}_q(\mu = 0.7 \text{ GeV})|_{n_f=3} = -0.201 \pm 0.003$$

$$\bar{C}_q(\mu = 1 \text{ GeV})|_{n_f=3} = -0.180 \pm 0.003$$

$$\bar{C}_q(\mu = 2 \text{ GeV})|_{n_f=3} = -0.163 \pm 0.003$$

$\overline{\text{MS}}$ scheme

$$\bar{C}_q(\mu)|_{n_f=3} \simeq -0.108 - 0.114 [\alpha_s(\mu)]^{\frac{50}{81}}$$

$$\bar{C}_q(0) \quad (= -\bar{C}_g(0)) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$\begin{aligned} A_q(0, \mu) &= \int_0^1 dx x [q(x, \mu) + \bar{q}(x, \mu)] \\ &= \frac{n_f}{4C_F + n_f} + \frac{4C_F A_q(0, \mu_0) + n_f (A_q(0, \mu_0) - 1)}{4C_F + n_f} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} + \dots \end{aligned}$$

$$\eta_{\mu\nu} T_q^{\mu\nu} = m\bar{\psi}\psi + \frac{\alpha_s}{4\pi} \left(\frac{4}{3} C_F m\bar{\psi}\psi + \frac{1}{3} n_f F^2 \right) + \dots$$

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$$\langle N(p) | T^{\mu\nu} | N(p) \rangle = 2p^\mu p^\nu$$

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$A_q(\mu_0) = \int_0^1 dx x [q(x, \mu_0) + \bar{q}(x, \mu_0)]$ global QCD analysis at NNLO

$$A_q(\mu_0 = 1.3 \text{ GeV}) = 0.613$$

CT18
(MMHT2014, NNPDF)

$$\langle N(p) | m \bar{\psi} \psi | N(p) \rangle = \langle N(p) | m_u \bar{u} u + m_d \bar{d} d + m_s \bar{s} s | N(p) \rangle = 2M (\sigma_{\pi N} + \sigma_s)$$

$$\sigma_{\pi N} = \frac{1}{2M} \langle N(p) | \frac{m_u + m_d}{2} (\bar{u} u + \bar{d} d) | N(p) \rangle = 59.1 \pm 3.5 \text{ MeV}$$

Hoferichter, Elvira, Kubis, Meißner, PRL115, 092301

$$\sigma_s = \frac{1}{2M} \langle N(p) | m_s \bar{s} s | N(p) \rangle = 45.6 \pm 6.2 \text{ MeV}$$

Alexandrou, et al., PRD102, 054517

$$A_q^\pi(\mu_0) = \int_0^1 dx x \left[q^\pi(x, \mu_0) + \bar{q}^\pi(x, \mu_0) \right]$$

global QCD analysis at **NLO**

$$A_q^\pi(\mu_0 = 1.3 \text{ GeV}) = \begin{cases} 0.70 \pm 0.02 & \mathbf{JAM ('18)} \\ 0.81 \pm 0.16 & \mathbf{xFitter ('20)} \\ 0.61 \pm 0.08 & \mathbf{JAM ('21)} \end{cases}$$

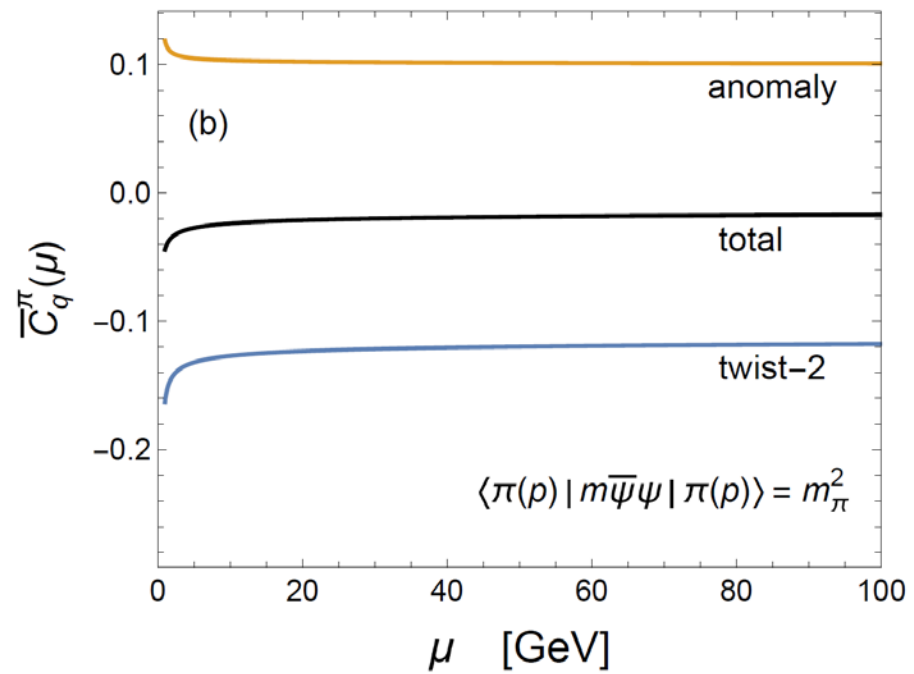
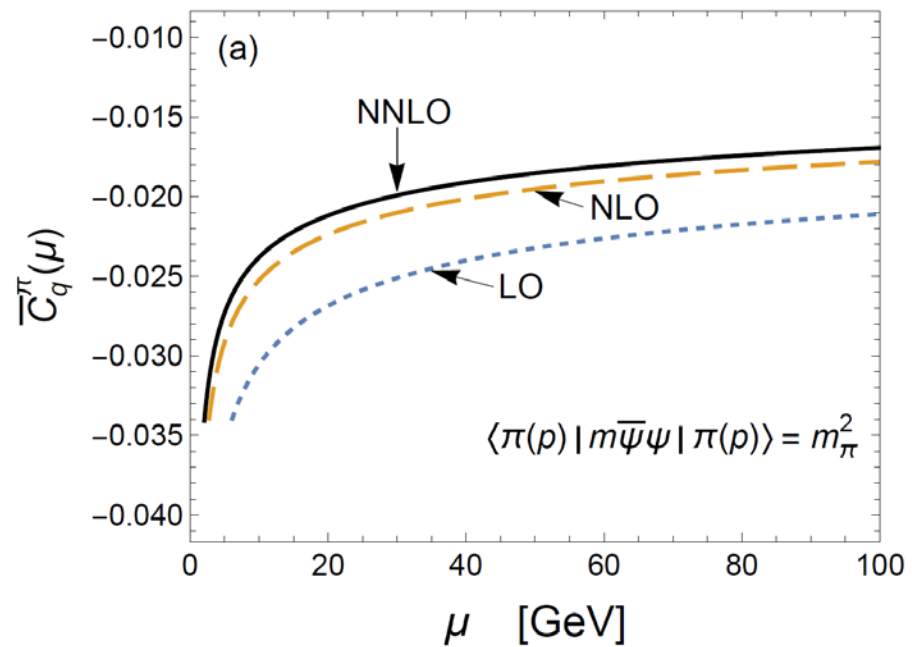
$$\langle \pi(p) | m \bar{\psi} \psi | \pi(p) \rangle$$

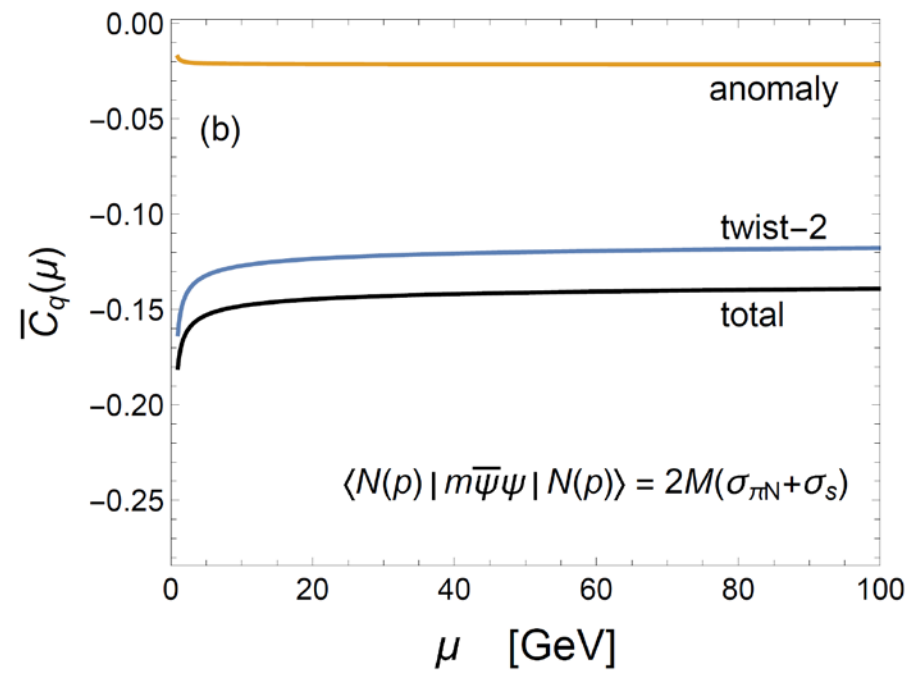
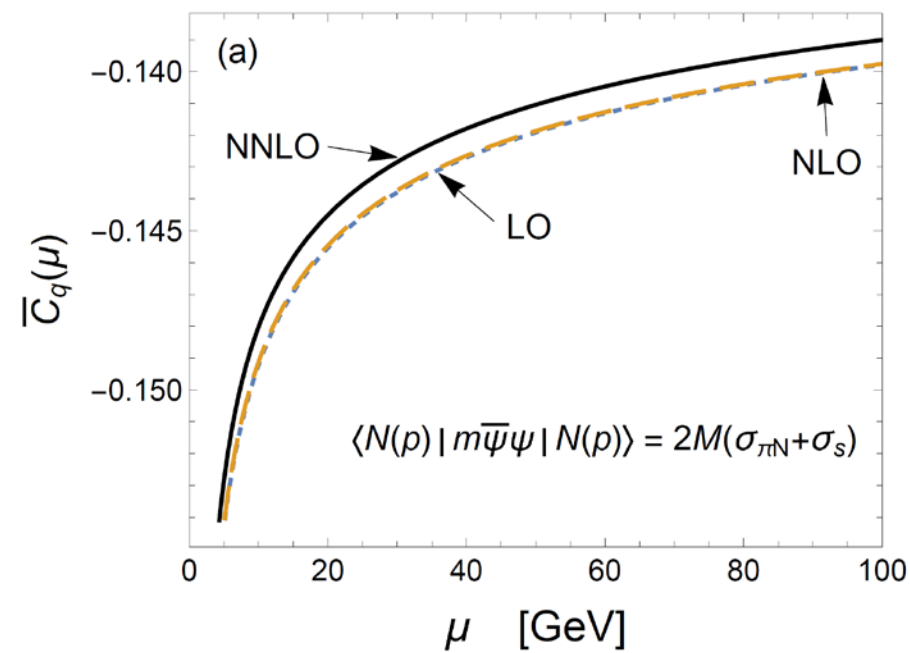
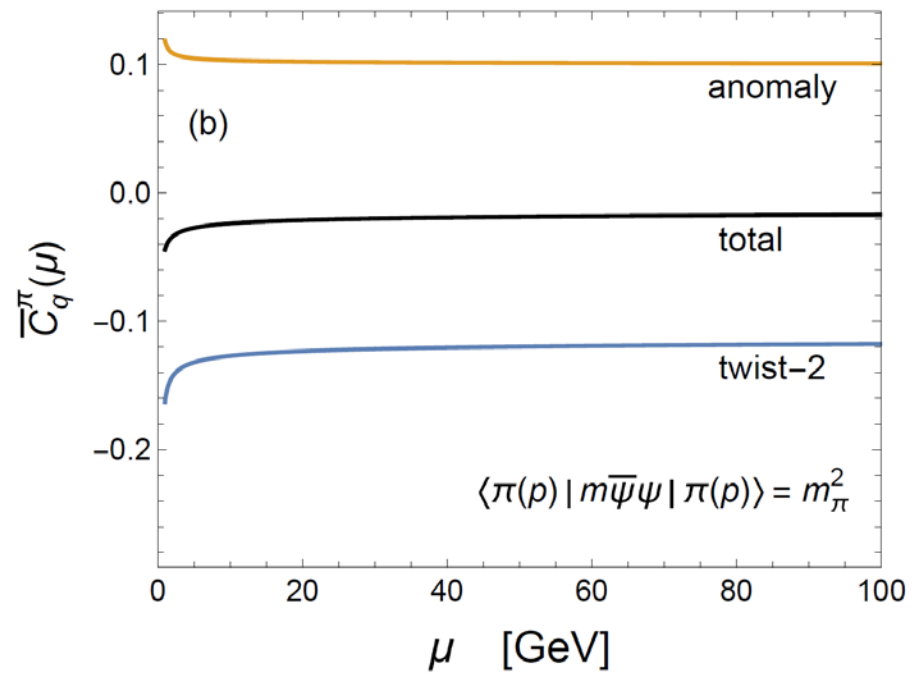
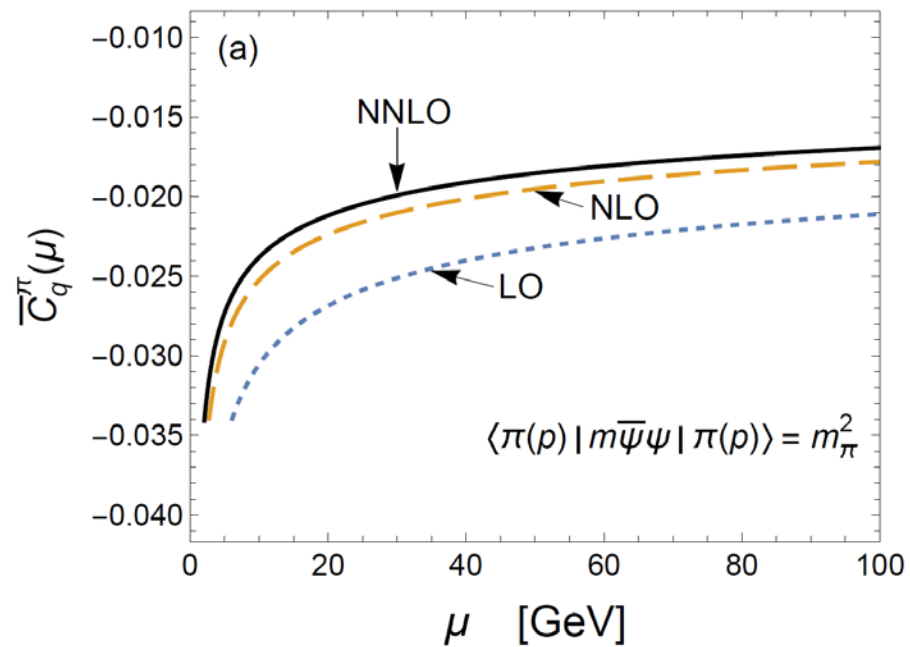
$$= m_\pi^2 + O(6\%)$$

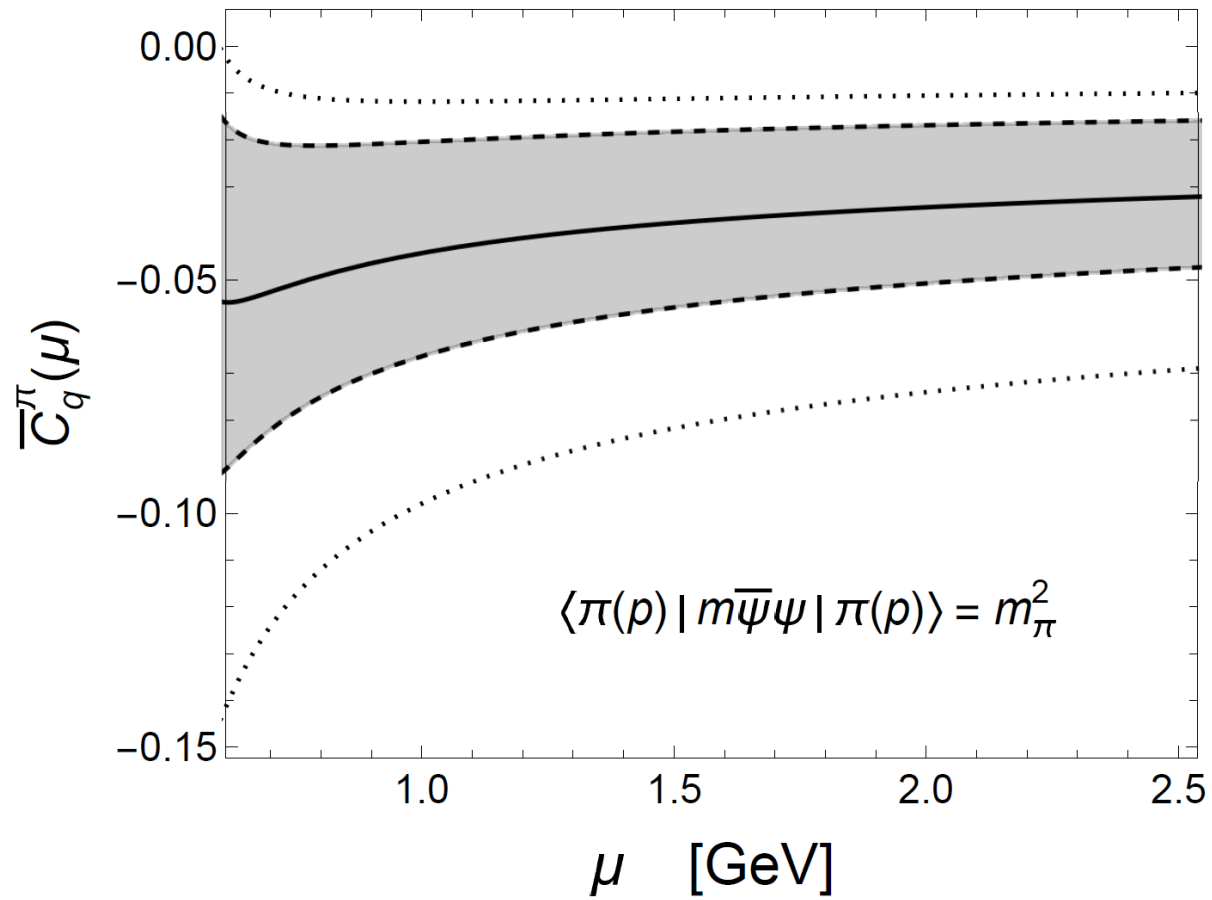
χ PT

Gasser, Leutwyler, *Annals Phys.* 158, 142

Colangelo, Gasser, Leutwyler, *PRL* 86, 5008







$$\bar{C}_q^\pi(\mu = 0.7 \text{ GeV})|_{n_f=3} = -0.05 \pm 0.03$$

$$\bar{C}_q^\pi(\mu = 1 \text{ GeV})|_{n_f=3} = -0.04 \pm 0.02$$

$$\bar{C}_q^\pi(\mu = 2 \text{ GeV})|_{n_f=3} = -0.03 \pm 0.02$$

$\overline{\text{MS}}$ scheme

at $t = 0$:

$$\bar{C}_q(0, \mu \sim 0.4 \text{ GeV}) = 0.25$$

Bag model [Ji, Melnitchouk, Song, PRD56, 5511 ('97)]

$$\bar{C}_q(0, \mu = 2 \text{ GeV}) \approx -0.11$$

Phenomenological [Lorce, EPJC78, 120 ('18)]

$$\bar{C}_q(0, \mu \sim 0.63 \text{ GeV}) = 0.014$$

Instanton [Polyakov, Son, JHEP09, 156 ('18)]

$$\bar{C}_q(0, \mu = 1 \text{ GeV}) = -0.021 \pm 0.008$$

LCSR [Azizi, Ozdem, EPJC80, 104 ('20)]

$$\bar{C}_q(0, \mu \rightarrow \infty) \simeq -0.15$$

Trace anomaly [Hatta, Rajan, KT, JHEP12, 008 ('18)]

nucleon

$$\bar{C}_q(0, \mu \sim 0.4 \text{ GeV}) = 0.25$$

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$$\bar{C}_q(0, \mu \rightarrow \infty) \simeq -0.15$$

Trace anomaly [Hatta, Rajan, KT, JHEP12, 008 ('18)]

$$\bar{C}_q(0, \mu = 1 \text{ GeV}) = -0.180 \pm 0.003$$

NNLO QCD [KT, JHEP03, 013 ('23)]

$$\bar{C}_q(0, \mu = 2 \text{ GeV}) = -0.163 \pm 0.003$$

pion

$$\bar{C}_q^\pi(0, \mu = 1 \text{ GeV}) = -0.04 \pm 0.02 \text{ NNLO QCD with NLO input [KT, JHEP03, 013 ('23)]}$$

① mass decomposition

Ji, PRD52 271 ('95)

Lorce, Moutarde, Trawinski, EPJC79, 89 ('19)

Metz, Pasquini, Rodini, PRD102, 114042 ('20)

Ji, Liu, Schafer, NPB971, 115537 ('21)

② pressure

$$-\bar{C}_{q,g} \frac{M}{V}$$

Lorce, EPJC78, 120 ('18)

Liu, PRD104, 076010 ('21)

③ nucleon's transverse spin sum rule

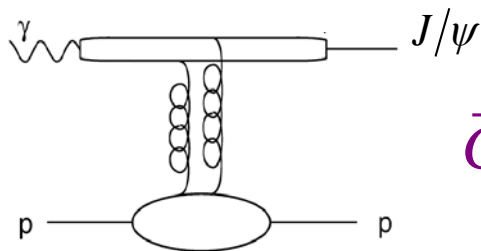
Hatta, KT, Yoshida,

JHEP 02 ('13) 003

$$J_{q,g} = \frac{1}{2} (A_{q,g} + B_{q,g}) + \frac{p^3}{2(p^0 + M)} \bar{C}_{q,g}$$

④ $\gamma p \rightarrow J/\psi p$ near threshold

JLab, EIC



$$\bar{C}_g (= -\bar{C}_q)$$

Y. Hatta, D. Yang, PRD98, 074003

Y. Hatta, A. Rajan, D. Yang,
PRD100, 014032

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$\bar{C}_q(0) \left(= -\bar{C}_g(0) \right) = -\frac{1}{4} A_q(0) + \frac{1}{8M^2} \langle N(p) | \eta_{\mu\nu} T_q^{\mu\nu} | N(p) \rangle$$

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2 \equiv T_q^{\mu\nu} + T_g^{\mu\nu}$$

$$\eta_{\mu\nu} T^{\mu\nu} = \frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \quad \left(\beta(g) \equiv \mu \frac{dg}{d\mu}, \gamma_m(g) = -\frac{\mu}{m} \frac{dm}{d\mu} \right)$$

$$\int_0^1 dx x [q(x, \mu) + \bar{q}(x, \mu)] = A_q(t=0, \mu)$$

$$A_q(0, \mu) = \frac{n_f}{4C_F + n_f} + \frac{4C_F A_q(0, \mu_0) + n_f (A_q(0, \mu_0) - 1)}{4C_F + n_f} \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{8C_F + 2n_f}{3\beta_0}} + \dots$$

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^{(0)} + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = M_q + M_g$$

$$M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^{(0)} + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

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$$M = M_q + M_g \quad M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

$$M = \left(M_q - M_m \right) + M_g + M_m \quad M_m = \left\langle \int d^3 x m \bar{\psi} \psi \right\rangle = \frac{\sigma_{\pi N} + \sigma_s}{M} M$$

$$M = \left(M_q - M_m - \Delta M_q \right) + \left(M_g - \Delta M_g \right) + M_m + \left(\Delta M_q + \Delta M_g \right)$$

$$\Delta M_q + \Delta M_g = \frac{1}{4} \left\langle \int d^3 x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M - M_m)$$

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^0 + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = M_q + M_g \quad M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

0.6, 0.4

$$M = \left(M_q - M_m \right) + M_g + M_m \quad M_m = \left\langle \int d^3 x m \bar{\psi} \psi \right\rangle = \frac{\sigma_{\pi N} + \sigma_s}{M} M$$

0.1

$$M = \left(M_q - M_m - \Delta M_q \right) + \left(M_g - \Delta M_g \right) + M_m + \left(\Delta M_q + \Delta M_g \right)$$

$$\Delta M_q + \Delta M_g = \frac{1}{4} \left\langle \int d^3 x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M - M_m)$$

0.2

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^0 + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = M_q + M_g \quad M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

$$0.6, 0.4 \quad -0.18, 0.18$$

$$M = \left(M_q - M_m \right) + M_g + M_m \quad M_m = \left\langle \int d^3 x m \bar{\psi} \psi \right\rangle = \frac{\sigma_{\pi N} + \sigma_s}{M} M$$

$$0.1$$

$$M = \left(M_q - M_m - \Delta M_q \right) + \left(M_g - \Delta M_g \right) + M_m + \left(\Delta M_q + \Delta M_g \right)$$

$$\Delta M_q + \Delta M_g = \frac{1}{4} \left\langle \int d^3 x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M - M_m)$$

$$0.2$$

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^{(0)} + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = \underbrace{M_q}_{0.4} + \underbrace{M_g}_{0.6} \qquad M_{q,g} = \left(\underbrace{A_{q,g}(0)}_{0.6, 0.4} + \underbrace{\bar{C}_{q,g}(0)}_{-0.18, 0.18} \right) M$$

$$M = \left(\underbrace{M_q}_{0.3} - \underbrace{M_m}_{0.1} \right) + \underbrace{M_g}_{0.6} + \underbrace{M_m}_{0.1} \qquad M_m = \left\langle \int d^3 x m \bar{\psi} \psi \right\rangle = \frac{\sigma_{\pi N} + \sigma_s}{M} M$$

$$M = \left(\underbrace{M_q}_{0.3} - \underbrace{M_m}_{0.1} - \underbrace{\Delta M_q}_{0.2} \right) + \left(\underbrace{M_g}_{0.6} - \underbrace{\Delta M_g}_{0.2} \right) + \underbrace{M_m}_{0.1} + \left(\underbrace{\Delta M_q}_{0.2} + \underbrace{\Delta M_g}_{0.2} \right)$$

$$\underbrace{\Delta M_q + \Delta M_g}_{0.2} = \frac{1}{4} \left\langle \int d^3 x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M - M_m)$$

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^{(0)} + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = \underbrace{M_q}_{0.4} + \underbrace{M_g}_{0.6} \qquad M_{q,g} = \left(\underbrace{A_{q,g}(0)}_{0.6, 0.4} + \underbrace{\bar{C}_{q,g}(0)}_{-0.18, 0.18} \right) M$$

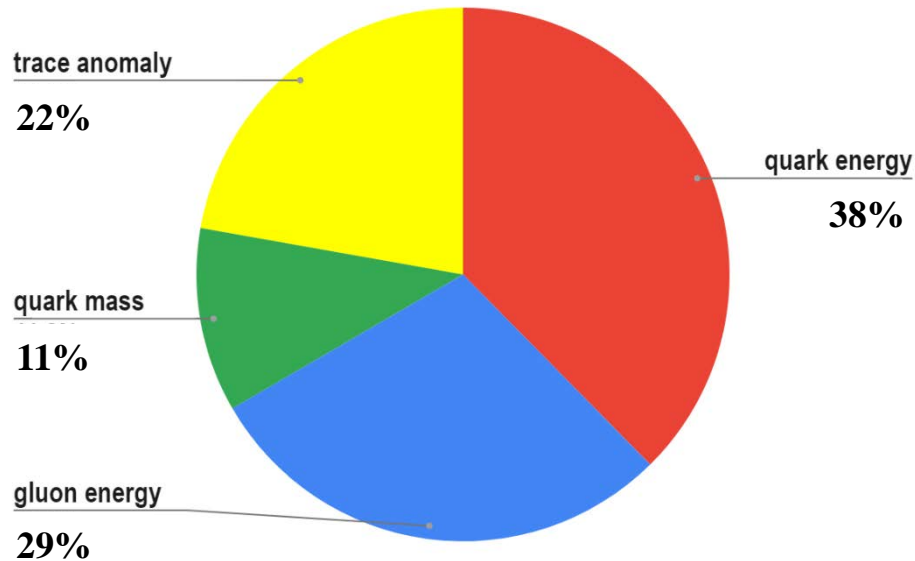
$$M = \left(\underbrace{M_q}_{0.3} - \underbrace{M_m}_{0.1} \right) + \underbrace{M_g}_{0.6} + \underbrace{M_m}_{0.1} \qquad M_m = \left\langle \int d^3 x m \bar{\psi} \psi \right\rangle = \frac{\sigma_{\pi N} + \sigma_s}{M} M$$

$$M = \left(\underbrace{M_q}_{0.4} - \underbrace{M_m}_{0.1} - \underbrace{\Delta M_q}_{0.2} \right) + \left(\underbrace{M_g}_{0.3} - \underbrace{\Delta M_g}_{0.2} \right) + \underbrace{M_m}_{0.1} + \left(\underbrace{\Delta M_q}_{0.2} + \underbrace{\Delta M_g}_{0.2} \right)$$

$$\underbrace{\Delta M_q + \Delta M_g}_{0.2} = \frac{1}{4} \left\langle \int d^3 x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M - M_m)$$

"Ji's decomposition" in NNLO QCD

proton mass



$$\mu = 1.3 \text{ GeV}$$

$$\langle \pi(p) | T_{q,g}^{00} | \pi(p) \rangle = 2A_{q,g}^\pi(0) p^0 p^0 + 2(M^\pi)^2 \bar{C}_{q,g}^\pi(0) \eta^{00}$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle \pi | \int d^3x T_{q,g}^{00} | \pi \rangle}{\langle \pi | \pi \rangle} = M^\pi \left(A_{q,g}^\pi(0) + \bar{C}_{q,g}^\pi(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3x \psi^\dagger (-i\mathbf{D} \cdot \boldsymbol{\alpha} + m\beta) \psi + \int d^3x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M^\pi = M_q^\pi + M_g^\pi \qquad M_{q,g}^\pi = \left(A_{q,g}^\pi(0) + \bar{C}_{q,g}^\pi(0) \right) M^\pi$$

$$M^\pi = \left(M_q^\pi - M_m^\pi \right) + M_g^\pi + M_m^\pi \qquad M_m^\pi = \left\langle \int d^3x m \bar{\psi} \psi \right\rangle$$

$$M^\pi = \left(M_q^\pi - M_m^\pi - \Delta M_q^\pi \right) + \left(M_g^\pi - \Delta M_g^\pi \right) + M_m^\pi + \left(\Delta M_q^\pi + \Delta M_g^\pi \right)$$

$$\Delta M_q^\pi + \Delta M_g^\pi = \frac{1}{4} \left\langle \int d^3x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M^\pi - M_m^\pi)$$

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$$M^\pi = M_q^\pi + M_g^\pi \qquad M_{q,g}^\pi = \left(A_{q,g}^\pi(0) + \bar{C}_{q,g}^\pi(0) \right) M^\pi$$

0.6, 0.4

$$M^\pi = \left(M_q^\pi - M_m^\pi \right) + M_g^\pi + M_m^\pi \qquad M_m^\pi = \left\langle \int d^3x m \bar{\psi} \psi \right\rangle$$

0.5

$$M^\pi = \left(M_q^\pi - M_m^\pi - \Delta M_q^\pi \right) + \left(M_g^\pi - \Delta M_g^\pi \right) + M_m^\pi + \left(\Delta M_q^\pi + \Delta M_g^\pi \right)$$

$$\Delta M_q^\pi + \Delta M_g^\pi = \frac{1}{4} \left\langle \int d^3x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M^\pi - M_m^\pi)$$

0.1

$$\langle \pi(p) | T_{q,g}^{00} | \pi(p) \rangle = 2A_{q,g}^\pi(0) p^0 p^0 + 2(M^\pi)^2 \bar{C}_{q,g}^\pi(0) \eta^{00}$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle \pi | \int d^3x T_{q,g}^{00} | \pi \rangle}{\langle \pi | \pi \rangle} = M^\pi \left(A_{q,g}^\pi(0) + \bar{C}_{q,g}^\pi(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3x \psi^\dagger (-i\mathbf{D} \cdot \boldsymbol{\alpha} + m\beta) \psi + \int d^3x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M^\pi = M_q^\pi + M_g^\pi \qquad M_{q,g}^\pi = \left(A_{q,g}^\pi(0) + \bar{C}_{q,g}^\pi(0) \right) M^\pi$$

0.6, 0.4 -0.04, 0.04

$$M^\pi = \left(M_q^\pi - M_m^\pi \right) + M_g^\pi + M_m^\pi \qquad M_m^\pi = \left\langle \int d^3x m \bar{\psi} \psi \right\rangle$$

0.5

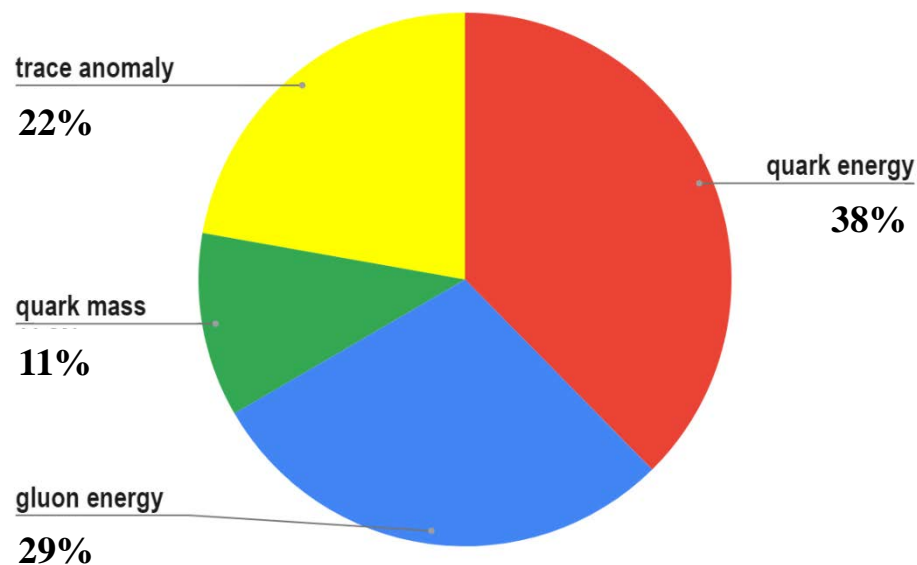
$$M^\pi = \left(M_q^\pi - M_m^\pi - \Delta M_q^\pi \right) + \left(M_g^\pi - \Delta M_g^\pi \right) + M_m^\pi + \left(\Delta M_q^\pi + \Delta M_g^\pi \right)$$

$$\Delta M_q^\pi + \Delta M_g^\pi = \frac{1}{4} \left\langle \int d^3x \left(\frac{\beta(g)}{2g} F^2 + \gamma_m(g) m \bar{\psi} \psi \right) \right\rangle = \frac{1}{4} (M^\pi - M_m^\pi)$$

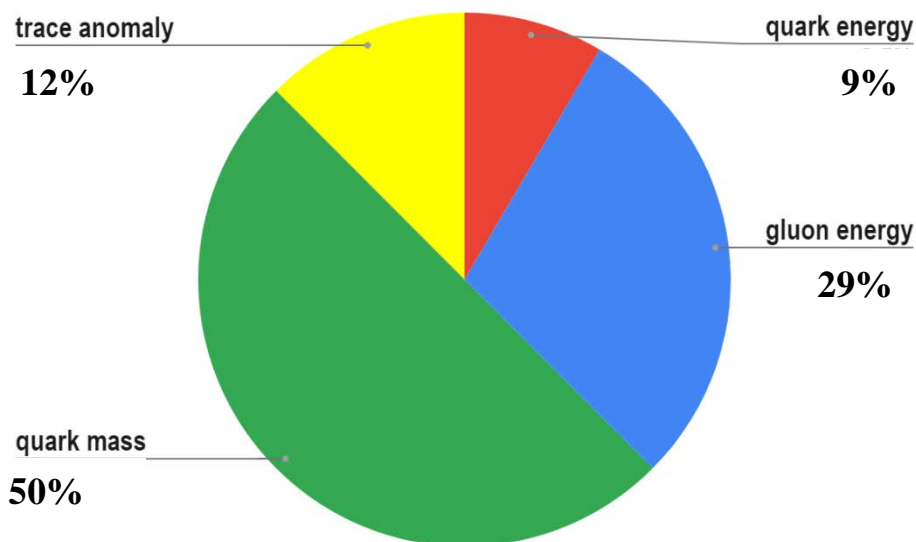
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"Ji's decomposition" in NNLO QCD

proton mass



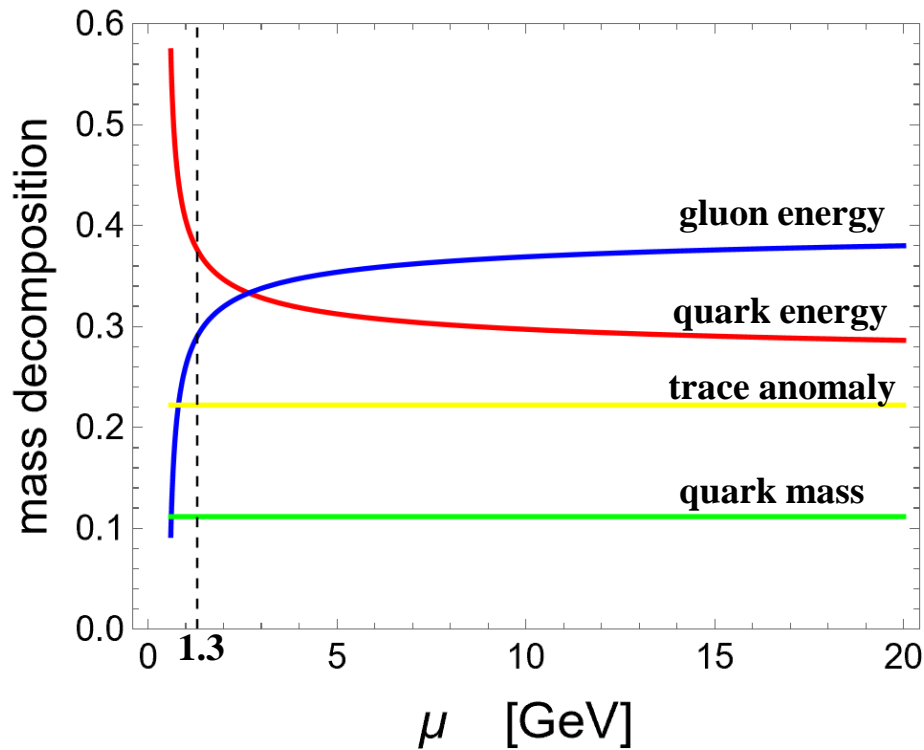
pion mass



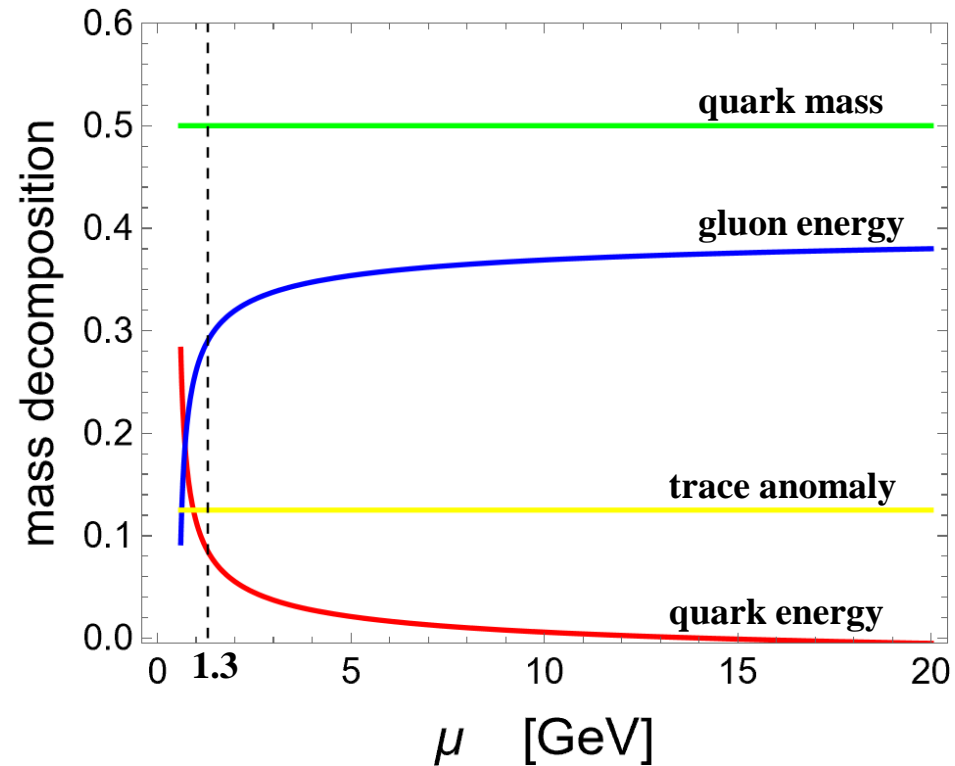
$$\mu = 1.3 \text{ GeV}$$

"Ji's decomposition" in NNLO QCD

proton mass



pion mass



$\overline{\text{MS}}$ scheme

$$\langle N(p) | T_{q,g}^{00} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(0)} p^{(0)} + \bar{C}_{q,g}(0) M \eta^{00} \right] u(p)$$

$$\langle \hat{H}_{q,g} \rangle = \frac{\langle N | \int d^3 x T_{q,g}^{00} | N \rangle}{\langle N | N \rangle} = M \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right)$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3 x \psi^\dagger (-i \mathbf{D} \cdot \boldsymbol{\alpha} + m \beta) \psi + \int d^3 x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2)$$

$$M = M_q + M_g$$

$$M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

Trace anomaly decomposition in NNLO QCD

$$\langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p) \left[A_{q,g}(0) \gamma^{(\mu} p^{\nu)} + \bar{C}_{q,g}(0) M \eta^{\mu\nu} \right] u(p)$$

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$2M^2 = \eta_{\mu\nu} \langle N | T^{\mu\nu} | N \rangle = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle$$

$$M = \tilde{M}_q + \tilde{M}_g \qquad \tilde{M}_{q,g} = \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right) M$$

Trace anomaly decomposition in NNLO QCD

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$$2M^2 = \eta_{\mu\nu} \langle N | T^{\mu\nu} | N \rangle = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle$$

$$M = \tilde{M}_q + \tilde{M}_g \qquad \tilde{M}_{q,g} = \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right) M$$

-0.1

1.1

0.6, 0.4

-0.18, 0.18

proton

0.4

0.6

0.6, 0.4

-0.04, 0.04

pion

$$\mu = 1.3 \text{ GeV}$$

Trace anomaly decomposition in NNLO QCD

$$\eta_{\mu\nu} \langle N(p) | T_{q,g}^{\mu\nu} | N(p) \rangle = 2M^2 \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right)$$

$$2M^2 = \eta_{\mu\nu} \langle N | T^{\mu\nu} | N \rangle = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle$$

$$M = \tilde{M}_q + \tilde{M}_g \qquad \tilde{M}_{q,g} = \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right) M$$

	-0.1	1.1		0.6, 0.4	-0.18, 0.18	
	0.4	0.6		0.6, 0.4	-0.04, 0.04	proton
						pion

$$\mu = 1.3 \text{ GeV}$$

$$\hat{H} = \hat{H}_q + \hat{H}_g = \int d^3x \psi^\dagger (-i\mathbf{D} \cdot \boldsymbol{\alpha} + m\beta) \psi + \int d^3x \frac{1}{2} (\mathbf{E}^2 + \mathbf{B}^2) \qquad \int d^3x T_{q,g}^{00} = \hat{H}_{q,g}$$

$$M = M_q + M_g \qquad M_{q,g} = \left(A_{q,g}(0) + \bar{C}_{q,g}(0) \right) M$$

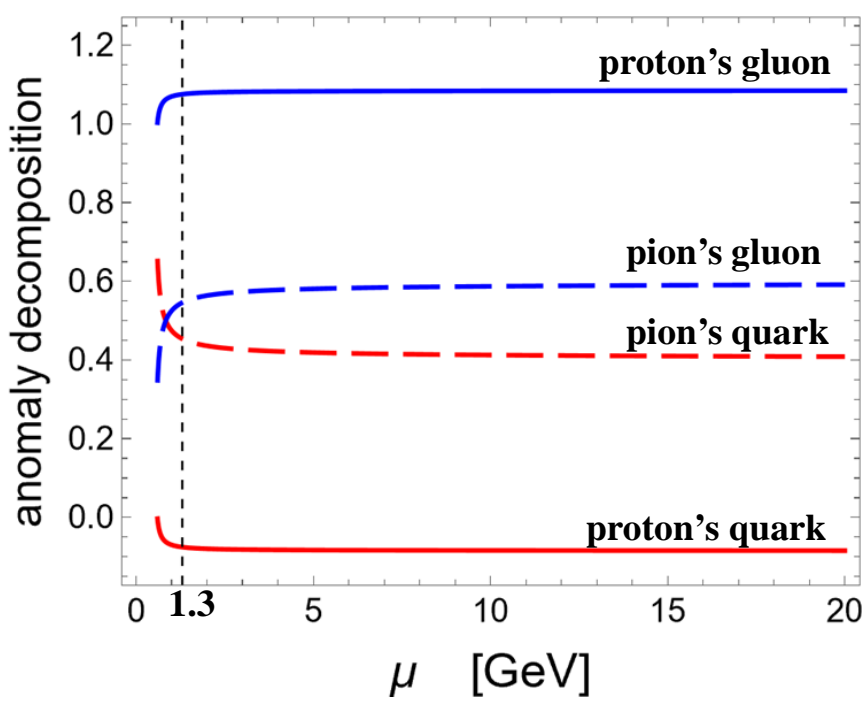
	0.4	0.6		0.6, 0.4	-0.18, 0.18	
	0.6	0.4		0.6, 0.4	-0.04, 0.04	proton
						pion

Trace anomaly decomposition in NNLO QCD

$$2M^2 = \eta_{\mu\nu} \langle N | T^{\mu\nu} | N \rangle = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle$$

$$M = \tilde{M}_q + \tilde{M}_g \quad \tilde{M}_{q,g} = \left(A_{q,g}(0) + 4\bar{C}_{q,g}(0) \right) M$$

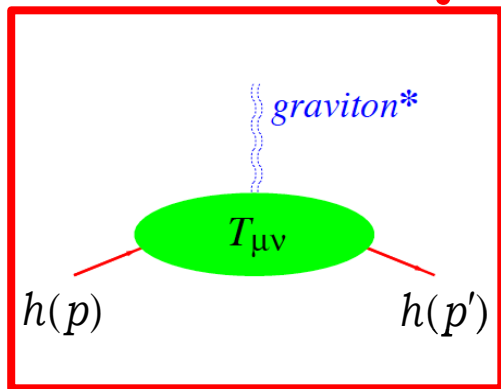
-0.1	1.1	0.6, 0.4	-0.18, 0.18	proton
0.4	0.6	0.6, 0.4	-0.04, 0.04	pion



$\overline{\text{MS}}$ scheme

Summary

Gravitational form factors can be accessed in hard processes @ JLab, EIC, ...



mass & energy distribution

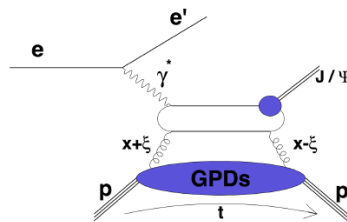
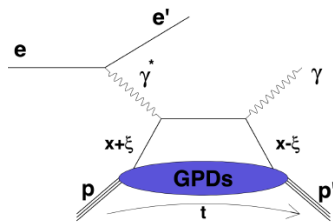
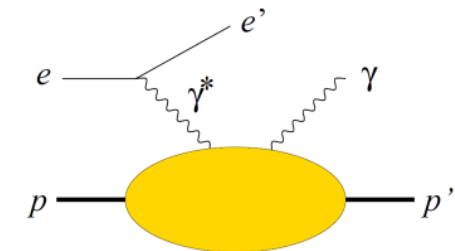
spin distribution

$$\langle N(p') | T_{q,g}^{\mu\nu} | N(p) \rangle = \bar{u}(p') \left[A_{q,g}(t) \gamma^{(\mu} P^{\nu)} + B_{q,g}(t) \frac{P^{(\mu} i\sigma^{\nu)\alpha} \Delta_\alpha}{2M} \right.$$

$$\left. + D_{q,g}(t) \frac{\Delta^\mu \Delta^\nu - g^{\mu\nu} \Delta^2}{M} + \bar{C}_{q,g}(t) M \eta^{\mu\nu} \right] u(p)$$

force & pressure distribution

mass & pressure distribution



$$\int_{-1}^1 dx x H^q(x, \eta, t) = A_q(t) + 4\eta^2 D_q(t)$$

$$\int_{-1}^1 dx x E^q(x, \eta, t) = B_q(t) - 4\eta^2 D_q(t)$$

$$\bar{C}_q(0, \mu) = -\bar{C}_g(0, \mu) = \text{LO} + \text{NLO} + \text{NNLO}$$

$$A_q(\mu_0) = \int_0^1 dx x [q(x, \mu_0) + \bar{q}(x, \mu_0)]$$

• trace anomaly for each q/g part of EMT

$$\langle N(p) | m \bar{\psi} \psi | N(p) \rangle$$

• NNLO term is $\sim 1\%$ level

$$\bar{C}_q^N(0, \mu = 1 \text{ GeV}) = -\bar{C}_g^N(0, \mu = 1 \text{ GeV}) = -0.180 \pm 0.003$$

NNLO QCD

$$\bar{C}_q^\pi(0, \mu = 1 \text{ GeV}) = -\bar{C}_g^\pi(0, \mu = 1 \text{ GeV}) = -0.04 \pm 0.02$$

NNLO QCD with NLO input

Using the NNLO values for $A_{q,g}(0), \bar{C}_{q,g}(0)$,

$$T^{00} = T_q^{00} + T_g^{00}$$

$$M = M_q + M_g$$

0.4	0.6
0.6	0.4

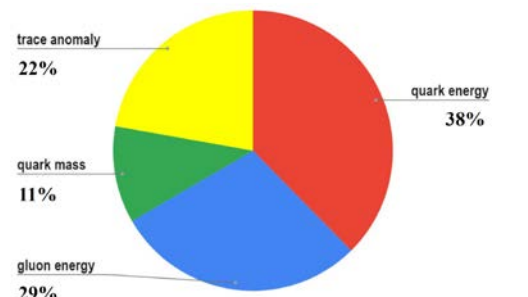
$$M_{q,g} = (A_{q,g}(0) + \bar{C}_{q,g}(0)) M$$

0.6, 0.4	-0.18, 0.18
0.6, 0.4	-0.04, 0.04

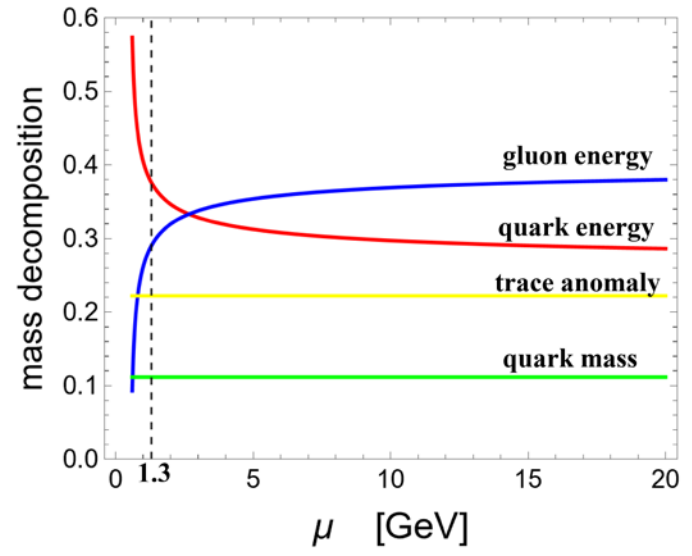
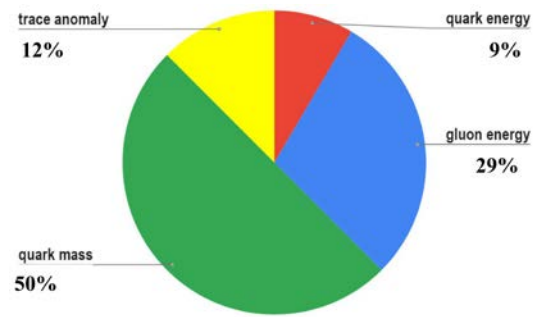
strong μ dep.

proton
pion

proton



pion



proton

$$\eta_{\mu\nu} T^{\mu\nu} = \eta_{\mu\nu} T_q^{\mu\nu} + \eta_{\mu\nu} T_g^{\mu\nu}$$

$$M = \tilde{M}_q + \tilde{M}_g$$

-0.1	1.1
0.4	0.6

$$\tilde{M}_{q,g} = (A_{q,g}(0) + 4\bar{C}_{q,g}(0)) M$$

0.6, 0.4	-0.18, 0.18
0.6, 0.4	-0.04, 0.04

weak μ dep.

proton
pion

Using the NNLO values for $A_{q,g}(0), \bar{C}_{q,g}(0),$

$$T^{00} = T_q^{00} + T_g^{00}$$

$$M = M_q + M_g$$

0.4	0.6
0.6	0.4

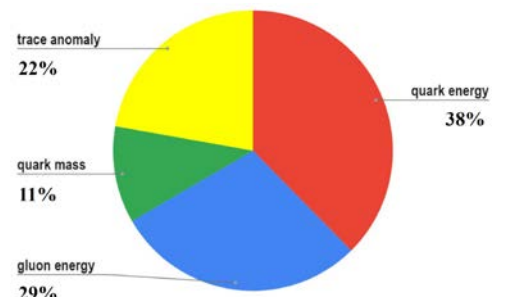
$$M_{q,g} = (A_{q,g}(0) + \bar{C}_{q,g}(0)) M$$

0.6, 0.4	-0.18, 0.18
0.6, 0.4	-0.04, 0.04

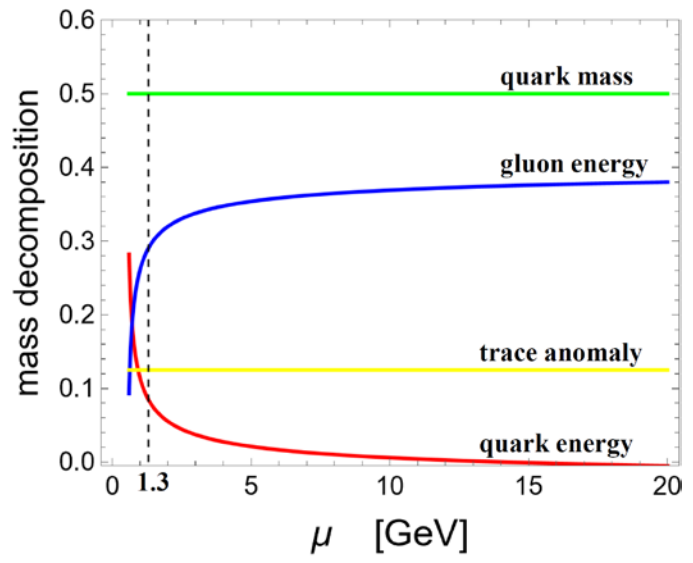
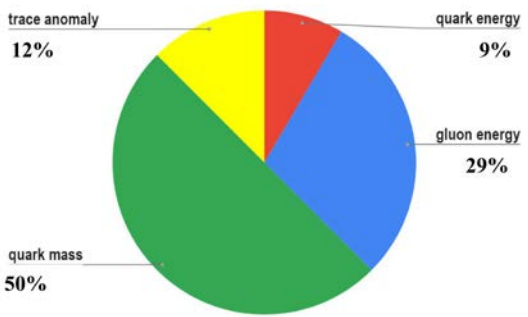
strong μ dep.

proton
pion

proton



pion



pion

$$\eta_{\mu\nu} T^{\mu\nu} = \eta_{\mu\nu} T_q^{\mu\nu} + \eta_{\mu\nu} T_g^{\mu\nu}$$

$$M = \tilde{M}_q + \tilde{M}_g$$

-0.1	1.1
0.4	0.6

$$\tilde{M}_{q,g} = (A_{q,g}(0) + 4\bar{C}_{q,g}(0)) M$$

0.6, 0.4	-0.18, 0.18
0.6, 0.4	-0.04, 0.04

weak μ dep.

proton
pion

backup

$$\begin{aligned}
\bar{C}_q(0, \mu) \Big|_{n_f=3} &= -0.145556 + 0.305556 \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} \\
&+ (0.09 - 0.25A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} \\
&+ \alpha_s(\mu) \left[0.00553609 + 0.0803962 \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} \right. \\
&+ \left. (0.0127684 - 0.0354678A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} - (0.0279651 - 0.0354678A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{31}{81}} \right] \\
&+ (\alpha_s(\mu))^2 \left[0.00174426 + 0.0312256 \frac{\langle N(p) | m\bar{\psi}\psi | N(p) \rangle}{2M^2} \right. \\
&- (0.0059729 - 0.0165914A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} \\
&- (0.00396745 - 0.00503187A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{31}{81}} \\
&+ \left. (0.0237481 - 0.0216233A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{112}{81}} \right]
\end{aligned}$$

$$\begin{aligned}
\bar{C}_q(0, \mu) \Big|_{n_f=3} &= -0.145556 + 0.305556 \frac{\langle N(p) | m \bar{\psi} \psi | N(p) \rangle}{2M^2} \\
&+ (0.09 - 0.25 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} \\
&+ \alpha_s(\mu) \left[0.00553609 + 0.0803962 \frac{\langle N(p) | m \bar{\psi} \psi | N(p) \rangle}{2M^2} \right. \\
&+ (0.0127684 - 0.0354678 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} - (0.0279651 - 0.0354678 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{31}{81}} \left. \right] \\
&+ (\alpha_s(\mu))^2 \left[0.00174426 + 0.0312256 \frac{\langle N(p) | m \bar{\psi} \psi | N(p) \rangle}{2M^2} \right. \\
&- (0.0059729 - 0.0165914 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{50}{81}} \\
&- (0.00396745 - 0.00503187 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{31}{81}} \\
&+ (0.0237481 - 0.0216233 A_q(\mu_0)) \left(\frac{\alpha_s(\mu)}{\alpha_s(\mu_0)} \right)^{\frac{112}{81}} \left. \right]
\end{aligned}$$

$$A_q(\mu_0) = \int_0^1 dx x [q(x, \mu_0) + \bar{q}(x, \mu_0)]$$

$$\langle N(p) | m \bar{\psi} \psi | N(p) \rangle$$

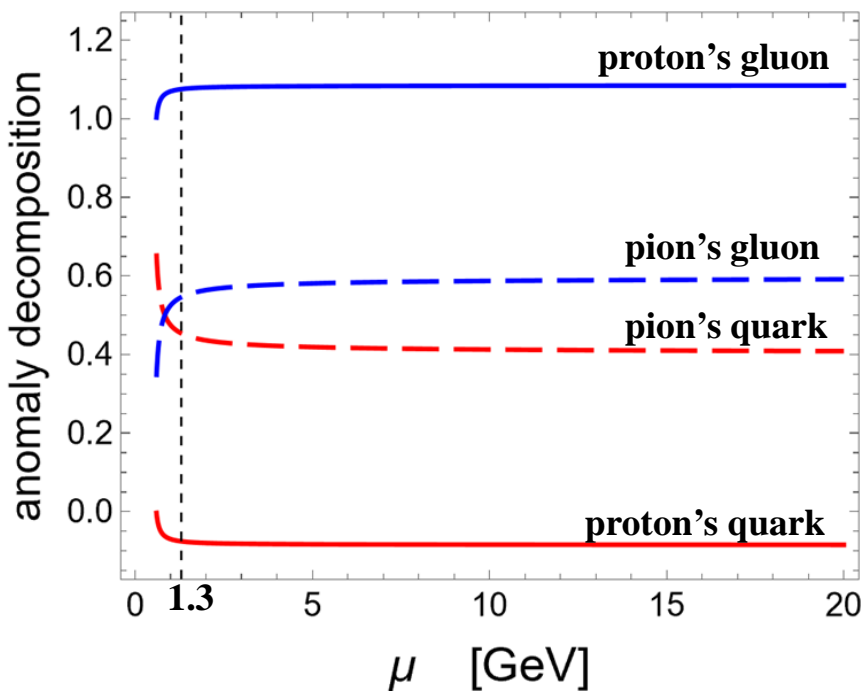
$$= \langle N(p) | m_u \bar{u} u + m_d \bar{d} d + m_s \bar{s} s | N(p) \rangle$$

Trace anomaly decomposition in NNLO QCD

$$2M^2 = \eta_{\mu\nu} \langle N | T^{\mu\nu} | N \rangle = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle$$

$$M = \tilde{M}_q + \tilde{M}_g$$

-0.1	1.1	$\tilde{M}_{q,g} = \left(\begin{matrix} A_{q,g}(0) & 4\bar{C}_{q,g}(0) \end{matrix} \right) M$	proton
0.4	0.6		



$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2$$

$$= T_q^{\mu\nu} + T_g^{\mu\nu}$$

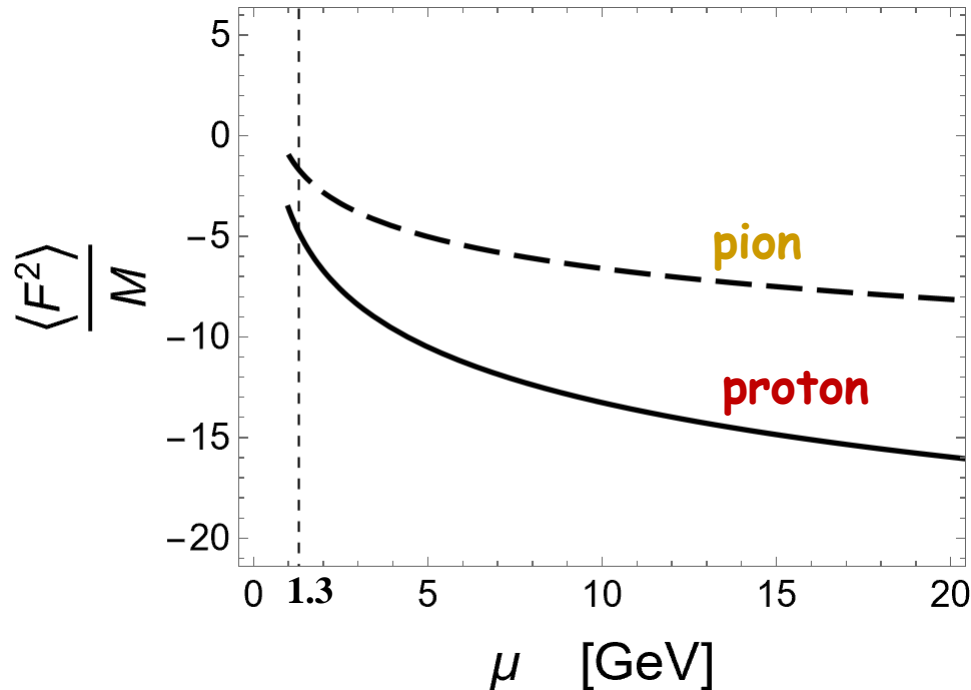
$$A_{q,g}(0) + 4\bar{C}_{q,g}(0) = \frac{\langle N(p) | \eta_{\mu\nu} T_{q,g}^{\mu\nu} | N(p) \rangle}{2M^2}$$

Hatta, Rajan, KT, JHEP 12 ('18) 008
 KT, JHEP 01 ('19) 120

$\overline{\text{MS}}$ scheme

Gluon field in NNLO QCD

$$\frac{\langle F^2 \rangle}{M} = \frac{2g}{\beta(g)} \left(1 - (1 + \gamma_m(g)) \frac{\langle m\bar{\psi}\psi \rangle}{M} \right) = \begin{cases} -4.8 & \text{proton} \\ -1.7 & \text{pion} \end{cases} \quad (\mu = 1.3 \text{ GeV})$$



$\overline{\text{MS}}$ scheme

$$T^{\mu\nu} = \frac{1}{2} \bar{\psi} \gamma^{(\mu} i \vec{D}^{\nu)} \psi + F^{\mu\rho} F_{\rho}{}^{\nu} + \frac{\eta^{\mu\nu}}{4} F^2$$

$$\langle p' | T_q^{\mu\nu} | p \rangle = \frac{1}{4} \langle p' | \bar{\psi} \left(-i \vec{\partial}^{\mu} + i \vec{\partial}^{\nu} + 2gA^{\mu} \right) \gamma^{\nu} \psi | p \rangle + (\mu \leftrightarrow \nu)$$

$$-i\partial^{\mu} \psi(x) = \left[\hat{P}^{\mu}, \psi(x) \right] \quad A^{\mu}(z^{-}) = \frac{1}{2} \int_{-\infty}^{\infty} dz'^{-} \text{sgn}(z'^{-} - z^{-}) F^{\mu+}(z'^{-})$$

intermediated states: "partonic"

$$A_q(t) + \eta^2 D_q(t) \simeq \langle x \rangle F_1^q(t), \quad B_q(t) - \eta^2 D_q(t) \simeq \langle x \rangle F_2^q(t)$$

$$\begin{aligned} & \frac{P^+ \eta \Delta_{\perp}^{\mu} \bar{u}(p') u(p)}{M} D_q(t) \simeq \langle N(p') | \bar{\psi} g A_{\perp}^{\mu} \gamma^+ \psi | N(p) \rangle \\ & = \frac{1}{2} \int_{-\infty}^{\infty} d\lambda \text{sgn}(\lambda) n_{\alpha} \langle N(p') | g F_a^{\mu\alpha}(\lambda n) \bar{\psi}(0) t^a \gamma^+ \psi(0) | N(p) \rangle \end{aligned}$$

$$A_q(t) + \eta^2 D_q(t) \simeq \langle x \rangle F_1^q(t), \quad B_q(t) - \eta^2 D_q(t) \simeq \langle x \rangle F_2^q(t)$$

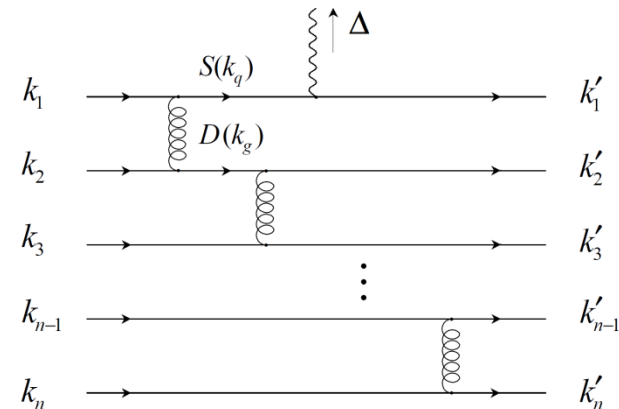
$$-\frac{P^+ \eta \Delta_{\perp}^{\mu} \bar{u}(p') u(p)}{M} D_q(t) \simeq \langle N(p') | \bar{\psi} g A_{\perp}^{\mu} \gamma^+ q \psi | N(p) \rangle$$

$$= \frac{1}{2} \int_{-\infty}^{\infty} d\lambda \operatorname{sgn}(\lambda) n_{\alpha} \langle N(p') | g F_a^{\mu\alpha}(\lambda n) \bar{\psi}(0) t^a \gamma^+ \psi(0) | N(p) \rangle$$

$t \rightarrow \infty$

$$A_q(t) \sim \frac{1}{t^2},$$

$$D_q(t) \sim \frac{1}{t^3}$$



$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

1-loop

$$\frac{\alpha_s}{4\pi} \frac{n_f}{3} F^2$$

$$\frac{\alpha_s}{4\pi} \left(-\frac{11C_A}{6} F^2 \right)$$

$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

1-loop

$$\frac{\alpha_s}{4\pi} \frac{n_f}{3} F^2$$

$$\frac{\alpha_s}{4\pi} \left(-\frac{11C_A}{6} F^2 \right)$$

2-loop

$$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2$$

$$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2$$

3-loop

$$\begin{aligned} & \left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\frac{52\zeta(3)}{9} - \frac{401}{324} \right) C_A C_F \right. \right. \\ & \left. \left. + \left(\frac{134}{27} - 4\zeta(3) \right) C_A^2 + \left(\frac{2407}{1458} - \frac{16\zeta(3)}{9} \right) C_F^2 \right\} \right. \\ & \left. + n_f^2 \left(-\frac{697C_A}{729} - \frac{169C_F}{1458} \right) \right] F^2 \end{aligned}$$

$$\begin{aligned} & \left(\frac{\alpha_s}{4\pi} \right)^3 \left[n_f \left(\left(\frac{1123}{162} - \frac{52\zeta(3)}{9} \right) C_A C_F \right. \right. \\ & \left. \left. + \left(4\zeta(3) + \frac{293}{36} \right) C_A^2 + \frac{16}{729} (81\zeta(3) - 98) C_F^2 \right) \right. \\ & \left. + n_f^2 \left(\frac{655C_A}{2916} - \frac{361C_F}{729} \right) - \frac{2857C_A^3}{108} \right] F^2 \end{aligned}$$

$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

1-loop

$$\frac{\alpha_s}{4\pi} \frac{n_f}{3} F^2$$

$$\frac{\alpha_s}{4\pi} \left(-\frac{11C_A}{6} F^2 \right)$$

2-loop

$$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2$$

$$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2$$

3-loop

$$\begin{aligned} & \left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\frac{52\zeta(3)}{9} - \frac{401}{324} \right) C_A C_F \right. \right. \\ & \left. \left. + \left(\frac{134}{27} - 4\zeta(3) \right) C_A^2 + \left(\frac{2407}{1458} - \frac{16\zeta(3)}{9} \right) C_F^2 \right\} \right. \\ & \left. + n_f^2 \left(-\frac{697C_A}{729} - \frac{169C_F}{1458} \right) \right] F^2 \end{aligned}$$

$$\begin{aligned} & \left(\frac{\alpha_s}{4\pi} \right)^3 \left[n_f \left(\left(\frac{1123}{162} - \frac{52\zeta(3)}{9} \right) C_A C_F \right. \right. \\ & \left. \left. + \left(4\zeta(3) + \frac{293}{36} \right) C_A^2 + \frac{16}{729} (81\zeta(3) - 98) C_F^2 \right) \right. \\ & \left. + n_f^2 \left(\frac{655C_A}{2916} - \frac{361C_F}{729} \right) - \frac{2857C_A^3}{108} \right] F^2 \end{aligned}$$

nucleon

-1

:

5

$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

1-loop	$\frac{\alpha_s}{4\pi} \frac{n_f}{3} F^2$	$\frac{\alpha_s}{4\pi} \left(-\frac{11C_A}{6} F^2 \right)$
2-loop	$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2$	$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2$
3-loop	$\left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\frac{52\zeta(3)}{9} - \frac{401}{324} \right) C_A C_F \right. \right.$ $\left. + \left(\frac{134}{27} - 4\zeta(3) \right) C_A^2 + \left(\frac{2407}{1458} - \frac{16\zeta(3)}{9} \right) C_F^2 \right\}$ $\left. + n_f^2 \left(-\frac{697C_A}{729} - \frac{169C_F}{1458} \right) \right] F^2$	$\left(\frac{\alpha_s}{4\pi} \right)^3 \left[n_f \left(\left(\frac{1123}{162} - \frac{52\zeta(3)}{9} \right) C_A C_F \right. \right.$ $\left. + \left(4\zeta(3) + \frac{293}{36} \right) C_A^2 + \frac{16}{729} (81\zeta(3) - 98) C_F^2 \right]$ $\left. + n_f^2 \left(\frac{655C_A}{2916} - \frac{361C_F}{729} \right) - \frac{2857C_A^3}{108} \right] F^2$
nucleon	-1	5

$$2m_\pi^2 = \eta_{\mu\nu} \langle \pi | T_q^{\mu\nu} | \pi \rangle + \eta_{\mu\nu} \langle \pi | T_g^{\mu\nu} | \pi \rangle$$

$$2M^2 = \langle N | \left(\frac{\beta(g)}{2g} F^2 + (1 + \gamma_m(g)) m \bar{\psi} \psi \right) | N \rangle \simeq \langle N | \frac{\beta(g)}{2g} F^2 | N \rangle$$

$$2M^2 = \eta_{\mu\nu} \langle N | T_q^{\mu\nu} | N \rangle + \eta_{\mu\nu} \langle N | T_g^{\mu\nu} | N \rangle$$

1-loop	$\frac{\alpha_s}{4\pi} \frac{n_f}{3} F^2$	$\frac{\alpha_s}{4\pi} \left(-\frac{11C_A}{6} F^2 \right)$
2-loop	$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{17C_A n_f}{27} + \frac{49C_F n_f}{54} \right) F^2$	$\left(\frac{\alpha_s}{4\pi} \right)^2 \left(\frac{28C_A n_f}{27} - \frac{17C_A^2}{3} + \frac{5C_F n_f}{54} \right) F^2$
3-loop	$\left(\frac{\alpha_s}{4\pi} \right)^3 \left[\left\{ n_f \left(\frac{52\zeta(3)}{9} - \frac{401}{324} \right) C_A C_F \right. \right.$ $\left. + \left(\frac{134}{27} - 4\zeta(3) \right) C_A^2 + \left(\frac{2407}{1458} - \frac{16\zeta(3)}{9} \right) C_F^2 \right\}$ $\left. + n_f^2 \left(-\frac{697C_A}{729} - \frac{169C_F}{1458} \right) \right] F^2$	$\left(\frac{\alpha_s}{4\pi} \right)^3 \left[n_f \left(\left(\frac{1123}{162} - \frac{52\zeta(3)}{9} \right) C_A C_F \right. \right.$ $\left. + \left(4\zeta(3) + \frac{293}{36} \right) C_A^2 + \frac{16}{729} (81\zeta(3) - 98) C_F^2 \right]$ $\left. + n_f^2 \left(\frac{655C_A}{2916} - \frac{361C_F}{729} \right) - \frac{2857C_A^3}{108} \right] F^2$
nucleon	-1	5
pion	1	1

$$2m_\pi^2 = \eta_{\mu\nu} \langle \pi | T_q^{\mu\nu} | \pi \rangle + \eta_{\mu\nu} \langle \pi | T_g^{\mu\nu} | \pi \rangle$$