



LHCspin: a polarized fixed-target experiment at the LHC



L. L. Pappalardo (pappalardo@fe.infn.it)

- World top energies
- high luminosity
- p and ion beams
- highly sophisticated state-of-the-art detectors

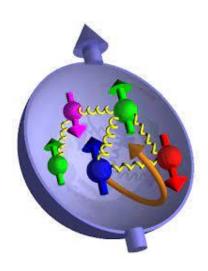


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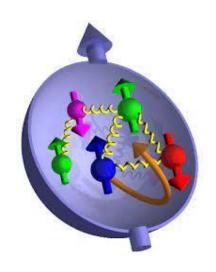


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- > ...which must not affect (degrade) the mainstream collider operation (beam quality, machine vacuum, detectors performance, background, etc.)

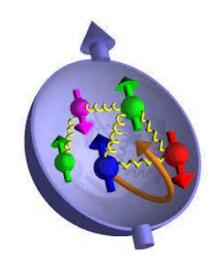
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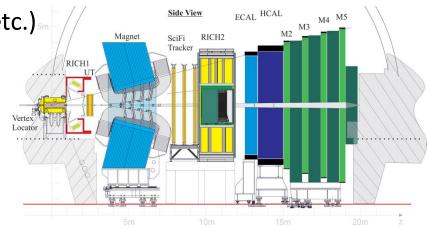
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- ...and the LHCb spectrometer is the perfect place to host it.

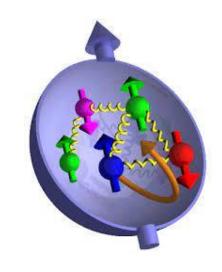


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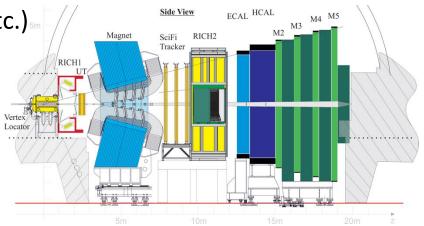




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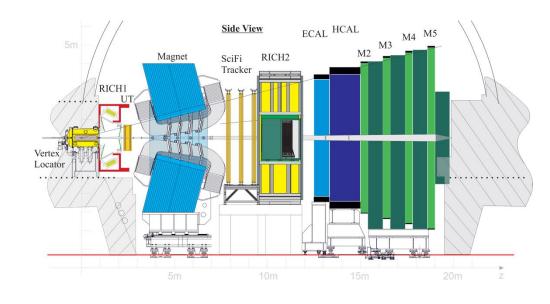






LHCb is a **general-purpose single-arm spectrometer**, fully instrumented in $2 < \eta < 5$ and optimised for detection of charmed and beauty hadrons

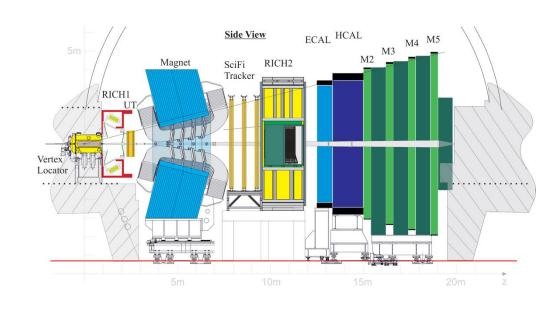
[JINST 3 (2008) S08005] [IJMPA 30 (2015)1530022]



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- Excellent particle identification and momentum resolution: $\sigma_p/p \leq 1.0 \% \ (p \in [2,200] \ {\rm GeV})$
- Precise primary and secondary vertex reconstruction (VELO)
- ➤ During LS2 major hardware upgrade to cope with the factor of 5 increase in luminosity starting from the Run 3 [http://arxiv.org/abs/2305.10515]

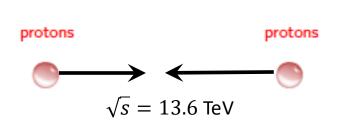


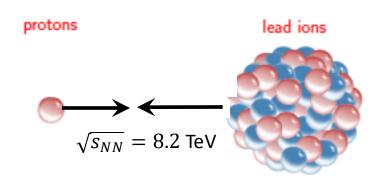
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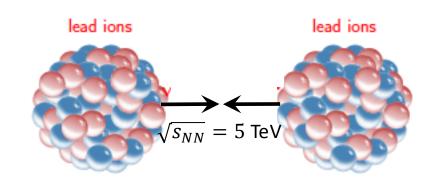
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Types of collisions (Collider mode):



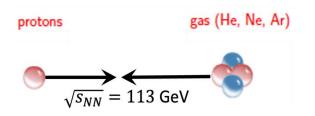


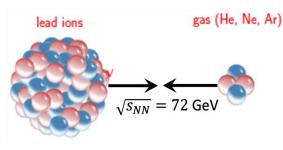


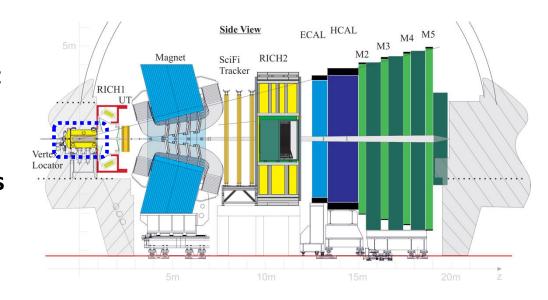
ECAL HCAL

RICH1

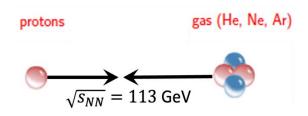
- ➤ Since 2015 LHCb can also be operated as a **fixed-target experiment** with the **SMOG system**, by injecting low pressure noble gases (He, Ne, Ar) into the VELO vessel.
- ➤ Unique opportunity to study pA/AA collisions on various targets exploiting the high-energy, high-intensity LHC beams!

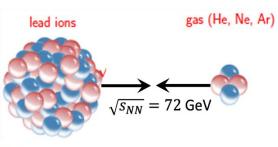


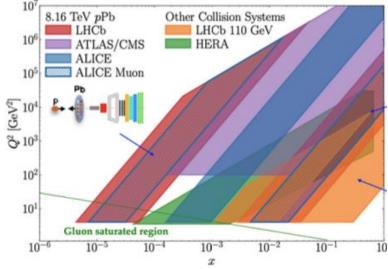




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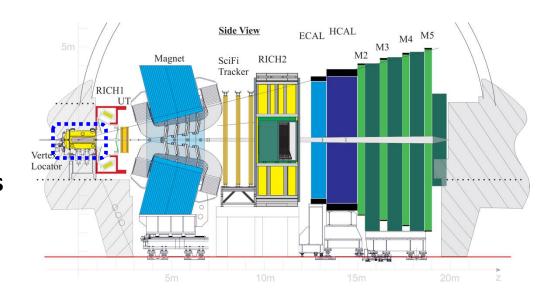






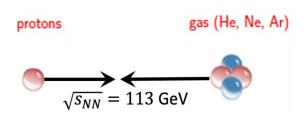
FT kinem. with $E_p = 6.8 \text{ TeV}$:

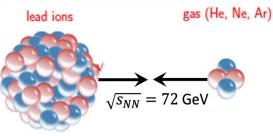
- $\sqrt{s_{NN}} \approx 115 \ GeV$
- $-3.0 \le y_{CM} \le 0$
- $x_F < 0$
- intermediate-large x_B
- intermediate Q^2

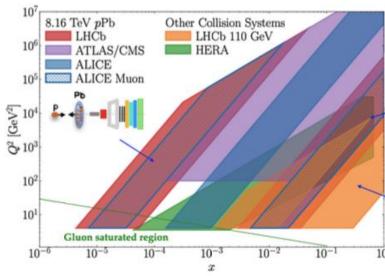


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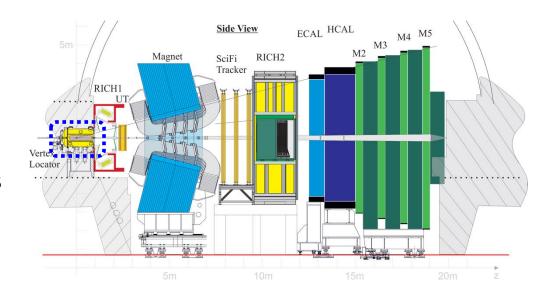


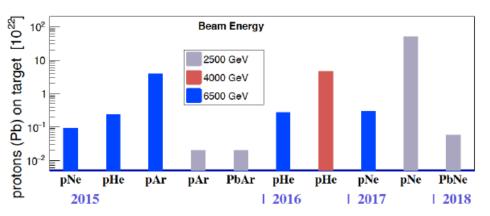




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Many interesting published analyses:

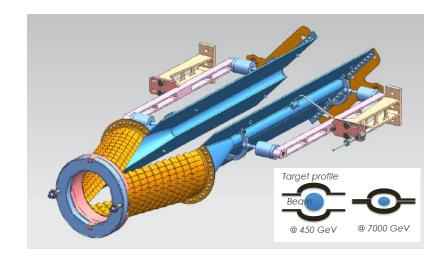
- Antiproton prod. cross section in p-He
- Charmonium production in p-Ne and Pb-Ne
- Open charm production in p-Ne and Pb-Ne

The SMOG2 upgrade

[SMOG2 TDR]



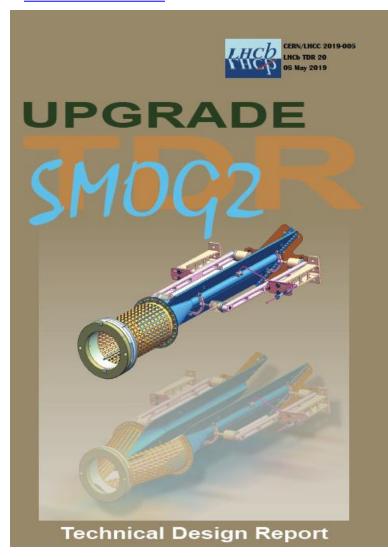
- > 20 cm storage cell for the target gas installed upstream of the VELO
- Brand new, more flexible and sophisticated Gas Feed System (GFS)



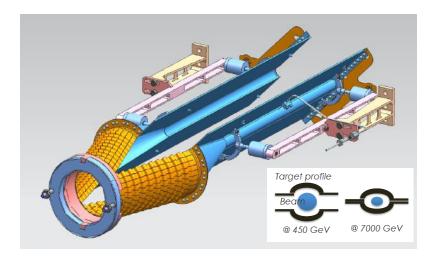


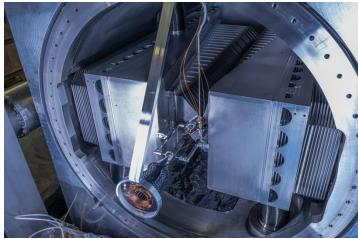
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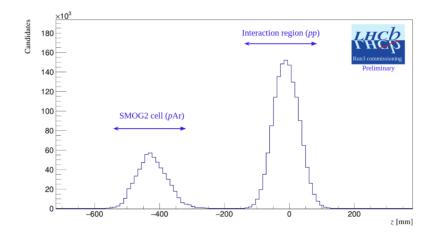


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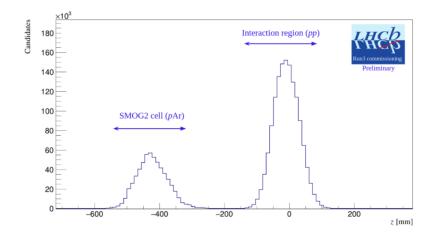




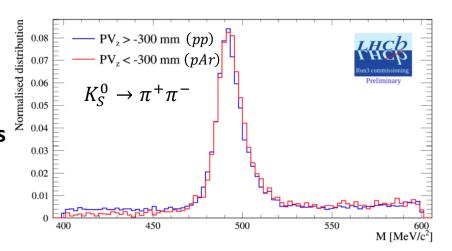
- more gas species: H_2 , D_2 , He, N_2 , O_2 , Ne, Ar (Kr and Xe to be tested)
- target density increased by large factor (up to 30)
- precise density (luminosity) determination
- negligible impact on LHC and LHCb performance
- can run un parallel with collider mode!



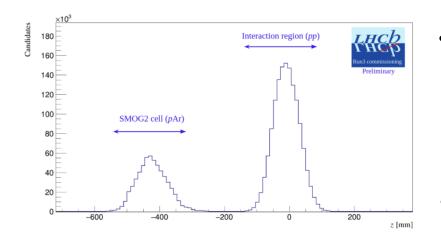
- FT and collider interaction regions are very well separated!
- LHCb is now the first (unique) LHC experiment with two simultaneous interaction regions!



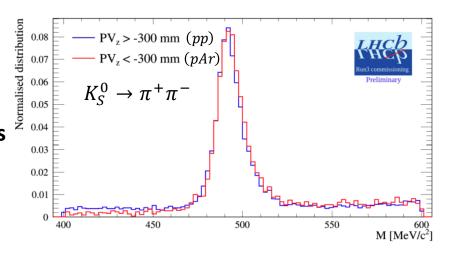
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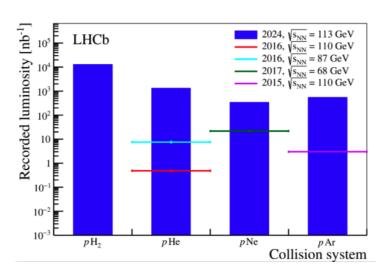
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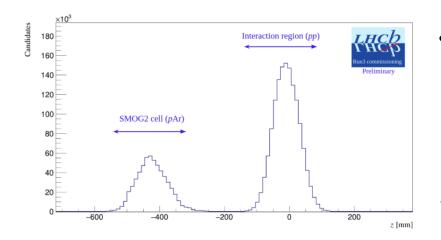


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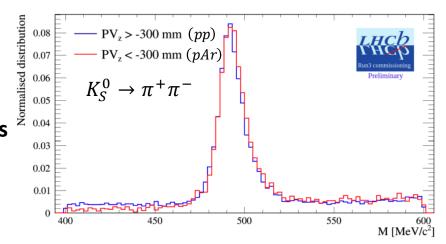


- SMOG2 is performing very well!
- Large beam-gas samples already collected in 2024 with all available gases!

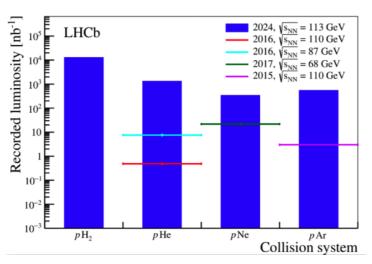


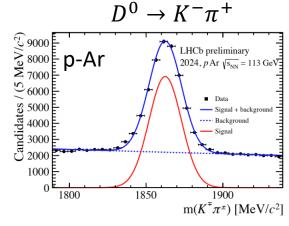


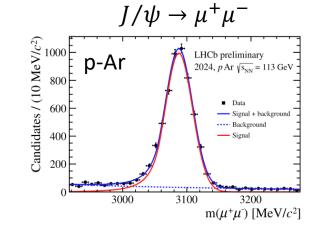
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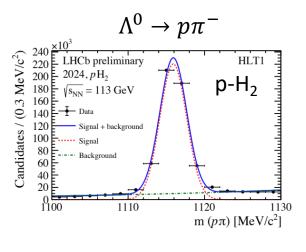


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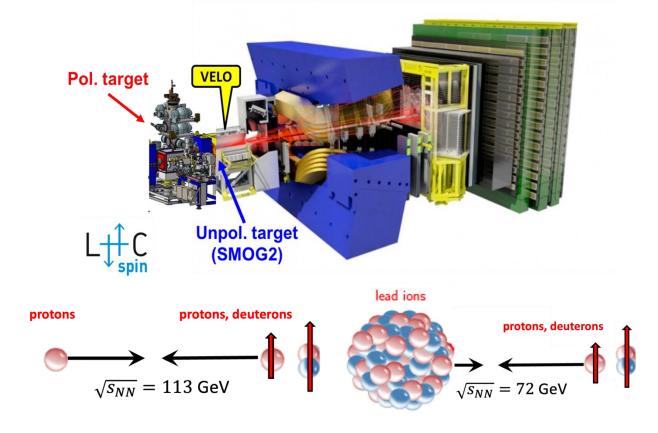






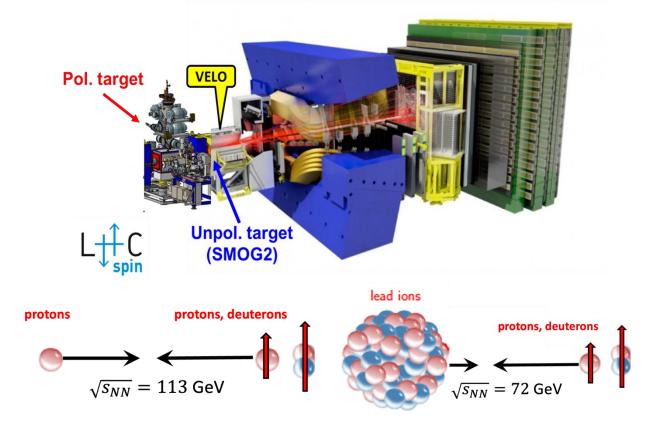
The LHCspin project

The LHCspin project represents the natural evolution of SMOG2 and will allow for the first time to perform spin physics measurements at the LHC through the implementation of a new-generation polarized gaseous target in the LHCb spectrometer.



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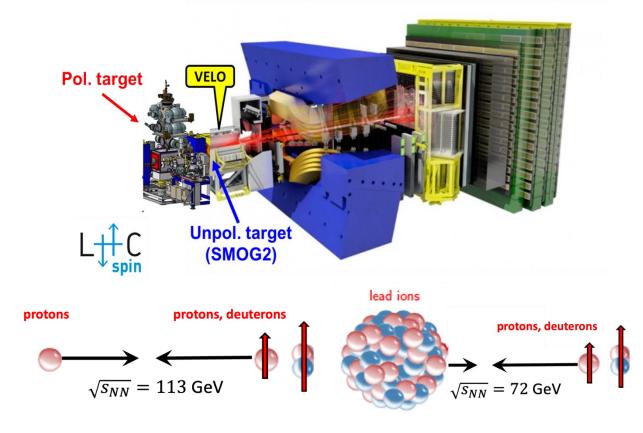


Physics Motivations

Study multi-dimensional nucleon structure at unique kinematic conditions (backward CM region, poorly explored large-x region at intermediate Q^2)

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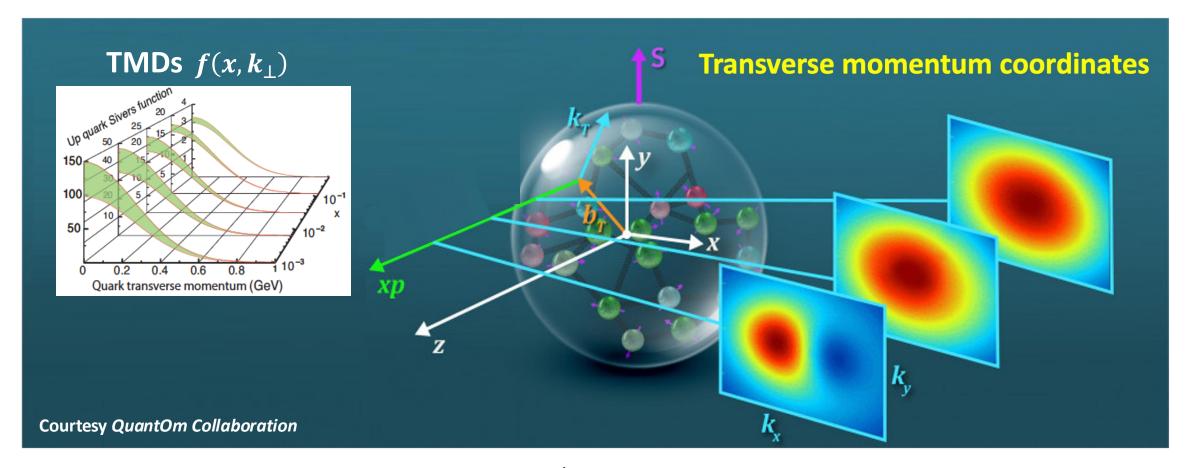
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Points of strenght

- ✓ use of well-established polarized gas target technology (HERMES @ DESY, ANKE @ COSY,...)
- ✓ marginal impact on LHC beam lifetime and LHCb mainstream physics program and performances
- ✓ can run in parallel with collider mode (well displaced interaction regions)
- ✓ can benefit from both protons and heavy-ion beams
- ✓ allows also injection of non-polarized gases (a-la SMOG2): H_2 , D_2 , He, N_2 , O_2 , Ne, Ar, Kr, Xe
- ✓ broad and unique physics program (next slides).

Nucleon tomography in momentum space: TMDs



- Describe **spin-orbit correlations of the form** $\vec{S} \cdot (\vec{p}_1 \times \vec{p}_2)$
- generate distorsions of the parton densities in transverse momentum plane (e.g. Sivers effect)
- can provide sensitivity to unknown parton OAM!

Quark and gluon TMDs

quark pol.

nucleon pol.

	U	L	Т
U	f_1		h_1^{\perp}
L		g_{1L}	h_{1L}^{\perp}
Т	f_{1T}^{\perp}	g_{1T}	$oldsymbol{h_1}, oldsymbol{h_{1T}}$

- 8 independent quark TMDs at leading-twist
- significant experimental progress in the last 20 years!
- main results from SIDIS (HERMES, COMPASS, JLAB, → EIC)
- many phenomenological extractions available from global analyses

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L		g_{1L}^g	$h_{1L}^{\perp g}$
Т	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$

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- Similar notation, but important differences!
- different naïve-time-reversal properties
- Experimental access still very limited!

	T-even	T-odd
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g	$h_1^{\perp g}$	h_1^g

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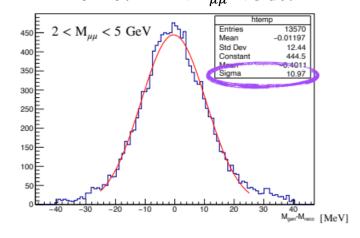
Polarized hadronic collisions with LHCspin offer a complementary approach

- Measure experimental observables sensitive to both quarks and gluons TMDs
- Make use of new probes (charmed and beauty mesons)
- Test non-trivial process dependence of quarks and (especially) gluons TMDs

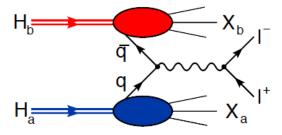
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Expected mass resolution $\sim 10~MeV$ in $2 < M_{\mu\mu} < 5~GeV$



Transv. polarized Drell-Yan



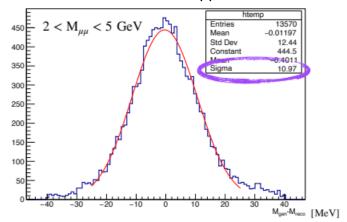
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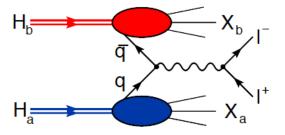
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Sensitive to unpol. and BM TMDs

$$d\sigma_{UU}^{DY} \propto f_1^{\bar{q}} \otimes f_1^q + \cos 2\phi \ h_1^{\perp,\bar{q}} \otimes h_1^{\perp,q}$$

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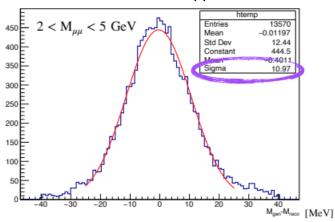
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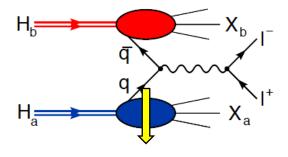
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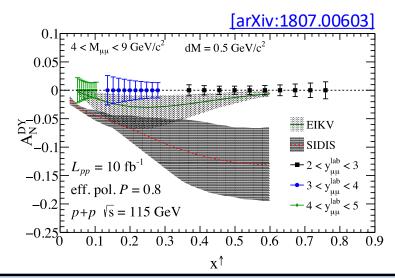
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Sensitive to quark TMDs through TSSAs

$$A_{N}^{DY} = \frac{1}{P} \frac{\sigma_{DY}^{\uparrow} - \sigma_{DY}^{\downarrow}}{\sigma_{DY}^{\uparrow} + \sigma_{DY}^{\downarrow}} \implies A_{UT}^{\sin\phi_{S}} \sim \frac{f_{1}^{q} \otimes f_{1T}^{\perp q}}{f_{1}^{q} \otimes f_{1}^{q}}, \quad A_{UT}^{\sin(2\phi - \phi_{S})} \sim \frac{h_{1}^{\perp q} \otimes h_{1}^{q}}{f_{1}^{q} \otimes f_{1}^{q}}, \dots$$



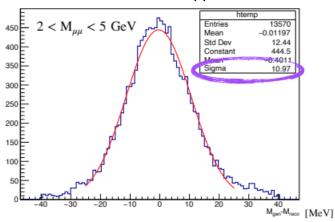
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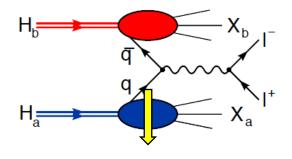
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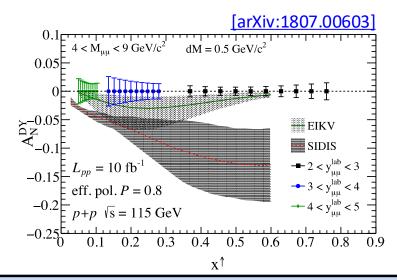
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- Extraction of qTMDs does not require knowledge of FF
- Verify sign change of Sivers func. wrt SIDIS

$$f_{1T}^{\perp}\big|_{DY} = -f_{1T}^{\perp}\big|_{SIDIS}$$

29

Test flavour sensitivity using both H and D targets

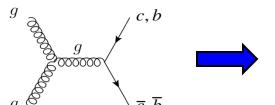
Gluon TMDs

nucleon pol.

gluon pol.

	U	Circularly	Linearly
U	$m{f}_1^g$		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
Т	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:



the most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

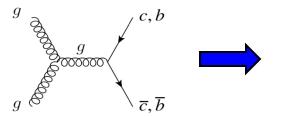
Gluon TMDs

gluon pol.

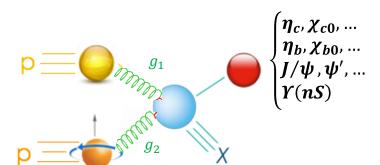
		U	C
pol.	U	f_1^g	
nucleon pol	L		
nnc	Т	$f_{1T}^{\perp g}$	

	U	Circularly	Linearly
U	$\boldsymbol{f_1^g}$		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
T	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$

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Polarized gTMDs can be accessed through TSSAs in **inclusive** heavy meson production

$$A_N = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \propto \left[f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \right] \sin \phi_S + \cdots$$

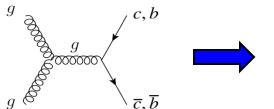
Gluon TMDs

nucleon pol

gluon pol.

UCircularlyLinearlyU f_1^g $h_1^{\perp g}$ L g_{1L}^g $h_{1L}^{\perp g}$ T $f_{1T}^{\perp g}$ g_{1T}^g h_1^g , $h_{1T}^{\perp g}$

In high-energy hadron collisions, heavy quarks are dominantly produced through gg fusion:

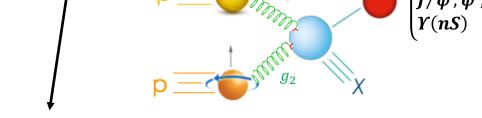


 η_c, χ_{c0}, \dots

the most efficient way to access the gluon dynamics inside the proton at LHC is to **measure heavy-quark observables**

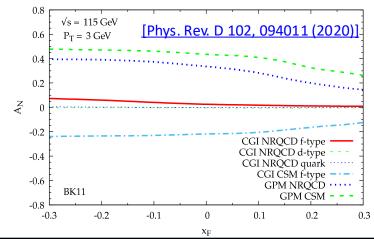
Polarized gTMDs can be accessed through TSSAs in **inclusive heavy meson production**

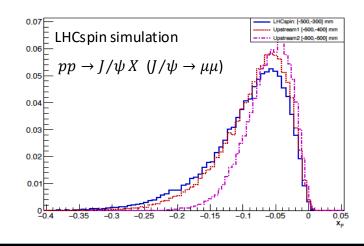
$$A_N = \frac{1}{P} \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}} \propto \left[f_{1T}^{\perp g}(x_a, k_{\perp a}) \otimes f_g(x_b, k_{\perp b}) \right] \sin \phi_S + \cdots$$



Gluon Sivers function:

- Sheds light on spin-orbit correlations of unpol. gluons inside a transv. pol. proton
- is sensitive to gluon OAM





TMD factorization requires $q_T(Q) \ll M_Q$. Can look at **associate quarkonia production**, where only the relative q_T needs to be small, e.g.: $pp^{(\uparrow)} \rightarrow J/\psi + J/\psi + X$

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$$\frac{\mathrm{d}\sigma}{\mathrm{d}M_{\mathcal{Q}\mathcal{Q}}\mathrm{d}Y_{\mathcal{Q}\mathcal{Q}}\mathrm{d}^2\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\mathrm{d}\Omega} = \frac{\sqrt{M_{\mathcal{Q}\mathcal{Q}}^2 - 4M_{\mathcal{Q}}^2}}{(2\pi)^28s\,M_{\mathcal{Q}\mathcal{Q}}^2}$$

$$\left\{ F_1(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\text{CS}}) \, \mathcal{C}\left[f_1^g f_1^g\right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) + F_2(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\text{CS}}) \, \mathcal{C}\left[w_2 h_1^{\perp g} h_1^{\perp g}\right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) \right\}$$

$$+ \left(F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right]\!\left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3'(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3' h_1^{\perp g} f_1^g\right]\!\left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) \right) \cos 2\phi_{\mathrm{CS}} \\ \left. + F_4(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_4 h_1^{\perp g} h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) \cos 4\phi_{\mathrm{CS}}\right) \right\} \\ \left. + \left(F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \right. \\ \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \left(x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}}) \,\mathcal{C}\!\left[w_3 f_1^g h_1^{\perp g}\right] \right) + \left. F_3(M_{\mathcal{Q}\mathcal{Q}, \theta_{\mathrm{CS}$$

gluon pol.

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0.3 0.25 $1/Q \cdot dQ/dp \frac{di \cdot J/\psi}{Q} = 0.15$ 0.15 0.1

0.05

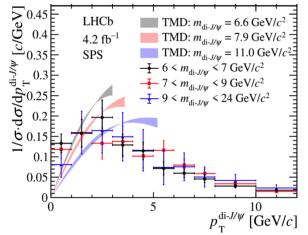
LHCb

 4.2 fb^{-1}

gluon pol.

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U	f_1^g		$h_1^{\perp g}$
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 $p_{\mathrm{T}}^{\mathrm{di}\text{-}J/\psi}$ [GeV/c]



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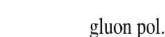
$$\left\{F_1(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\left[f_1^g f_1^g\right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}^T}) + F_2(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\mathrm{CS}}) \,\mathcal{C}\left[w_2 h_1^{\perp g} h_1^{\perp g}\right](x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}^T})\right\}$$

$$+\left(F_3(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\!\left[w_3f_1^gh_1^{\perp g}\right]\!\left(x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right)+\right.\left.F_3'(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\!\left[w_3'h_1^{\perp g}f_1^g\right]\!\left(x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right)\right)\cos2\phi_{\mathrm{CS}}\right.\\\left.+F_4(M_{\mathcal{Q}\mathcal{Q}},\theta_{\mathrm{CS}})\,\mathcal{C}\!\left[w_4h_1^{\perp g}h_1^{\perp g}\right]\left(x_{1,2},\boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}\right)\cos4\phi_{\mathrm{CS}}\right\}$$

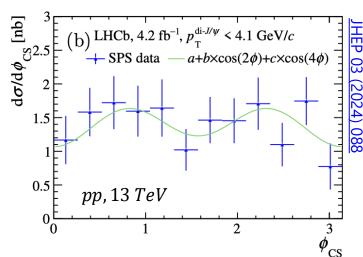
 $\int_{\phi/f^{-1}}^{0.25} dp/\omega p \cdot \rho/I$

LHCb

 4.2 fb^{-1}



	U	Circularly	Linearly
U	f_1^g		$h_1^{\perp g}$
L		g_{1L}^g	$h_{1L}^{\perp g}$
Т	$f_{1T}^{\perp g}$	g_{1T}^g	$h_1^g,h_{1T}^{\perp g}$



$$\left\langle \cos 2\phi_{\text{CS}} + F_4(M_{\mathcal{Q}\mathcal{Q}}, \theta_{\text{CS}}) \, \mathcal{C}\left[w_4 h_1^{\perp g} h_1^{\perp g}\right] (x_{1,2}, \boldsymbol{P}_{\mathcal{Q}\mathcal{Q}_T}) \cos 4\phi_{\text{CS}} \right\rangle$$

$$\left\langle \cos 2\phi_{\text{CS}} \right\rangle = -0.029 \pm 0.050 \, \left(\text{stat}\right) \pm 0.009 \, \left(\text{syst}\right)$$

 $p_{\mathrm{T}}^{\mathrm{di}\text{-}J/\psi}$ [GeV/c]

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LHCb

 4.2 fb^{-1}

SPS

TMD: $m_{\text{di-}J/\psi} = 7.9 \text{ GeV}/c^2$

TMD: $m_{\text{di-}J/\psi} = 11.0 \text{ GeV}/c^2$

 $p_{\mathrm{T}}^{\mathrm{di}\text{-}J/\psi} \stackrel{10}{[\mathrm{GeV}/c]}$

- $\frac{1}{1}$ 6 < $m_{\text{di-}J/\psi}$ < 7 GeV/ c^2

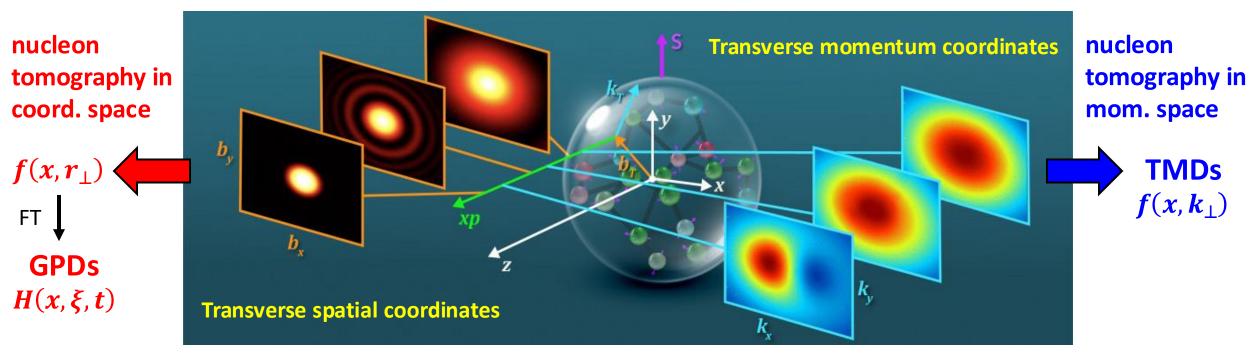
 $\frac{1}{1} 7 < m_{\text{di-J/}\psi} < 9 \text{ GeV/}c^2$ $\frac{1}{1} 9 < m_{\text{di-J/}\psi} < 24 \text{ GeV/}c^2$

azimuthal amplitudes consistent with zero

 $\langle \cos 4\phi_{\rm CS} \rangle = -0.087 \pm 0.052 \text{ (stat)} \pm 0.013 \text{ (syst)}$

- a few-% asymmetry cannot be excluded
- uncertainties statistically dominated
- But very challenging at fixed-target kinematics

GPDs: a complementary approach to the nucleon tomography



GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \ ilde{m{H}_T}$

Courtesy QuantOm Collaboration

U nucleon pol. f_1 g_{1L} f_{1T}^{\perp} g_{1T}

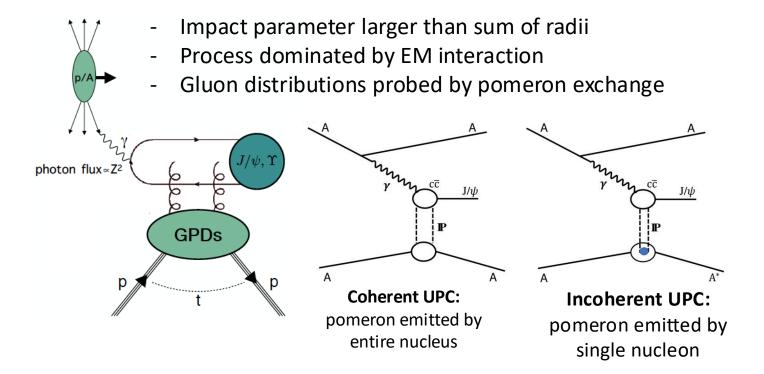
quark	nol
quain	por.
1	

L. L. Pappalardo

GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \ ilde{m{H}_T}$

Gluon GPDs can be accessed at LHC in **Ultra-Peripheral collisions (UPC)** where a quasi-real photon is emitted by the relativistic beam particle [PRD 85 (2012), 051502]

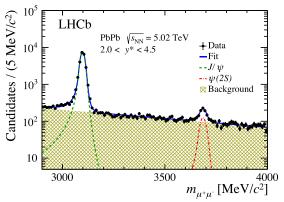
At LHC energies, these photons are energetic enough to trigger the production of hard dileptons and charmonia and bottomonia.



GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \; ilde{H}_T$

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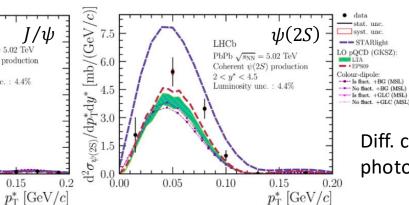
PbPb $\sqrt{s_{\rm NN}} = 5.02 \text{ TeV}$

Coherent J/ψ production

0.15

J. High Energ. Phys. 2023

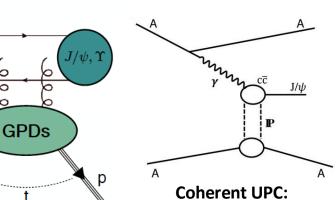
Pb-Pb collisions $\sqrt{s_{NN}} = 5.02 \, TeV$



Impact parameter larger than sum of radii

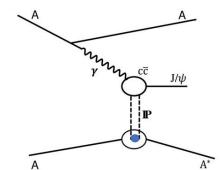
Process dominated by EM interaction

Gluon distributions probed by pomeron exchange



pomeron emitted by

entire nucleus



Incoherent UPC: pomeron emitted by single nucleon

Diff. cross section vs. p_T for coherent J/ψ and $\psi(2S)$ photoproduction, compared with models

0.05

0.10

0.00

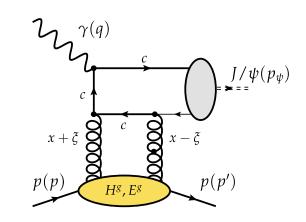
 $d^2\sigma_{J/\psi}/dp_T^*dy^* [mb/(GeV/c)]$

photon flux∝Z²

GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \ ilde{ ilde{H}}_T$

With LHCspin photo-production of J/ψ in polarized UPC of proton (or lead) beams with H^\uparrow target can be studied, providing constraints to the essentially unknown gluon GPD E_q which plays a crucial role in the Ji sum rule:

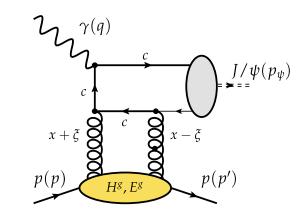
$$J^{g} = \frac{1}{2} \int_{0}^{1} dx \Big(H^{g}(x, \xi, 0) + E^{g}(x, \xi, 0) \Big)$$



GPD	$oldsymbol{U}$	L	T
$oldsymbol{U}$	H		\mathcal{E}_T
L		$ ilde{H}$	$ ilde{E}_T$
T	E	$ ilde{E}$	$H_T, \ ilde{m{H}_T}$

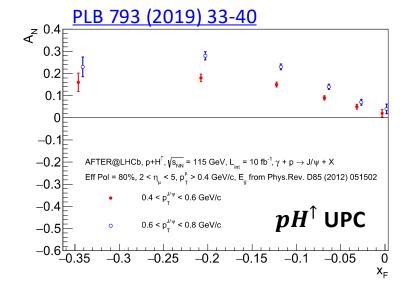
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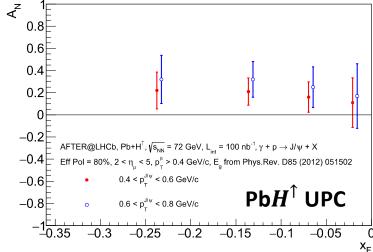
$$J^{g} = \frac{1}{2} \int_{0}^{1} dx \Big(H^{g}(x, \xi, 0) + E^{g}(x, \xi, 0) \Big)$$



$$A_{N} = \frac{\sigma^{h_{A}h_{B}^{\downarrow}} - \sigma^{h_{A}h_{B}^{\uparrow}}}{\sigma^{h_{A}h_{B}^{\downarrow}} + \sigma^{h_{A}h_{B}^{\uparrow}}} = \frac{\int dk \frac{dn_{A}}{dk} A_{N}^{\gamma} \sigma^{\gamma h_{B}}}{\int dk \left[\frac{dn_{A}}{dk} \sigma^{\gamma h_{B}} + \frac{dn_{B}}{dk} \sigma^{\gamma h_{A}}\right]}$$

The hadronic STSA A_N can be parametrized in terms of the photonic STSA A_N^γ which incorporates the GPDs H^g and E^g through their gluonic CFFs \mathcal{H}^g and \mathcal{E}^g

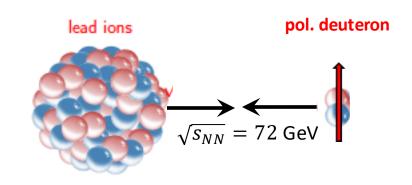




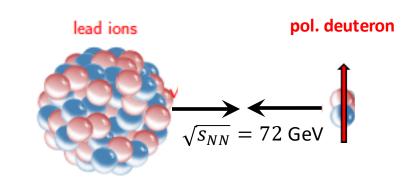
- Extraction based on models for the GPD H^g (Goloskokov-Kroll) and E^g (PRD 85, 051502 (2012))
- AFTER model-dependent predictions very promising for pH^{\uparrow} UPC

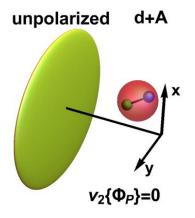
L. L. Pappalardo

- probe collective phenomena in heavy-light systems through ultrarelativistic collisions of heavy nuclei with trasv. pol. deuterons
- polarized light target nuclei offer a unique opportunity to control the orientation of the formed fireball by measuring the elliptic flow relative to the polarization axis (ellipticity).



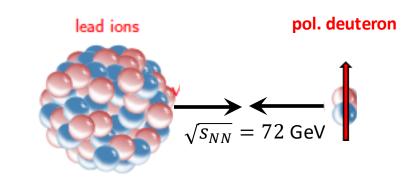
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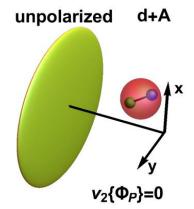




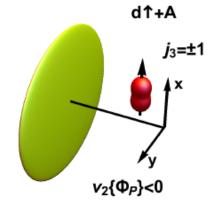
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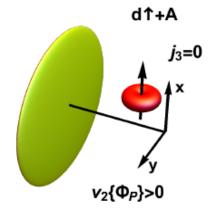




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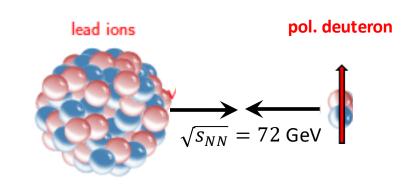


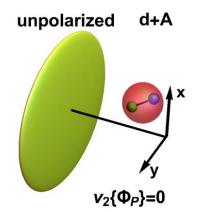
 $j_3 = \pm 1 \rightarrow \text{prolate fireball}$ stretched along the pol. axis, corresponds to $v_2 < 0$



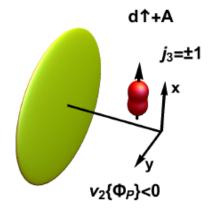
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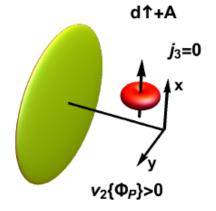




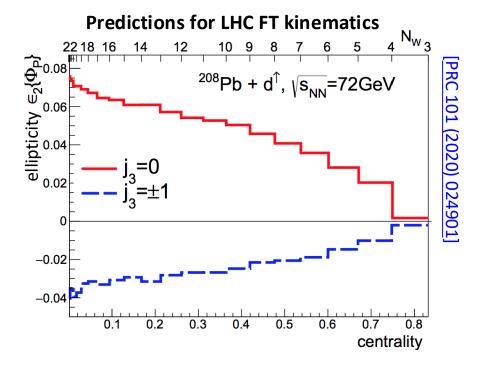
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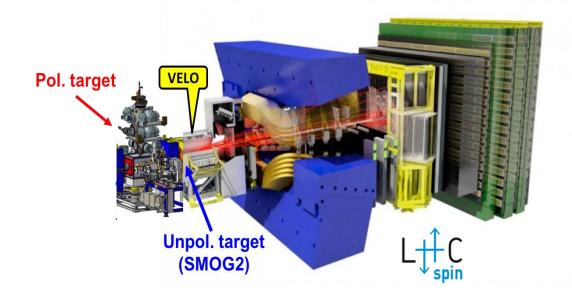
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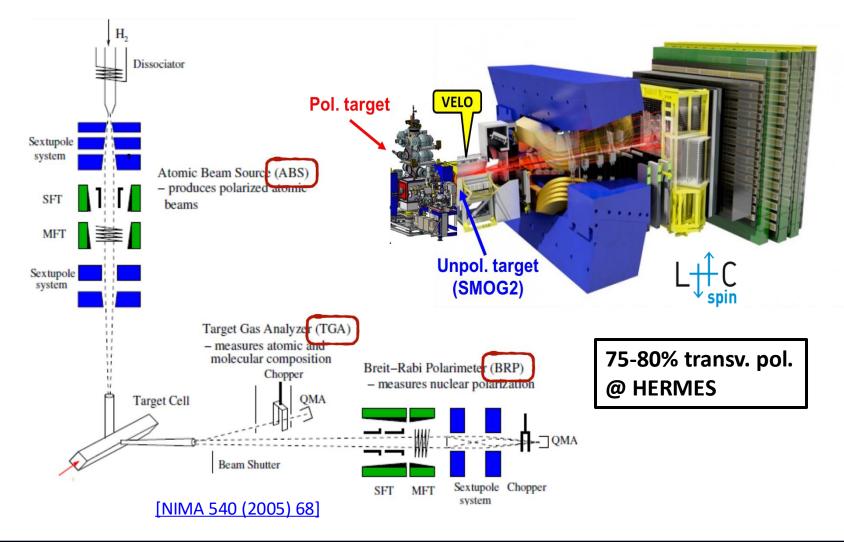
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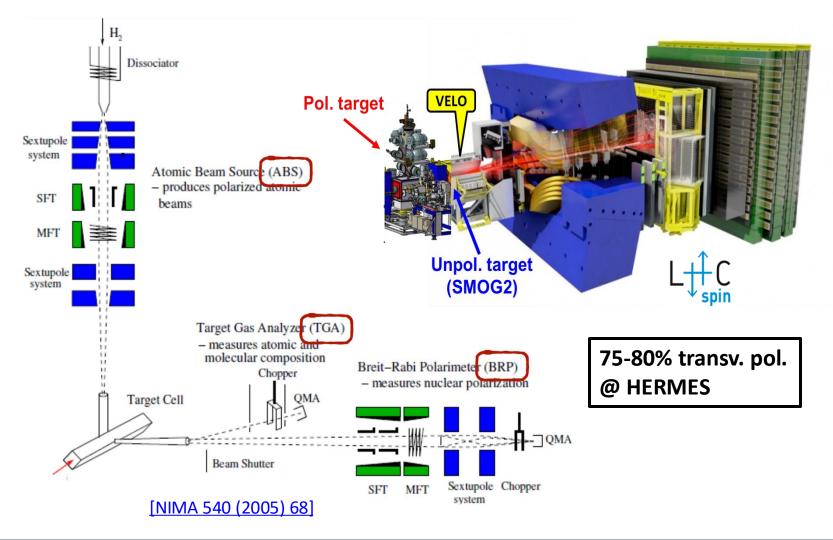
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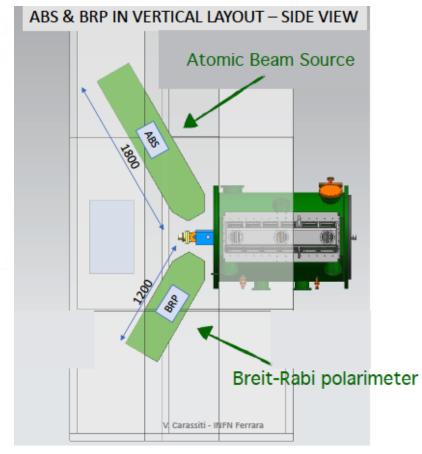


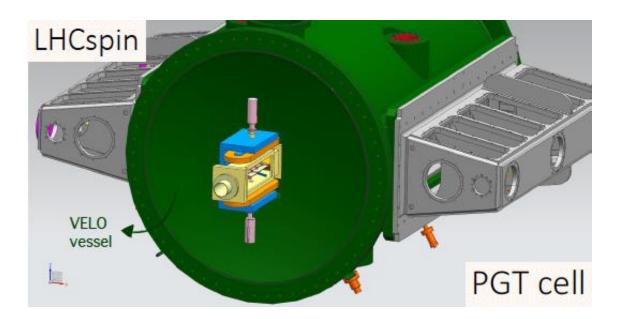
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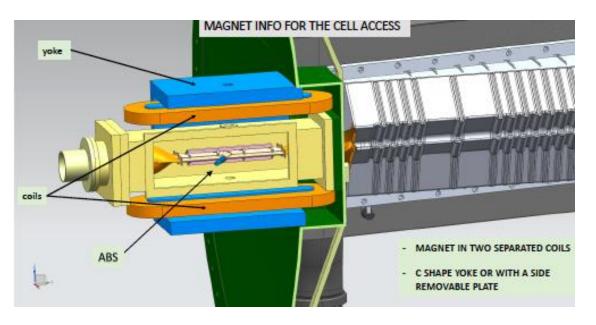


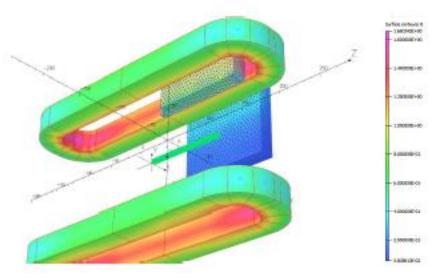
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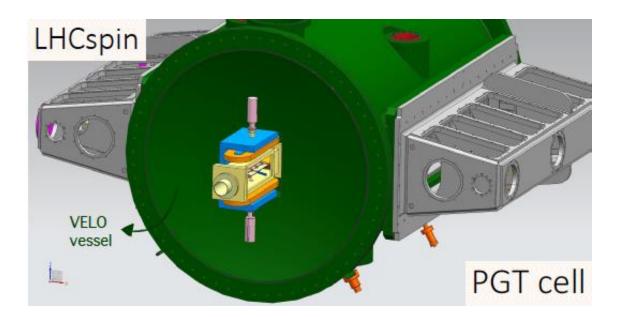


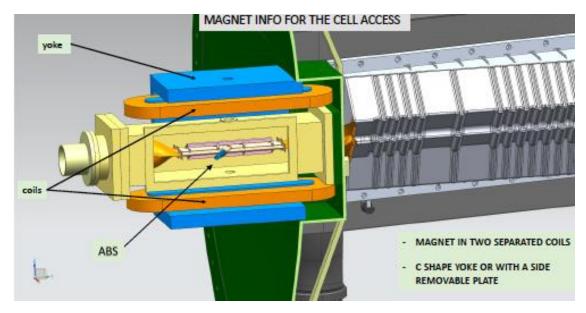


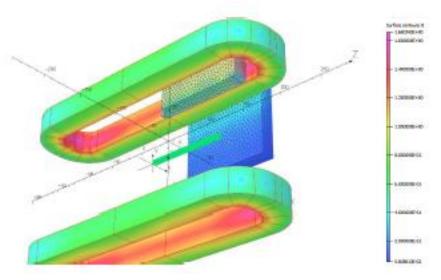




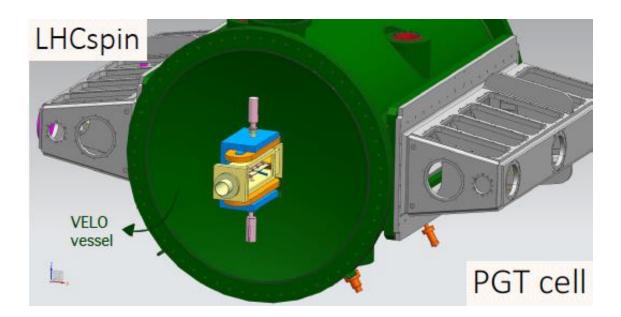
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- Required B=300~mT with $\Delta B/B \sim 10\%$

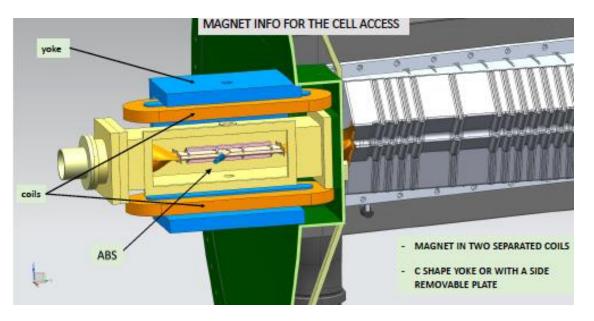


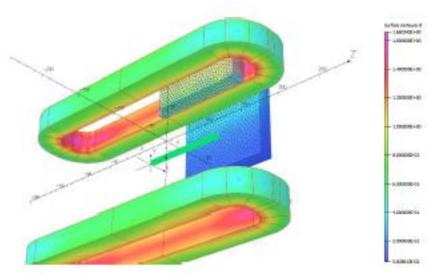




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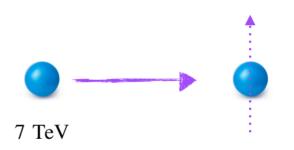






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- Possibility to switch from dipole magnet to solenoid to realize a Longitudinal polarized target in a future phase

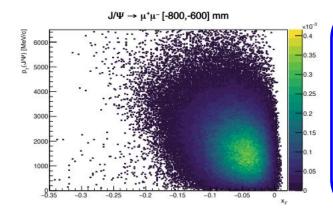
Kinematic coverage



$$\sqrt{s} = \sqrt{2m_N E_p} = 115 \text{ GeV}$$

Using
$$x_F = 2E_T/\sqrt{s_{NN}} \sinh(y^*)$$

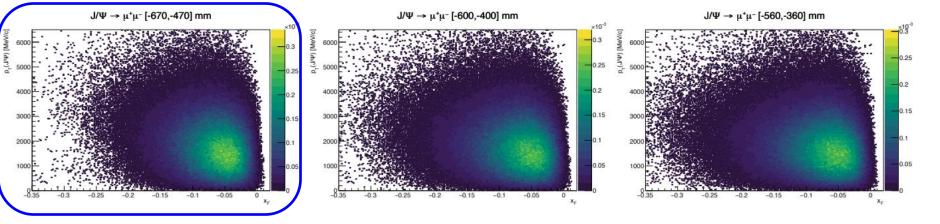
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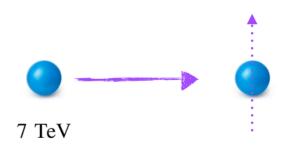
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- x_B/x_F ranges shrink towards larger backward distances
- SMOG2 cell position: [-560,-360] mm
- Probable position of LHCspin cell: [-670,-470] mm



Kinematic coverage

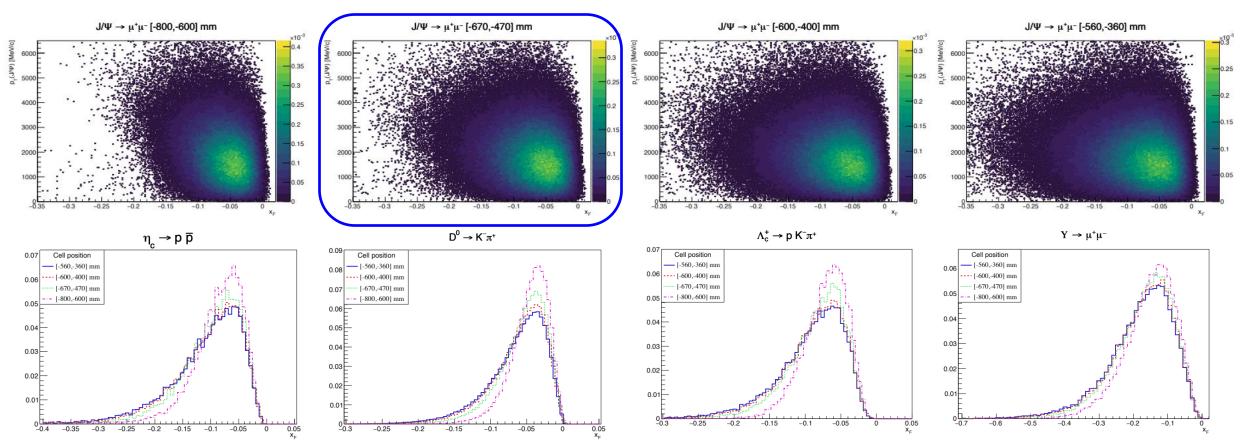


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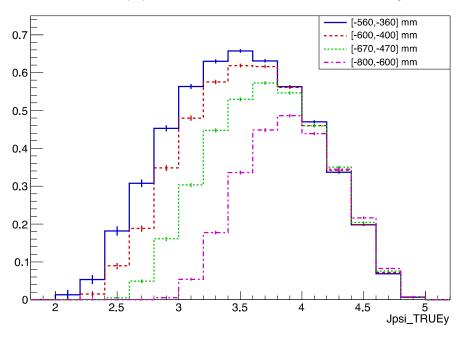
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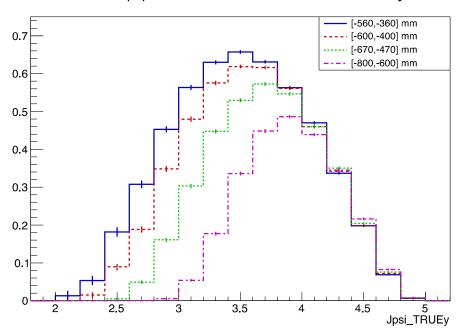
Expected performance

 $J/\Psi \rightarrow \mu^{+}\mu^{-}PV X$ track reconstruction efficiency



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Target

•
$$I_0 = 6.5 \cdot 10^{16} s^{-1}$$
 (HERMES)

•
$$C_{tot} = 17.4 \text{ l/s} (20 \text{ cm cell})$$

•
$$\theta$$
= 3.7 ·10¹³ atoms/cm²

Beam (Run4)

- $2.2 \cdot 10^{11}$ p/bunch
- 2760 bunches

$$I_{beam} = 6.8 \cdot 10^{18} \ p/s$$

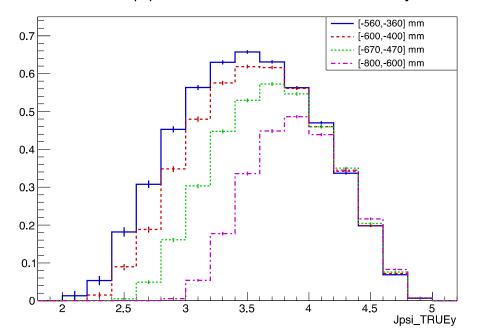


 $\mathcal{L}_{pH}(300 \, K) \approx 2.5 \cdot 10^{32} \, \text{cm}^{-2} \text{s}^{-1}$

 $L_{pH}(Run) \approx 5 fb^{-1}$

Expected performance

 $J/\Psi \rightarrow \mu^{+}\mu^{-}PV X$ track reconstruction efficiency



Channel	Events / week	Total yield
$J/\psi \to \mu^+\mu^-$	1.3×10^{7} !!	1.5×10^{9}
$D^0 o K^-\pi^+$	6.5×10^{7}	7.8×10^{9}
$\psi(2S) \to \mu^+\mu^-$	2.3×10^{5}	2.8×10^{7}
$J/\psi J/\psi \rightarrow \mu^+ \mu^- \mu^+ \mu^-$ (DPS)	8.5	1.0×10^{3}
$J/\psi J/\psi \to \mu^+\mu^-\mu^+\mu^- \text{ (SPS)}$	2.5×10^{1}	3.1×10^{3}
Drell Yan (5 $< M_{\mu\mu} < 9 \text{ GeV}$)	7.4×10^{3}	8.8×10^{5}
$\Upsilon ightarrow \mu^+ \mu^-$	5.6×10^{3}	6.7×10^{5}
$\Lambda_c^+ \to p K^- \pi^+$	1.3×10^{6}	1.5×10^{8}

- > fully reconstructed and selected events!
- based on extrapolation of current SMOG2 performance

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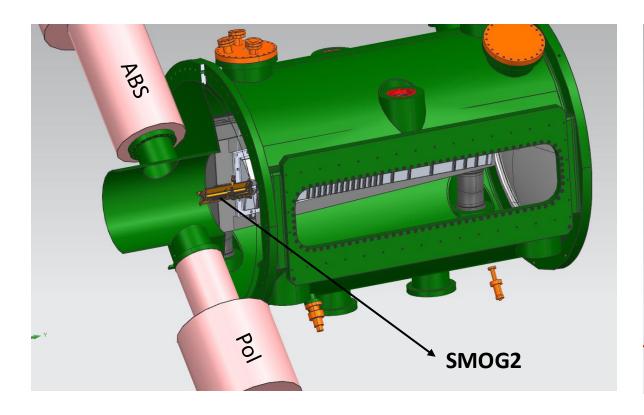
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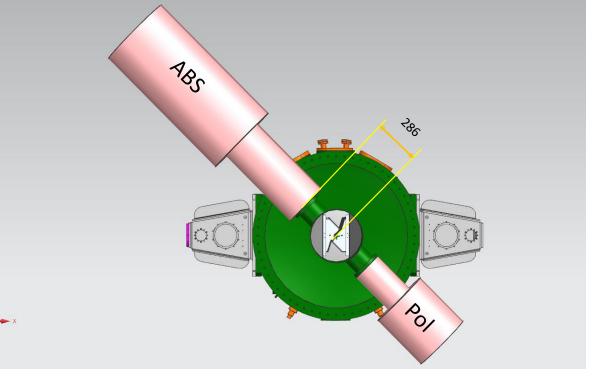
$$L_{pH}(Run) \approx 5 fb^{-1}$$

The jet target option

Alternative solution with **jet target** also under evaluation:

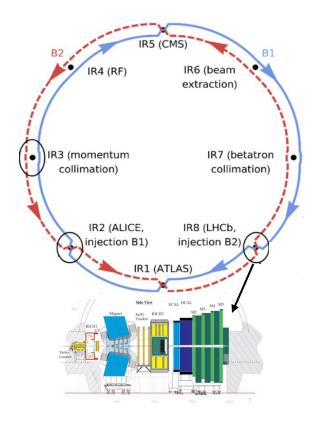
- lower density ($\sim 10^{12} \text{ atoms}/cm^2$) \rightarrow about a factor of 40 smaller
- higher polarization (up to 90%)
- lower systematics in P measurement (virtually close to 0)
- Compatible with SMOG2 setup





Necessary pre-requisites for approval of the project at LHCb (Run5)

- R&D campaign for the apparatus towards the final setup for LHCb
- feasibility studies in a dedicated exp. area served by LHC beams

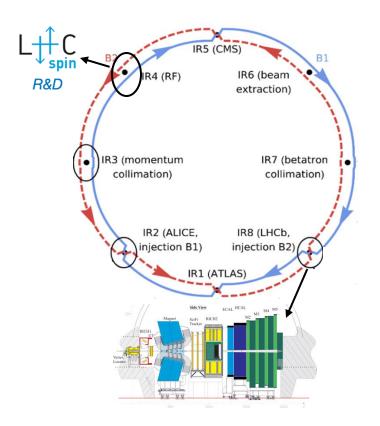


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IR4 is the ideal place for our R&D:

- Lots of free space for our instrumentation
- Rails, cables and racks already available in-situ









ABS and BR-polarimeter are presently at COSY (Julich) and will soon be moved to Ferrara for first tests and optimization prior to full installation at IR4.

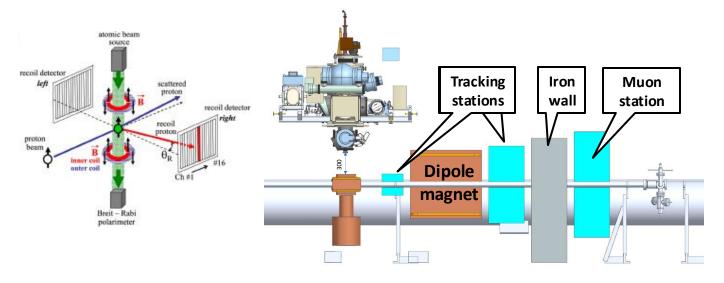


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LS3 (2026-30):

- Installation at IR4 of existing setup (ABS + BR polarimeter)
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- Possibility to implement a minimal spectrometer

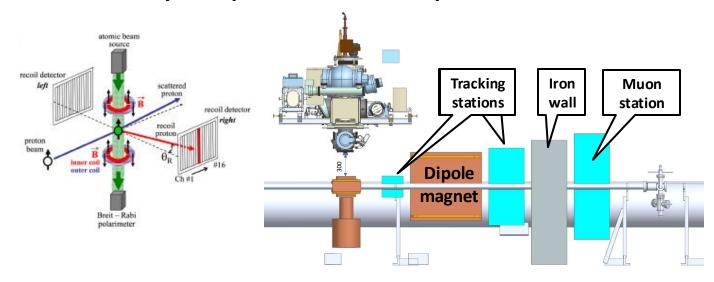


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Run4 (2030-33):

- In-beam polarimetry studies
- proof-of-principle prototype experiment: first pol. meas. at the LHC (SSAs in inclusive light hadron production)

Timeline of the project

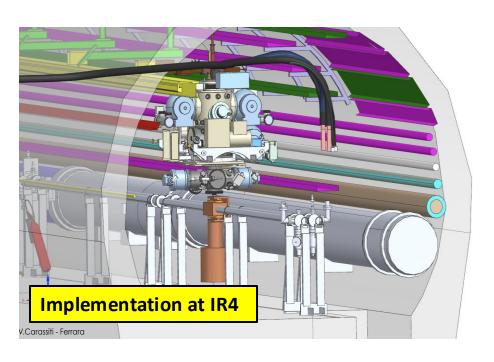


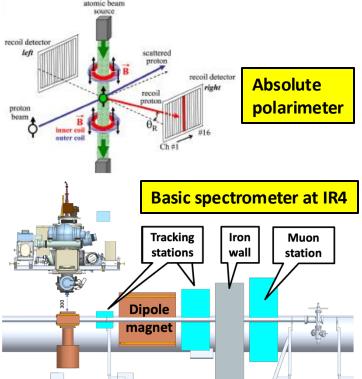
First tests and characterization in Ferrara

Installation at IR4 (LHC):

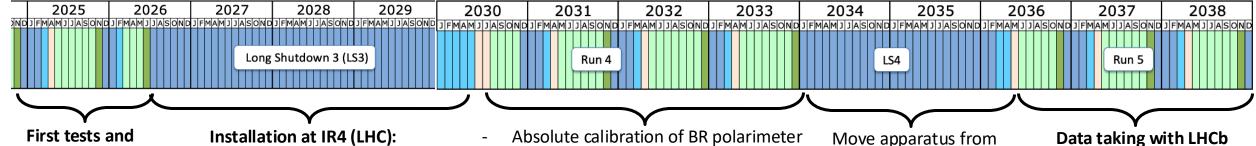
- ABS + Breit-Rabi polarimeter
- Absolute polarimeter
- minimal spectrometer?

- Absolute calibration of BR polarimeter
- Polarimetry measurements
- First physics measurements (SSAs)





Timeline of the project

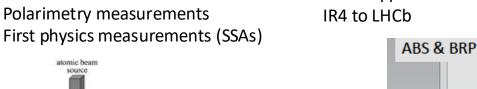


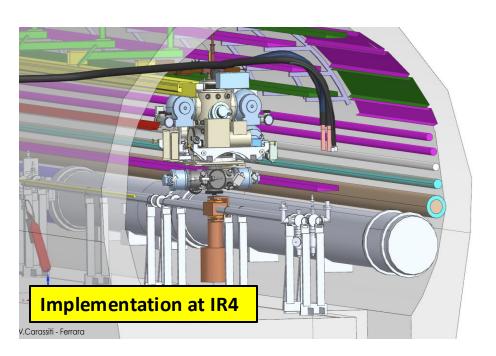
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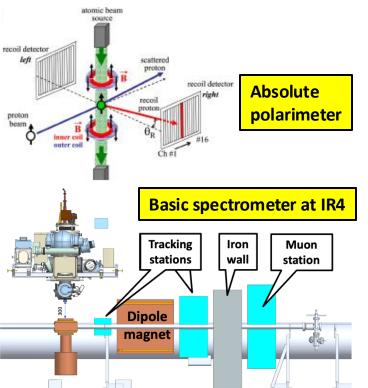
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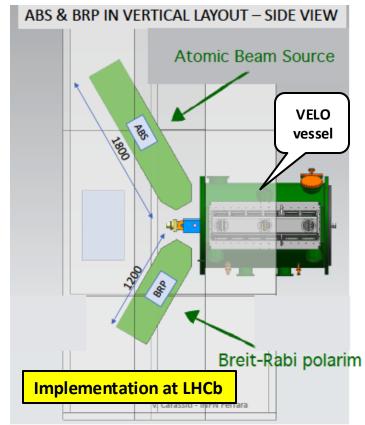
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- Absolute calibration of BR polarimeter









Conclusions

- > The Fixed-Target program at LHCb is active since Run 2, now greatly enriched with SMOG2
- LHCspin is the natural evolution: a polarized fixed target at LHCb will bring spin-physics for the first time at the LHC and will open the way to a broad and unique physics program
- > Novel approaches and reactions will be exploited for studies of the 3D nucleon structure
- First insights into the yet unknown gluon TMDs (such as the GSF) will be possible thanks to the excellent capabilities of LHCb in reconstructing quarkonia states and heavy mesons.
- Cutting-edge unpolarized physics will also be at reach (cold nuclear matter effects, intrinsic charm, QGP studies, etc.)

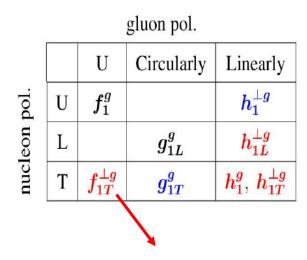
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If approved, LHCspin will make LHCb the first experiment simultaneously running in collider and fixed-target mode with unpolarized and polarized targets, opening a whole new range of explorations at the LHC!

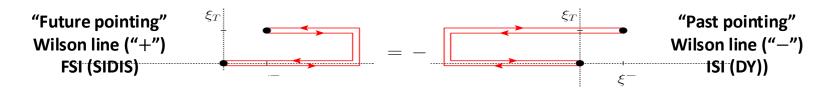
Backup

gluon TMDs



Theory framework well consolidated ...but experimental access still extremely limited!

Gluon correlator depends on 2 path-dependent gauge links, different for ISI and FSI:



$$[+,+] \longleftarrow f_{1T}^{\perp g\left[e\,p^{\uparrow}\to e'\,Q\bar{Q}X\right]}\left(x,p_{T}^{2}\right) = -f_{1T}^{\perp g\left[p^{\uparrow}\,p\to\gamma\,\gamma\,X\right]}\left(x,p_{T}^{2}\right) \longleftarrow [-,-] \qquad \text{Sign-change relation expected for the other T-odd gTMDs } h_{1}^{g} \text{ and } h_{1T}^{\perp g}!$$

- Depending on their combinations, there are 2 independent versions of each gTMD that can be probed in different processes and can have different magnitude and widths and different x and k_T dependencies!
- E.g. there are 2 types of f_1^g and $h_1^{\perp g}$: [++]=[--] Weizsacker-Williams (WW); [+-]=[-+] DiPole (DP)
- 2 indep. GSF: $f_{1T}^{\perp g[+,+]}$ "f-type" \rightarrow antisymm. colour structure; $f_{1T}^{\perp g[+,-]}$ "d-type" \rightarrow symm. colour structure

L. L. Pappalardo

A synergic attack to gTMDs

[D. Boer: Few-body Systems 58, 32 (2017)]

	DIS	DY	SIDIS	$pA \to \gamma \operatorname{jet} X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$		$pp \to J/\psi \gamma X$ $pp \to \Upsilon \gamma X$
$f_1^{g[+,+]}$ (WW)	×	×	×	×	\checkmark	$\sqrt{}$	\checkmark
$f_1^{g[+,-]}$ (DP)	\checkmark	\checkmark	\checkmark		×	×	×

	$pp \to \gamma \gamma X$	$pA \to \gamma^* \text{ jet } X$	$e p \to e' Q \overline{Q} X$ $e p \to e' j_1 j_2 X$		$pp \to J/\psi \gamma X$ $pp \to \Upsilon \gamma X$
$h_1^{\perp g [+,+]} (WW)$	\checkmark	×	\checkmark		\checkmark
$h_1^{\perp g [+,-]} (\mathrm{DP})$	×		×	×	×

	DY	SIDIS	$p^{\uparrow} A \to h X$	$p^{\uparrow}A \to \gamma^{(*)} \text{ jet } X$		$e p^{\uparrow} \rightarrow e' Q \overline{Q} X$ $e p^{\uparrow} \rightarrow e' j_1 j_2 X$
$f_{1T}^{\perp g [+,+]} (WW)$	×	×	×	×	\checkmark	\checkmark
$f_{1T}^{\perp g [+,-]} (DP)$	√	√	\checkmark	\checkmark	×	×

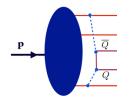
Can be	measured	at the	EIC
Can be	measured	at the	•

Can be measured at RHIC & LHC
(including LHCb+SMOG2/LHCspin

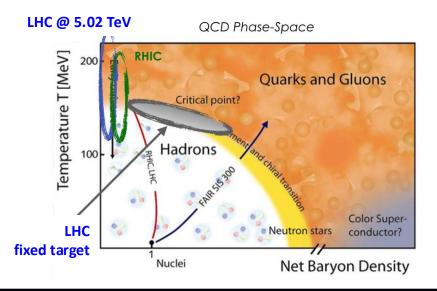
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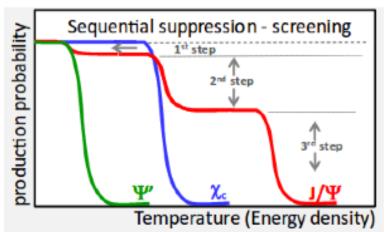
More physics reach with unpolarized FT reactions

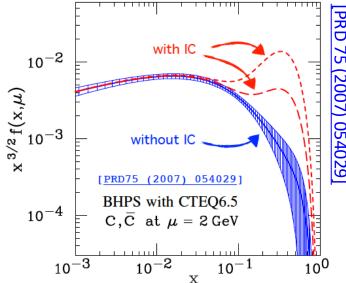
- Intrinsic heavy-quark [S.J. Brodsky et al., Adv. High Energy Phys. 2015 (2015) 231547]
 - 5-quark Fock state of the proton may contribute at high x!
 - charm PDFs at large x could be larger than obtained from conventional fits

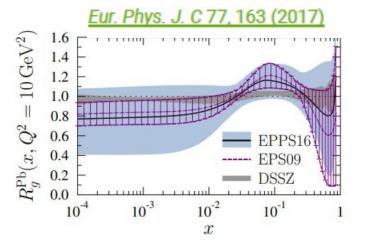


- **pA collisions** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - constraints on nPDFs (e.g. on poorly understood gluon antishadowing at high x)
 - studies of parton energy-loss and absorption phenomena in the cold medium
 - reactions of interest for cosmic-ray physics and DM searches
- **PbA collisions at** $\sqrt{s_{NN}} \approx 72$ **GeV** (using unpolarized gas: He, N, Ne, Ar, Kr, Xe)
 - Study of QGP formation (search for predicted sequential quarkonium suppression)









 $c\overline{c}$ states: J/ψ , χ_c , ψ' ,...

Different binding energies, different dissociation temperatures \rightarrow **medium** thermometer

A preliminary analysis tool for pseudo-data

A pseudo-data set based on a Transversely Pol. H target has been generated to study the interplay between statistical and systematic uncertainties (due to the measurement of the polarization).

Similar approach used at HERMES (Appendix C of [JHEP, 12:010, 2020]):

- Use official LHCb MC data for inclusive production of $J/\psi \to \mu^+\mu^-$ in fixed-target configuration (PYTHIA8 + EPOS)
- Introduce a spin-dependence in the simulation: assign to each simulated event a target polarization state (↑ or ↓) using a random extraction modulated with a model for the cross section
- The model assumes a dominant $\sin \phi$ modulation (e.g. sensitive to the gluon Sivers) plus a suppressed $\sin 2\phi$ modulation (to account e.g. for possible higher-twist contributions). Both terms depend mildly on the kinematics (x, p_T) :

$$\rho = \frac{1}{2} \left[1 + \left(a_1 + a_2 \frac{x - \overline{x}}{x_{max}} + a_3 \frac{p_T - \overline{p_T}}{p_{T\ max}} \right) \sin \phi + \left(b_1 + b_2 \frac{x - \overline{x}}{x_{max}} + b_3 \frac{p_T - \overline{p_T}}{p_{T\ max}} \right) \sin 2\phi \right]$$

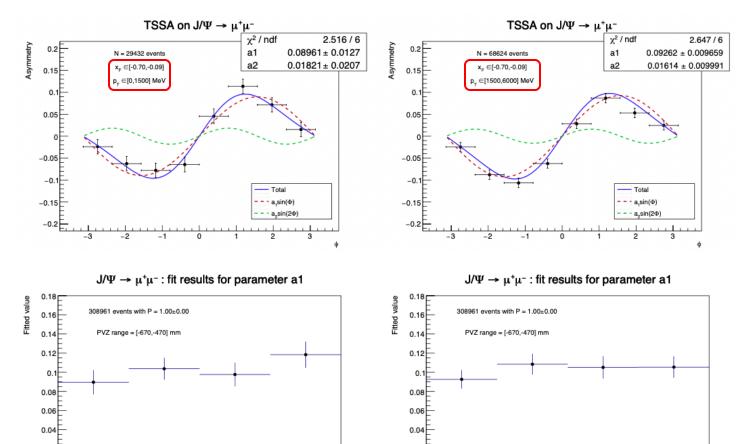
Using these pseudo-data the TSSA is computed in the usual way:

$$A_N = \frac{1}{P} \frac{N^{\uparrow} - N^{\downarrow}}{N^{\uparrow} + N^{\downarrow}}$$

and the uncertainties on $N^{\uparrow(\downarrow)}$ (Poisson) and P (systematic) propagated accordingly.

A preliminary analysis tool for pseudo-data

• The data points are binned in x_F and p_T (2D binning), represented vs. ϕ and fitted with $f = a_1 \sin \phi + a_2 \sin 2\phi$ where the free parameters a_1 and a_2 represent the amplitude of the corresponding azimuthal modulation



- The extracted parameters a_1 and a_2 are consistent with those used to generate the model (no bias is observed)
- With the available MC statistics (corresponding to 2 weeks of data-taking) there is no sensitivity for the $\sin 2\phi$ term
- The amplitudes a_1 are the reported vs. x_F in bins of p_T (and vice-versa)
- A mild kinematic dependence is observed consistent with the model

 $x_F \in [-0.06, -0.04]$

x_F ∈ [-0.09,-0.06]

 $x_F \in [-0.06, -0.04]$

Statistical vs Systematics uncertainties

• The analysis tool described above allows to study the interplay between statistical uncertainties and systematic uncertainties (due to the measurement of the polarization) under different data-taking scenarios

$p_T \; ({ m MeV})$	x_F	$a_1 \ (\Delta P = 0\%)$	$a_1 \ (\Delta P = 5\%)$	$a_1 \ (\Delta P = 20\%)$	$a_1 \ (\Delta P = 50\%)$
[0,1500]	[-0.70, -0.09]	0.090 ± 0.013	0.089 ± 0.013	0.087 ± 0.014	0.087 ± 0.022
[0,1500]	[-0.09, -0.06]	0.104 ± 0.011	0.104 ± 0.012	0.103 ± 0.016	0.100 ± 0.027
[0,1500]	[-0.06, -0.04]	0.098 ± 0.012	0.098 ± 0.013	0.097 ± 0.016	0.094 ± 0.027
[0,1500]	[-0.04, 0.05]	0.118 ± 0.014	0.117 ± 0.014	0.114 ± 0.017	0.113 ± 0.030
$[1500,\!6000]$	[-0.70, -0.09]	0.093 ± 0.010	0.092 ± 0.010	0.090 ± 0.013	0.089 ± 0.023
$[1500,\!6000]$	[-0.09, -0.06]	0.108 ± 0.011	0.108 ± 0.011	0.108 ± 0.015	0.107 ± 0.027
$[1500,\!6000]$	[-0.06, -0.04]	0.105 ± 0.012	0.105 ± 0.012	0.104 ± 0.015	0.103 ± 0.026
$[1500,\!6000]$	[-0.04, 0.05]	0.105 ± 0.011	0.105 ± 0.012	0.102 ± 0.015	0.102 ± 0.026

- A 5% systematic uncertainty on P has no impact on the total uncertainty on a_1
- For $\Delta P = 20\%$ the systematic uncertainty amounts to 30-40% of the statistical uncertainty
- For $\Delta P = 50\%$ the systematic uncertainty approximately equals the statistical uncertainty
- We expect $\Delta P pprox 10-15\%$ for the storage cell hypothesis (and close to 0 for the jet target hypothesis)

Plan:

- Installation of existing setup (ABS + polarimeter from COSY) + minimal detection apparatus during LS3 (2026-28)
- R&D and proof-of-principle experiment during Run4 (2029-2032)
- Three levels of complexity (depending on effective manpower, involved expertise and fundings):

Option 1 (minimal): pol. get targt (no storage cell) for polarimetry measurements with pH^{\uparrow} elastic scattering

Option 2 (intermediate): implementation of an openable storage cell along the beam-pipe (a-la SMOG2) for:

- polarimetry measurements (pH^{\uparrow} elastic scattering)
- study of recombination/depolarization at the cell walls
- study of beam-target interactions

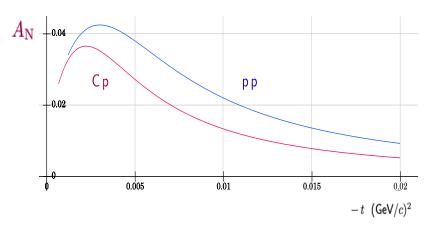
Option 3 (full small-scale experiment): implementation of a minimal spectrometer (in conjunction with gettarget or storage cell) for simple (but unique!) physics measurements (spin asymmetries).

Option 1 (minimal)

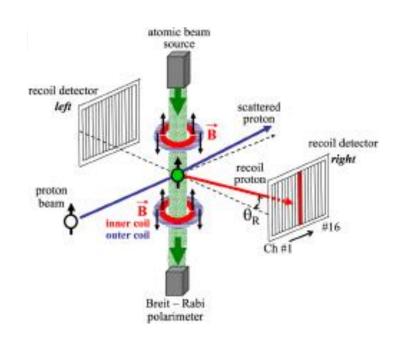
Under **Coulomb-Nuclear Interference** (CNI) conditions one can measure L-R asymmetries in elastic pH^{\uparrow} scattering:

$$A_N(t) = \frac{\mu_p - 1}{m_p} \sqrt{-3t_e} \frac{(t/t_e)^{3/2}}{3(t/t_e)^2 + 1} \qquad t_e = -\frac{8\pi\sqrt{3}\alpha_{EM}|ZZ'|}{\sigma_{tot}(s)}$$

$$t_e = -\frac{8\pi\sqrt{3}\alpha_{EM}|ZZ'|}{\sigma_{tot}(s)}$$



> Technique used for polarimetry at RHIC with p-C scattering (A. Poblaguev et al., PoS PSTP 2017 (2018) 022)



Analyzing power is maximal (4-5%) for $t = t_e$

For a 7 TeV proton beam:

• $\sigma_{tot} \approx 47 \ mb$

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- proton recoil energies: 1.7 4.6 MeV
- proton recoil angles $87^{\circ} < \theta_{lab} < 89^{\circ}$
- Find more here: https://www.maths.tcd.ie/~nhb/talks/2019 07 16 nhb.pdf

Option 1 (minimal)

Goals:

- First absolute polarimetry measurements with elastic beam-target scattering at LHC energies
- test model predictions for Coulomb-Nuclear Interference
- allow for absolute calibration of the Breit-Rabi polarimeter

Tasks and challenges:

- detect small (<5%) Left-Right asymmetry of elastically scattered low-energy protons at large angles (\sim 90°)
- due to very small recoil energies (< 5 MeV), need to place detectors in primary vacuum

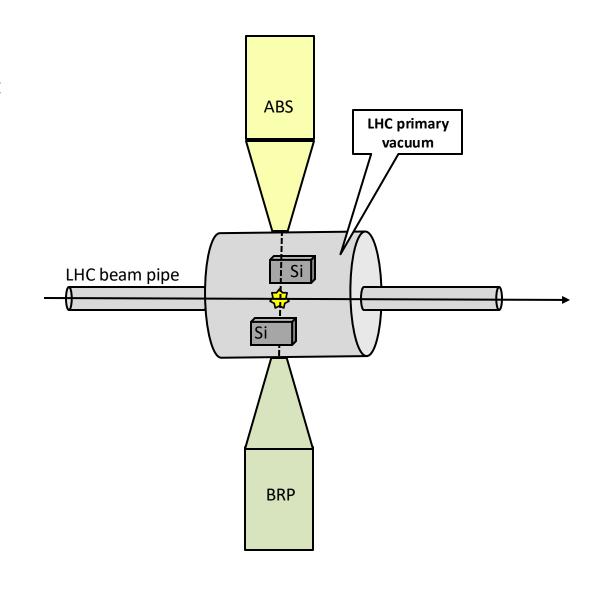
Apparatus:

- target system: jet target (ABS) + Breit-Rabi Polarimeter (BRP)
- detection system: pairs of small-area Si (strip/pixel) detectors in LHC primary vacuum (angular coverage: 30-50 mrad around 90°

Needed expertise (apart from pol. target):

- Si detectors (to be operated in primary vacuum)
- electronics + DAQ

- ...



Option 2 (intermediate)

Goals:

- those of Option 1 (but using a storage cell)
- study of recombination/depolarization at the cell walls (coatings)
- study of beam-target interactions (under realistic conditions for future use at LHCb): beam-induced depolarization, aperture, impedance, etc.

Apparatus:

- same as previous + storage cell
- 300 mT transverse magnet to maintain polarization inside the cell

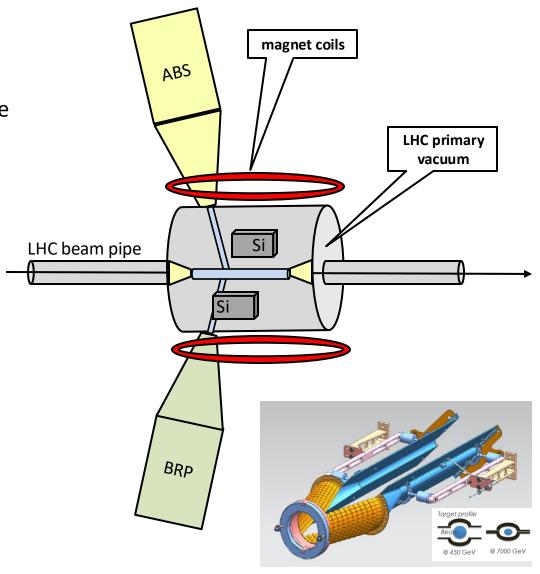
Challenges:

- cell must be openable (with independent motor) to avoid interference with beam during injection/tuning
- cell lateral walls must be very thin to allow passage of low-energy recoiling protons (for polarimetry measurement)

Needed expertise (apart from pol. target):

- Si detectors (to be operated in primary vacuum)
- electronics + DAQ
- transverse magnet
- cryogenics (for SC magnet)





Option 3 (full small-scale experiment)

Goals:

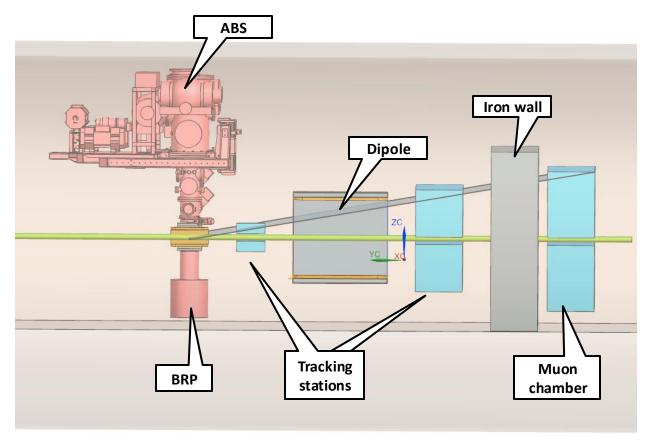
- Those of option 1 (polarimetry measurements)
- measurement of single-spin asymmetries in inclusive hadron (e.g. $J/\psi \to \mu^+\mu^-$) production in $pH^\uparrow, pD^\uparrow, AH^\uparrow, AD^\uparrow$
- proof of principle of the future (large-scale) experiment with LHCb.

Needed expertise (apart from pol. target):

- Si detectors (to be operated in primary vacuum)
- dipole magnet
- tracking detectors (SciFi, Si strip, drift chamber?)
- muon chambers (MWPC?)
- DAQ + electronics
- slow control
- tracking/reconstruction algorithms
- ...

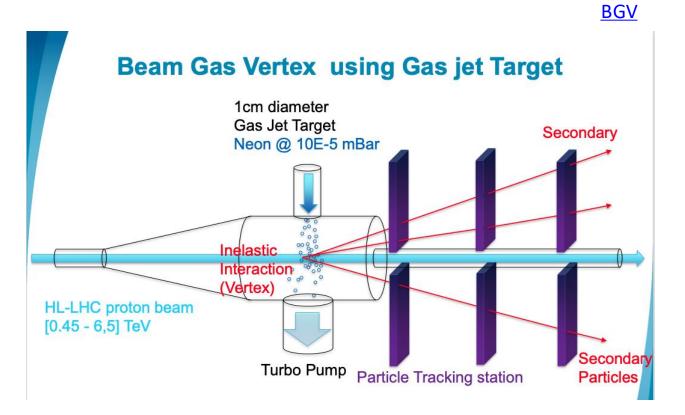
Apparatus:

- jet-target (but could be done also with storage cell)
- full (minimal) spectrometer
 (dipole magnet, tracking stations, muon system)
- PID detectors (Calo, RICH)?



Instrumentation at IR4

- A Beam-Gas Vertex (BGV) apparatus, based on a gas get target and used in the past for beam emittance measurements (currently not in use), could be replaced by our apparatus
- LHCspin could also serve as a new BGV





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